### MODELING EXAMPLES

1.	Prepare your geometry	1
2.	ThermoElectric	2
2.1.	General presentation of the files	2
2.2.	Examples	5
2.3.	Validity	17
3.	Magnetostatic	27
3.1.	General presentation of the files	27
3.2.	Examples	29
3.3.	Validity	29
3.4.	Cartesian model	38
4. Elasticity		40
4.1.	General presentation of the files	40
4.2.	Examples	41
4.3.	Validity	41

Contents

# 1. Prepare your geometry

First of all, you need to create the 3D model of your conductor with Salome or gmsh and to make the mesh. When the mesh is created, with specific names for the volumes and surfaces, export it in .med (example: example\_mesh.med) with Salome. You also need to partition the mesh if you want to do the calculations in parallel, running this command below.

example:

feelpp\_mesh\_partitioner --ifile example\_mesh.med --part 4 --nochdir

Where part indicate the number of processors you want to use. This will create a new file, call example\_mesh\_p4.json that you will use in your configure file.

### 2. ThermoElectric

2.1. General presentation of the files. You can use ether the linear or the nonlinear model for your calculation, but for our example, we use the linear model, so we select the thermoelectric-linear (respectively thermoelectric-nonlinear) model in the Json file (details in the Example json).

When all the files (detailed below) are created, you can run your calculation with this command :

1

 $\verb|mpirun -np "number_of_processor_you_have_chosen_in_the_partition" feelpp_hfm\_thermoel| \\$ 

### $2.1.1.\ Material.$

$\alpha$	alpha		
$\sigma_0 / \sigma$	sigma0 / sigma		
$k_0 / k$	k0 / k		
$T_0$	T0		

There is a dedicated section in the Json file, named Materials, to configure the magnet properties. The structure of a json file is as follows (conditions are details in Condition)

```
{
    "Name": "ThermoElectric",
    "ShortName": "TE",
    "Model": "thermoelectric-linear",
    "Materials":
    {
   "Name_of_the_first_volume":
     "name": "material_of_this_volume",
      "alpha":"_",
      "TO":"_",
      "sigma0":"_",
      "k0":" ".
      "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
      "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
    },
  "other_volume":
   {
   }
    "BoundaryConditions":
   "first_condition(like potential or temperature)":
     "type_of_condition":
      "Surface_concerned_by_the_condition":
          "expr1":"_",
          "expr2":"_"
        },
      "Surface_concerned_by_the_condition":
```

```
{
    "expr1":"_",
    "expr2":"_"
}
},
"other_condition":
{
    ...
}
},
"PostProcess":
    {
    "Fields":["temperature","potential","current"]
}
```

WARNING: The name of the volumes and surfaces must be the same as defined in Salome. Be careful to the units of the material properties. They need to be consistent with the length unit used for the rest. For instance the length unit is in mm in . It is also necessary to create a file to configure the calculation, call .cfg file (example thermoelectric\_3D\_V1T1\_N1\_cvgcfg (for the T1V1 model)). It will configure which file you will use in your calculations and which type of solver you use (here we use the Krylov method to solve both electro and thermal problem).

```
dim=3
geofile="name_of_the_file_created_by_the_partition.json" (or .msh)
geofile-path=$cfgdir

conductor_volume="name_of_your_volume"

[thermoelectric]
model_json=$cfgdir/"name_of_your_file.json"
weakdir=false

[electro]
pc-type=gamg
#ksp-monitor=true
ksp-rtol="relative_convergence_tolerance"
ksp-atol="absolute_convergence_tolerance"
ksp-maxit="maximum_number_of_iterations"
ksp-use-initial-guess-nonzero=1

[thermal]
pc-type=gamg
```

```
#ksp-monitor=true
ksp-rtol="relative_convergence_tolerance"
ksp-atol="absolute_convergence_tolerance"
ksp-use-initial-guess-nonzero=1
```

There are few differences between the linear and the nonlinear calculation. For the nonlinear model, just add this lines in the section thermoelectric:

```
nonlinear
eps_potential=1.e-4
eps_temperature=1.e-4
resolution=picard
itmax_picard=10
update_intensity=true
marker_intensity="the_surface"
target_intensity="the_intensity" (be careful of the sign)
eps_intensity=1.e-2
verbosity=2
```

We can define the current I using the Ohm's law, defining the voltage ine the json file.

# 2.1.2. Condition. There are three type of conditions:

### 1 Dirichlet

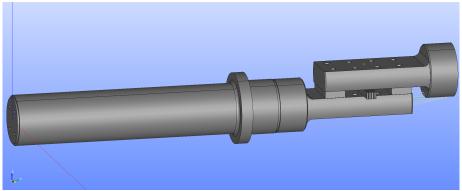
### 2 Neumann

},

**WARNING**: Your have to set a condition for each surfaces you have defined. For those where there is no conditions, set an homogeneous Neumann condition ("expr":"0")

### 2.2. Examples.

2.2.1. Current sensor. Here we want to model a current sensor placed in the connection of a magnet. This sensor consist in 6 constantan plates placed between the connectors as shown in the images. The goal is to allow the users to directly see the evolution of the current by measuring the potential in the connection. The main problem is the temperature reached by the sensor, due to the high current, the fusion of the constantan being around 1500 K, but we don't want to exceed an elevation of 200 degree. Because an excessive increase of the temperature will We model all of this with salome and do the mesh. The mesh is more precise on the sensor (the 6 constantan plates)



Parameters. There are 8 volumes, 2 connectors in copper and the 6 constantan plates which constitute the sensor. { "Name": "ThermoElectric", "ShortName": "TE", "Model": "thermoelectric-linear", "Materials": { "A1": { "name":"A1", "alpha": "3.35e-3", "T0": "293", "sigma0":"58e+3", "k0":"0.38", "sigma":"sigma0/(1+alpha\*(T-T0)):sigma0:alpha:T:T0", "k":"k0\*T/((1+alpha\*(T-T0))\*T0):k0:T:alpha:T0" }, "CurrentLead\_1": { "name": "CurrentLead\_1", "alpha":"3.35e-3", "T0":"293", "sigma0":"58e+3", "k0":"0.38", "sigma":"sigma0/(1+alpha\*(T-T0)):sigma0:alpha:T:T0", "k":"k0\*T/((1+alpha\*(T-T0))\*T0):k0:T:alpha:T0"

```
},
"Constantan_1":
{
    "name": "Constantan_1",
    "alpha":"0",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma": "sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k": "k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
"Constantan 2":
{
    "name": "Constantan_2",
    "alpha":"0",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k": "k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
},
"Constantan_3":
    "name": "Constantan_3",
    "alpha":"0",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma": "sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k": "k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
},
"Constantan_4":
    "name": "Constantan_4",
    "alpha":"0",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma": "sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k": "k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
},
"Constantan_5":
{
    "name": "Constantan 5",
    "alpha":"0",
```

```
"T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
"Constantan_6":
{
    "name": "Constantan_6",
    "alpha":"0",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"0.019",
    "sigma": "sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
}
},
Conditions.
"BoundaryConditions":
"potential":
{
  "Dirichlet":
"VO":
{
    "expr1":"0",
    "expr2":"A1"
},
"V1":
{
    "expr1":"0.1",
    "expr2": "CurrentLead_1"
}
 }
},
"temperature":
{
  "Dirichlet":
  {
"VO":
{
    "expr1":"293",
    "expr2":"A1"
```

```
},
"V1":
{
    "expr1":"293",
    "expr2": "CurrentLead 1"
 }
}
},
"PostProcess":
"Fields":["temperature", "potential", "current"]
}
There is also
a specific file
to study only one plate, to be more precise on the temperature reach. We can use
here a nonlinear model for the thermoelectric study.
{
    "Name": "ThermoElectric",
    "ShortName": "TE",
    "Model": "thermoelectric-nonlinear",
    "Materials":
"Constantan_1":
{
    "name": "Constantan_1",
    "alpha": "2.e-5",
    "T0":"293",
    "sigma0":"2.04e+3",
    "k0":"19.5e-3",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
}
    },
    "BoundaryConditions":
"potential":
    "Dirichlet":
"Interface_0":
    "expr1":"0",
```

```
"expr2": "Constantan_1"
},
"Interface_1":
{
    "expr1":"0.1153",
    "expr2": "Constantan_1"
}
    },
    "Neumann":
"Fixer":
{
    "expr":"0"
},
"Free_edge":
{
    "expr":"0"
}
    }
},
"temperature":
{
    "Robin":
"Free_edge":
    "expr1":"50.e-6", //the heat transfer coefficient
    "expr2":"293"
}
    },
    "Neumann":
"Fixer":
{
    "expr":"0"
"Interface_0":
}
    "expr":"0"
},
"Interface_1":
{
    "expr":"0"
}
    }
```

```
}
    },
    "PostProcess":
    {
    "Fields":["temperature","potential","current","joules"]
    }
}
```

Results. First we can see the potential we should be measuring.

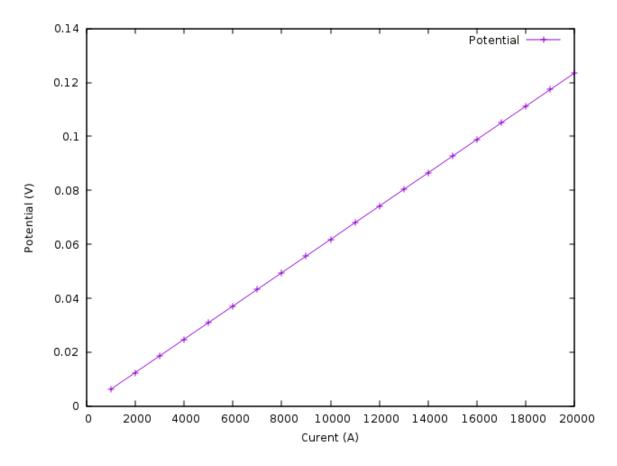


FIGURE 1. Potential as a function of current

Therefore, we need a voltmeter which can be precise between 0.01 and 0.15 Volt.

Next, we want to control the temperature reached by the sensor, the purpose being not to destroy the sensor. The melting temperature of the Constantan is near 1500 K, but we want to stop well before reaching this point to avoid deformations of the sensor. The main parameter that we can control is the heat transfer coefficient  $\mathbf{h}$  ( $W.m^{-2}.K^{-1}$ ). This coefficient  $\mathbf{h}$  can be control by displaying or not a ventilator to be in natural or forced convection.

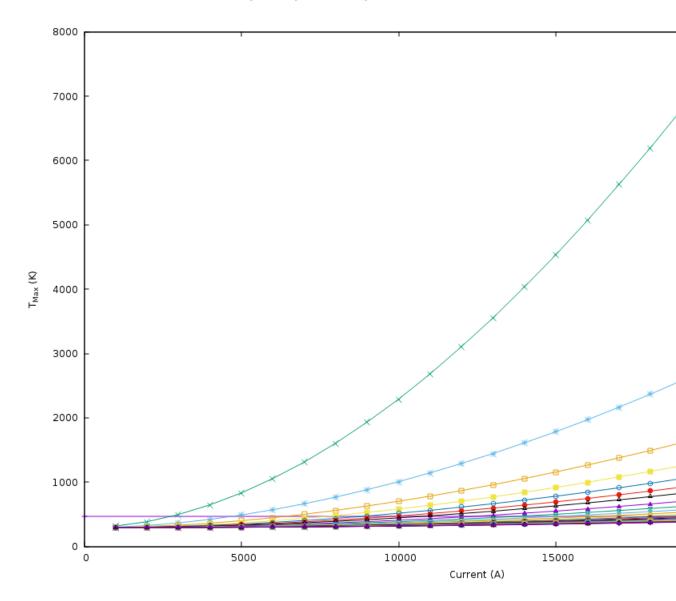


FIGURE 2. Temperature max in 1 plate as a function of current

Here we see that in natural convection (h=15), the temperature reached is by far too high.

A heat transfer coefficient maximum allow to set a higher current but is more difficult to set up.

2.2.2. Double Helix. For this piece, the purpose is to see the elevation of the temperature in the conductor around the double helix. We model a double helix on salome then do the mesh.

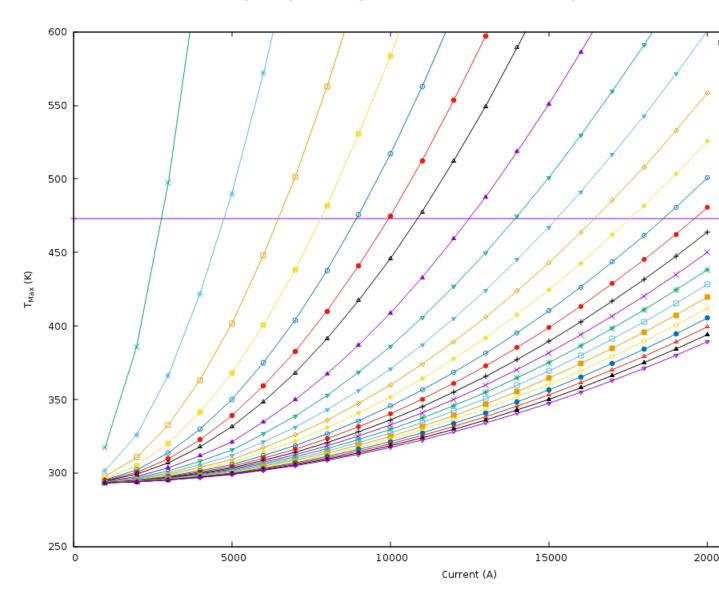


FIGURE 3. Temperature max in 1 plate as a function of current, forced convection only

```
Parameters.
```

```
{
    "Name": "ThermoElectric",
    "ShortName": "TE",
    "Model":"thermoelectric-nonlinear",
    "Materials":
    {
    "Cu":
```



FIGURE 4. Double Helix on 'Paraview'

```
{
    "name": "copper",
    "alpha": "3.75e-3",
    "T0":"293",
    "sigma0":"56.e+3",
    "k0":"0.4",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
},
"Glue0":
{
    "name": "glue0",
    "alpha":"0",
    "T0":"293",
    "sigma0":"0",
    "k0":"1.2e-3",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
},
"Glue1":
{
    "name": "glue1",
    "alpha":"0",
    "T0":"293",
```

```
"sigma0":"0",
    "k0":"1.2e-3",
    "sigma":"sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
    "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
}
    },
Conditions.
"BoundaryConditions":
"potential":
 "Dirichlet":
 {
"V1":
{
    "expr1":"9",
    "expr2":"Cu"
},
"VO":
{
    "expr1":"0",
    "expr2":"Cu"
}
 },
 "Neumann":
\verb|"Interface0_1":
    "expr":"0"
},
"Interface0_2":
{
    "expr":"0"
},
"Rint":
    "expr":"0"
},
"Rext":
    "expr":"0"
"iRint1":
```

```
"expr":"0"
},
"iRext1":
{
   "expr":"0"
},
"iRint2":
{
    "expr":"0"
},
"iRext2":
{
    "expr":"0"
}
 }
},
"temperature":
  "Robin":
 {
"Rint":
{
    "expr1":"85000.e-6",
    "expr2":"293"
},
"Rext":
{
    "expr1":"85000.e-6",
    "expr2":"293"
},
"iRint1":
{
    "expr1":"85000.e-6",
    "expr2":"293"
},
"iRext1":
{
    "expr1":"85000.e-6",
    "expr2":"293"
},
"iRint2":
{
    "expr1": "85000.e-6",
    "expr2":"293"
},
```

```
"iRext2":
    "expr1": "85000.e-6",
    "expr2":"293"
}
  },
  "Neumann":
"Interface0_1":
    "expr":"0"
},
"Interface0_2":
{
    "expr":"0"
},
"VO":
{
    "expr":"0"
},
"V1":
{
    "expr":"0"
}
  }
}
},
"PostProcess":
"Fields":["temperature", "potential", "current"]
}
}
```

Results. We can see the repartition of the temperature in the helix.

We can note that the peak temperature is inside the double helix.

2.3. Validity. In this example, we approximate the magnet with an axisymmetric copper torus. Thus we can consider only a quarter of this torus for our study. The torus is modeled thanks to the file geo, which also name the volume (omega) and each surface.

Equations. First of all, we start with the standard heat equation, with the heat from the Joule effect :

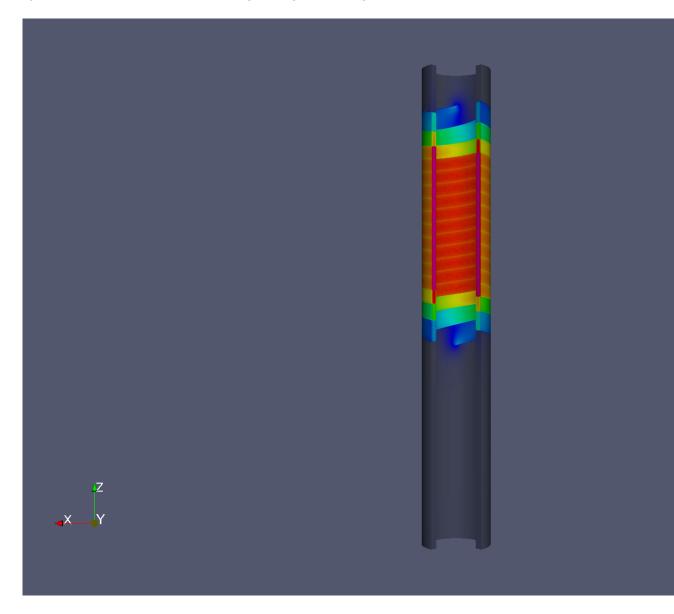


Figure 5. Temperature's repartition on paraview

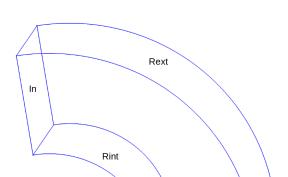
$$\rho C_p \frac{\partial T}{\partial t} - \nabla \cdot (k \nabla T) = \sigma (\frac{U}{2\pi r})^2$$

Coefficients  $\sigma$  and k are in fact, temperature-dependent, shown in this equations :

• 
$$\sigma = \frac{\sigma \ 0}{1 + \alpha (T - T_{ref})}$$

• 
$$k = k_0 \frac{T}{(1 + \alpha(T - T_{ref}))T_{ref}}$$

Volume omega



ν \<u>z</u> x

Out

FIGURE 6. Model of the quarter of torus on 'Gmsh'

But, for the next, we consider that  $\alpha = 0$  and  $T=T_{ref}$ . Thereby, we are in a linear problem that we can solve with the thermoelectric-linear model in the Json file.

In our case, we consider that T only depends on the radius, so  $\frac{\partial T}{\partial t} = 0$ . We can now consider this equation :

$$T = A\log(r) - \frac{\sigma}{2k}(\frac{U}{2\pi})^2\log^2(r) + B$$

The constants A and B are determined by the boundary conditions (Dirichlet and Robin).

Finally, we have this equation:

$$T = -a\log(\frac{r}{r_0})^2 + T_{max}$$

• 
$$T_{max} = \frac{2ak}{h_1r_1 + h_2r_2}\log(\frac{r_2}{r_1}) + \frac{h_1r_1T_{w1} + h_2r_2T_{w2}}{h_1r_1 + h_2r_2} + a\frac{h_1r_1log(\frac{r_1}{r_0})^2 + h_2r_2log(\frac{r_2}{r_0})^2}{h_1r_1 + h_2r_2}$$

$$\bullet \ r_0 = e^{\frac{T_{w2} - T_{w1}}{b} + \frac{ac}{b}}$$

• 
$$a = \frac{\sigma_0}{2k} (\frac{U}{2\pi})^2$$

• 
$$b = k(\frac{1}{h_1 r_1} + \frac{1}{h_2 r_2}) + \log(\frac{r_2}{r_1})$$

• 
$$c = \log(\frac{r_2}{r_1})\log(r_1r_2) + 2k(\frac{\log(r_1)}{h_1r_1} + \frac{\log(r_2)}{h_2r_2})$$

 $\mathbf{r}_0$  is the radius for which the temperature is at its maximum.

Parameters. In our case, we choose the parameters like this :

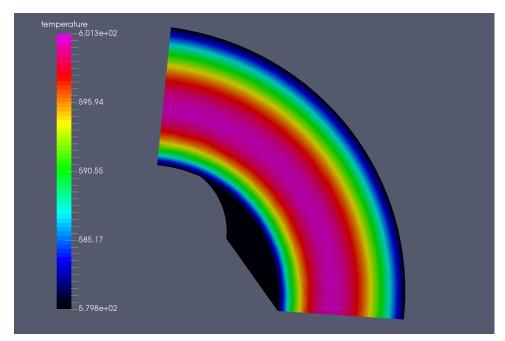
$\sigma_0$	electrical conductivity at $T_0$	$[52.10^6; 58.10^6]$	$58.10^{6}$	$S.m^{-1}$
k	thermal conductivity	[360;380]	380	$W.m^{-1}.K^{-1}$
$r_1$	internal radius	$1.10^{-3}$	$1.10^{-3}$	m
$r_2$	internal radius	$2.10^{-3}$	$2.10^{-3}$	m
$T_{w1}$	water cooling temperature on radius $r_1$	[293;310]	293	K
$T_{w2}$	water cooling temperature on radius $r_2$	[293;310]	293	K
$h_2$	heat transfer coefficient	[70000;90000]	80000	$W.m^{-2}.K^{-1}$
$h_1$	heat transfer coefficient	$h_1 = h_2 \frac{r_2}{r_1}$	$W.m^{-2}.K^{-1}$	'
$V_0$	electrical potential		0.3	V

Conditions. We set 2 conditions :

- Dirichlet for the potential :  $V_{in}=0$  V and  $V_{out}=V_0\frac{1}{4}=0.075V$  because we consider only one quarter of the torus.
- Robin for the temperature :

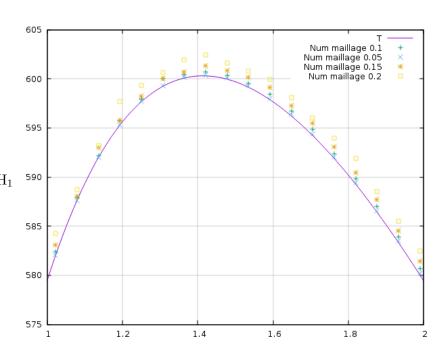
– On 
$$\mathbf{r}_{int}$$
:  $h_1 = 80000 \frac{r_2}{r_1} = 160000 W.m^{-2}.K^{-1}$  and  $\mathbf{T}_{w1} = 293~\mathrm{K}$ 

– On 
$$r_{ext}$$
 :  $h_2 = 80000 \ W.m^{-2}.K^{-1}$  and  $T_{w2} = 293 \ K$ 



Results. We can vary the degree of the finite element from 1 (linear) to 2 (quadratic).

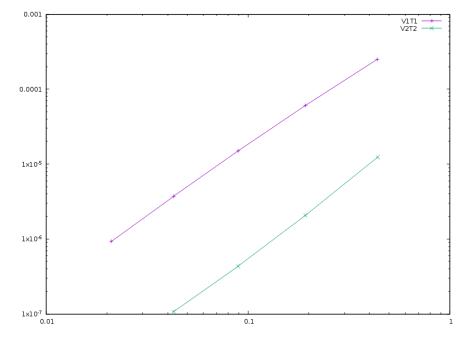
To prove the convergence towards the theory, we plot the difference between  $L_2/H_1$ and the theoretical formulas for T and V. The scale is logarithmic, to see directly the slope and note that it is directly linked to the degree of the finite element used.

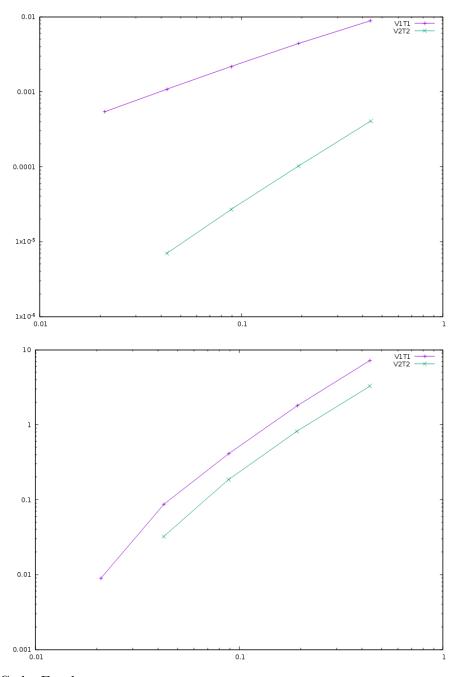


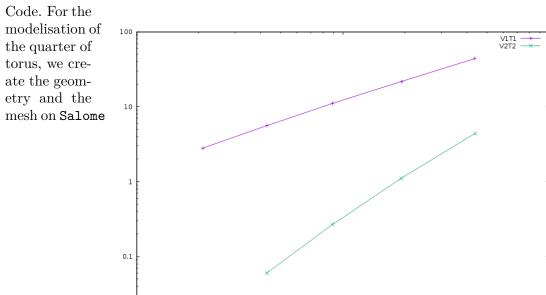
$$T = -a\log(\frac{r}{r_0})^2 + T_{max}$$

$$V = \frac{0.3*atan2(x,y)}{2\pi}$$

- $\bullet$  For L2, we have directly the output on the table obtained whether for T or for V
- $\bullet$  For H1, we need to do a quadratical mean between the H1 and L2 of the table  $(\sqrt{H1^2+L2^2})$  for T and V







```
and export the
file in .geo
// Define Main params
Unit = 1.e-3;
lc = 1*Unit;
lc_ext = 3*lc;
lc_inf = 1*lc;
h=0.2;
r1=1;
r2=2;
L=2*r2;
Mesh.ElementOrder = 1;
Point(1) = \{0, 0, -L, h\};
Point(2) = \{r1, 0, -L, h\};
Point(3) = \{r2, 0, -L, h\};
Point(4) = \{0, r1, -L, h\};
Point(5) = \{0, r2, -L, h\};
Circle(1) = \{2, 1, 4\};
Circle(2) = \{3, 1, 5\};
Line(3) = \{4, 5\};
Line(4) = \{2, 3\};
Line Loop(5) = \{3, -2, -4, 1\};
Plane Surface(1) = {5};
out[] = Extrude {0,0,L} {Surface{1};};
Physical Volume("omega") = {out[1]};
Physical Surface("top") = {out[0]};
Physical Surface("bottom") = {1};
Physical Surface("Rint") = {out[5]};
Physical Surface("Rext") = {out[3]};
Physical Surface("in") = {out[2]};
Physical Surface("out") = {out[4]};
Next step is
to create a
file.json which
define the model
we will use,
the material
```

```
and sets the
conditions.
{
    "Name": "ThermoElectric",
    "ShortName": "TE",
    "Model": "thermoelectric-linear",
    "Materials":
        "omega":
        {
            "name": "copper",
            "alpha": "3.35e-3",
            "T0":"293",
            "sigma0":"58e+3",
            "k0":"0.38",
            "sigma": "sigma0/(1+alpha*(T-T0)):sigma0:alpha:T:T0",
            "k":"k0*T/((1+alpha*(T-T0))*T0):k0:T:alpha:T0"
        }
    },
    "BoundaryConditions":
        "potential":
            "Dirichlet":
                "in":
                     "expr1":"0",
    "expr2":"omega"
                "out":
                     "expr1": "0.3/4.", //since we only consider a quarter of the torus
    "expr2":"omega"
        },
        "temperature":
            "Robin":
                 "Rext":
                 {
                     "expr1":"0.08",
                     "expr2":"293"
```

```
},
                 "Rint":
                  }
                      "expr1":"0.08*(2./1.)",
                      "expr2":"293"
             }
         }
    },
    "PostProcess":
    {
         "Fields":["temperature", "potential", "current"]
    }
}
Lastly, we cre-
ate a file.cfg
to configure what we will calculate.
\begin{array}{l} thermoelectric\_3D\_V1T1\_N1\_cvg.cfg \ (for \ the \ T1V1 \ model) \end{array}
geofile=quarter-turn3D.geo
geofile-path=$cfgdir
gmsh.hsize=0.2
conductor_volume=omega
[convergence]
max_iter=5
[functions]
#V_exact
f=0.3*atan2(x,y)/(2*pi):x:y:z
#T_exact
g=600.312-58.e+3/(2*0.38)*(0.3/(2*pi))^2*log(sqrt(x*x+y*y)/sqrt(1*2))^2:x:y:z
[thermoelectric]
model_json=$cfgdir/quarter-turn3D.json
weakdir=false
[electro]
pc-type=gamg
#ksp-monitor=true
ksp-rtol=1e-7
ksp-atol=1e-5
ksp-maxit=2000
```

ksp-use-initial-guess-nonzero=1

```
[thermal]
pc-type=gamg
#ksp-monitor=true
ksp-rtol=1e-8
ksp-atol=1e-6
ksp-use-initial-guess-nonzero=1
```

Finally, to execute our program, run this command:

```
to study the convergence mpirun -np 4 feelpp_test_convergence_3D_V1T1_N1 --config-file thermoelectric_3D_V1T1_N1_cvg
```

It will create a table with all the informations you need. For our example, we showed the convergence using L2 and H1 (in section Results).

```
to apply for a real case (theory not known) mpirun -np 4 feelpp_hfm_thermoelectric_model_3D_V1T1_N1 --config-file thermoelectric_3D_V1T
```

This command will create files in /feel/hifimagnet/ThermoElectricModel/... You can see the results with Paraview or Ensight opening Thermics.case or Electrics.case in the software of your choice.

# 3. Magnetostatic

3.1. **General presentation of the files.** There are several applications we can use, but all this applications follow the same form:

The kind of your test can be:

**biotsavart:** calculate **B** and **A** analytically and with BiotSavart, depending of **j** given.

biotsavart\_num: same than biotsavart but use the thermoelectric program to calculate j.

You either set the dimension (1, 2 or 3D) of the box inside the conductor.

3.1.1. *Material*. There is no json file for the magnetostatic studies (except for the biot-savart\_num which use the termoelectric program). Instead, the configuration file (.cfg) use a .d file, which define the geometry and the material of the conductor, like this:

```
#Power[MW] Current[A]
  "_"
#Helices N_Elem
"_"
```

```
\#N R1[m] R2[m] HalfL[m] Rho[Ohm.m] Alpha[1/K] E_Max[Pa] K[W/(m.K)] h[W/(m^2.K)]
<T_Water>[°C] T_Max[°C]
# Bitter I=j1*r1*log(r2/r1)*2*L=11767.657994358965
#Type R1
                R2 Z1 Z2 J
                                                        Rho
                                                                 N turns
                       11 11
                                11 11 11 11
                11 11
                                                        11 11
                                                                 11 11
# Supra
#Bz(0)[tesla] Power[MW] Bz_total(0)[tesla]
         н н н
   BO_H[t] Sum_BO[t] Power_H[MW] Sum_Power[MW]
#H
MARGE DE SECURITE CONTRAINTES= 8.0 %
As there is different applications, there is different configure files, the biotsavart
dim=3
units=mm
geofile="name_of_the_file_created_by_the_partition.json" (.msh or .geo)
geofile-path=$cfgdir
gmsh.hsize="_"
conductor_volume="name_of_your_volume"
[convergence]
max_iter="_"
[functions]
#theoretical function of j
j=\{-58.e+3*(0.5/(2*Pi))*y/(x^2+y^2),58.e+3*(0.5/(2*Pi))*x/(x^2+y^2),0\}:x:y:z
[biot_savart]
conductor="name_of_your_volume"
box=box"dimension_of_your_box(1,2_or_3D)"
[magnetic_field-bmap]
geo-data="name_of_your_file.d"
geo-path=$cfgdir
helix-intensity="_"
bitter-intensity="_"
supra-intensity="_"
```

```
dim=3
units=mm
geofile=torus3D.geo
geofile-path=$cfgdir
gmsh.hsize=10
conductor_volume=omega
[convergence]
max_iter=1
[functions]
j=\{-58.e+3*(0.5/(2*Pi))*y/(x^2+y^2),58.e+3*(0.5/(2*Pi))*x/(x^2+y^2),0\}:x:y:z
[biot_savart]
conductor=omega
box=box3D
[magnetic_field-bmap]
geo-data=torus3D.d
geo-path=$cfgdir
helix-intensity=0
bitter-intensity=11767.7
supra-intensity=0
[thermoelectric]
model_json=$cfgdir/biotsavart.json
weakdir=false
[electro]
pc-type=gamg
#ksp-monitor=true
ksp-rtol=1e-7
ksp-atol=1e-5
ksp-maxit=2000
ksp-use-initial-guess-nonzero=1
[thermal]
pc-type=gamg
#ksp-monitor=true
ksp-rtol=1e-8
ksp-atol=1e-6
ksp-use-initial-guess-nonzero=1
```

# 3.2. Examples.

3.3. Validity. In this example, we approximate the magnet with an axisymmetric copper torus. This torus is modeled with a file .geo, which name the volume of the torus (omega), each surface and a volume inside the torus (box3D), also known as a sphere.

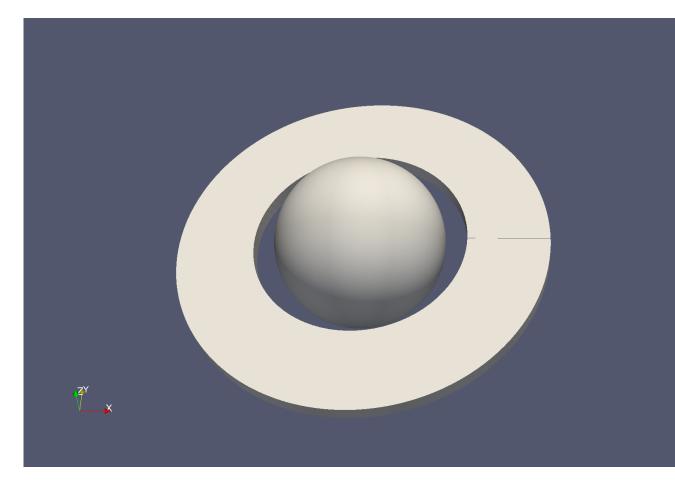


FIGURE 7. Model of the torus with the 3D box

Equations. We use the Biot & Savart's law to prove the validity. First, we need to set the conditions to use this law,  $\Omega_{cond} \cap \Omega_{mgn} = \emptyset$ , with  $\Omega_{cond}$  the conductor in which the current is passing and  $\Omega_{mgn}$  the domain in which we want to know the magnetic field. With this, using the magnetostatic equation, we can write the magnetic potential  $\mathbf{A}$  as

$$\nabla^2 \mathbf{A} = -\mu \mathbf{j}$$

The general solution of this equation is:

$$A(\mathbf{r}) = -\mu \int_{\Omega_{cond}} G(\mathbf{r}, \mathbf{r'}) \mathbf{j}(\mathbf{r'})$$

with  $G(\mathbf{r}, \mathbf{r}')$  the Green's function defined as:

$$G(\mathbf{r}, \mathbf{r'}) = \frac{-1}{4\pi} \frac{1}{|\mathbf{r} - \mathbf{r'}|}$$

with  $\mathbf{r} \in \Omega_{mqn}$  and  $\mathbf{r'} \in \Omega_{cond}$ 

Then we rewrite the expression of the magnetic potential A:

$$A(\mathbf{r}) = \frac{\mu_0}{4\pi} \int_{\Omega_{cond}} \frac{\mathbf{j}(\mathbf{r'})}{|\mathbf{r} - \mathbf{r'}|} d\mathbf{r'} \qquad \mathbf{r} \in \Omega_{mgn}$$

As  $\mathbf{B}$  is defined as the curl of  $\mathbf{A}$ , we can write the so called Biot & Savart's law:

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int_{\Omega_{cond}} \frac{\mathbf{j}(\mathbf{r'}) \times (\mathbf{r} - \mathbf{r'})}{|\mathbf{r} - \mathbf{r'}|^3} d\mathbf{r'} \qquad \mathbf{r} \in \Omega_{mgn}$$

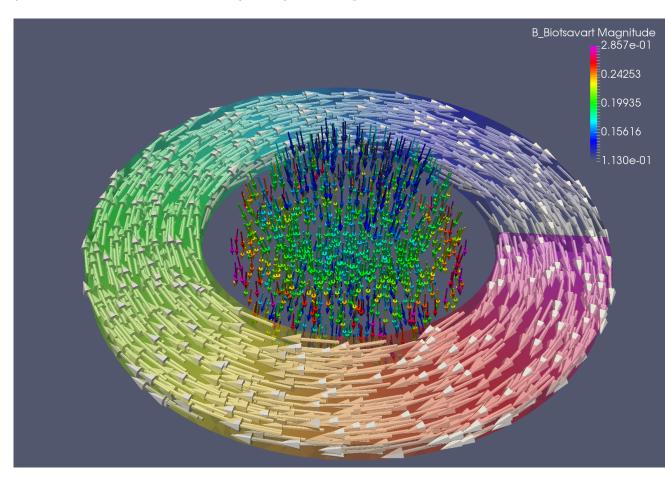
Parameters. We use here the same parameters as in the ThermoElectric section for the .json file. We select the biter-intensity (shown in biotsavart\_num\_box3D\_3D\_V1T1\_N1\_cvgcfg) We also need a file .d to define the characteristics of the conductors.

Results. On the next figure, the current is represented by the white arrow in the torus, the color of the torus corresponding to the potential. The box3D (the sphere inside the torus) is represented with the magnetic induction's vector (**B**).

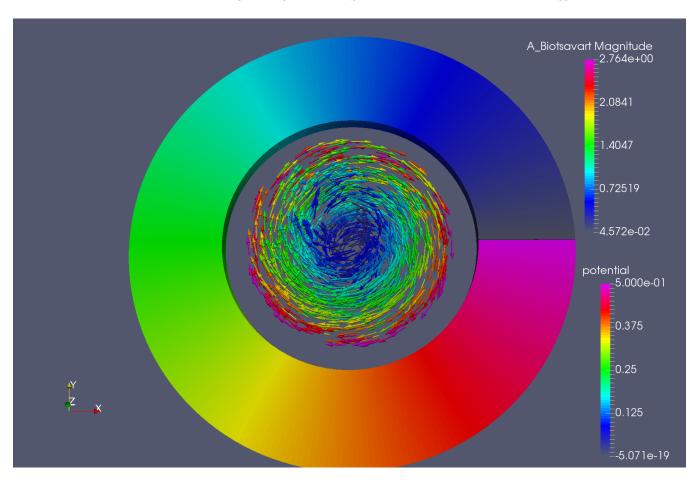
We can see that the magnetic potential  $(\mathbf{A})$  is in the same way that the current in the torus. The scale and the direction of the magnetic field  $(\mathbf{B})$  are coherent.

Code. For the modelisation of the torus and his box inside, we create the geometry with this file .geo

```
torus3D_box3D.geo
// Define Main params
Unit = 1.e-3;
h=5;
r1=61.2*0.5;
r2=106.4*0.5;
L=4.61/2.;
eps=0.1;
theta1=Asin( eps/(2*r1) );
theta2=Asin( eps/(2*r2) );
```



```
// 1st quarter
Point(1) = \{0, 0, -L, h\};
Point(2) = \{r1*Cos(theta1), eps/2., -L, h\};
Point(3) = \{r2*Cos(theta2), eps/2., -L, h\};
Point(4) = \{0, r1, -L, h\};
Point(5) = \{0, r2, -L, h\};
Point(6) = \{-r1, 0, -L, h\};
Point(7) = \{-r2, 0, -L, h\};
Point(8) = \{0, -r1, -L, h\};
Point(9) = \{0, -r2, -L, h\};
Point(10) = \{r1*Cos(-theta1), -eps/2., -L, h\};
Point(11) = \{r2*Cos(-theta2), -eps/2., -L, h\};
Circle(1) = \{2, 1, 4\};
Circle(2) = \{4, 1, 6\};
Circle(3) = \{6, 1, 8\};
Circle(4) = \{8, 1, 10\};
```



```
Circle(5) = {3, 1, 5};
Circle(6) = {5, 1, 7};
Circle(7) = {7, 1, 9};
Circle(8) = {9, 1, 11};

Line(9) = {2, 3};
Line(10) = {10, 11};

dL=newl; Line Loop(dL) = {1:4, 10, -8, -7, -6, -5, -9};
S=news; Plane Surface(S) = {dL};

out[] = Extrude {0,0,2*L} {Surface{S};};

Physical Volume("omega") = {out[1]};
Physical Surface("top") = {out[0]};
Physical Surface("bottom") = {S};
Physical Surface("Rint") = {out[2], out[3], out[4], out[5]};
```

```
Physical Surface("Rext") = {out[7], out[8], out[9], out[10]};
Physical Surface("in") = {out[6]};
Physical Surface("out") = {out[11]};
// Define BiotSavart box
Boxdim=3;
hs=1;
np=10;
z0=-0.8*r1;
z1 = -z0;
C0=newp; Point(C0) = \{0, 0, 0, hs\};
P0=newp; Point(P0) = \{0, 0, z0, hs\};
P1=newp; Point(P1) = \{0, 0, z1, hs\};
Q0=newp; Point(Q0) = \{0, z1, 0, hs\};
R0=newp; Point(R0) = \{z1, 0, 0, hs\};
COPO=newl; Line(COPO) = \{CO, PO\};
POP1=newl; Line(POP1) = {P0, P1};
BSO=newl; Circle(BSO) = {PO, CO, QO};
BS1=newl; Circle(BS1) = {PO, CO, RO};
BS2=newl; Circle(BS2) = {Q0, C0, R0};
BS3=newl; Circle(BS3) = {Q0, C0, P1};
BS4=newl; Circle(BS4) = {RO, CO, P1};
Sb_Sph=newl; Line Loop(Sb_Sph)={BS0, BS2, -BS1};
S_Sph=newl; Ruled Surface(S_Sph)={Sb_Sph};
S2_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, Pi/2. } { Duplicata{ Surface{S_Sph}; } };
S3_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, 2*Pi/2. } { Duplicata{ Surface{S_Sph}; }
S4_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, 3*Pi/2. } { Duplicata{ Surface{S_Sph}; }
Nb_Sph=newl; Line Loop(Nb_Sph)={BS2, BS4, -BS3};
N_Sph=newl; Ruled Surface(N_Sph)={Nb_Sph};
N2_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, Pi/2. } { Duplicata{ Surface{N_Sph}; } };
N3_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, 2*Pi/2. } { Duplicata{ Surface{N_Sph}; } }
N4_Sph = Rotate { { 0, 0, 1 }, { 0, 0, 0 }, 3*Pi/2. } { Duplicata{ Surface{N_Sph}; }
SLoop=news; Surface Looop(SLoop)={S_Sph, N_Sph, S2_Sph, N2_Sph, S3_Sph, N3_Sph, S4_Sph
RMN=newv; Volume(RMN)={SLoop};
If (Boxdim == 1)
Physical Line("box1D") = {POP1};
```

```
{\tt EndIf}
```

```
If ( Boxdim == 2 )
Physical Surface("box2D") = {S_Sph, N_Sph, S2_Sph, N2_Sph, S3_Sph, N3_Sph, S4_Sph, N4_Sph}
EndIf
If ( Boxdim == 3 )
 Physical Volume("box3D") = {RMN};
EndIf
the next step is to make a file .d which fix some parameters on the torus
torus3D.d
#Power[MW] Current[A]
12.5
       31000.
#Helices N_Elem
#N R1[m] R2[m] HalfL[m] Rho[Ohm.m] Alpha[1/K] E_{Max}[Pa] K[W/(m.K)] h[W/(m^2.K)]
<T_Water>[°C] T_Max[°C]
0.
# Bitter I=j1*r1*log(r2/r1)*2*L=11767.657994358965
                         Z1
                                     Z2
#Type R1
                R2
                                                  J
                                                                  Rho N_turns
        30.6e-3
                   53.2e-3
                              -2.305e-3
                                           2.305e-3
                                                       150833116.00212305
                                                                              1
                                                                                         1
                    53.2e-3
                             -2.305e-3
         30.6e-3
                                          2.305e-3
                                                        124827406.34658459 1
#1
                                                                                        1
# Supra
#Bz(0)[tesla] Power[MW] Bz_total(0)[tesla]
 22.7526804266798 12.500000000 22.7526804266798
     BO_H[t] Sum_BO[t] Power_H[MW] Sum_Power[MW]
#H
MARGE DE SECURITE CONTRAINTES= 8.0 %
Finally we can use the biotsavart_num (in which we use the thermoelectric model to
calculate j) or the biotsavart (in which we specify manually j)
dim=3
units=mm
geofile=biotsavart_box3D.geo
geofile-path=$cfgdir
gmsh.hsize=10
conductor_volume=omega
```

```
[convergence]
max_iter=1
[functions]
j=\{-58.e+3*(0.5/(2*Pi))*y/(x^2+y^2),58.e+3*(0.5/(2*Pi))*x/(x^2+y^2),0\}:x:y:z
[biot_savart]
conductor=omega
box=box3D
[magnetic_field-bmap]
geo-data=torus3D.d
geo-path=$cfgdir
helix-intensity=0
bitter-intensity=11767.7
supra-intensity=0
dim=3
units=mm
geofile=torus3D_box3D.geo
geofile-path=$cfgdir
gmsh.hsize=10
conductor_volume=omega
 [convergence]
max_iter=3
[functions]
j=\{58.e+3*(0.5/(2*Pi))*y/(x^2+y^2),-58.e+3*(0.5/(2*Pi))*x/(x^2+y^2),0\}:x:y:z
 u=0.5*atan2(y,x)/(2*Pi)*(atan2(y,x)>0)+(0.5*(atan2(y,x)+2*Pi)/(2*Pi))*(atan2(y,x)<0): (atan2(y,x)+2*Pi)/(2*Pi))*(atan2(y,x)<0): (atan2(y,x)+2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/(2*Pi)/
t=362.156146169164-58.e+3/(2*0.38)*(0.5/(2*Pi))^2*log(sqrt(x*x+y*y)/39.4354779237947
[biot_savart]
conductor=omega
box=box3D
[magnetic_field-bmap]
geo-data=torus3D.d
geo-path=$cfgdir
helix-intensity=0
bitter-intensity=-11767.7
supra-intensity=0
[thermoelectric]
```

model\_json=\$cfgdir/biotsavart.json

weakdir=false

# [electro] pc-type=gamg #ksp-monitor=true ksp-rtol=1e-7 ksp-atol=1e-5 ksp-maxit=2000 ksp-use-initial-guess-nonzero=1 [thermal] pc-type=gamg #ksp-monitor=true ksp-rtol=1e-8 ksp-atol=1e-6 ksp-use-initial-guess-nonzero=1 For the numerical file, we use a json file like in the thermoelectric section. { "Name": "ThermoElectric", "ShortName": "TE", "Model": "thermoelectric-linear", "Materials": { "omega": { "name": "copper", "alpha":"3.35e-3", "T0":"293", "sigma0":"58e+3", "k0":"0.38", "sigma":"sigma0/(1+alpha\*(T-T0)):sigma0:alpha:T:T0", "k":"k0\*T/((1+alpha\*(T-T0))\*T0):k0:T:alpha:T0" } }, "BoundaryConditions": "potential": "Dirichlet": "in": "expr1":"0.5", "expr2":"omega"

```
},
                  "out":
                      "expr1":"0",
    "expr2": "omega"
             }
         },
         "temperature":
             "Robin":
             {
                  "Rext":
                  {
                      "expr1":"0.08",
                      "expr2":"293"
                 },
                  "Rint":
                  {
                      "expr1":"0.08",
                      "expr2":"293"
                 }
             }
        }
    },
    "PostProcess":
    {
         "Fields": ["temperature", "potential", "current"]
    }
}
```

3.4. Cartesian model. Until now, we only focus on the study in the center of the torus. Indeed, the program can only calculate the magnetic potential and field outside the torus, therefore to be able to calculate  $\bf A$  and  $\bf B$  around the torus that implies to create a really complex geometry. Instead of this, we use the Cartesian model which allow to calculate  $\bf A$  and  $\bf B$  inside the torus, this allow us to calculate inside the torus, then we model a large sphere to represent the air all around the conductor.

With a simple torus, we can calculate the potential analytically like before.

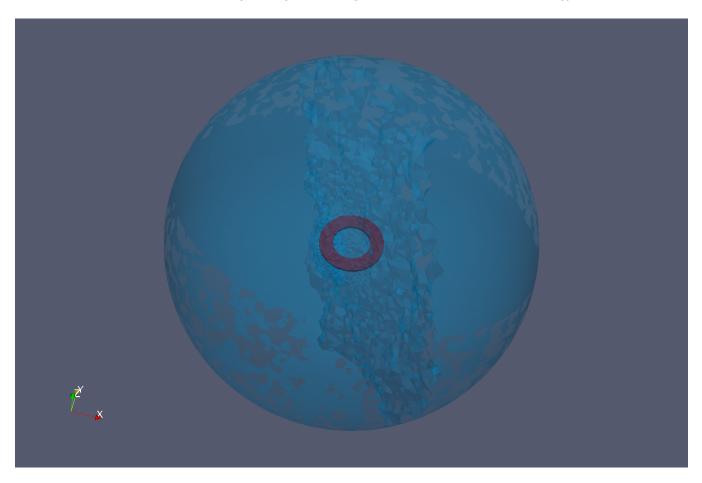
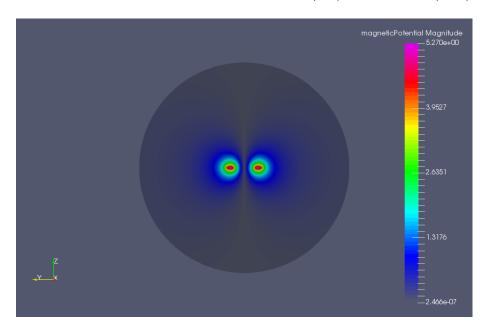
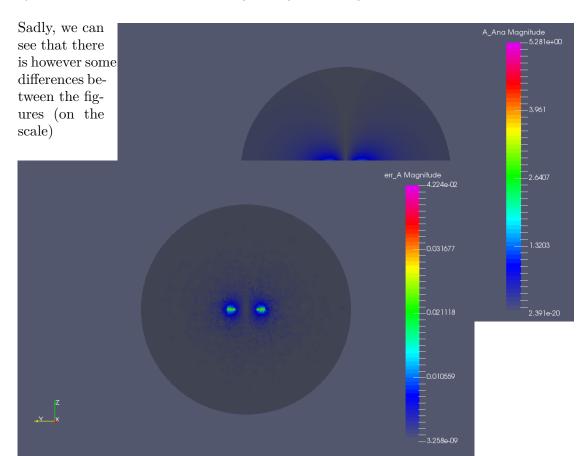


FIGURE 8. model of the torus (red) with the box (blue).

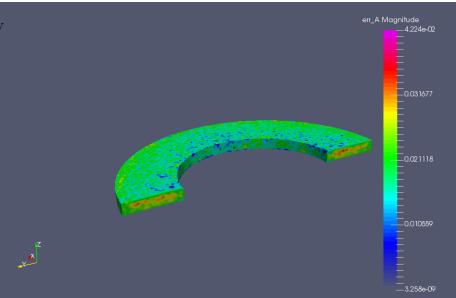




The maximum of the error is approximately 0.8

# 4. Elasticity

# 4.1. General presentation of the files. In the continuity of the study of the magnet, we use the file already used in Thermoelectric and Magnetostatic. We create a



```
new file .json
to set up the
conditions for
elasticity.
{
    "Name": "CoupledCart",
    "ShortName": "MSC",
    "Model": "Elasticity",
    "Materials":
    {
        "name_of_the_volume":
             "name": "material",
             "E": "Young_modulus",
             "nu": "Poisson's_ratio",
             "alphaT": "linear_dilatation_coefficient",
             "rho": "density"
        }
    },
    "BoundaryConditions":
        "condition(like displacement_x y or z)":
        {
             "type_of_condition (Dirichlet, Neumann or Robin)":
                 "surface_concerned":
                 {
                     "expr":"_"
                 },
            }
        },
        "other_condition":
        },
    },
    "PostProcess":
        "Fields":["displacement","Von-Mises","tresca","principal-stresses"]
    }
}
```

We also need to set up the configuration file (.cfg) to compute the elasticity part, which consist in adding the elasticity section.

[elasticity]
filename=\$cfgdir/quarter-torus3D-elasticity.json
on.type=elimination\_symmetric
thermal\_dilatation=false
do\_export\_all=true

# # precondtioner config
pc-type=gamg #lu,gasm,ml
ksp-monitor=true
# ksp-converged-reason=1

## 4.2. Examples.

### 4.3. Validity.

Conditions. We consider the conductor as a solenoid with finite thickness and infinite length. This allow us to ignore the z components in our equations. We admit that there is only a radial expansion.

Equations. Taking back the equations in the Maths for Hifimagnet, we consider:

$$div\sigma + \mathbf{j} \times \mathbf{b} = 0$$

With the conditions set in the previous section, we have:

$$-\sigma_{\theta} + \frac{\partial}{\partial r}(r\sigma_r) = -rj_{\theta}b_z$$

Parameters. In our case (a coil of copper), we choose the parameters like this:

E	Young modulus	$[124.10^9; 128.10^9]$	$128.10^9$	$Pa = kg.m^{-1}.s^{-2}$
$\nu$	Poisson's ratio	0.33	0.33	-
$\alpha_T$	linear dilatation coefficient	$[16, 6.10^{-6}; 18.10^{-6}]$	$18.10^{-6}$	$K^{-1}$
$\rho$	density	[8920; 8960]	8950	$kg.m^{-3}$

Results. As we can see in this coarse mesh, the scale is coherent (he unit being in Pa).

