

Mathematical Methods for Graduate Students in Physics

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Non-constant coefficient ODE

- Missing a dependent function in the ODE
- it can be reduced to an ODE of one-order less. For example

$$ty'' + 4y' = t^2 \quad (1)$$

- Set $y' = z$ which implies that $y'' = z'$,
- It leads to a linear first order ODE.

$$z' + \frac{4}{t}z = t \quad (2)$$

Non-constant coefficient ODE: Exact Equations

- If the ODE can be written as a derivative of another ODE, we say it is exact

$$a_n(x) \frac{d^n y}{dx^n} + \cdots + a_1(x) \frac{dy}{dx} + a_0(x)y = f(x) \quad (3)$$

- is exact if the LHS can be written as

$$a_n(x) \frac{d^n y}{dx^n} + \cdots + a_0(x)y = \frac{d}{dx} \left[b_{n-1}(x) \frac{d^{n-1} y}{dx^{n-1}} + \cdots + b_0(x)y \right] \quad (4)$$

- This holds if

$$a_0(x) - a_1'(x) + a_2''(x) - \cdots + (-1)^n a_n^{(n)}(x) = 0 \quad (5)$$

Example

- Find the general solution to the ODE

$$(1 - x^2) \frac{d^2 y}{dx^2} - 3x \frac{dy}{dx} - y = 1 \quad (6)$$

- solution**

- Comparing to equation (3), we find $a_2 = 1 - x^2$, $a_1 = -3x$, $a_0 = 1$.
- Check that $a_0 - a_1' + a_2'' = 0$, hence

$$\frac{d}{dx} \left[b_1(x) \frac{dy}{dx} + b_0(x)y \right] = 0 \quad (7)$$

- Matching

$$\frac{d}{dx} \left[b_1(x) \frac{dy}{dx} + b_0(x)y \right] = b_1 \frac{d^2 y}{dx^2} + (b_1' + b_0) \frac{dy}{dx} + b_1' y \quad (8)$$

- We can then compute b_1 and b_0 by comparing with the original equation

$$b_1 = 1 - x^2, \quad b_1' + b_0 = -3x, \quad b_1' = -1 \quad (9)$$

- This implies that $b_1 = 1 - x^2$ and $b_0 = -x$.

$$\frac{d}{dx} \left[(1 - x^2) \frac{dy}{dx} - xy \right] = 1 \quad (10)$$

- Integrating give

$$\frac{dy}{dx} - \left(\frac{x}{1 - x^2} \right) y = \frac{x + c_1}{(1 - x^2)} \quad (11)$$

- Solving further gives

$$y = \frac{c_1 \sin^{-1} x + c_2}{\sqrt{1 - x^2}} - 1 \quad (12)$$

Solution to Laplace equation[Spherical Harmonics]

- In spherical coordinates, the Laplacian is given by

$$\nabla^2 f = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 f}{\partial \phi^2} = 0 \quad (13)$$

- Using the method of separation of variables

$$f(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi) \quad (14)$$

- Inserting this decomposing in equation (13)

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{1}{\Theta} \frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \frac{1}{\Phi} \frac{1}{\sin^2 \theta} \frac{d^2 \Phi}{d\phi^2} = 0 \quad (15)$$

Laplace equation

- The Φ evolves independently

$$+ \frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -\frac{\sin^2 \theta}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) - \frac{\sin^2 \theta}{\Theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) = -m^2 \quad (16)$$

- We can write

$$\frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} = -m^2, \quad (17)$$

- where m^2 is the separation constant, which is chosen to be negative so that the solution for $\Phi(\phi)$ are periodic in ϕ

$$\Phi(\phi) = \begin{cases} e^{im\phi} \\ e^{-im\phi} \end{cases} \quad \text{for} \quad m = 0, 1, 2, \dots \quad (18)$$

Euler equation

- The radial component becomes

$$\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) = \lambda, \quad (19)$$

- Imposing regularity, $\lambda = \ell(\ell + 1)$ Leading to a well-known Euler equation

$$r^2 \frac{d^2 R}{dr^2} + 2r \frac{dR}{dr} - \ell(\ell + 1)R = 0 \quad (20)$$

- This implies that the solution has the following form $R = r^s$. Putting the solution back into equation (20)

$$R(r) = \begin{cases} r^\ell \\ r^{-\ell-1} \end{cases} \quad (21)$$

Associated Legendre equation

- Now consider the angular component

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \left[\ell(\ell + 1) - \frac{m^2}{\sin^2 \theta} \right] \Theta = 0 \quad (22)$$

- Changing variables to $x = \cos \theta$ and $y = \Theta(\theta)$, the above equation becomes

$$(1 - x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + \left[\ell(\ell + 1) - \frac{m^2}{\sin^2 \theta} \right] y = 0 \quad (23)$$

- The solution this equation leads to associated Legendre polynomials [details later]

$$y(x) = P_\ell^m(x), \quad \ell = 0, 1, 2, 3 \text{ and } -\ell \leq m \leq \ell \quad (24)$$

- Combining the Θ and Φ solutions leads to the Spherical Harmonics.

$$Y_\ell^m(\theta, \varphi) = \sqrt{\frac{(2\ell + 1)(\ell - m)!}{4\pi(\ell + m)!}} P_\ell^m(\cos \theta) e^{im\varphi} \quad (25)$$

Legendre equations: Series solution to ODE

- Legendre's equation has the form

$$(1 - x^2)y'' - 2xy' + \ell(\ell + 1)y = 0 \quad (26)$$

where ℓ is a constant. Show that $x = 0$ is an ordinary point and $x = \pm 1$ are regular singular points of this equations

Solution

- Since $x = 0$ is an ordinary point in the Legendre equation, we expect to find two linearly independent solutions of the form

$$y(x) = \sum_{n=0}^{\infty} a_n x^n \quad (27)$$

$$y'(x) = \sum_{n=1}^{\infty} n a_n x^{n-1} \quad (28)$$

$$y''(x) = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2} \quad (29)$$

- This implies

$$2xy' = \sum_{n=1}^{\infty} 2x n a_n x^{n-1} = \sum_{n=0}^{\infty} 2n a_n x^n \quad (30)$$

$$(1-x^2)y'' = \sum_{n=0}^{\infty} [(n+2)(n+1)a_{n+2} - n(n-1)a_n] x^n \quad (31)$$

- Substituting all in equation (26) gives

$$\sum_{n=0}^{\infty} [n(n-1)a_n x^{n-2} - n(n-1)a_n x^n - 2na_n x^n + \ell(\ell+1)a_n x^n] = 0 \quad (32)$$

- Shifting the index of the first term and collecting like terms

$$\sum_{n=0}^{\infty} \{(n+2)(n+1)a_{n+2} - [n(n+1) - \ell(\ell+1)] a_n\} x^n = 0 \quad (33)$$

- The recurrence relation

$$a_{n+2} = \frac{[n(n+1) - \ell(\ell+1)] a_n}{(n+2)(n+1)} \quad (34)$$

Solution

- For the first few integer values $n = 0, 1, 2, \dots$ for the even

$$a_2 = -\frac{\ell(\ell+1)}{2 \cdot 1} a_0$$

$$a_4 = -\frac{(\ell-2)(\ell-3)}{3 \cdot 4} a_2 = (-1)^2 \frac{\ell(\ell-2)(\ell+1)(\ell+3)}{4!} a_0$$

$$a_n = (-1)^n \frac{\ell(\ell-2) \cdots (\ell-2n+2) \cdot (\ell+1)(\ell+3) \cdots (\ell+2+2\ell-1)}{(2n)!}$$

Similarly for the odd number sin terms of a_1

$$a_3 = -\frac{(-1)(\ell+2)}{2 \cdot 3} a_1$$

$$a_5 = -\frac{(\ell-3)(\ell-4)}{4 \cdot 5} a_3 = (-1)^2 \frac{(\ell-1)(\ell-3)(\ell+2)(\ell+4)}{5!} a_1$$

$$a_{2n+1} = (-1)^n \frac{(\ell-1)(\ell-3) \cdots (\ell-2n+1)(\ell+2) \cdots (\ell+2n)}{(2n+1)!} a_1$$

Legendre polynomial

- If we choose $a_0 = 1$ and $a_1 = 0$, we obtain

$$y_1(x) = 1 - \ell(\ell+1)\frac{x^2}{2!} + (\ell-2)\ell(\ell+1)(\ell+3)\frac{x^4}{4!} - \dots$$

$$y_2(x) = x - (\ell-1)(\ell+2)\frac{x^3}{3!} + (\ell-3)(\ell-1)(\ell+2)(\ell+4)\frac{x^5}{5!} - \dots$$

The two solutions are linearly independent

$$y(x) = c_1 y_1(x) + c_2 y_2(x) \quad (35)$$

In general

$$y_1(x) = 1 + \sum_{n=1}^{\infty} (-1)^n \frac{\ell(\ell-2)\cdots(\ell-2n+2) \cdot (\ell+1)(\ell+3)\cdots(\ell+2n-1)}{(2n)!}$$

$$y_2(x) = x + \sum_{n=1}^{\infty} (-1)^n \frac{(\ell-1)(\ell-3)\cdots(\ell-2n+1)(\ell+2)\cdots(\ell+2n)}{(2n+1)!}$$

Legendre polynomial

- These solution constitute the Legendre polynomials

$$P_n(x) = \frac{1}{2^n} \sum_{k=0}^{\lfloor \frac{n}{2} \rfloor} (-1)^k \binom{n}{k} \binom{2n-2k}{n} x^{n-2k}, \quad (36)$$

And for the first few terms in the series

$$P_0 = 1 \quad (37)$$

$$P_1 = x \quad (38)$$

$$P_2 = \frac{1}{2} (3x^2 - 1) \quad (39)$$

$$P_3 = \frac{1}{2} (5x^3 - 3x) \quad (40)$$

$$P_4 = \frac{1}{8} (35x^4 - 30x^2 + 3) \quad (41)$$

$$P_5 = \frac{1}{8} (63x^5 - 70x^3 + 15x) \quad (42)$$

$$P_6 = \frac{1}{16} (231x^6 - 315x^4 + 105x^2 - 5) \quad (43)$$

Find the solution to the following

① Bessel

$$z^2 y'' + zy' + (z^2 - \nu^2)y = 0 \quad (44)$$

② Associated Legendre

$$(1 - z^2)y'' - 2zy' + \left[\ell(\ell + 1) - \frac{m^2}{1 - z^2} \right] y = 0 \quad (45)$$

③ Hypergeometric

$$z(1 - z)y'' + [c - (a + b + 1)z]y' - aby = 0 \quad (46)$$