

Intermittent symmetry breaking of the Agmon - Hörmander maximizers

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Fourier Transform of Functions | ... and of Measures

$$\hat{F}(x) = \int_{\mathbb{R}^d} F(y) e^{ix \cdot y} dy.$$

$$\hat{\mu}(x) = \int_{\mathbb{R}^d} e^{ix \cdot y} \mu(dy).$$

Plancherel Identity:
(functions only!)

$$\int_{\mathbb{R}^d} |\hat{F}(x)|^2 \frac{dx}{(2\pi)^d} = \int_{\mathbb{R}^d} |F(y)|^2 dy.$$

σ : uniform measure on S^2 . $\hat{\sigma} = 4\pi \frac{\sin(r)}{r}$.

$$\int_{\mathbb{R}^3} |\hat{\sigma}(x)|^2 \frac{dx}{(2\pi)^3} = \frac{2}{\pi} |S^2| \int_0^\infty \sin^2(r) dr = \infty !!!$$

However, $\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{\sigma}(x)|^2 \frac{dx}{(2\pi)^3} = \frac{2}{\pi} |S^2| \frac{1}{R} \int_0^R \sin^2(r) dr = \frac{1}{\pi} |S^2| + O\left(\frac{1}{R}\right)$. $(R \rightarrow \infty)$.

S. Agmon & L. Hörmander, 1976. There is $C_d(R) > 0$

such that

$$\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}_\sigma(x)|^2 \frac{dx}{(2\pi)^d} \leq C_d(R) \int_{S^{d-1}} |f|^2 d\sigma, \quad \forall f \in L^2(S^{d-1}).$$

Moreover,

$$\lim_{R \rightarrow \infty} \frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}_\sigma(x)|^2 \frac{dx}{(2\pi)^d} = \frac{1}{\pi} \int_{S^{d-1}} |f|^2 d\sigma.$$

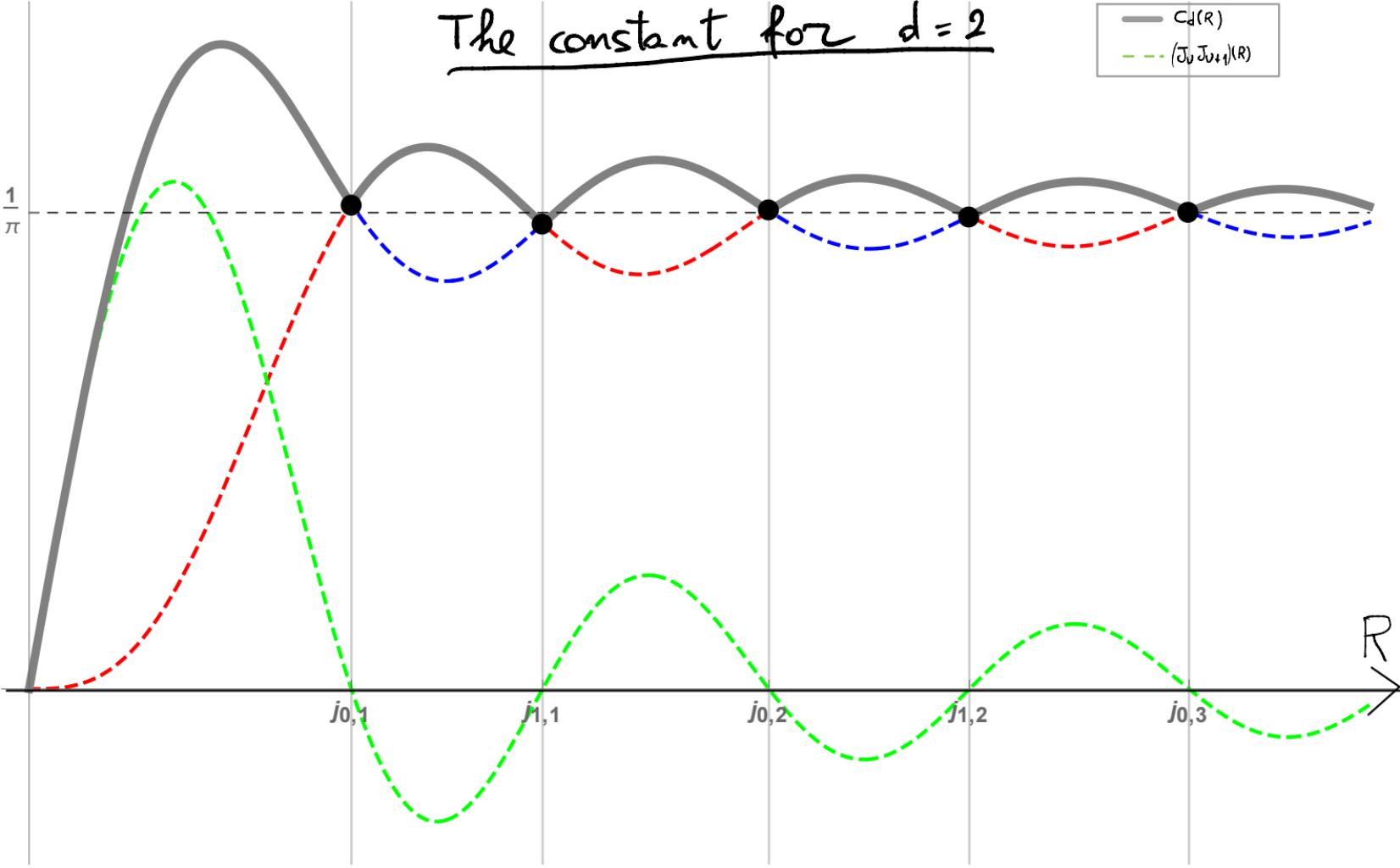
$d \geq 2$

Our result, pt. 1. The best constant is

$$C_d(R) = \begin{cases} \frac{R}{2} J_\nu^2(R) - \frac{R}{2} J_{\nu-1}(R) J_{\nu+1}(R), & \text{if } J_\nu(R) J_{\nu+1}(R) > 0, \\ \frac{R}{2} J_{\nu+1}^2(R) - \frac{R}{2} J_\nu(R) J_{\nu+2}(R), & \text{if } J_\nu(R) J_{\nu+1}(R) \leq 0. \end{cases}$$

$\nu := \frac{d}{2} - 1$, J_α = Bessel function of the first kind.

The constant for $d = 2$



Our result, pt. 2.

$$\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}(x)|^2 \frac{dx}{(2\pi)^d}$$

$$\text{Recall: } C_d(R) = \max_{0 \neq f \in L^2(S^{d-1})} \frac{\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}(x)|^2 \frac{dx}{(2\pi)^d}}{\|f\|_{L^2(S^{d-1})}^2}.$$

The max is attained if and only if:

i) when $(J_v J_{v+1})(R) > 0$, for $f(\omega) = c$.

$$\omega = (\omega_1, \dots, \omega_d) \in S^{d-1}$$

ii) when $(J_v J_{v+1})(R) < 0$, for $f(\omega) = \sum_{j=1}^d \alpha_j \omega_j$.

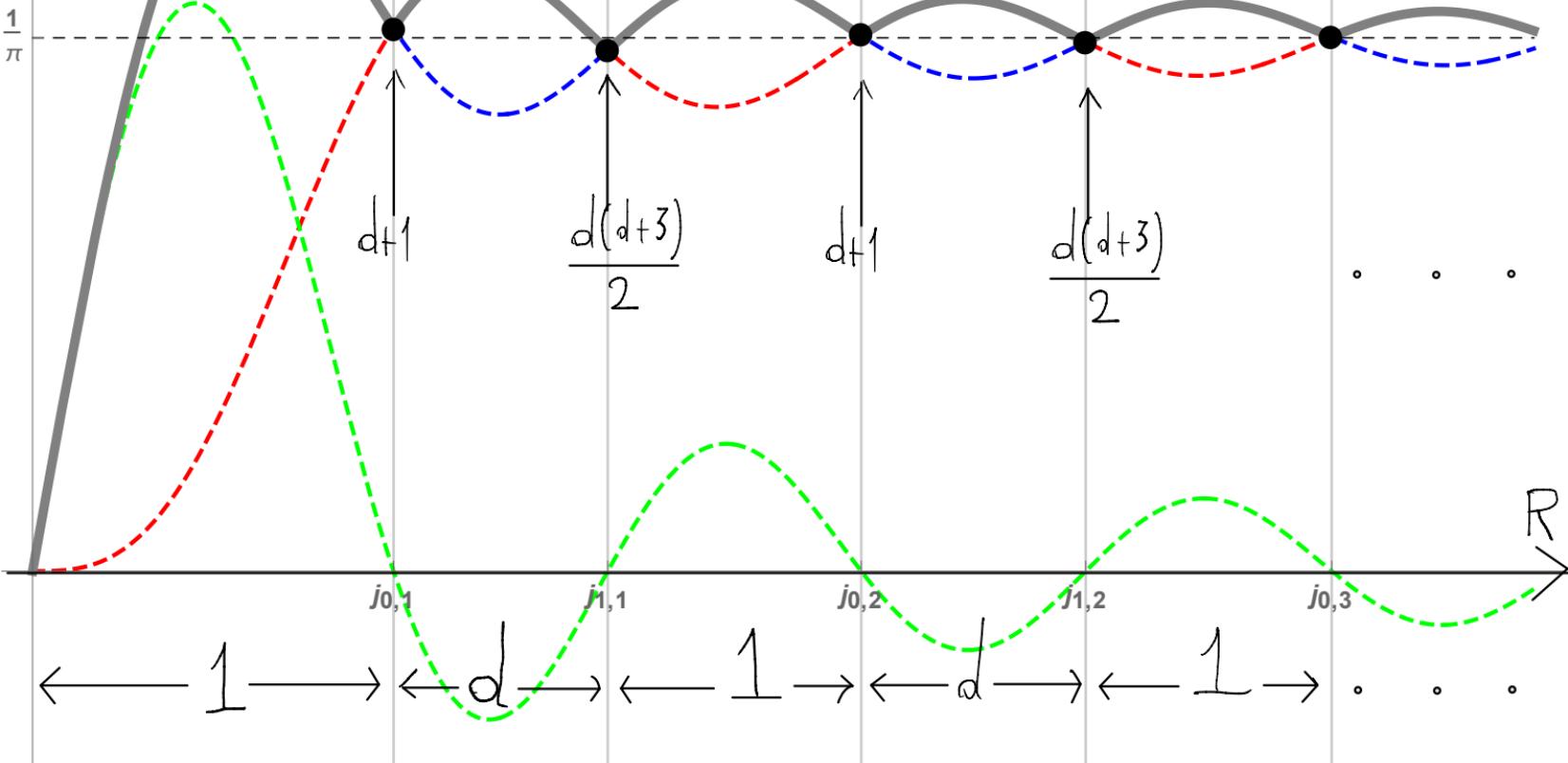
iii-a when $J_v(R) = 0 \neq J_{v+1}(R)$, for $f(\omega) = c + \sum_{j=1}^d \alpha_j \omega_j$.

iii-b when $J_v(R) \neq 0 = J_{v+1}(R)$, for $f(\omega) = H_2(\omega)$.

$c \in \mathbb{C}$, $(\alpha_1, \alpha_2, \dots, \alpha_d) \in \mathbb{C}^d$, $H_2(\omega)$ a HARMONIC POLY, $\deg H_2 \leq 2$.

Dimension of maximizer space
as a function of $R > 0$

— $C_d(R)$
- - - $(J_{\nu} J_{\nu+1})(R)$



Motivation. (or Why should we care?)

Ag.-Hör. is the simplest example of a Fourier restriction estimate, i.e.

$$\|\widehat{f_0}\|_{X(\mathbb{R}^d)} \leq C \|f\|_{L^2(S^{d-1})}.$$

Best constant is known for $X(\mathbb{R}^d) = L^{2\kappa}(\mathbb{R}^d)$, $\kappa \geq 2$, $3 \leq d \leq 7$.

In all these cases, maximizers are constants,

i.e. fully symmetric.

This is our case for some values of R ,

but not for all: symmetry breaking.

Proof

We want

$$C_d(R) = \max_{f \neq 0} \frac{\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}_0(x)|^2 \frac{dx}{(2\pi)^d}}{\|f\|_{L^2(S^{d-1})}^2}.$$

Spherical harmonics : $Y_k(\omega)$ homogeneous harmonic poly
deg. = k .

Decompose $f = \sum_{k=0}^{\infty} Y_k$ so that $\|f\|_{L^2(S^{d-1})}^2 = \sum_{k=0}^{\infty} \|Y_k\|_{L^2(S^{d-1})}^2$.

KEY FACT : $\frac{1}{R} \int_{\{|x| \leq R\}} |\hat{f}_0(x)|^2 \frac{dx}{(2\pi)^d} = \sum_{k=0}^{\infty} C_{k,d}(R) \|Y_k\|_{L^2(S^{d-1})}^2$

Thus,

$$C_d(R) = \max \frac{\sum_{k=0}^{\infty} \Lambda_{k,d}(R) \|Y_k\|^2}{\sum_{k=0}^{\infty} \|Y_k\|^2}$$

$$= \max \left\{ \Lambda_{0,d}(R), \Lambda_{1,d}(R), \Lambda_{2,d}(R), \dots \right\}.$$

The Λ 's are Lommel integrals : $v = \frac{d}{2} - 1$.

$$\Lambda_{k,d}(R) = \frac{1}{R} \int_0^R J_{v+k}^2(r) r dr = \frac{R}{2} J_{v+k}^2(R) - \frac{R}{2} \left(J_{v+k-1} J_{v+k+1} \right)(R).$$

Claim: $\max \{ \Lambda_{0,d}(R), \Lambda_{1,d}(R) \} \geq \Lambda_{k,d}(R)$ for all k .

proof:

$$\begin{aligned} R(\Lambda_{k,d}(R) - \Lambda_{k+2,d}(R)) &= \int_0^R [J_{v+k}^2(r) - J_{v+k+2}^2(r)] r dr \\ &= \int_0^R (J_{v+k}(r) + J_{v+k+2}(r))(J_{v+k}(r) - J_{v+k+2}(r)) r dr \\ (\text{Bessel recursions}) \quad \cong \int_0^R J_{v+k+1}(r) \frac{d}{dr} (J_{v+k+1}(r)) dr &\cong J_{v+k+1}^2(R) \geq 0. \end{aligned}$$

We conclude that

$\Lambda_{0,d}(R) \geq \Lambda_{2,d}(R) \geq \Lambda_{4,d}(R) \geq \dots$	$\Lambda_{1,d}(R) \geq \Lambda_{3,d}(R) \geq \Lambda_{5,d}(R) \geq \dots$
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□

End of proof.

We saw that $C_d(R) = \max \{ \Lambda_{0,d}(R), \Lambda_{1,d}(R) \}$.

Turns out that $\Lambda_{0,d}(R) - \Lambda_{1,d}(R) = (J_v J_{v+1})(R)$.

Conclusion: $C_d(R) = \begin{cases} \Lambda_{0,d}(R) & , (J_v J_{v+1})(R) > 0, \\ \Lambda_{1,d}(R) & , (J_v J_{v+1})(R) < 0, \\ \Lambda_{0,1}(R) = \Lambda_{1,d}(R), & (J_v J_{v+1})(R) = 0. \end{cases}$