1.1 Introduction to package qm

The qm package was written by Eric Majzoub, University of Missouri (email: majzoube-at-umsystem.edu), with help from Robert Dodier and Barton Willis.

This purpose of this package is to provide computational tools for solving quantum mechanics problems in a finite-dimensional Hilbert space. It was written with students in mind and is appropriate for upper-level undergraduate quantum mechanics at the level of Townsend's A Modern Introduction to Quantum Mechanics. Please report any errors or unexpected behavior by submitting an issue on the Github page for this project.

The package is loaded with: load(qm);

The qm package provides functions and standard definitions to solve quantum mechanics problems in a finite dimensional Hilbert space. For example, one can calculate the outcome of Stern-Gerlach experiments using built-in definitions of the Sx, Sy, and Sz operators for arbitrary spin, e.g. $s=\{1/2, 1, 3/2, ...\}$. For spin-1/2 the standard basis kets in the x, y, and z-basis are available as $\{xp,xm\}$, $\{yp,ym\}$, and $\{zp,zm\}$, respectively. One can create general ket vectors with arbitrary but finite dimension and perform standard computations such as expectation value, variance, etc. The angular momentum $|j,m\rangle$ representation of kets is also available. Tensor product states for multiparticle systems can be created to perform calculations such as computing the Clebsh-Gordon coefficients.

Let us consider a simple example involving spin-1/2 particles. A bra vector in the **z**-basis may be written as

$$< psi| = a < z+| + b < z-|.$$

The bra <psi| will be represented in Maxima by the row vector [a b], where the basis vectors are

$$$$

and

$$<_{z-|} = [0 1].$$

This bra vector can be created with the mbra command

or taking the quantum mechanical dagger of the corresponding ket. In a Maxima session this looks like the following. The basis kets {zp,zm} are transformed into bras using the dagger function.

1.1.1 Types of kets and bras

There are two types of kets and bras available in the qm package, the first type is given by a matrix representation, as returned by the functions mbra and mket. mkets are column vectors and mbras are row vectors, and their components are entered as Maxima lists in the mbra and mket functions. The second type of bra or ket is abstract; there is no matrix representation. Abstract bras and kets are entered using the bra and ket functions, while also using Maxima lists for the elements. These general kets are displayed in Dirac notation. Abstract bras and kets are used for both the (j,m) representation of states and also for tensor products. For example, a tensor product of two ket vectors |a> and |b> is input as ket([a,b]) and displayed as

```
|a,b> (general ket)
```

Note that abstract kets and bras are assumed to be orthonormal. These general bras and kets may be used to build arbitrarily large tensor product states.

The following examples illustrate some of the basic capabilities of the qm package. Here both abstract, and concrete (matrix representation) kets are shown. The last example shows how to contstruct an entangled Bell pair.

```
(%i1) ket([a,b])+ket([c,d]);
                                   |c, d\rangle + |a, b\rangle
(%o1)
(%i2) mket([a,b]);
                                         [a]
(\%02)
                                         Γ
                                         ГъЪ
(%i3) mbra([a,b]);
(\%03)
                                       [a b]
(%i4) bell:(1/sqrt(2))*(ket([u,d])-ket([d,u]));
                                   |u, d> - |d, u>
(\%04)
                                       sqrt(2)
(%i5) dagger(bell);
                                   <u, d| - <d, u|
(\%05)
                                       sqrt(2)
```

Note that ket([a,b]) is treated as tensor product of states a and b as shown below.

```
(%i1) braket(bra([a1,b1]),ket([a2,b2]));
(%o1) kron_delta(a1, a2) kron_delta(b1, b2)
```

Constants that multiply kets and bras must be declared complex by the user in order for the dagger function to properly conjugate such constants. The example below illustrates this behavior.

The following shows how to declare a ket with both real and complex components in the matrix representation.

1.1.2 Special ket types

Some kets are difficult to work with using either the matrix representation or the general ket representation. These include tensor products of (j,m) kets used in the addition of angular momentum computations. For this reason there are a set of tpkets and associated tpXX functions defined in section (j,m)-kets and bras.

1.1.3 Types of spin operators

When working with kets and bras in the matrix representation, use the spin operators Sxx. When working with abstract kets and bras in the (j,m) representation use the operators Jxx. The family of Sxx operators are represented as matrices in Maxima, while the family of Jxx operators are rule based or function based.

1.2 Functions and Variables for qm

hbar

Planck's constant divided by 2*%pi. hbar is not given a floating point value, but is declared to be a real number greater than zero.

```
\text{ket }([k_1,k_2,\ldots]) [Function]
```

ket creates a general state ket, or tensor product, with symbols k_i representing the states. The state kets k_i are assumed to be orthonormal.

ketp (abstract ket)

[Function]

ketp is a predicate function for abstract kets. It returns true for abstract kets and false for anything else.

```
bra ([b_1, b_2, ...])
                                                                                [Function]
      bra creates a general state bra, or tensor product, with symbols b<sub>i</sub> representing the
      states. The state bras b_i are assumed to be orthonormal.
      (%i1) k:ket([u,d]);
                                                 lu. d>
      (\%01)
      (%i2) b:bra([u,d]);
      (\%02)
                                                 <u, d|
      (\%i3) b . k;
      (\%03)
                                                    1
brap (abstract bra)
                                                                                [Function]
      brap is a predicate function for abstract bras. It returns true for abstract bras and
      false for anything else.
mket ([c_1, c_2, ...])
                                                                                [Function]
      mket creates a column vector of arbitrary finite dimension. The entries c_i can be any
      Maxima expression. The user must declare any relevant constants to be complex.
      For a matrix representation the elements must be entered as a list in [...] square
      brackets.
      (%i1) declare([c1,c2],complex);
      (%o1)
                                                  done
      (%i2) mket([c1,c2]);
                                                 [ c1 ]
      (\%02)
                                                       ]
                                                 [ c2 ]
      (%i3) facts();
      (%o3) [kind(hbar, real), hbar > 0, kind(c1, complex), kind(c2, complex)]
mketp (ket)
                                                                                [Function]
      mketp is a predicate function that checks if its input is an mket, in which case it returns
      true, else it returns false. mketp only returns true for the matrix representation of
      a ket.
      (%i1) k:ket([a,b]);
      (\%01)
                                                 |a, b>
      (%i2) mketp(k);
      (\%02)
                                                  false
      (%i3) k:mket([a,b]);
                                                  [a]
      (\%03)
                                                       ]
                                                  [ b ]
      (%i4) mketp(k);
```

mbra $([c_1,c_2,\ldots])$ [Function] mbra creates a row vector of arbitrary finite dimension. The entries c_i can be any Maxima expression. The user must declare any relevant constants to be complex.

For a matrix representation the elements must be entered as a list in [...] square

true

brackets.

(%04)

```
(%i1) kill(c1,c2);
      (%o1)
                                                done
      (%i2) mbra([c1,c2]);
                                             [ c1 c2 ]
      (\%02)
      (%i3) facts();
      (%o3)
                                   [kind(hbar, real), hbar > 0]
mbrap (bra)
                                                                             [Function]
     mbrap is a predicate function that checks if its input is an mbra, in which case it returns
      true, else it returns false. mbrap only returns true for the matrix representation of
      a bra.
      (%i1) b:mbra([a,b]);
      (\%01)
                                              [a b]
      (%i2) mbrap(b);
      (\%02)
                                                true
   Two additional functions are provided to create kets and bras in the matrix representat-
ion. These functions conveniently attempt to automatically declare constants as complex.
For example, if a list entry is a*sin(x)+b*cos(x) then only a and b will be declare-d
complex and not x.
autoket ([a_1, a_2, \ldots])
                                                                             [Function]
      autoket takes a list [a1, a2, ...] and returns a ket with the coefficients ai declare-
      d complex. Simple expressions such as a*sin(x)+b*cos(x) are allowed and will
      declare only the coefficients as complex.
      (%i1) autoket([a,b]);
                                                [ a ]
                                                [ ]
      (%o1)
                                                [ b ]
      (%i2) facts();
              [kind(hbar, real), hbar > 0, kind(a, complex), kind(b, complex)]
      (%i1) autoket([a*sin(x),b*sin(x)]);
                                            [ a sin(x) ]
      (%o1)
                                            [ b sin(x) ]
      (%i2) facts();
              [kind(hbar, real), hbar > 0, kind(a, complex), kind(b, complex)]
autobra (|a_1, a_2, \ldots|)
                                                                             [Function]
      autobra takes a list [a_1, a_2, \ldots] and returns a bra with the coefficients a_i declare-
      d complex. Simple expressions such as a*sin(x)+b*cos(x) are allowed and will
      declare only the coefficients as complex.
      (%i1) autobra([a,b]);
                                              [a b]
      (%o1)
```

[kind(hbar, real), hbar > 0, kind(a, complex), kind(b, complex)]

(%i2) facts();

(%o2)

```
(%i1) autobra([a*sin(x),b]);
     (%o1)
                                       [a sin(x) b]
     (%i2) facts();
     (%02) [kind(hbar, real), hbar > 0, kind(a, complex), kind(b, complex)]
dagger (vector)
                                                                      [Function]
     dagger is the quantum mechanical dagger function and returns the conjugate
     transpose of its input. Arbitrary constants must be declare-d complex for dagger
     to produce the conjugate.
     (%i1) dagger(mbra([%i,2]));
                                          [ - %i ]
                                          [
     (%o1)
                                                 1
                                          [2]
braket (psi,phi)
                                                                       [Function]
     Given a bra psi and ket phi, braket returns the quantum mechanical bracket
     <psi|phi>.
     (%i1) declare([a,b,c],complex);
     (%o1)
                                            done
     (%i2) braket(mbra([a,b,c]),mket([a,b,c]));
     (\%02)
                                        c + b + a
     (%i3) braket(dagger(mket([a,b,c])),mket([a,b,c]));
                     c conjugate(c) + b conjugate(b) + a conjugate(a)
     (%i4) braket(bra([a1,b1,c1]),ket([a2,b2,c2]));
     (\%04)
                 kron_delta(a1, a2) kron_delta(b1, b2) kron_delta(c1, c2)
norm (psi)
                                                                       [Function]
     Given a ket or bra psi, norm returns the square root of the quantum mechanical
     bracket <psi|psi>. The vector psi must always be a ket, otherwise the function
     will return false.
     (%i1) declare([a,b,c],complex);
     (%o1)
                                            done
     (%i2) norm(mket([a,b,c]));
     (%o2)
                  sqrt(c conjugate(c) + b conjugate(b) + a conjugate(a))
magsqr (c)
                                                                       [Function]
     magsqr returns conjugate(c)*c, the magnitude squared of a complex number.
     (%i1) declare([a,b,c,d],complex);
     (%o1)
                                            done
     (%i2) A:braket(mbra([a,b]),mket([c,d]));
     (%02)
                                          bd+ac
     (%i3) P:magsqr(A);
     (%o3) (b d + a c) (conjugate(b) conjugate(d) + conjugate(a) conjugate(c))■
```

1.2.1 Spin-1/2 state kets and associated operators

Spin-1/2 particles are characterized by a simple 2-dimensional Hilbert space of states. It is spanned by two vectors. In the z-basis these vectors are {zp,zm}, and the basis kets in the z-basis are {xp,xm} and {yp,ym} respectively.

zp	Return the $ z+\rangle$ ket in the z-basis.		[Function]
zm	Return the $ z-\rangle$ ket in the z-basis.		[Function]
хр	Return the $ x+\rangle$ ket in the z-basis.		[Function]
xm	Return the $ x-\rangle$ ket in the z-basis.		[Function]
ур	Return the $ y+\rangle$ ket in the z-basis.		[Function]
ym	Return the $ y\rangle$ ket in the z-basis.		[Function]
	(%i1) zp;		
	(%o1)	[1] [] [0]	
	(%i2) zm;		
	(%o2)	[0] [] [1]	
	(%i1) yp;		
	(%o1)	[1] [] [sqrt(2)] [%i] [] [sqrt(2)]	
	(%i2) ym;	[1]	
	(%o2)	[] [sqrt(2)] [%i] [] [sqrt(2)]	

Switching bases is done in the following example where a z-basis ket is constructed and the x-basis ket is computed.

1.2.2 Pauli matrices and Sz, Sx, Sy operators

[Function]

Returns the Pauli x matrix.

sigmay [Function]

Returns the Pauli y matrix.

sigmaz [Function]

Returns the Pauli z matrix.

[Function]

Returns the spin-1/2 Sx matrix.

Sy Returns the spin-1/2 Sy matrix. [Function]

Sz [Function]

Returns the spin-1/2 Sz matrix.

(%i2) Sy;

commutator (X,Y)

[Function]

Given two operators \mathtt{X} and $\mathtt{Y},$ return the commutator \mathtt{X} . \mathtt{Y} - \mathtt{Y} . $\mathtt{X}.$

(%i1) commutator(Sx,Sy);

(%o1)

]
]
]
]
]
2]
]
-]
]

anticommutator (X,Y)

[Function]

Given two operators X and Y, return the commutator X . Y + Y . X.

[1 0] [0 1]

1.2.3 SX, SY, SZ operators for any spin

SX (s) [Function]

SX(s) for spin s returns the matrix representation of the spin operator Sx. Shortcuts for spin-1/2 are Sx,Sy,Sz, and for spin-1 are Sx1,Sy1,Sz1.

SY (s) [Function]

SY(s) for spin s returns the matrix representation of the spin operator Sy. Shortcuts for spin-1/2 are Sx,Sy,Sz, and for spin-1 are Sx1,Sy1,Sz1.

SZ (s) [Function]

SZ(s) for spin s returns the matrix representation of the spin operator Sz. Shortcuts for spin-1/2 are Sx,Sy,Sz, and for spin-1 are Sx1,Sy1,Sz1.

Example:

(%i1) SY(1/2); %i hbar] 0 2] (%o1)] [%i hbar] Γ 2] (%i2) SX(1); hbar Г 0 0] sqrt(2)] hbar hbar (%o2) 0 sqrt(2) sqrt(2)]]] hbar] 0 0 ٦ sqrt(2)

1.2.4 Expectation value and variance

expect (0,psi) [Function]

Computes the quantum mechanical expectation value of the operator O in state psi, <psi|O|psi>.

qm_variance (0,psi)

[Function]

Computes the quantum mechanical variance of the operator O in state psi, $sqrt(\langle psi|O^2|psi \rangle - \langle psi|O|psi \rangle^2)$.

1.2.5 Angular momentum and ladder operators in the matrix representation

SP (s) [Function]

SP is the raising ladder operator S_{+} for spin s.

SM (s) [Function]

SM is the raising ladder operator S_{-} for spin s.

Examples of the ladder operators:

(%i1) SP(1);				
	[0 s	qrt(2) hba	ır 0]
	[]
(%o1)	[0	0	sqrt(2)	hbar]
	[]
	[0	0	0]
(%i2) SM(1);				
	[0	0	0]
	[]
(%o2)	[sqrt	(2) hbar	0	0]
	[]
	[0	sqrt(2) hba	ar 0]

1.2.6 Rotation operators

RX (s,t) [Function]

 $\mathtt{RX}(\mathtt{s})$ for spin \mathtt{s} returns the matrix representation of the rotation operator \mathtt{Rx} for rotation through angle \mathtt{t} .

RY (s,t) [Function]

RY(s) for spin s returns the matrix representation of the rotation operator Ry for rotation through angle t.

RZ (s,t) [Function]

RZ(s) for spin s returns the matrix representation of the rotation operator Rz for rotation through angle t.

(%i1) RY(1,t);

Proviso: assuming 4*t # 0

Proviso: assumi	ng 4*t # U		
	$[\cos(t) + 1]$	sin(t)	1 - cos(t)
	[]
	[2	sqrt(2)	2]
	[]
	[sin(t)		sin(t)]
(%o1)	[cos(t)]
	[sqrt(2)		sqrt(2)]
	[]
	[1 - cos(t)]	sin(t)	cos(t) + 1
	[]
	[2	sqrt(2)	2]

1.2.7 Time-evolution operator

UU (H,t) [Function]

UU(H,t) is the time evolution operator for Hamiltonian H. It is defined as the matrix exponential matrixexp(-%i*H*t/hbar).

1.3 Angular momentum representation of kets and bras

1.3.1 Matrix representation of (j,m)-kets and bras

The matrix representation of kets and bras in the qm package are represented in the z-basis. To create a matrix representation of of a ket or bra in the (j,m)-basis one uses the $spin_mket$ and $spin_mbra$ functions.

spin_mket $(s, m_s, [1,2])$ [Function] spin_mket returns a ket in the z-basis for spin s and z-projection m_s , for axis 1=X or 2=Y.

spin_mbra $(s, m_s, [1,2])$ [Function] spin_mbra returns a bra in the z-basis for spin s and z-projection m_s , for axis 1=X or 2=Y.

1.3.2 Angular momentum (j,m)-kets and bras

To create kets and bras in the $|j,m\rangle$ representation you use the abstract ket and bra functions with j,m as arguments, as in ket([j,m]) and bra([j,m]).

Some convenience functions for making the kets are the following:

```
(\%i1) jmbot(3/2);
                                               3
      (%o1)
                                                                             [Function]
jmket (j,m)
     jmket creates a (j,m)-ket.
     (%i1) jmket(3/2,1/2);
                                               |-, ->
      (\%01)
jmketp (jmket)
                                                                             [Function]
     jmketp checks to see that the ket has an m-value that is in the set \{-j,-j+1,\ldots,+j\}.
     (%i1) jmketp(ket([j,m]));
      (%o1)
                                                false
      (%i2) jmketp(ket([3/2,1/2]));
      (%o2)
                                                true
jmbrap (jmbra)
                                                                             [Function]
     jmbrap checks to see that the bra has an m-value that is in the set \{-j,-j+1,\ldots,+j\}.
jmcheck (j,m)
                                                                              [Function]
     jmcheck checks to see that m is one of \{-j, \ldots, +j\}.
      (\%i1) jmcheck(3/2,1/2);
      (%o1)
                                                true
Jp (jmket)
                                                                              [Function]
     Jp is the J_+ operator. It takes a jmket jmket(j,m) and returns sqrt(j*(j+1)-
     m*(m+1))*hbar*jmket(j,m+1).
Jm (jmket)
                                                                             [Function]
     Jm is the J_ operator. It takes a jmket jmket(j,m) and returns sqrt(j*(j+1)-m*(m-
     1))*hbar*jmket(j,m-1).
Jsqr (jmket)
                                                                             [Function]
     Jsqr is the J^2 operator. It takes a jmket jmket(j,m) and returns
     (j*(j+1)*hbar^2*jmket(j,m).
Jz (jmket)
                                                                              [Function]
     Jz is the J<sub>z</sub> operator. It takes a jmket jmket(j,m) and returns m*hbar*jmket(j,m).
```

These functions are illustrated below.

1.3.3 Addition of angular momentum in the (j,m)-representation

Addition of angular momentum calculations can be performed in the (j,m)-representation using the function definitions below. The internal representation of kets and bras for this purpose is the following. Given kets $|j1,m1\rangle$ and $|j2,m2\rangle$ a tensor product of (j,m)-kets is instantiated as:

and the corresponding bra is instantiated as:

where the factor of 1 is the multiplicative factor of the tensor product. We call this the *common factor* (cf) of the tensor product. The general form of a tensor product in the (j,m) representation is:

Using the function definitions below one must be careful to avoid errors produced by Maxima's automatic list arithmetic. For example, do not use (J1z+J2z), and instead use the defined function Jtz. Similarly for any of the operators that are added together, one should always use the total Jtxx defined function.

tpket (jmket1,jmket2)

[Function]

tpket instantiates a tensor product of two (j,m)-kets.

[Function]

tpbra instantiates a tensor product of two (j,m)-bras.

tpbraket (tpbra,tpket)

[Function]

tpbraket returns the bracket of a tpbra and a tpket.

```
(%i1) k:tpket(jmtop(1),jmbot(1));
     (%o1)
                                [tpket, 1, |1, 1>, |1, - 1>]
     (%i2) K:Jtsqr(k);
     (%o2) [tpket, 2 hbar , |1, 1>, |1, - 1>] + [tpket, 2 hbar , |1, 0>, |1, 0>]■
     (%i3) B:tpdagger(k);
     (\%03)
                               [tpbra, 1, <1, 1|, <1, -1|]
     (%i4) tpbraket(B,K);
     (\%04)
                                           2 hbar
tpcfset (cf,tpket)
                                                                       [Function]
     tpcfset manually sets the common factor cf of a tpket.
tpscmult (a, tpket)
                                                                       [Function]
     tpscmult multiplies the tensor product's common factor by a.
     (%i1) k1:tpket(ket([1/2,1/2]),ket([1/2,-1/2]));
                                [tpket, 1, |-, ->, |-, - ->]
     (\%01)
     (%i2) tpscmult(c,k1);
     (%02)
                                [tpket, c, |-, ->, |-, - ->]
tpadd (tpket,tpket)
                                                                       [Function]
     tpadd adds two tpkets. This function is necessary to avoid trouble with Maxima's
     automatic list arithmetic.
     (%i1) k1:tpket(ket([1/2,1/2]),ket([1/2,-1/2]));
                                            1 1 1
                               [tpket, 1, |-, ->, |-, -->]
     (%o1)
     (%i2) k2:tpket(ket([1/2,-1/2]),ket([1/2,1/2]));
                               [tpket, 1, |-, -->, |-, ->]
     (\%02)
                                                2
     (%i3) tpadd(k1,k2);
                                   1 1
               [tpket, 1, |-, ->, |-, -->] + [tpket, 1, |-, -->, |-, ->]
```

tpdagger (tpket or tpbra)

(%o3)

[Function]

tpdagger takes the quantum mechanical dagger of a tpket or tpbra.

2 2 2

(%i1)
$$k1:tpket(ket([1/2,1/2]),ket([1/2,-1/2]));$$

(%i2) tpdagger(k1);

J1z (tpket) [Function]

J1z returns the tensor product of a tpket with Jz acting on the first ket.

J2z (tpket) [Function]

J2z returns the tensor product of a tpket with Jz acting on the second ket.

$$(\%i1)$$
 k:tpket(ket([3/2,3/2]),ket([1/2,1/2]));

(%i2) J1z(k);

(%i3) J2z(k);

Jtz (tpket) [Function]

Jtz is the total z-projection of spin operator acting on a tpket and returning $(J_{1z}+J_{2z})$.

(%i1)
$$k:tpket(ket([3/2,3/2]),ket([1/2,1/2]));$$

(%i2) Jtz(k);

J1sqr (tpket) [Function]

J1sqr returns Jsqr for the first ket of a tpket.

J2sqr (tpket) [Function]

J2sqr returns Jsqr for the second ket of a tpket.

J1p (tpket) [Function]

J1p returns J_+ for the first ket of a tpket.

[Function]

J2p (tpket)[Function] J2p returns J_+ for the second ket of a tpket. Jtp (tpket) [Function] Jtp returns $(J_{1+}+J_{2+})$ for the tpket. J1m (tpket) [Function] J1m returns J₋ for the first ket of a tpket. J2m (tpket)[Function] J2m returns $J_{=}$ for the second ket of a tpket. Jtm (tpket) [Function] Jtm returns $(J_{1-}+J_{2-})$ for the tpket. J1p2m (tpket)[Function] J1p2m returns $(J_{1+}J_{2-})$ for the tpket. (%i1) k:tpket(ket([3/2,1/2]),ket([1/2,1/2]));3 1 1 1 [tpket, 1, |-, ->, |-, ->] (%o1) (%i2) b:tpdagger(k); 3 1 [tpbra, 1, <-, -|, <-, -|] (%02)2 2 2 2 (%i3) J1p2m(k);3 3 1 [tpket, sqrt(3) hbar , |-, ->, |-, -->] (%03)(%i4) J1m2p(k);(%04)0 J1m2p (tpket)

[Function] J1zJ2z (tpket) J1zJ2z returns $(J_{1z}J_{2z})$ for the tpket.

J1m2p returns $(J_{1}-J_{2+})$ for the tpket.

Jtsqr (tpket) [Function] Jtsqr returns $(J_1^2+J_2^2+J_{1+}J_{2-}+J_{1-}J_{2+}+J_{1z}J_{2z})$ for the tpket.

(%i1)
$$k:tpket(ket([3/2,-1/2]),ket([1/2,1/2]));$$

(%i2) B:tpdagger(k);

(%i3) K2:Jtsqr(k);

(%i4) tpbraket(B,K2);

get_j is a convenience function that computes j from j(j+1)=q where q is a rational number. This function is useful after using the function Jtsqr.

(%o1)
$$j = -2$$

1.3.3.1 Example computations

For the first example, let us see how to determine the total spin state $|j,m\rangle$ of the two-particle state $|1/2,1/2;1,1\rangle$.

(%i2) Jtsqr(k);

(%i3) get_j(15/4);

This is an eigenket of Jtsqr, thus $|3/2,3/2\rangle = |1/2,1/2;1,1\rangle$, and it is also the top state. One can now apply the lowering operator to find the other states: $|3/2,1/2\rangle$, $|3/2,-1/2\rangle$, and $|3/2,-3/2\rangle$.

Let us see how to compute the matrix elements of the operator (J1z-J1z) in the z-basis for two spin-1/2 particles. Note that we use the tpadd and tpscmult functions to add the two operators. First, we form the four basis kets $\{phi_1,phi_2,phi_3,phi_4\}$ of the form $|j_1,m_1;j_2,m_2\rangle$. The next four entries are for the operator acting on the basis kets. We skip taking the braket below; the common factor is the resulting matrix element.

```
(%i1) phi1:tpket(ket([1/2,1/2]),ket([1/2,1/2]));
(%o1)
                          [tpket, 1, |-, ->, |-, ->]
                                               2 2
(%i2) phi2:tpket(ket([1/2,1/2]),ket([1/2,-1/2]));
                                     1 1
(%o2)
                         [tpket, 1, |-, ->, |-, - ->]
                                     2 2
                                              2
(%i3) phi3:tpket(ket([1/2,-1/2]),ket([1/2,1/2]));
(\%03)
                         [tpket, 1, |-, -->, |-, ->]
                                          2
                                                2 2
(\%i4) phi4:tpket(ket([1/2,-1/2]),ket([1/2,-1/2]));
                                          1
                                               1
(\%04)
                        [tpket, 1, |-, - ->, |-, - ->]
                                    2
                                          2 2
(%i5) tpadd(J1z(phi1),tpscmult(-1,J2z(phi1)));
(\%05)
(%i6) tpadd(J1z(phi2),tpscmult(-1,J2z(phi2)));
                        [tpket, hbar, |-, ->, |-, - ->]
(\%06)
(%i7) tpadd(J1z(phi3),tpscmult(-1,J2z(phi3)));
                       [tpket, - hbar, |-, - ->, |-, ->]
(\%07)
(%i8) tpadd(J1z(phi4),tpscmult(-1,J2z(phi4)));
(%08)
```

In the example below we calculate the Clebsh-Gordon coefficients of the two-particle state with two spin-1/2 particles. We begin by defining the top rung of the ladder and stepping down.

```
(%i1) top:tpket(jmtop(1/2),jmtop(1/2));
                                       1
(%o1)
                         [tpket, 1, |-, ->, |-, ->]
                                    2 2
(%i2) Jtsqr(top);
                                   2
                                      1 1
(\%02)
                      [tpket, 2 hbar , |-, ->, |-, ->]
                                       2 2
                                               2 2
(%i3) get_j(2);
(\%03)
                                   j = 1
(%i4) Jtz(top);
                                      1 1
(\%04)
                        [tpket, hbar, |-, ->, |-, ->]
                                      2 2
(%i5) JMtop:ket([1,1]);
(\%05)
                                  |1, 1>
(%i6) mid:Jtm(top);
                     1 1
                           1
                                 1
      [tpket, hbar, |-, ->, |-, -->] + [tpket, hbar, |-, -->, |-, ->]
                     2 2
                            2
                                 2
(%i7) Jm(JMtop);
(\%07)
                            sqrt(2) |1, 0> hbar
(%i8) mid:tpscmult(1/(sqrt(2)*hbar),mid);
                      1 1
                            1 1
                                                    1
                1
(%08) [tpket, -----, |-, ->, |-, -->] + [tpket, -----, |-, -->, |-, ->]
                     2 2 2 2
             sqrt(2)
                                                 sqrt(2)
(%i9) bot:Jtm(mid);
(\%09)
                  [tpket, sqrt(2) hbar, |-, - ->, |-, - ->]
(%i10) Jm(ket([1,0]));
(%o10)
                            sqrt(2) |1, - 1> hbar
(%i11) bot:tpscmult(1/(sqrt(2)*hbar),bot);
                       [tpket, 1, |-, - ->, |-, - ->]
(%o11)
                                  2 2 2
```

1.4 General tensor products

Tensor products are represented as lists in the qm package. The ket tensor product |z+,z+> can be represented as ket([u,d]), for example, and the bra tensor product <a,b| is represented as bra([a,b]) for states a and b. For a tensor product where the identity is one of the elements of the product, substitute the string Id in the ket or bra at the desired location. See the examples below for the use of the identity in tensor products.

Examples below show how to create abstract tensor products that contain the identity element Id and how to take the bracket of these tensor products.

In the next example we construct the state function for an entangled Bell pair, construct the density matrix, and then trace over the first particle to obtain the density submatrix for particle 2.

```
(%i1) bell:(1/sqrt(2))*(ket([u,d])-ket([d,u]));
                            |u, d> - |d, u>
(%o1)
                               sqrt(2)
(%i2) rho:bell . dagger(bell);
     |u, d> . <u, d| - |u, d> . <d, u| - |d, u> . <u, d| + |d, u> . <d, u|
(%o2) ------
                                  2
(%i3) assume(not equal(u,d));
(\%03)
                           [notequal(u, d)]
(\%i4) trace1:bra([u,Id]) . rho . ket([u,Id])+bra([d,Id]) . rho . ket([d,Id]);
                    |-, u> . <-, u| |-, d> . <-, d|
(\%04)
                                          2
                          2
```

1.5 Quantum harmonic oscillator

The qm package can perform simple quantum harmonic oscillator calculations involving the ladder operators a⁺ and a⁻. These are referred to in the package as ap and am respectively. For computations with arbitrary states to work you must declare the harmonic oscillator state, say n, to be both scalar and integer, as shown in the examples below.

```
ap ap is the raising operator a<sup>+</sup> for quantum harmonic oscillator states.

[Function]

am

a is the lowering operator a<sup>-</sup> for quantum harmonic oscillator states.
```

A common problem is to compute the 1st order change in energy of a state due to a perturbation of the harmonic potential, say an additional factor $V(x) = x^2 + g*x^4$ for small g. This example is performed below, ignoring any physical constants in the problem.

Another package that handles quantum mechanical operators is operator_algebra written by Barton Willis.

Appendix A Function and Variable index

\mathbf{A}	Jtsqr
am	Jtz 17 Jz 14
ap 23 autobra 5 autoket 5	K
В	ket 3 ketp 3
bra 4 braket 6 brap 4	M
C	magsqr. 6 mbra 4 mbrap 5
commutator	$\begin{array}{llllllllllllllllllllllllllllllllllll$
D	
dagger6	N
\mathbf{E}	norm
expect	Q
\mathbf{G}	qm_variance
get_j	
	\mathbf{R}
J	RX
J1m	RY
J1m2p	RZ
J1p	
J1p2m	C
J1sqr	\mathbf{S}
J1z	sigmax 8
J1zJ2z	sigmay 8
J2p	$\mathtt{sigmaz}8$
J2sqr	SM
J2z	spin_mbra
Jm	spin_mket
jmbot	SP
jmbrap	Sx
jmcheck	Sy 8
jmket	SY
jmketp	Sz
jmtop	SZ
Jsqr	
Jtm	

\mathbf{T}	\mathbf{X}
tpadd 16 tpbra 15 tpbraket 15	xm 7 xp 7
tpcfset 16 tpdagger 16 tpket 15 tpscmult 16	Y ym
U	Z zm
hhar 3	