

3.19.4 A Tutorial Type Example of a MR Calculation

Perhaps, the most important use of the MR-CI module is for the calculation of transition energies and optical spectra. Let us first calculate the first excited singlet and triplet state of the formaldehyde molecule using the MR-CI method together with the Davidson correction to approximately account for the effect of unlinked quadruple substitutions. We deliberately choose a somewhat small basis set for this calculation which is already reasonable since we only look at a valence excited state and want to demonstrate the principle.

Suppose that we already know from a ground state calculation that the HOMO of H_2CO is an oxygen lone pair orbitals and the LUMO the π^* MO. Thus, we want to calculate the singlet and triplet $n \rightarrow \pi^*$ transitions and nothing else. Consequently, we only need to correlate two electrons in two orbitals suggesting a CAS(2,2) reference space.

```
# A simple MRCI example
! def2-SVP def2-SVP/C UseSym

%method frozencore fc_ewin
    end

%mrcki ewin -3,1000
    CIType MRCI
    EUnselOpt FullMP2
    DavidsonOpt Davidson1
    UseIVOs true
    tsel 1e-6
    tpre 1e-2
    MaxMemInt 256
    MaxMemVec 32
    IntMode FullTrafo
    AllSingles true
    Solver Diag
    # ground state 1A1
    NewBlock 1 0
        NRoots 1
        Excitations cisd
        Refs CAS(2,2) end
    End
    # HOMO LUMO transition 1A2
    NewBlock 1 1
        NRoots 1
        Excitations cisd
        Refs CAS(2,2) end
    End
    # HOMO LUMO triplet transition 3A2
    NewBlock 3 1
        NRoots 1
        Excitations cisd
        Refs CAS(2,2) end
    end
end

* int 0 1
C 0 0 0 0.000000 0.000 0.000
O 1 0 0 1.200371 0.000 0.000
H 1 2 0 1.107372 121.941 0.000
H 1 2 3 1.107372 121.941 180.000
*
```

This input – which is much more than what is really required - needs some explanations: First of all, we choose a standard RHF calculation with the SVP basis set and we assign the SV/C fitting basis although it is not used in the SCF procedure at all. In the %mrcki block all details of the MR-CI procedure are specified. First, EWin(%method frozencore fc_ewin) selects the MOs within the given orbital energy range to be included in the correlation treatment. The CIType variable selects the type of multireference treatment. Numerous choices are possible and

MRCI is just the one selected for this application.

Note

The CItYpe statement selects several default values for other variables. So it is a very good idea to place this statement at the beginning of the MR-CI block and possibly overwrite the program selected defaults later. If you place the CItYpe statement after one of the values which it selects by default your input will simply be overwritten!

The variables EUnselOpt and DavidsonOpt control the corrections to the MR-CI energies. EUnselOpt specifies the way in which the MR-CI energies are extrapolated to zero threshold T_{Sel} . Here we choose a full MR-MP2 calculation of the missing contributions to be done *after* the variational step, i.e. using the relaxed part of the reference wavefunction as a 0th order state for MR-PT. The DavidsonOpt controls the type of estimate made for the effect of higher substitutions. Again, multiple choices are possible but the most commonly used one (despite some real shortcomings) is certainly the choice Davidson1. The flag UseIVOs instructs the program to use “improved virtual orbitals”. These are virtual orbitals obtained from a diagonalization of the Fock operator from which one electron has been removed in an averaged way from the valence orbitals. Thus, these orbitals “see” only a $N - 1$ electron potential (as required) and are not as diffuse as the standard virtual orbitals from Hartree-Fock calculations. If you input DFT orbitals in the MR-CI module (which is perfectly admissible and also recommended in some cases, for example for transition metal complexes) then it is recommended to turn that flag off since the DFT orbitals are already o.k. in this respect. The two thresholds T_{Sel} and T_{Pre} are already explained above and represent the selection criteria for the first order interacting space and the reference space respectively. T_{Sel} is given in units of Eh and refers to the second order MR-MP2 energy contribution from a given excited CSF. 10^{-6} Eh is a pretty good value. Reliable results for transition energies start with $\approx 10^{-5}$; however, the total energy is converging pretty slowly with this parameter and this is one of the greatest drawbacks of individually selecting CI procedures! (see below). T_{Pre} is dimensionless and refers to the *weight* of a given initial reference after diagonalization of the given initial reference space (10^{-4} is a pretty good value and there is little need to go much lower. Aggressive values such as 10^{-2} only select the truly leading configurations for a given target state which can be time saving. Intermediate values are not really recommended). The parameters MaxMemInt and MaxMemVec tell the program how much memory (in MB) it is allowed to allocate for integrals and for trial and sigma-vectors respectively.

The flag IntMode tells the program to perform a full integral transformation. This is possible for small cases with less than, say, 100–200 MOs. In this case that it is possible it speeds up the calculations considerably. For larger molecules you *have to* set this flag to RITrafo which means that integrals are recomputed on the fly using the RI approximation which is more expensive but the only way to do the calculation. To switch between the possible modes use:

```
%mrci  IntMode FullTrafo # exact 4 index transformation
          RITrafo      # use auxiliary basis sets
```

For small molecules or if high accuracy in the total energies is required it is much better to use the exact four index transformation. The limitations are that you will run out of disk space or main memory with more than ca. 200–300 MOs.

The variable Solver can be diag (for Davidson type diagonalization) or DIIS for multirrot DIIS type treatments.

```
%mrci  Solver Diag      # Davidson solver
          DIIS          # Multirrot DIIS
```

For CI methods, the diag solver is usually preferable. For methods like ACPF that contain nonlinear terms, DIIS is imperative.

Next in the input comes the definition of what CI matrices are to be constructed and diagonalized. Each multiplicity defines a *block* of the CI matrix which is separately specified. Here we ask for two blocks – singlet and triplet. The general syntax is:

```
NewBlock Multiplicity Irrep
          NRroots 1      # Number of roots to determine
          Excitations cisd # Type of excitations
```

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```

Refs CAS (NEl,NOrb) end # Reference space def.
end # Finalize the block

```

Now that all input is understood let us look at the outcome of this calculation:

The first thing that happens after the SCF calculation is the preparation of the frozen core Fock matrix and the improved virtual orbitals by the program `orca_ciprep`. From the output the energies of the IVOs can be seen. In this case the LUMO comes down to -8.2 eV which is much more reasonable than the SCF value of $+3.4$ eV. Concomitantly, the shape of this MO will be much more realistic and this important since this orbital is in the reference space!

```

-----
ORCA CI-PREPARATION
-----

Reading the GBW file          ... done
Symmetry usage                ... ON

One-Electron Matrix          ... test1.H.tmp
GBW-File                     ... test1.gbw
Improved virtual orbitals     ... test1.ivo
First MO in the CI            ... 2
Internal Fock matrix          ... test1.fi.tmp
LastInternal Orbital          ... 6
Integral package used         ... LIBINT
Reading the GBW file          ... done
Symmetry usage                ... ON

Reading the one-electron matrix ... done
Forming inactive density       ... done
Forming averaged valence density ...
Scaling the occupied orbital occupation numbers
First MO                      ... 2
Last MO                       ... 7
Number of electrons in the range ... 12
Scaling factor                 ... 0.917

done
Forming internal density       ... done
Forming Fock matrix/matrices   ...
Nuclear repulsion             ... 31.371502
Core repulsion                 ... 31.371502
One-electron energy           ... -114.942080
Fock-energy                     ... -94.993431
Final value                    ... -73.596254
done
Modifying virtual orbitals     ...
Last occupied MO               ... 7
Total number of MOs            ... 38
Number of virtual MOs          ... 30
Doing diagonalization with symmetry
The improved virtual eigenvalues:
  0: -0.2955 au  -8.041 eV  2- B2
  1: -0.0701 au  -1.907 eV  6- A1
  2: -0.0176 au  -0.479 eV  3- B1
  3:  0.0064 au   0.175 eV  7- A1
  4:  0.2922 au   7.951 eV  8- A1
  5:  0.2948 au   8.021 eV  3- B2
  6:  0.3836 au  10.439 eV  4- B1
  7:  0.4333 au  11.790 eV  9- A1
  8:  0.4825 au  13.128 eV  5- B1

```

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```

  9:  0.5027 au   13.680 eV   10- A1
 10:  0.7218 au   19.642 eV   11- A1
 11:  0.8351 au   22.723 eV    4- B2
 12:  0.9371 au   25.501 eV    6- B1
 13:  1.0265 au   27.933 eV    1- A2
 14:  1.1141 au   30.317 eV   12- A1
 15:  1.2869 au   35.017 eV    5- B2
 16:  1.4605 au   39.743 eV    7- B1
...
done
Transforming integrals          ... done
Storing passive energy          ... done (   -73.59625384 Eh)
Transforming internal FI        ... done
      .... done with the Frozen Core Fock matrices

```

The next step is to transform the electron-electron repulsion integrals into the MO basis:

```

-----
SHARK HALF TRANSFORMATION
-----

Number of basis functions          ...    38
Number of operators                ...     1
  Operator 0:    2- 37

Integral generator used            ... SHARK
Contraction scheme used           ... SEGMENTED CONTRACTION
MaxCore in resort                 ...   256 MB

Half transformed integrals for op= 0 ... test1.SHARK_MNPQ0.tmp
Resorted half transformed integrals ... test1.JAO.tmp
Starting integral generation + half trafo...
Half trafo (segmented) done. Total time =  0.1 sec. integrals=  0.0 sec trafo=  0.
↪0 sec
Starting integral resorting        ... done (  0.0 sec)

SHARK half integral transformation done. Total time =  0.1 sec.

-----
FULL TRANSFORMATION
-----

Processing MO  10
Processing MO  20
Processing MO  30
Full transformation done
Number of integrals made          ...   222111
Number of integrals stored        ...    59070
Timings:
Time for first half transformation ...    0.068 sec
Time for second half transformation ...    0.014 sec
Total time                       ...    0.086 sec

```

This will result in a few additional disk files required by `orca_mrci`. The program then tells you which multiplicities will be treated in this MRCI run:

```

-----
CI-BLOCK STRUCTURE
-----

```

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```

Number of CI-blocks          ... 3

=====
CI BLOCK 1
=====
Multiplicity                ... 1
Irrep                      ... 0
Number of reference defs ... 1
    Reference    1: CAS(2,2)

Excitation type              ... CISD
Excitation flags for singles:
    1 1 1 1
Excitation flags for doubles:
    1 1 1 / 1 1 1 / 1 1 1

=====
CI BLOCK 2
=====
Multiplicity                ... 1
Irrep                      ... 1
Number of reference defs ... 1
    Reference    1: CAS(2,2)

Excitation type              ... CISD
Excitation flags for singles:
    1 1 1 1
Excitation flags for doubles:
    1 1 1 / 1 1 1 / 1 1 1

=====
CI BLOCK 3
=====
Multiplicity                ... 3
Irrep                      ... 1
Number of reference defs ... 1
    Reference    1: CAS(2,2)

Excitation type              ... CISD
Excitation flags for singles:
    1 1 1 1
Excitation flags for doubles:
    1 1 1 / 1 1 1 / 1 1 1

-----
----- ALL SETUP TASKS ACCOMPLISHED -----
----- (      0.139 sec) -----
-----

```

Now that all the setup tasks have been accomplished the MRCI calculation itself begins.

```

#####
#                                     #
#                               M R C I                               #
#                                     #
# TSel      =  1.000e-06 Eh      #
# TPre      =  1.000e-02      #
# TIntCut   =  1.000e-10 Eh      #
# Extrapolation to unselected MR-CI by full MP2      #
# DAVIDSON-1 Correction to full CI      #
#

```

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```

#                                     #
#####

-----
INTEGRAL ORGANIZATION
-----

Reading the one-Electron matrix          ... done
E0 read was -73.596253835266
Reading the internal Fock matrix          ... assuming it to be equal to the one-
↪electron matrix!!!
done
Preparing the integral list              ... done
Loading the full integral                 ... done
Making the simple integrals              ... done

          *****
          *                CI-BLOCK 1                *
          *****

Configurations with insufficient # of SOMOs WILL be rejected
Building a CAS(2,2) for multiplicity 1 and irrep=A1
Reference Space:
  Initial Number of Configurations :      2
  Internal Orbitals :      2 -      6
  Active Orbitals :      7 -      8
  External Orbitals :      9 -     37
The number of CSFs in the reference is 2
Calling MRPT_Selection with N(ref)=2

```

In the first step, the reference space is diagonalized. From this CI, the most important configurations are selected with Tpre:

```

-----
REFERENCE SPACE CI
-----

Pre-diagonalization threshold           : 1.000e-02
Warning: Setting NGuessMat to 512
N(ref-CFG)=2 N(ref-CSF)=2

          ****Iteration      0****
Lowest Energy       : -113.779221580786
Maximum Energy change : 113.779221580786 (vector 0)
Maximum residual norm : 0.000000000000

          *** CONVERGENCE OF RESIDUAL NORM REACHED ***
Reference space selection using TPre= 1.00e-02

... found 1 reference configurations (1 CSFs)
... now redoing the reference space CI ...

Warning: Setting NGuessMat to 512
N(ref-CFG)=1 N(ref-CSF)=1

          ****Iteration      0****
Lowest Energy       : -113.778810013503
Maximum Energy change : 113.778810013503 (vector 0)
Maximum residual norm : 0.000000000000

```

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*** CONVERGENCE OF RESIDUAL NORM REACHED ***

In this case, the CAS space only has 2 correctly symmetry adapted CSFs one of which (the closed-shell determinant) is selected. In general, larger CAS spaces usually carry around a lot of unnecessary CSFs which are not needed for anything and then the selection is important to reduce the computational effort. The result of the second reference space CI is printed:

```
-----
CI-RESULTS
-----
```

```
The threshold for printing is 0.30 percent
The weights of configurations will be printed. The weights are summed over
all CSFs that belong to a given configuration before printing
```

```
STATE 0: Energy= -113.778810014 Eh RefWeight= 1.0000 0.00 eV 0.0 cm**-1
1.0000 : h---h---[20]
```

Energy is the total energy in Eh. In the present case we can compare to the SCF energy -113.778810014 Eh and find that the reference space CI energy is identical, as it has to be since the lowest state coincides with the reference space. **RefWeight** gives the weight of the reference configurations in a CI state. This is 1.0 in the present case since there were only reference configurations. The number 1.000 is the weight of the following configuration in the CI vector. The description of the configuration **h-h-[20]p-p-** is understood as follows:² The occupation of the active orbitals is explicitly given in square brackets. Since the HOMO orbitals is number 7 from the SCF procedure, this refers to MOs 7 and 8 in the present example since we have two active orbitals. The 2 means doubly occupied, the 0 means empty. Any number (instead of -) appearing after an h gives the index of an internal orbital in which a hole is located. Similarly, any number after a p gives the index of an virtual (external) MO where a particle is located. Thus h-h-[20] is a closed shell configuration and it coincides with the SCF configuration—this was of course to be expected. The second root (in CI-Block 2) h-h-[11] by comparison refers to the configuration in which one electron has been promoted from the HOMO to the LUMO and is therefore the desired state that we wanted to calculate. Things are happy therefore and we can proceed to look at the output.

The next step is the generation of excited configurations and their selection based on Tsel:

```
-----
MR-PT SELECTION Tsel= 1.00e-06
-----
```

```
Setting reference configurations WITH use of symmetry
Building active patterns WITH use of symmetry
```

```
Selection will be done from 1 spatial configurations
```

```
Selection will make use of spatial symmetry
```

```
( 0) Refs      : Sel:      1CFGs/      1CSFs Gen:      1CFGs/      _
↪ 1CSFs (      0.000 sec)
Building active space densities      ...      0.002 sec
Building active space Fock operators  ...      0.000 sec
( 1) (p,q)->(r,s): Sel:      1CFGs/      1CSFs Gen:      1CFGs/      _
↪ 1CSFs (      0.000 sec)
( 2) (i,-)->(p,-): Sel:      1CFGs/      1CSFs Gen:      1CFGs/      _
↪ 1CSFs (      0.000 sec)
( 3) (i,j)->(p,q): Sel:      8CFGs/      8CSFs Gen:      8CFGs/      _
↪ 8CSFs (      0.000 sec)
( 4) (i,p)->(q,r): Sel:      0CFGs/      0CSFs Gen:      1CFGs/      _
↪ 1CSFs (      0.000 sec)
( 5) (p,-)->(a,-): Sel:      8CFGs/      8CSFs Gen:      8CFGs/      _
↪ 8CSFs (      0.000 sec)
```

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² Note that for printing we always sum over all linearly independent spin couplings of a given spatial configuration and only print the summed up weight for the configuration rather than for each individual CSF of the configuration.

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(6) (i,-)→(a,-): Sel:	52CFGs/	52CSFs Gen:	52CFGs/	↳
↳ 52CSFs (0.000 sec)			
(7) (i,j)→(p,a): Sel:	95CFGs/	166CSFs Gen:	96CFGs/	↳
↳ 167CSFs (0.000 sec)			
(8) (i,p)→(q,a): Sel:	21CFGs/	42CSFs Gen:	22CFGs/	↳
↳ 44CSFs (0.000 sec)			
(9) (p,q)→(r,a): Sel:	3CFGs/	3CSFs Gen:	5CFGs/	↳
↳ 5CSFs (0.000 sec)			
(10) (i,p)→(a,b): Sel:	555CFGs/	1082CSFs Gen:	584CFGs/	↳
↳ 1139CSFs (0.001 sec)			
(11) (p,q)→(a,b): Sel:	124CFGs/	124CSFs Gen:	148CFGs/	↳
↳ 148CSFs (0.000 sec)			
(12) (i,j)→(a,b): Sel:	1688CFGs/	2685CSFs Gen:	1887CFGs/	↳
↳ 2947CSFs (0.001 sec)			
Selection results:				
Total number of generated configurations:	2814			
Number of selected configurations	:	2557 (90.9%)		
Total number of generated CSFs	:	4522		
Number of selected CSFs	:	4173 (92.3%)		
The selected tree structure:				
Number of selected Internal Portions	:	11		
Number of selected Singly External Portions:		27		
average number of VMOs/Portion	:	6.39		
percentage of selected singly externals	:	22.83		
Number of selected Doubly External Portions:		21		
average number of VMOs/Portion	:	107.59		
percentage of selected doubly externals	:	27.76		

Here, the program loops through classes of excitations. For each excitation it produces the excited configurations (CFGs) and from it the linearly independent spin functions (CSFs) which are possible within the configuration. It then calculates the interaction with the contracted 0th order roots and includes all CSFs belonging to a given CFG in the variational space if the largest second order perturbation energy is larger or equal to Tsel. In the present case ≈136,000 CSFs are produced of which 25% are selected. For larger molecules and basis sets it is not uncommon to produce 10⁹–10¹⁰ configurations and then there is no choice but to select a much smaller fraction than 20%. For your enjoyment, the program also prints the total energies of each state after selection:

Diagonal second order perturbation results:				
State	E(tot)	E(0)+E(1)	E2(sel)	E2(unsel)
	Eh	Eh	Eh	Eh
-----	-----	-----	-----	-----
0	-114.108350270	-113.778810014	-0.329433	-0.000107

You can ignore this output if you want. In cases that the perturbation procedure is divergent (not that uncommon!) the total energies look strange—don't worry—the following variational calculation is still OK. The second order perturbation energy is here divided into a selected part E2(sel) and the part procedure by the unselected configurations E2(unsel). Depending on the mode of EUnselOpt this value may already be used later as an estimate of the energetic contribution of the unselected CSFs.³

Now we have ≈ 4,200 CSFs in the variational space of CI block 1 and proceed to diagonalize the Hamiltonian over these CSFs using a Davidson or DIIS type procedure:

```
-----
DAVIDSON-DIAGONALIZATION
-----
```

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³ In this case the maximum overlap of the 0th order states with the final CI vectors is computed and the perturbation energy is added to the “most similar root”. This is of course a rather crude approximation and a better choice is to recompute the second order energy of the unselected configurations rigorously as is done with EUnselOpt = FullMP2.

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```

Dimension of the eigenvalue problem      ... 4173
Number of roots to be determined         ... 1
Maximum size of the expansion space      ... 4
Maximum number of iterations             ... 35
Convergence tolerance for the residual   ... 1.000e-06
Convergence tolerance for the energies   ... 1.000e-06
Orthogonality tolerance                  ... 1.000e-14
Level Shift                             ... 0.000e+00
Constructing the preconditioner           ... o.k.
Building the initial guess                ... o.k.
Number of trial vectors determined        ... 4

      ****Iteration 0****
Size of expansion space: 3
Lowest Energy      : -113.937028067251
Maximum Energy change : 113.937028067251 (vector 0)
Maximum residual norm : 0.741727830968

      ****Iteration 1****
Size of expansion space: 4
Lowest Energy      : -114.082265676116
Maximum Energy change : 0.145237608865 (vector 0)
Maximum residual norm : 0.012707561344
Rebuilding the expansion space

      ****Iteration 2****
Size of expansion space: 2
Lowest Energy      : -114.085350429118
Maximum Energy change : 0.003084753001 (vector 0)
Maximum residual norm : 0.002880697397

      ****Iteration 3****
Size of expansion space: 3
Lowest Energy      : -114.086043274125
Maximum Energy change : 0.000692845007 (vector 0)
Maximum residual norm : 0.000098595378

      ****Iteration 4****
Size of expansion space: 4
Lowest Energy      : -114.086074300143
Maximum Energy change : 0.000031026018 (vector 0)
Maximum residual norm : 0.000004959126
Rebuilding the expansion space

      ****Iteration 5****
Size of expansion space: 2
Lowest Energy      : -114.086076038587
Maximum Energy change : 0.000001738444 (vector 0)
Maximum residual norm : 0.000000572348

      *** CONVERGENCE OF RESIDUAL NORM REACHED ***

Storing the converged CI vectors          ... test1.mrci.vec

      *** DAVIDSON DONE ***
Returned from DIAG section

```

The procedure converges on all roots simultaneously and finishes after six iterations which is reasonable. Now the program calculates the Davidson correction (DavidsonOpt) which is printed for each root.

Davidson type correction:

Root= 0 W= 0.912 E0= -113.778810014 ECI= -114.086076039 DE=-0.026913

Already in this small example the correction is pretty large, ca. 27 mEh for the ground state (and ≈ 36 mEh for the excited state, later in the output). Thus, a contribution of ≈ 9 mEh = 0.25 eV is obtained for the transition energy which is certainly significant. Unfortunately, the correction becomes unreliable as the reference space weight drops or the number of correlated electrons becomes large. Here 0.912 and 0.888 are still OK and the system is small enough to expect good results from the Davidson correction.

The next step is to estimate the correction for the unselected configurations:

Unselected CSF estimate:

Full relaxed MR-MP2 calculation ...

Selection will be done from 1 spatial configurations

Selection will make use of spatial symmetry

Selection will make use of spatial symmetry

Selection will make use of spatial symmetry

done

Selected MR-MP2 energies ...

Root= 0 E(unsel)= -0.000106931

In the present case this is below 1 mEh and also very similar for all three states such that it is not important for the transition energy.

CI-RESULTS

The threshold for printing is 0.30 percent

The weights of configurations will be printed. The weights are summed over all CSFs that belong to a given configuration before printing

```
STATE 0: Energy= -114.113096211 Eh RefWeight= 0.9124 0.00 eV 0.0 cm**-1
      0.9124 : h---h---[20]
      0.0114 : h 6h 6[22]
```

The final ground state energy is -114.113096211 which is an estimate of the full CI energy in this basis set. The leading configuration is still the closed-shell configuration with a weight of $\approx 91\%$. However, a double excitation *outside* the reference space contributes some 1%. This is the excitation MO6,MO6 \rightarrow LUMO,LUMO. This indicates that more accurate results are expected once MO6 is also included in the reference space (this is the HOMO-1). The excited state is dominated by the HOMO-LUMO transition (as desired) but a few other single- and double- excitations also show up in the final CI vector.

Now that all CI vectors are known we can order the states according to increasing energy and print (vertical) transition energies:

TRANSITION ENERGIES

The lowest energy is -114.113096211 Eh

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	1	A1	0	0	0.000	0.000	0.0
1	3	A2	0	2	134.086	3.649	29428.4
2	1	A2	0	1	148.499	4.041	32591.8

This result is already pretty good and the transition energies are within ≈ 0.1 eV of their experimental gas phase values (≈ 3.50 and ≈ 4.00 eV) and may be compared to the CIS values of 3.8 and 4.6 eV which are considerably in error.

In the next step the densities and transition densities are evaluated and the absorption and CD spectra are calculated (in the dipole length formalism) for the spin-allowed transitions together with state dipole moments:

ABSORPTION SPECTRUM VIA TRANSITION ELECTRIC DIPOLE MOMENTS							
Transition	Energy	Energy	Wavelength	fosc(D2)	D2	DX	
DY DZ	(eV)	(cm-1)	(nm)		(au**2)	(au)	
(au) (au)							
0-1A1 -> 0-1A2	4.040866	32591.8	306.8	0.0000000000	0.000000	-0.000000	
0.000000 0.000000							
CD SPECTRUM VIA TRANSITION ELECTRIC DIPOLE MOMENTS							
Transition	Energy	Energy	Wavelength	R	MX	MY	
MZ	(eV)	(cm-1)	(nm)	(1e40*cgs)	(au)	(au)	
(au)							
0-1A1 -> 0-1A2	4.040866	32591.8	306.8	-0.000000	-0.000000	-0.000000	0.
59348							
STATE DIPOLE MOMENTS							
Root	Block	TX (Debye)	TY (Debye)	TZ (Debye)	T (Debye)		
0	0	0.000000	-0.000000	2.33244	2.33244		
0	2	0.000000	-0.000000	1.45831	1.45831		
0	1	0.000000	-0.000000	1.58658	1.58658		

Here the transition is symmetry forbidden and therefore has no oscillator strength. The state dipole moment for the ground state is 2.33 Debye which is somewhat lower than 2.87 Debye from the SCF calculation. Thus, the effect of correlation is to reduce the polarity consistent with the interpretation that the ionicity of the bonds, which is always overestimated by HF theory, is reduced by the correlation. Finally, you also get a detailed population analysis for each generated state density which may be compared to the corresponding SCF analysis in the preceding part of the output.

This concludes the initial example on the use of the MR-CI module. The module leaves several files on disk most of which are not yet needed but in the future will allow more analysis and restart and the like. The `.ivo` file is a standard `.gbw` type file and the orbitals therein can be used for visualization. This is important in order to figure out the identity of the generated IVOs. Perhaps they are not the ones you wanted and then you need to re-run the MR-CI with the IVOs as input, `NoIter` and the IVO feature in the new run turned off! We could use the IVOs as input for a state averaged CASSCF calculation:

```
! moread UseSym KDIIS
%moinp "Test-SYM-MRCI-H2CO.ivo"

%casscf nel      2
        norb     2
        irrep    0,1,1
```

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```

mult      1,1,3
nroots    1,1,1
end

```

If we based a MR-ACPF calculation on this reference space we will find that the calculated transition energies are slightly poorer than in the MRCI+Q calculation. This is typical of approximate cluster methods that usually require somewhat larger reference spaces for accurate results. A similar result is obtained with SORCI.

```

%mrcki  CItType          SORCI
        tsel              1e-6
        tpre              1e-4
        tnat              1e-5
        AllSingles        true
        doNatOrbs          true
        IntMode            FullTrafo
        # ground state 1A1
        NewBlock 1 0
          NRoots 1
          Excitations cisd
          Refs CAS(2,2) end
        End
        # HOMO LUMO transition 1A2
        NewBlock 1 1
          NRoots 1
          Excitations cisd
          Refs CAS(2,2) end
        End
        # HOMO LUMO triplet transition 3A2
        NewBlock 3 1
          NRoots 1
          Excitations cisd
          Refs CAS(2,2) end
        End
end

```

This gives:

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	1	A1	0	0	0.000	0.000	0.0
1	3	A2	0	2	144.563	3.934	31728.0
2	1	A2	0	1	161.179	4.386	35374.7

This is systematically 0.4 eV too high. But let us look at the approximate average natural orbital (AANOs) occupation numbers:

----- AVERAGE NATURAL ORBITALS -----

```

Trace of the density to be diagonalized = 12.000000
Sum of eigenvalues = 12.000000
Natural Orbital Occupation Numbers:
N[ 2] ( A1)=  1.99831062
N[ 3] ( A1)=  1.99761604
N[ 4] ( A1)=  1.99479313
N[ 5] ( B1)=  1.99016881
N[ 6] ( B2)=  1.95818285
N[ 7] ( B1)=  1.33014178
N[ 8] ( B2)=  0.70688423
N[ 9] ( B1)=  0.00988561
N[10] ( A1)=  0.00436843

```

This shows that there is a low-occupancy orbital (MO6) that has not been part of the reference space. Thus, we try the same calculation again but now with one more active orbital and two more active electrons:

```
! moread
%moinp "Test-SYM-MRCI-H2CO.gbw"

%casscf nel 4
      norb 3
      irrep 0,1,1
      mult 1,1,3
      nroots 1,1,1
      end

%mrcki CItpe          SORCI
      tsel            1e-6
      tpre            1e-4
      tnat            1e-5
      AllSingles      true
      doNatOrbs       true
      IntMode          FullTrafo
      # ground state 1A1
      NewBlock 1 0
        NRoots 1
        Excitations cisd
        Refs CAS(4,3) end
      End
      # HOMO LUMO transition 1A2
      NewBlock 1 1
        NRoots 1
        Excitations cisd
        Refs CAS(4,3) end
      End
      # HOMO LUMO triplet transition 3A2
      NewBlock 3 1
        NRoots 1
        Excitations cisd
        Refs CAS(4,3) end
      End
end
```

This gives:

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	1	A1	0	0	0.000	0.000	0.0
1	3	A2	0	2	145.494	3.959	31932.3
2	1	A2	0	1	162.222	4.414	35603.6

Which is now fine since all essential physics has been in the reference space. Inspection of the occupation numbers show that there is no suspicious orbital any more. Note that this is still a much more compact calculation than the MRCI+Q.

Likewise, we get an accurate result from MRACPF with the extended reference space.

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	1	A1	0	0	0.000	0.000	0.0
1	3	A2	0	2	134.985	3.673	29625.8
2	1	A2	0	1	148.330	4.036	32554.6

However, the SORCI calculation is *much* more compact. For larger molecules the difference becomes more and more pronounced and SORCI or even MRDDCI2 (with or without +Q) maybe the only feasible methods—if at all.

3.19.5 Excitation Energies between Different Multiplicities

As an example for a relatively accurate MRCI+Q calculation consider the following job which calculates the triplet-ground and as the first excited singlet states of O₂.

```
! ano-pVQZ RI-AO cc-pVQZ/JK VeryTightSCF NoPop Conv UseSym RI-MP2 PModel
%mp2 density relaxed natorbs true end
%base "O2"
* xyz 0 3
O 0 0 0
O 0 0 1.2
*

$new_job
! ano-pVQZ RI-AO cc-pVQZ/JK VeryTightSCF NoPop Conv UseSym KDIIS
! moread
%moinp "O2.mp2nat"
%casscf nel      8
      norb      6
      irrep     1,0,1
      nroots    1,2,1
      mult      3,1,1
      trafostep  ri
      switchstep nr
      end

%mrcki  citype mrcki
      tsel 1e-7
      tpre 1e-5
      newblock 3 1 nroots 1 refs cas(8,6) end end
      newblock 1 0 nroots 2 refs cas(8,6) end end
      newblock 1 1 nroots 1 refs cas(8,6) end end
      end

* xyz 0 3
O 0 0 0
O 0 0 1.2
*
```

Note that the linear molecule is run in D_{2h}. This creates a slight problem as the CASSCF procedure necessarily breaks the symmetry of the ¹Δ state.

```
LOWEST ROOT (ROOT 0, MULT 3, IRREP B1g) = -149.765383866 Eh -4075.323 eV
```

STATE	ROOT	MULT	IRREP	DE/a.u.	DE/eV	DE/cm** ⁻¹
1:	0	1	B1g	0.033334	0.907	7316.0
2:	0	1	Ag	0.033650	0.916	7385.3
3:	1	1	Ag	0.062381	1.697	13691.1

The result of the MRCI+Q is:

```
-----
TRANSITION ENERGIES
-----
```

The lowest energy is -150.176905551 Eh

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	3	B1g	0	0	0.000	0.000	0.0
1	1	B1g	0	2	36.971	1.006	8114.2
2	1	Ag	0	1	38.021	1.035	8344.7
3	1	Ag	1	1	62.765	1.708	13775.2

These excitation energies are accurate to within a few hundred wavenumbers. Note that the ≈ 200 wavenumber splitting in the degenerate $^1\Delta$ state is due to the symmetry breaking of the CAS and the individual selection. Repeating the calculation with the MP2 natural orbitals gives an almost indistinguishable result and a ground state energy that is even lower than what was found with the CASSCF orbitals. Thus, such natural orbitals (that might often be easier to get) are a good substitute for CASSCF orbitals and at the same time the symmetry breaking due to the use of symmetry appears to be difficult to avoid.

----- TRANSITION ENERGIES -----

The lowest energy is -150.177743426 Eh

State	Mult	Irrep	Root	Block	mEh	eV	1/cm
0	3	B1g	0	0	0.000	0.000	0.0
1	1	B1g	0	2	37.369	1.017	8201.5
2	1	Ag	0	1	38.237	1.040	8392.1
3	1	Ag	1	1	62.731	1.707	13767.9

3.19.6 Correlation Energies

The logic we are following here is the following: CID minus SCF gives the effect of the doubles; going to CISD gives the effect of the singles; QCISD(=CCD) minus CID gives the effect of the disconnected quadruples. QCISD minus QCID gives simultaneously the effect of the singles and the disconnected triples. They are a bit difficult to separate but if one looks at the singles alone and compares with singles + disconnected triples, a fair estimate is probably obtained. Finally, QCISD(T) minus QCISD gives the effect of the connected triples. One could of course also use CCSD instead of QCISD but I felt that the higher powers of T_1 obscure the picture a little bit—but this is open to discussion of course.

First $\text{H}_2\text{O}/\text{TZVPP}$ at its MP2/TZVPP equilibrium geometry ($T_{\text{pre}} = 10^{-6}$ and $T_{\text{sel}} = 10^{-9}$ Eh for the MRCI and MRACPF calculations):

Excitation class	Energy (Eh)	Delta-Energy (mEh)
None (RHF)	-76.0624	
Doubles (CID)	-76.3174	255
+Singles (CISD)	-76.3186	1
+Disconnected Quadruples (QCID)	-76.3282	11
+Disconnected Triples (QCISD)	-76.3298	2
+Connected Triples (QCISD(T))	-76.3372	7
CASSCF(8,6)	-76.1160	
CASSCF(8,6) + MRCI	-76.3264	210
CASSCF(8,6) + MRCI+Q	-76.3359	10
CASSCF(8,6) + MRACPF	-76.3341	218

One observes quite good agreement between single- and multireference approaches. In particular, the contribution of the disconnected triples and singles is very small. The estimate for the disconnected quadruples is fairly good from either the multireference Davidson correction or the ACPF and the agreement between CCSD(T) and these MR methods is 2-3 mEh in the total energy which is roughly within chemical accuracy.

In order to also have an open-shell molecule let us look at NH with a N-H distance of 1.0 Å using the TZVPP basis set.

Excitation class	Energy (Eh)	Delta-Energy (mEh)
None (UHF)	-54.9835	
Doubles (CID)	-55.1333	150
+Singles (CISD)	-55.1344	1
+Disconnected Quadruples (QCID)	-55.1366	3
+Disconnected Triples (QCISD)	-55.1378	1
+Connected Triples (QCISD(T))	-55.1414	4
CASSCF(6,5)	-55.0004	
CASSCF(6,5) + MRCI	-55.1373	137
CASSCF(6,5) + MRCI+Q	-55.1429	6
CASSCF(6,5) + MRACPF	-55.1413	141

Again, the agreement is fairly good and show that both single- and multiple reference approaches converge to the same limit.

3.19.7 Thresholds

Now we choose the CO molecule (1.128 Ångström) with the SVP basis set and study the convergence of the results with respect to the selection threshold. Comparison to high level single-reference approaches is feasible (The SCF energy is -112.645 946 Eh).

Reference Values for Total Energies

The single-reference values are:

```
BD: -112.938 48002
      CCSD: -112.939 79145
      QCISD: -112.941 95700
      BD (T): -112.950 17278
      CCSD (T): -112.950 63889
      QCISD (T): -112.951 37425
      MP4 (SDTQ): -112.954 80113
```

The calculations without connected triples (BD, CCSD, QCISD) are about the best what can be achieved without explicitly considering triple excitations. The CCSD is probably the best in this class. As soon as connected triples are included the CCSD(T), QCISD(T) and BD(T) values are close and from experience they are also close to the full CI values which is then expected somewhere between -112.950 and -112.952 Eh.

Convergence of Single Reference Approaches with Respect to T_{sel}

Next it is studied how these single reference methods converge with T_{sel} :

```
Closed-Shell ACPF:
  Tsel      Energy      (NCSF)      Energy      (NCSF)
  (Eh)      AllSingles=true  AllSingles=false
Tsel=0      -112.943 387 (5671)
Tsel=1e-14   -112.943 387 (2543)  -112.943 387 (2478)
Tsel=1e-10   -112.943 387 (2543)  -112.941 023 (2453)
Tsel=1e-08   -112.943 387 (2451)  -112.937 087 (2346)
Tsel=1e-06   -112.943 350 (2283)  -112.937 046 (2178)
Tsel=1e-05   -112.943 176 (1660)  -112.936 821 (1555)
Tsel=1e-04   -112.944 039 ( 782)  -112.938 381 ( 677)
```


It is clear that the convergence is erratic if the singles are not automatically included. This is the reason for making this the default from release 2.6.35 on. In the present case singles will only be selected due to round-off errors since by Brillouin's theorem the singles have zero-interaction with the ground state determinant. Thus, for individually selecting single-reference methods it is a good idea to automatically include all single-excitations in order to get converged results. The alternative would be a different singles selection procedure which has not yet been developed however. The selection of doubles appear to converge the total energies reasonably well. It is seen that the selection selects most CSFs between 10^{-5} and 10^{-7} Eh. Already a threshold of 10^{-6} Eh yields an error of less than 0.1 mEh which is negligible in relation to reaction energies and the like. Even 10^{-5} Eh gives an error of less than 0.1 kcal/mol.

Convergence of Multireference Approaches with Respect to T_{pre}

We next turn to multireference treatments. Here we want to correlate all valence electrons in all valence orbitals and therefore a CAS(10,8) is the appropriate choice. We first ask for the converged value of T_{pre} by using $T_{\text{sel}} = 10^{-14}$ and obtain for MRCI+Q:

```
TPre = 1e-1: -112.943 964
        1e-2: -112.952 963
        1e-3: -112.953 786
        1e-4: -112.954 019
        1e-5: -112.954 336
        1e-6: -112.954 416
        1e-7: -112.954 440
```

Thus, pretty good convergence is obtained for $T_{\text{pre}} = 10^{-4} - 10^{-6}$. Hence 10^{-4} is the default.

To show a convenient input consider the following:

```
#
# Here we calculate the CO ground state correlation energy with several methods
#
! Def2-SVP Def2-SV/C RI-MP2 CCSD(T)
%base "1"

%mp2 density relaxed
      donatorbs true
end

* int 0 1
C    0  0  0  0.000000      0.000      0.000
O    1  0  0  1.128        0.000      0.000
*

$new_job

! aug-SVP MRACPF
! moread
%moinp "1.mp2nat"
# the CASSCF is done with MP2 natural orbitals which is a good idea and
# secondly we use a large level shift in order to help convergence
%casscf nel      10
        norb      8
        mult      1
        nroots    1
        shiftup    2
        shiftdn    2
end

%mrcki tsel 1e-8
        tpre 1e-6
end
```

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```
* int 0 1
C      0  0  0  0.000  0.000  0.000
O      1  0  0  1.128  0.000  0.000
*
```

This job computes at the same time all of the below and demonstrates once more the agreement between consequent single- and multireference correlation methods

```
SCF           = -112.6459
RI-MP2        = -112.9330
CCSD          = -112.9398
CCSD(T)       = -112.9506
CASSCF(10,8)  = -112.7769
MRACPF        = -112.9514
```

3.19.8 Energy Differences - Bond Breaking

For the calculation of energy differences we start again with the reference CCSD(T) calculation; this method is one of the few which can claim chemical accuracy in practical applications:

```
Reference Total Energies for N2 at 1.0977 Angstr\ "{o}m with
The SVP basis
  E(CCSD)      = -109.163 497
  E(CCSD(T)) = -109.175 625
Nitrogen Atom (4S), SVP basis, unrestricted
  E(CCSD)      = -54.421 004
  E(CCSD(T)) = -54.421 7183
Energy Difference:
Delta-E(CCSD)  = -0.321 489 = 8.75 eV
Delta-E(CCSD(T)) = -0.332 188 = 9.04 eV
```

The basis set is of course not suitable for quantitative comparison to experimental values. However, this is not the point here in these calculations which are illustrative in nature. The SVP basis is just good enough to allow for a method assessment without leading to excessively expensive calculations.

This is now to be compared with the corresponding energy differences computed with some single-reference approaches. A typical input is (this is a somewhat old-fashioned example – in the present program version you would do a full valence CASSCF(10,8) or CASSCF(6,6) and invoke the MR-methods with a single keyword):

```
! HF def2-SVP def2-TZVPP/C VeryTightSCF NoPop

%base "1"

* xyz 0 1
N 0 0 0
N 0 0 1.0977
*
%method
  frozencore fc_ewin
end

%mrcki
  EWin          -3,1000
  CItType       MRACPF2a
  Solver        DIIS
  IntMode       FullTrafo
  UseIVOs       true
  AllSingles    true
  TSel          1e-14
```

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```

    TPre          1e-05
    TNat          0.0
    ETol          1e-10
    RTol          1e-10
    NewBlock 1 *
      NRoots 1
      Excitations CISD
      refs CAS(0,0) end
    end
  end
end

$new_job

! ROHF def2-SVP def2-TZVPP/C VeryTightSCF NoPop PModel

%base "2"

* xyz 0 4
N 0 0 0
*
%method
  frozencore fc_ewin
end

%mrcki
  EWin          -3,1000
  CITYPE        MRACPF2a
  IntMode       FullTrafo
  UseIVOs       true
  AllSingles    true
  TSel          1e-14
  TPre          1e-05
  TNat          0.0
  ETol          1e-10
  RTol          1e-10
  NewBlock 4 *
    NRoots 1
    Excitations CISD
    refs CAS(3,3) end
  end
end

```

The results are:

Single reference approaches:

Method	N2-Molecule	N-Atom	Delta-E
CISD+Q	: -109.167 904	-54.422 769	8.77 eV
ACPF	: -109.166 926	-54.421 783	8.80 eV
ACPF2	: -109.166 751	-54.421 333	8.82 eV
ACPF2a	: -109.166 730	-54.421 186	8.83 eV
CEPA1	: -109.159 721	-54.422 564	8.56 eV
CEPA2	: -109.172 888	-54.422 732	8.91 eV
CEPA3	: -109.161 034	-54.422 589	8.59 eV
AQCC	: -109.160 574	-54.420 948	8.67 eV
CEPA-0	: -109.174 924	-54.422 951	8.95 eV

With exception is CEPA1 and CEPA3, the results are OK. The reason for the poor performance of these methods is simply that the formalism implemented is only correct for closed shells – open shells require a different formalism which we do not have available in the MRCI module (but in the single reference MDCI module). Due to the simple approximations made in CEPA2 it should also be valid for open shells and the numerical results are in support of that.

Next we turn to the multireference methods and take a CAS(10,8) reference as for CO in order to correlate all valence

electrons.⁴

Multi reference approaches:

Method	N2-Molecule	N-Atom	Delta-E
MRCISD+Q:	-109.180 089	-54.422 667	9.11 eV
MRACPF :	-109.178 708	-54.421 685	9.12 eV
MRACPF2 :	-109.177 140	-54.421 236	9.11 eV
MRAQCC :	-109.175 947	-54.420 851	9.10 eV
SORCI :	-109.179 101	-54.422 703	9.08 eV

This test calculation pleasingly shows the high consistency of multireference approaches which all converge more or less to the same result which must be accurate.

3.19.9 Energy Differences - Spin Flipping

There are a number of interesting situations in which one is interested in a small energy difference which arises from two states of different multiplicity but same orbital configuration. This is the phenomenon met in diradicals or in magnetic coupling in transition metal complexes. As a primitive model for such cases one may consider the hypothetical molecule H-Ne-H in a linear configuration which will be used as a model in this section.

The reference value is obtained by a MR-ACPF calculation with all valence electrons active (again, this example is somewhat old fashioned – in the present program version you would do a CASSCF calculation followed by MR methods with a single keyword):

```
! ROHF def2-SVP def2-TZVPP/C VeryTightSCF NoPop
%basis
  NewAuxCGTO Ne "AutoAux" end
end
* xyz 0 3
H 0 0 0
Ne 0 0 2.0
H 0 0 4.0
*
%method frozencore fc_ewin
  end
%mrcki EWin -3,1000
  CIType MRACPF2a
  IntMode FullTrafo
  Solver DIIS
  UseIVOs true
  TSel 0
  TPre 1e-10
  ETol 1e-09
  RTol 1e-09
  DoDDCIMP2 true
  NewBlock 1 *
    NRoots 1
    Excitations CISD
    refs CAS(10,6) end
  end
  NewBlock 3 *
    NRoots 1
    Excitations CISD
    refs CAS(10,6) end
  end
end
```

⁴ Most of these results have been obtained with a slightly earlier version for which the MR energies are a little different from that what the present version gives. The energy differences will not be affected.