

Pauli matrices

$$\sigma_1 = \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_2 = \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_3 = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\sigma_1^2 = \sigma_2^2 = \sigma_3^2 = -i \sigma_1 \sigma_2 \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = I$$

Matrix exponentiating

$$e^{ia(\hat{n} \cdot \vec{\sigma})} = I \cos a + i(\hat{n} \cdot \vec{\sigma}) \sin a$$

$$e^{i\theta \hat{A}} = \cos \theta \hat{I} + i \sin \theta \hat{A}$$

$$\hat{H} = -\mu \mathbf{B} \cdot \hat{\mathbf{S}}, \quad i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H} |\psi(t)\rangle, \quad |\psi(t)\rangle = \exp \left[-\frac{i}{\hbar} \hat{H} t \right] |\psi(0)\rangle \equiv \hat{U}(t) |\psi(0)\rangle,$$

$$\mathbf{B} = B_0 \hat{\mathbf{x}} \quad \hat{U}(t) = \exp \left[\frac{i\omega_0 t}{2} \hat{\sigma}^x \right], \quad \omega_0 \equiv \mu B_0.$$

$$\begin{aligned} \hat{U}(t) &= \hat{I} \left(1 - \frac{(\omega_0 t/2)^2}{2!} + \frac{(\omega_0 t/2)^4}{4!} + \dots \right) \\ &\quad + i\hat{\sigma}^x \left(\omega_0 t/2 - \frac{(\omega_0 t/2)^3}{3!} + \frac{(\omega_0 t/2)^5}{5!} + \dots \right) \\ &= \cos \frac{\omega_0 t}{2} \hat{I} + i \sin \frac{\omega_0 t}{2} \hat{\sigma}^x \\ &\doteq \begin{pmatrix} \cos \frac{\tau}{2} & i \sin \frac{\tau}{2} \\ i \sin \frac{\tau}{2} & \cos \frac{\tau}{2} \end{pmatrix}, \quad \tau \equiv \omega_0 t. \end{aligned}$$

$$\hat{U}(\theta, \phi, \lambda) \doteq \begin{pmatrix} \cos \frac{\theta}{2} & -e^{i\lambda} \sin \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} & e^{i(\phi+\lambda)} \cos \frac{\theta}{2} \end{pmatrix}$$

$$u(\theta, \phi, \lambda)$$

$$\hat{U} = \hat{U} \left(\tau, +\frac{\pi}{2}, -\frac{\pi}{2} \right)$$

theoretical predictions of expectation values of the spin operators

$$\langle \hat{S}^\alpha \rangle \equiv \langle \psi | \hat{S}^\alpha | \psi \rangle$$

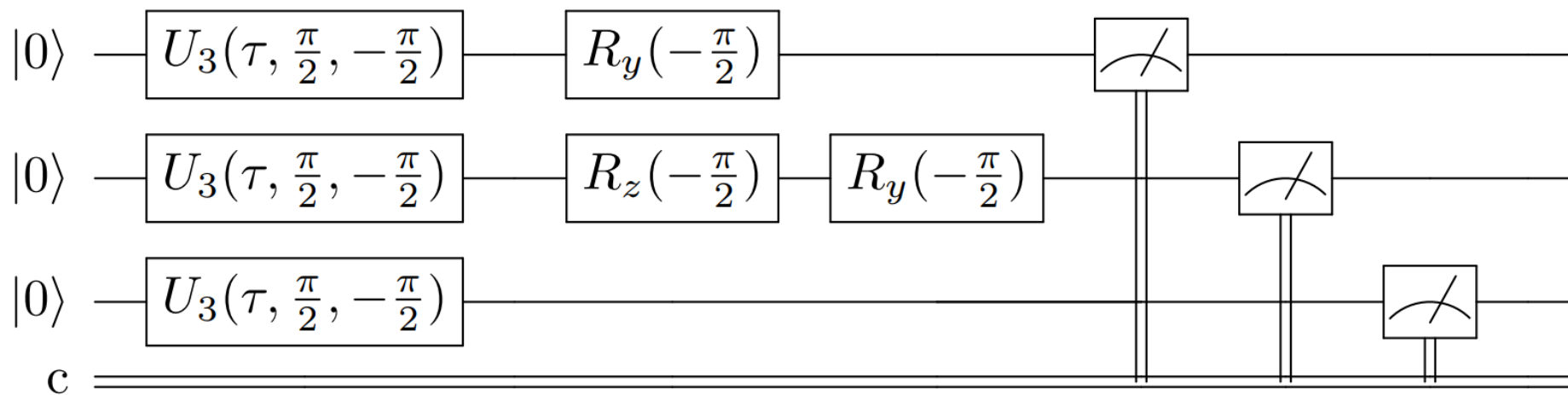
$$\langle \hat{S}^x(t) \rangle = 0$$

$$\langle \hat{S}^y(t) \rangle = \frac{\hbar}{2} \sin \omega_0 t$$

$$\langle \hat{S}^z(t) \rangle = \frac{\hbar}{2} \cos \omega_0 t.$$

$$\hat{H} = -\mu \mathbf{B} \cdot \hat{\mathbf{S}},$$

$$\mathbf{B} = B_0 \hat{\mathbf{x}}$$



$$\hat{H} = -\mu \mathbf{B} \cdot \hat{\mathbf{S}}, \qquad \hat{H} \rightarrow -\frac{\mu B_0 \hbar}{2} \sigma^y \doteq -\hbar \omega \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \qquad e^{-i\hat{H}t/\hbar} = e^{i\frac{\omega t}{2}\sigma^y}.$$

$$\mathbf{B} = B_0 \hat{y}$$

$$e^{i\omega t/2\sigma^y} = \hat{I} \cos(\omega t/2) + i\sigma^y \sin(\omega t/2) \doteq \begin{pmatrix} \cos(\omega t/2) & \sin(\omega t/2) \\ -\sin(\omega t/2) & \cos(\omega t/2) \end{pmatrix}$$

$$\hat{U}(\theta, \phi, \lambda) \doteq \begin{pmatrix} \cos \frac{\theta}{2} & -e^{i\lambda} \sin \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} & e^{i(\lambda+\phi)} \cos \frac{\theta}{2} \end{pmatrix} \qquad e^{-i\hat{H}t/\hbar} = \hat{U}(\omega t, \pi, \pi),$$

$$\hat{H} = \hat{A} + \hat{B},$$

$$[\hat{A}, \hat{B}] \neq 0.$$

$$\mathbf{B} = \frac{B_0}{\sqrt{2}} (\hat{\mathbf{x}} + \hat{\mathbf{y}})$$

$$\hat{H} = \frac{\omega_0}{\sqrt{2}} [\hat{\sigma}^x + \hat{\sigma}^y].$$

$$\hat{A} = \frac{\omega_0}{\sqrt{2}} \hat{\sigma}^x \text{ and } \hat{B} = \frac{\omega_0}{\sqrt{2}} \hat{\sigma}^y,$$

$$[\hat{A}, \hat{B}] = \frac{\omega_0^2}{2} [\hat{\sigma}^x, \hat{\sigma}^y] = i \frac{\omega_0^2}{2} \hat{\sigma}^z \neq 0.$$

$$\hat{U}(t) = \exp [-i\hat{H}t/\hbar] = \exp [-i\hat{A}t/\hbar - i\hat{B}t/\hbar].$$

$$\hat{U}_A(t) = \exp [-i\hat{A}t/\hbar], \quad \hat{U}_B(t) = \exp [-i\hat{B}t/\hbar]$$

As \hat{A} and \hat{B} do not commute, $\hat{U}(t) \neq \hat{U}_A(t)\hat{U}_B(t)$.

$$\begin{aligned} & \exp [i(\hat{A} + \hat{B})t/\hbar] \\ &= \lim_{n \rightarrow \infty} \left(\exp [-i\hat{A}t/n\hbar] \exp [-i\hat{B}t/n\hbar] \right)^n, \end{aligned}$$

or

$$\hat{U}(t) = \lim_{n \rightarrow \infty} \left(\hat{U}_A \left(\frac{t}{n} \right) \hat{U}_B \left(\frac{t}{n} \right) \right)^n.$$

Baker–Campbell–Hausdorff formula

$$Z = X + Y + \frac{1}{2}[X, Y] + \frac{1}{12}[X, [X, Y]] - \frac{1}{12}[Y, [X, Y]] + \cdots,$$

$$e^X e^Y = e^Z$$

Zassenhaus formula

$$e^{t(X+Y)} = e^{tX} e^{tY} e^{-\frac{t^2}{2}[X,Y]} e^{\frac{t^3}{6}(2[Y,[X,Y]]+[X,[X,Y]])} e^{\frac{-t^4}{24}([[[X,Y],X],X]+3[[[X,Y],X],Y]+3[[[X,Y],Y],Y])} \dots$$

As a corollary of this, the Suzuki–Trotter decomposition follows.

$$e^{A+B} = \lim_{n \rightarrow \infty} (e^{A/n} e^{B/n})^n.$$

Lie product formula

A first order formula consists of the approximation stated in the introduction, where the matrix exponential of a sum is approximated by a product of matrix exponentials:

$$e^{A+B} \approx e^A e^B$$

There exists a second-order formula, called the Suzuki-Trotter decomposition

$$e^{A+B} \approx e^{B/2} e^A e^{B/2}$$

$$e^{-it(XX+ZZ)} = e^{-it/2ZZ} e^{-itXX} e^{-it/2ZZ} + \mathcal{O}(t^3).$$

By means of recursions, higher-order approximations can be found

If all Hamiltonian terms commute, the task of simulating this Hamiltonian is reduced to implementing individually.

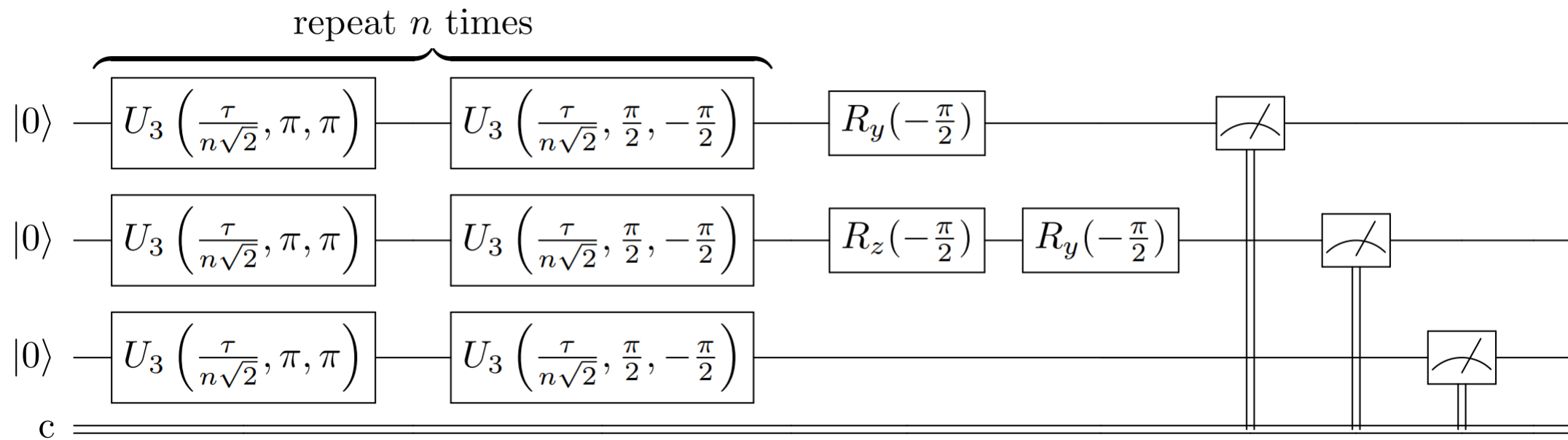
$$\exp \left[i \frac{\omega_0 t}{2\sqrt{2}} (\hat{\sigma}^x + \hat{\sigma}^y) t \right] = \hat{U}_3 \left(\tau, \frac{3\pi}{4}, -\frac{3\pi}{4} \right).$$

Theoretical spin expectation values (Time dependant)

$$\langle \hat{S}^x(t) \rangle = -\frac{\hbar}{2\sqrt{2}} \sin \omega_0 t,$$

$$\langle \hat{S}^y(t) \rangle = \frac{\hbar}{2\sqrt{2}} \sin \omega_0 t,$$

$$\langle \hat{S}^z(t) \rangle = \frac{\hbar}{2} \cos \omega_0 t.$$



one executes a circuit for some large number of shots

In Qiskit, the counts associated with executing a circuit over some number of shots

$$\langle \hat{S}^y \rangle = \frac{1}{N_{\text{shots}}} \left[(n_{000} + n_{001} + n_{100} + n_{101}) \left(+\frac{\hbar}{2} \right) + (n_{010} + n_{011} + n_{110} + n_{111}) \left(-\frac{\hbar}{2} \right) \right]$$

$$\langle \hat{S}^x \rangle = \frac{\hbar}{2} \frac{c_{000} + c_{001} + c_{010} + c_{011} - c_{100} - c_{101} - c_{110} - c_{111}}{N_{\text{shots}}}.$$