

Quantum Many-Body Simulations of Double Dot System

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Introduction

Quantum-Dot Model

- Schrödinger equation: $\mathcal{H}|\psi\rangle = E|\psi\rangle$, $\mathcal{H} = -\sum_i \frac{\nabla_i^2}{2} + f(\mathbf{r}) + V(R, \mathbf{r})$

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Quantum-Dot Model

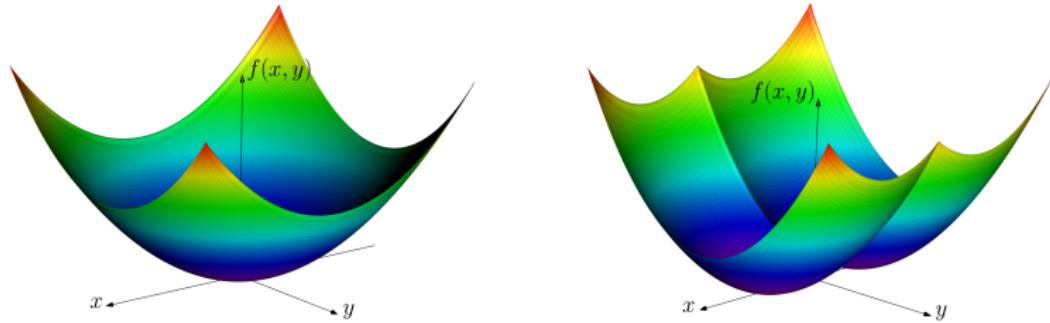
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- Confinement: Harmonic Oscillator¹, Double-Well²
 $V(\mathbf{r}) = \frac{1}{2}m\omega^2 r^2$ $V(\mathbf{r}) = \frac{1}{2}m\omega^2(r^2 - \delta R|x| + R^2)$



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Methods

Methods

**Hartree-Fock
Variational Monte-Carlo**

Methods: Variational Principle

$$E_0 \leq \frac{\langle \Psi | \mathcal{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

Slater Determinant and Energy Functional

Methods: Slater Determinant and Energy Functional

- Pauli Principle

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$$\bullet E[\Psi] = \frac{\langle \Psi | \mathcal{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \sum_p \langle p | \mathcal{H}_0 | p \rangle + \frac{1}{2} \sum_{p,q} [\langle pq | f_{12} | pq \rangle \pm \langle pq | f_{12} | qp \rangle]$$



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 - $\mathcal{J} \equiv \langle \psi_k^* | f_{12} | \psi_k \rangle = \int \psi_k^*(\mathbf{r}) f_{12} \psi_k(\mathbf{r}) d\mathbf{r}$
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- Expand: $\phi_i(\mathbf{r}) = \sum_{p=1}^L C_{pi} \chi_p(\mathbf{r})$
- Roothan-Hall: $\mathbf{F}\mathbf{C}_i = \boldsymbol{\varepsilon} S \mathbf{C}_i$
 - $F_{pq} = h_{pq} + \sum_{pq} \rho_{pq} (2\langle pq | f_{12} | rs \rangle - \langle pq | f_{12} | sr \rangle)$
 - $h_{pq} \equiv \langle p | h | q \rangle$
 - $\rho_{pq} \equiv \sum_{i=1}^{\frac{N}{2}} C_{pi} C_{qi}^*$
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- Poople-Nesbet: $\mathbf{F}^+ \mathbf{C}^+ = \boldsymbol{\epsilon} \mathbf{S} \mathbf{C}^+, \mathbf{F}^- \mathbf{C}^- = \boldsymbol{\epsilon}^- \mathbf{S} \mathbf{C}^-$

Variational Monte-Carlo

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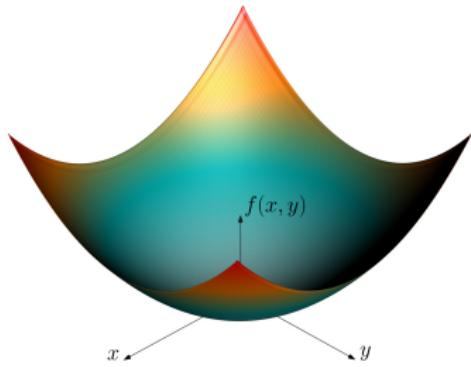
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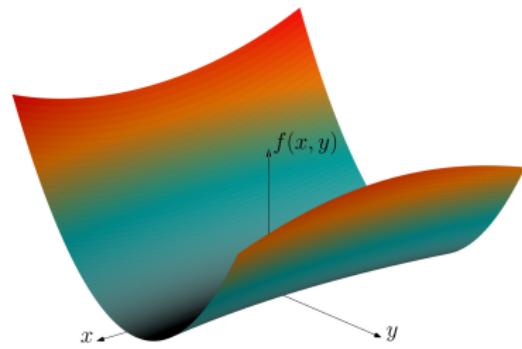
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 - Importance Sampling
 - Focker-Planck equation and Langevin equation
 - $r^{(b)} = r^{(a)} + D \Delta t F^{(a)} + \sqrt{\Delta t} \xi$
 - Quantum force: $F = \frac{2}{\Psi} \nabla \Psi$
 - $\frac{T(b, a, \Delta t)}{T(a, b, \Delta t)} = \text{Greensfunction ratio}$

Minimization

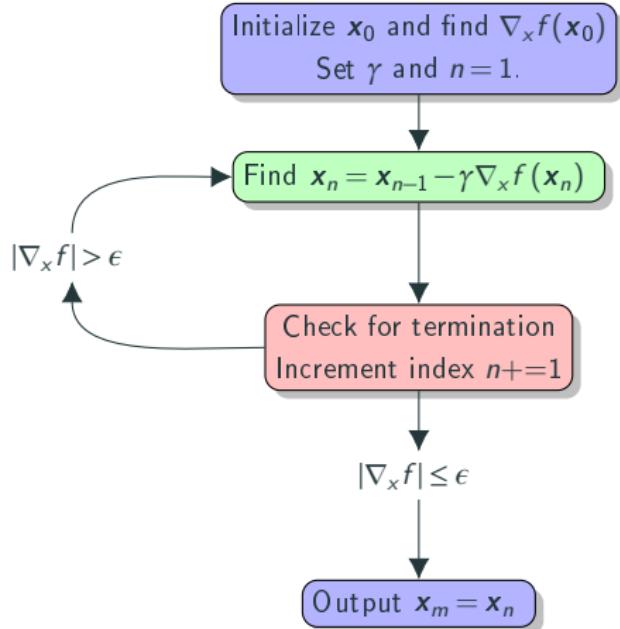
Single-Well



Rosenbrock



Minimization: Gradient Descent



Minimization: Gradient Descent

| x_0 | γ | Iterations | x_m | $f(x_m)$ |
|--------|----------|------------|---|-------------------------|
| (5, 5) | 0.9 | 20 | (-0.072, -0.072) | 0.010 |
| (5, 5) | 0.9 | 50 | (-8.920×10^{-5} , -8.920×10^{-5}) | 1.591×10^{-8} |
| (5, 5) | 0.9 | 100 | (-1.273×10^{-9} , -1.273×10^{-9}) | 3.242×10^{-18} |
| (5, 5) | 0.5 | 20 | (0.0, 0.0) | 0.0 |
| (5, 5) | 0.5 | 50 | (0.0, 0.0) | 0.0 |
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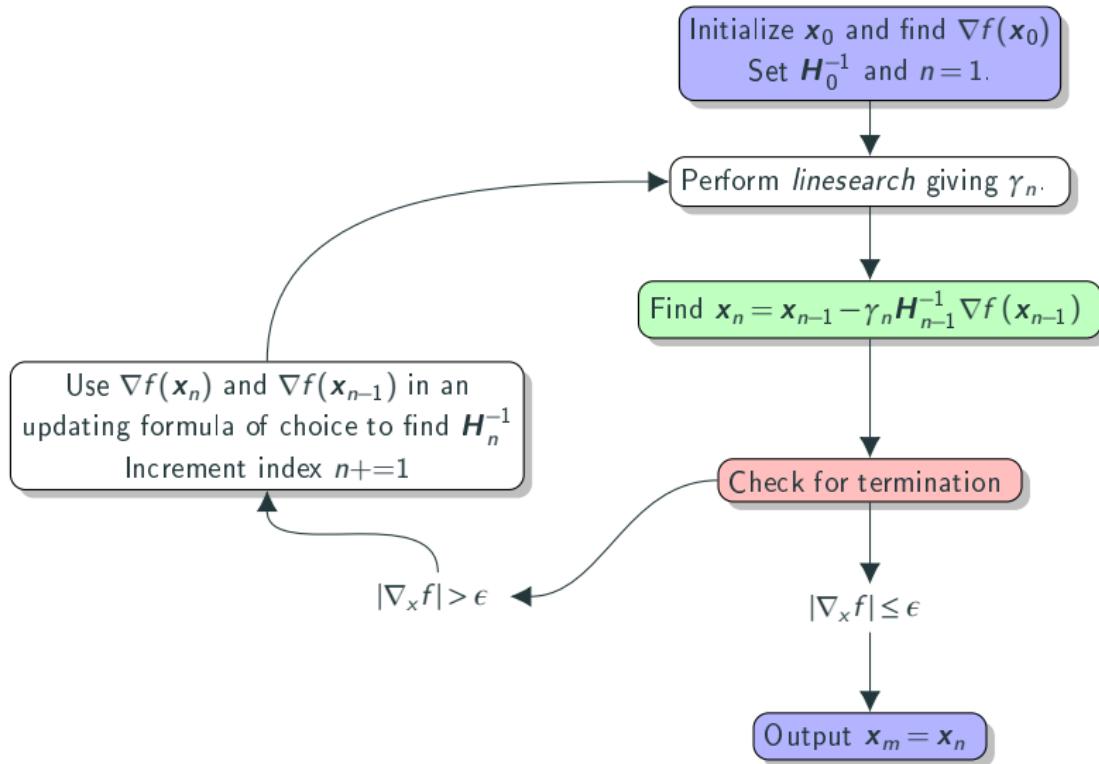
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|----------|----------|------------|----------------|------------------------|
| (0, 0.5) | 0.001 | 100 | (0.181, 0.030) | 0.034 |
| (0, 0.5) | 0.001 | 500 | (0.512, 0.258) | 0.327 |
| (0, 0.5) | 0.001 | 1000 | (0.675, 0.454) | 0.106 |
| (0, 0.5) | 0.001 | 100000 | (1.000, 1.000) | 0.0 |
| (0, 0.5) | 0.0001 | 100 | (0.027, 0.068) | 1.399 |
| (0, 0.5) | 0.0001 | 500 | (0.105, 0.009) | 0.801 |
| (0, 0.5) | 0.0001 | 1000 | (0.184, 0.031) | 0.666 |
| (0, 0.5) | 0.0001 | 100000 | (0.994, 0.989) | 3.131×10^{-5} |

Minimization: Gradient Descent

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Minimization: Quasi-Newton BFGS



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|-----------|------------|-----------------|----------|
| (1,1) | 1 | (-0.071,-0.071) | 1.000 |
| (-1,2) | 1 | (0.447,-0.894) | 1.000 |
| (1,1) | 2 | (0.000,0.000) | 0.000 |
| (-1,2) | 2 | (0.000,0.000) | 0.000 |
| (10,10) | 1 | (-0.071,-0.071) | 1.000 |
| (10,10) | 2 | (0.000,0.000) | 0.000 |
| (100,100) | 1 | (-0.071,-0.071) | 1.000 |
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| x_0 | Iterations | x_m | $f(x_m)$ |
|-------------|------------|----------------|----------|
| (-0.5,2.0) | 1 | (-0.706,0.708) | 7.280 |
| (-0.5,2.0) | 2 | (-0.780,0.649) | 3.342 |
| (-0.5,2.0) | 10 | (0.238,0.051) | 0.584 |
| (-0.5,2.0) | 30 | (1.000,1,000) | 0.000 |
| (5.5,-10.0) | 1 | (-0.996,0.091) | 85.214 |
| (5.5,-10.0) | 2 | (-0.908,1.087) | 10.549 |
| (5.5,-10.0) | 10 | (0.027,0.012) | 0.9613 |
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Minimization: Simulated Annealing

Wavefunction

Wavefunction

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Wavefunction

$$\phi_i(\mathbf{r}) = \sum_{p=1}^L C_{pi} \chi_p(\mathbf{r})$$

Wavefunction: Integral Elements

$$\langle \phi_i(\mathbf{r}) | \phi_j(\mathbf{r}) \rangle$$

$$\langle \phi_i(\mathbf{r}) | x_d^k | \phi_j(\mathbf{r}) \rangle$$

$$\langle \phi_i(\mathbf{r}) | \nabla^2 | \phi_j(\mathbf{r}) \rangle$$

$$\langle \phi_i(\mathbf{r}_1) \phi_j(\mathbf{r}_2) | f_{12} | \phi_k(\mathbf{r}_1) \phi_l(\mathbf{r}_2) \rangle$$

Wavefunction: Single-Well

- Hermite Function: $\psi_n(\mathbf{r}) \equiv \prod_d N_d H_{n_d}(\sqrt{\omega}x_d) \exp\left(-\frac{\omega}{2}x_d^2\right)$

Wavefunction: Single-Well

- Hermite Function: $\psi_n(\mathbf{r}) \equiv \prod_d N_d H_{n_d}(\sqrt{\omega}x_d) \exp\left(-\frac{\omega}{2}x_d^2\right)$
- Solution in polar³

³E. Anisimovas and A. Matulis. “Energy spectra of few-electron quantum dots”. In: *Journal of Physics: Condensed Matter* (1998), p. 601.

Wavefunction: Single-Well

- Hermite Function: $\psi_n(\mathbf{r}) \equiv \prod_d N_d H_{n_d}(\sqrt{\omega}x_d) \exp\left(-\frac{\omega}{2}x_d^2\right)$
- Solution in polar³
- $\psi_n(\mathbf{r}) = \prod_d N_d \sum_{l=1}^{n_d} C_{n_d l}^{\text{Hermite}} g_l\left(\frac{\omega}{2}, \mathbf{r}, \mathbf{0}\right), \quad g_l\left(\frac{\omega}{2}, \mathbf{r}, \mathbf{0}\right) = x_d^{(l)} \exp\left(-\frac{\omega^2}{2}x_d^2\right)$

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- Solution in Cartesian⁴

$$\langle g_i(\mathbf{r}) | g_j(\mathbf{r}) \rangle$$

$$\langle g_i(\mathbf{r}) | x_d' | g_j(\mathbf{r}) \rangle$$

$$\langle g_i(\mathbf{r}) | \nabla^2 | g_j(\mathbf{r}) \rangle$$

$$\langle g_i(\mathbf{r}_1) g_j(\mathbf{r}_2) | f_{12} | g_k(\mathbf{r}_1) g_l(\mathbf{r}_2) \rangle$$

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Wavefunction: Single-Well Integral Elements

$$\langle \psi_i^{\text{HO}} | \psi_j^{\text{HO}} \rangle = N_i \delta_{ij}$$

$$\langle \psi_i^{\text{HO}} | h^{\text{HO}} | \psi_j^{\text{HO}} \rangle = N_i \epsilon_i^{\text{HO}} \delta_{ij}$$

$$\langle \psi_i^{\text{HO}} \psi_j^{\text{HO}} | \frac{1}{r_{12}} | \psi_k^{\text{HO}} \psi_l^{\text{HO}} \rangle = \frac{aN_{ijkl}}{\sqrt{2\omega}} \sum_{tuvw}^{ijkl} H_{tuvw}^{ijkl} \sum_{pq}^{t+v, u+w} E_p^{tv} E_q^{uw} (-1)^q \xi_{p+q}(\frac{\omega}{2}, 0)$$

$$E_t^{i_d+1} = \frac{1}{2\omega} E_{t-1}^i \quad \xi_{t_d+1}^n = t_d \xi_{t_d-1}^{n+1}$$

$$E_0^0 = K_{AB} \quad \xi_0^n = (-b)^n \zeta_n(0)$$

$$\zeta_n^{\text{2D}}(x) = \int_{-1}^1 \frac{u^{2n}}{\sqrt{1-u^2}} e^{-u^2} du \quad \zeta_n^{\text{3D}}(x) = \int_{-1}^1 u^{2n} e^{-u^2} du$$

$$b = \begin{cases} \frac{\omega}{2}, & \text{2D} \\ \omega, & \text{3D} \end{cases}$$

Wavefunction: Double-Well

- Perturbation of harmonic oscillator: $U^{\text{DW}}(r) = V^{\text{HO}}(r) + V_n^{\text{DW}}(r)$

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 - $H_{ij}^{\text{DW}} = \epsilon_i^{\text{HO}} \delta_{ij} + \langle \psi_i^{\text{HO}} | V_n^{\text{DW}} | \psi_j^{\text{HO}} \rangle$

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- Integral-Elements

$$\langle \psi_p^{\text{DW}} \left| \psi_q^{\text{DW}} \right\rangle = \delta_{pq}$$

$$\langle \psi_p^{\text{DW}} \left| h^{\text{DW}} \right| \psi_q^{\text{DW}} \rangle = \epsilon_p^{\text{DW}} \delta_{pq}$$

$$\left\langle \psi_p^{\text{DW}} \psi_q^{\text{DW}} \left| \frac{1}{r_{12}} \right| \psi_r^{\text{DW}} \psi_s^{\text{DW}} \right\rangle = \sum_{tuvw}^{ijkl} C_{tp}^{\text{DW}} C_{uq}^{\text{DW}} C_{vr}^{\text{DW}} C_{ws}^{\text{DW}} \left\langle \psi_t^{\text{HO}} \psi_u^{\text{HO}} \left| \frac{1}{r_{12}} \right| \psi_v^{\text{HO}} \psi_w^{\text{HO}} \right\rangle$$

Wavefunction: Slater-Jastrow

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- Padé-NQS: $J = J_{\text{Padé}} J_{\text{NQS}}$

Implementation

Implementation

- C++ and Eigen

Implementation

- C++ and Eigen
 - Performance

Implementation

- C++ and Eigen
 - Performance
 - Generalization

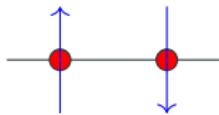
Implementation

- C++ and Eigen
 - Performance
 - Generalization
- Python

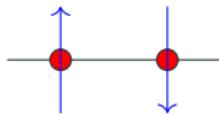
Implementation

- C++ and Eigen
 - Performance
 - Generalization
- Python
 - Generate C++ code

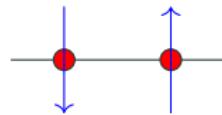
Implementation: Cartesian



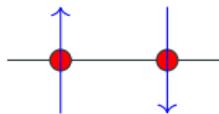
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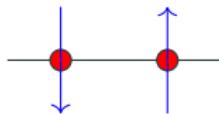
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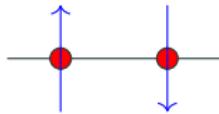
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(1,0)

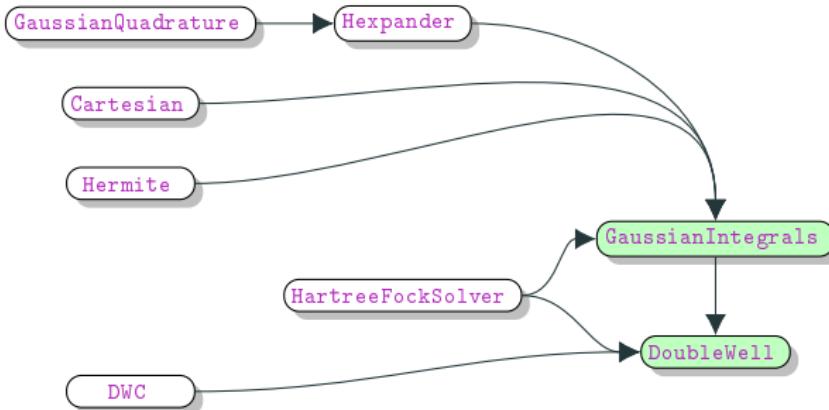


(0,1)

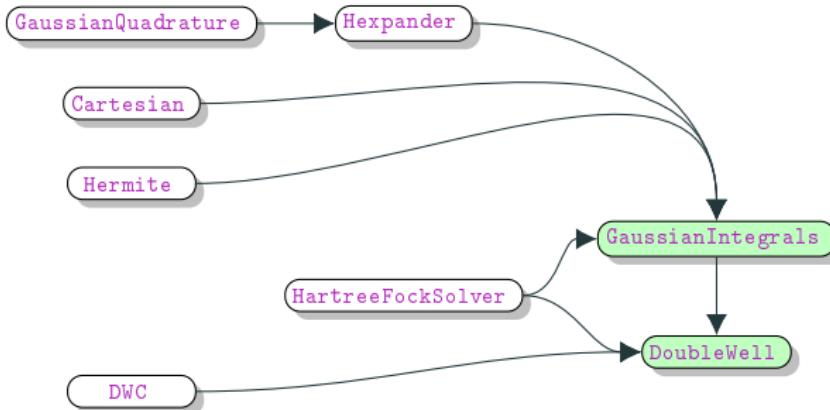


(0,0)

Implementation: Hartree-Fock

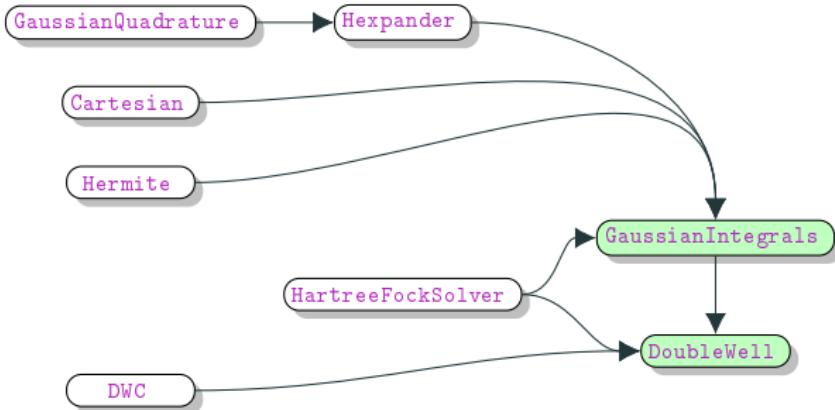


Implementation: Hartree-Fock



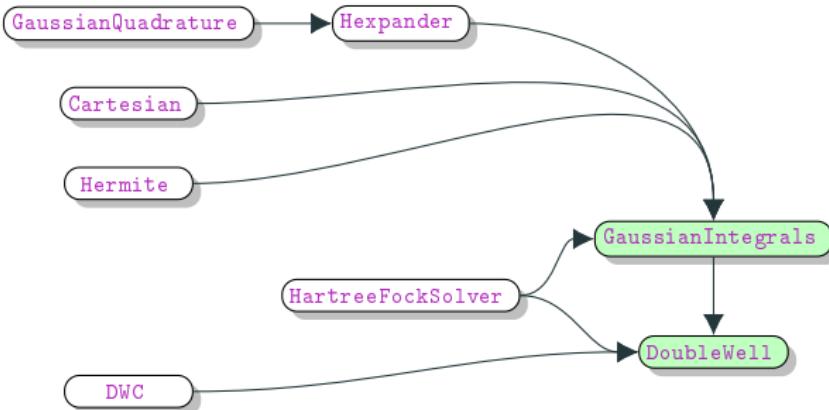
- Parallelization

Implementation: Hartree-Fock



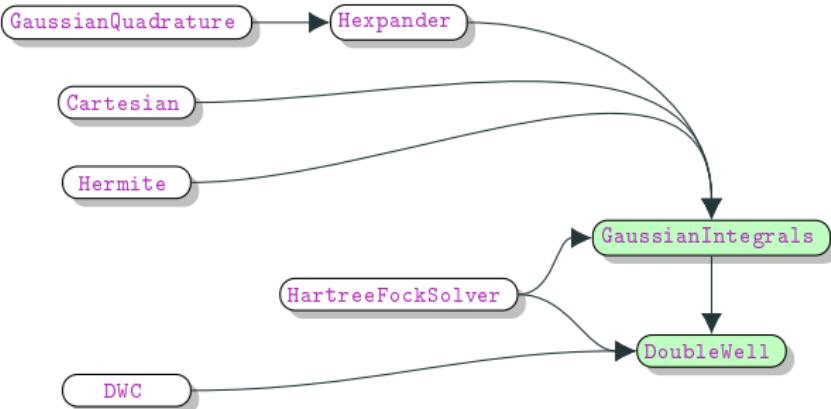
- Parallelization
 - Two-body element is computationally expensive
 - $S_i = \sum_{j=0}^{P_i} \prod_d (n_{jd} + 1)$

Implementation: Hartree-Fock



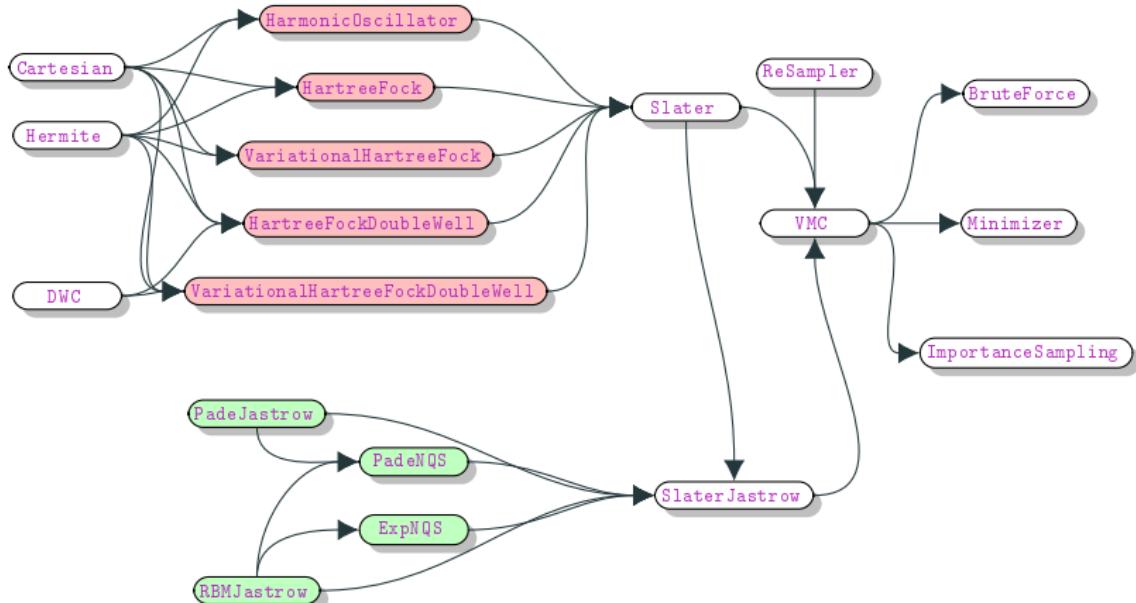
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Implementation: Hartree-Fock

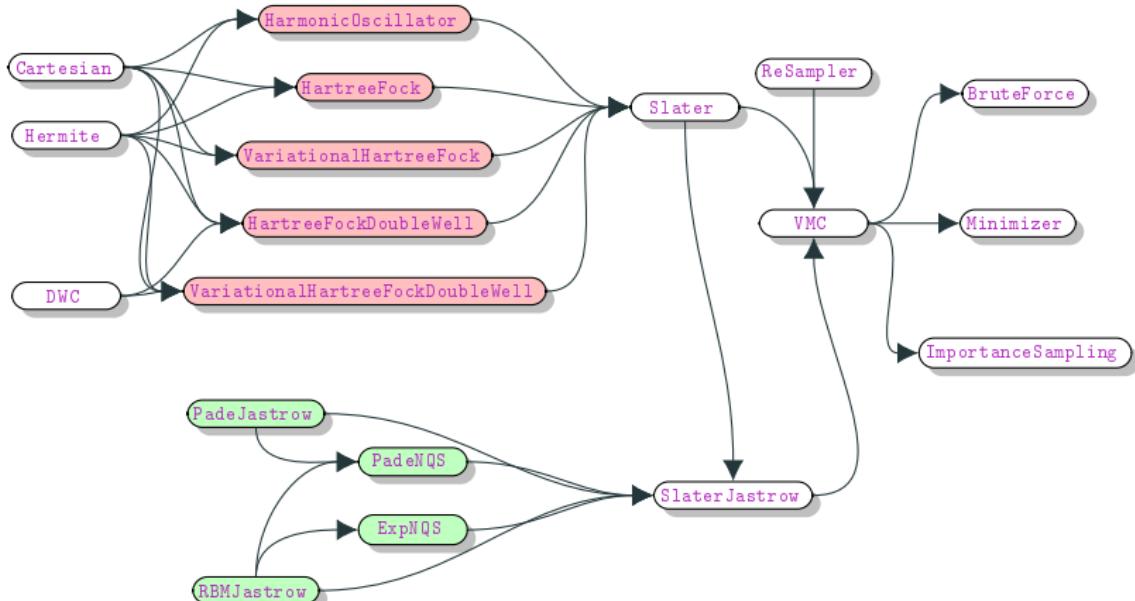


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 - Two-body element is computationally expensive
 - $S_i = \sum_{j=0}^{P_i} \prod_d (n_{jd} + 1)$
- Hartree-Fock algorithm only run on one process
- Tabulation of Two-Body matrix

Implementation: Variational Monte-Carlo

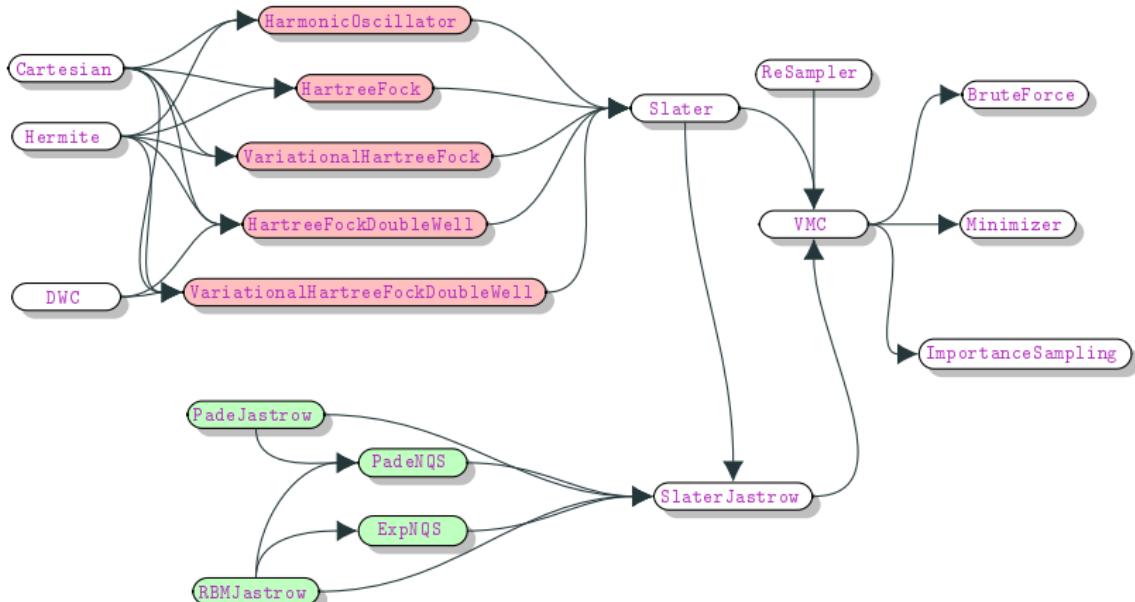


Implementation: Variational Monte-Carlo



- Hermite generated with Python and SymPy

Implementation: Variational Monte-Carlo

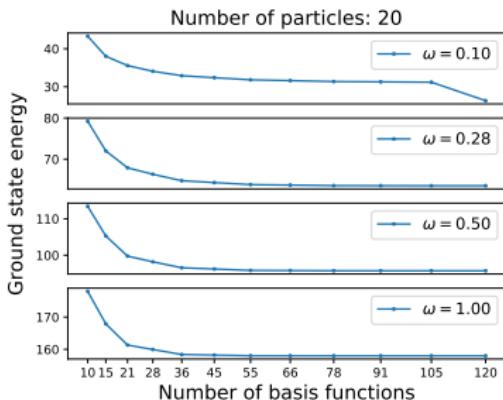
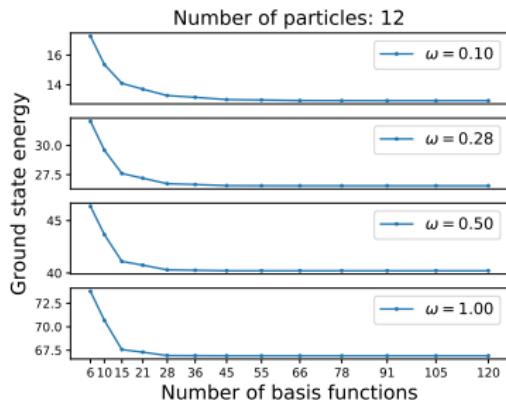
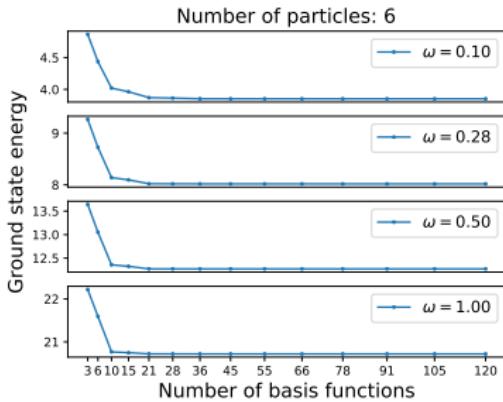
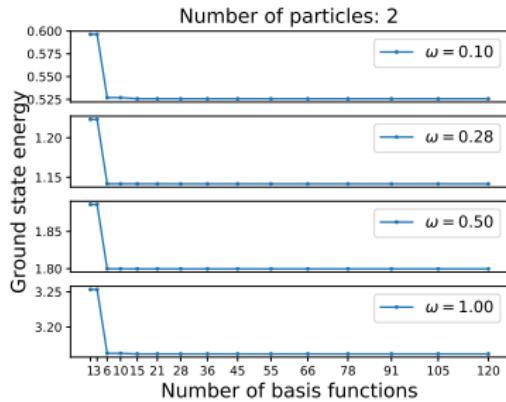


- **Hermite** generated with Python and SymPy
- Wavefunction class generated with Python

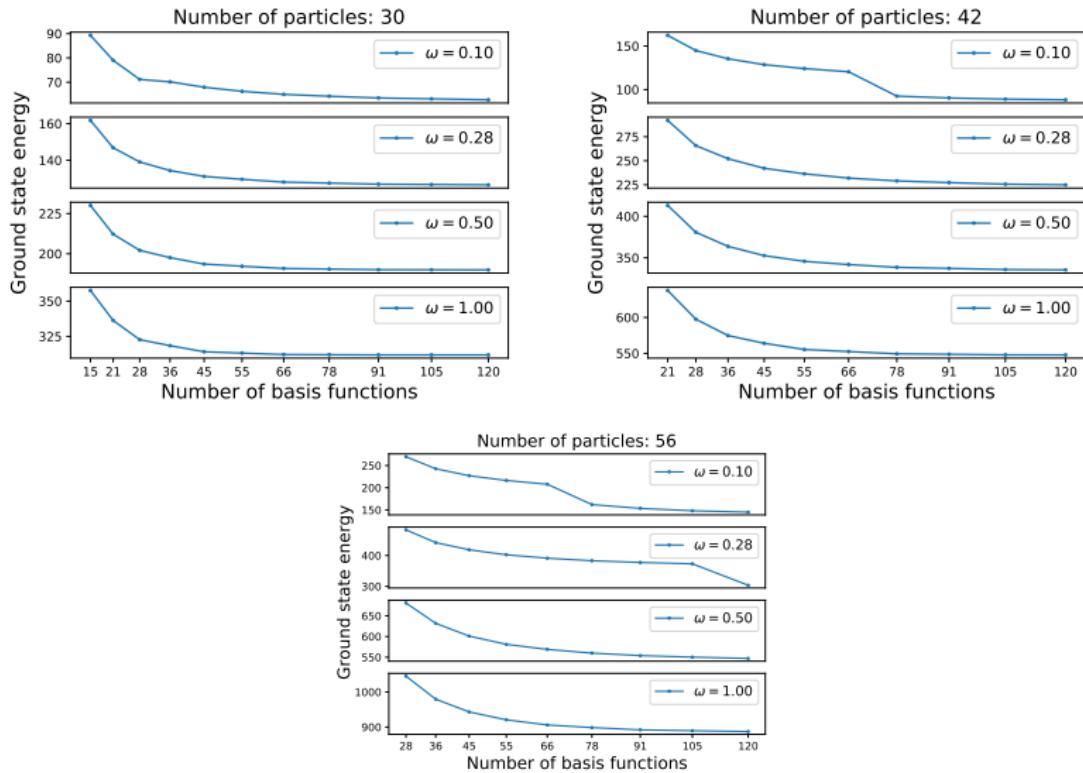
Results

Benchmark

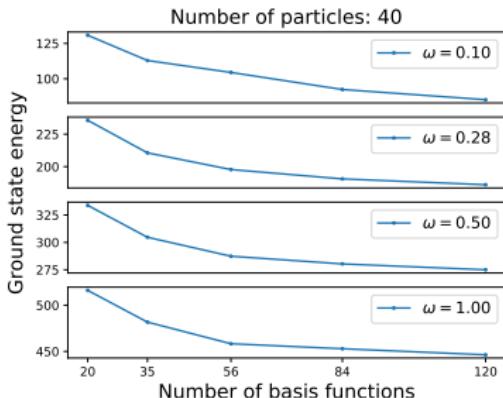
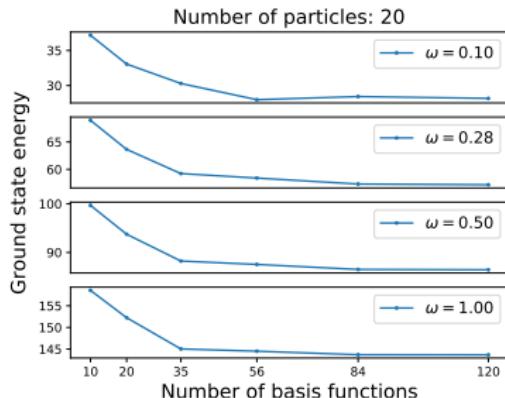
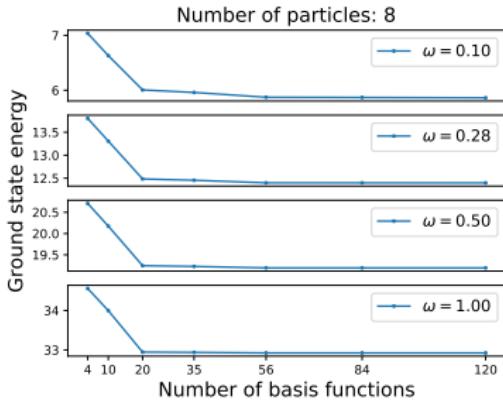
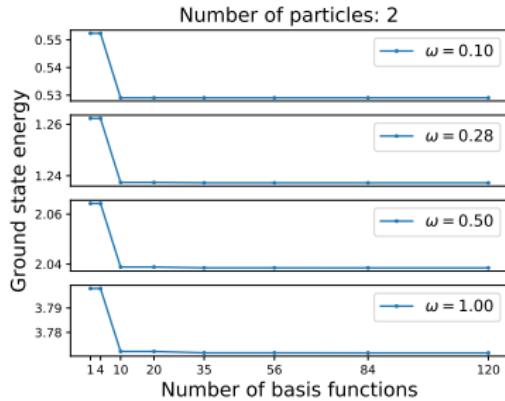
Results: Benchmark



Results: Benchmark



Results: Benchmark



Results: Benchmark

| ω [a.u] | N | | | |
|----------------|-----------|------------|------------|-------------|
| | 2 | 6 | 12 | 20 |
| 0.1 | 0.4407(4) | 3.5650(4) | 12.3164(4) | 30.0480(4) |
| 0.28 | 1.0020(4) | 7.6198(4) | 25.5948(3) | 61.8090(3) |
| 0.5 | 1.6650(4) | 11.8017(4) | 39.3166(3) | 93.9240(2) |
| 1.0 | 3.0000(5) | 20.2863(3) | 68.1465(3) | 156.2778(2) |

| ω [a.u] | N | |
|----------------|------------|-------------|
| | 2 | 8 |
| 0.1 | 0.50006(5) | 5.80479(4) |
| 0.28 | 1.20156(5) | 12.48178(4) |
| 0.5 | 2.00027(5) | 19.33356(4) |
| 1.0 | 3.72985(5) | 33.30958(4) |

$$\psi = \psi^{\text{HO}}(\sqrt{\alpha\omega}) J_{\text{Pad\'e}}$$

Results: Benchmark

| ω [a.u] | N | | | |
|----------------|----------------|-----------------|-----------------|------------------|
| | 2 | 6 | 12 | 20 |
| 0.1 | 0.46552(5){15} | 3.70137(4){36} | 12.64342(4){91} | - |
| 0.28 | 1.04939(4){6} | 7.89627(4){36} | 26.21301(4){66} | 62.93503(5){120} |
| 0.5 | 1.70130(4){6} | 12.02776(4){21} | 39.76442(3){45} | 95.21976(3){91} |
| 1.0 | 3.05625(4){6} | 20.45876(3){36} | 66.37115(3){45} | 157.41119(3){78} |

| ω [a.u] | N | | | |
|----------------|----------------|-----------------|-----------------|------------------|
| | 2 | 6 | 12 | 20 |
| 0.10 | 0.44473(5){15} | 3.63897(4){36} | 12.46408(4){91} | - |
| 0.28 | 1.04978(4){6} | 7.72929(4){36} | 25.96595(4){66} | 62.65652(3){120} |
| 0.50 | 1.66418(4){6} | 11.97781(4){21} | 39.57182(3){45} | 94.76303(3){91} |
| 1.00 | 3.00624(4){6} | 20.38811(3){36} | 66.28996(3){45} | 157.46167(3){78} |

$$\psi_p = \sum_l C_{lp} \psi_l^{\text{HO}}(\sqrt{\omega} r) J_{\text{Pad\'e}}, \quad \psi_p = \sum_l C_{lp} \psi_l^{\text{HO}}(\sqrt{\alpha\omega} r) J_{\text{Pad\'e}}$$

Results: Benchmark

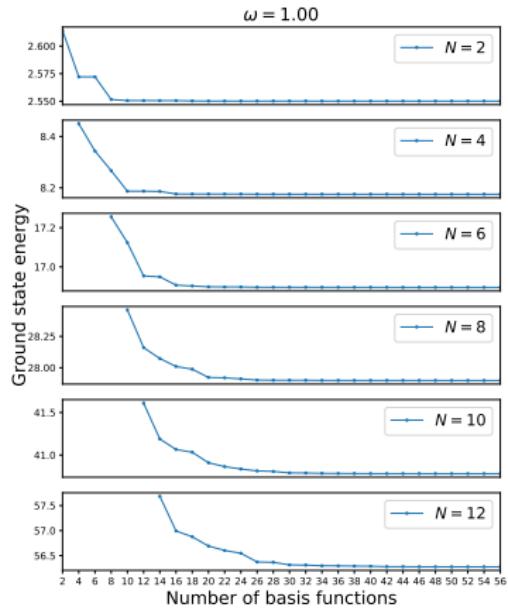
| ω | N | |
|----------|----------------|------------------|
| | 2 | 8 |
| 0.1 | 0.51122(5){70} | 5.87372(4){120} |
| 0.28 | 1.21844(5){70} | 12.36177(4){168} |
| 0.5 | 2.02030(4){20} | 19.15006(4){112} |
| 1.0 | 3.72918(5){20} | 33.58046(4){168} |

| ω | N | |
|----------|----------------|------------------|
| | 2 | 8 |
| 0.1 | 0.50751(5){70} | 5.84082(4){240} |
| 0.28 | 1.20320(5){20} | 12.37435(4){168} |
| 0.5 | 2.01439(4){20} | 19.09917(4){112} |
| 1.0 | 3.72959(5){70} | 33.04162(4){168} |

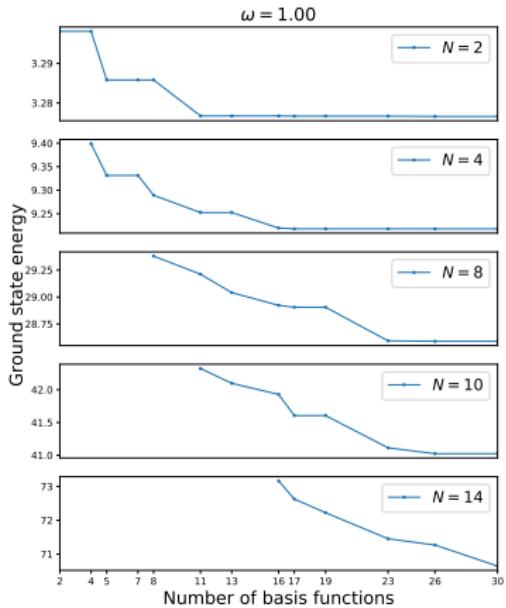
$$\psi_p = \sum_l C_{lp} \psi_l^{\text{HO}}(\sqrt{\omega} r) J_{\text{Pad\'e}}, \quad \psi_p = \sum_l C_{lp} \psi_l^{\text{HO}}(\sqrt{\alpha\omega} r) J_{\text{Pad\'e}}$$

Results: Double-Well Hartree-Fock

2D



3D



Results: Double-Well Variational Monte-Carlo

| ω | N | | | |
|----------|----------------|----------------|-----------------|-----------------|
| | 2 | 4 | 6 | 8 |
| 1.0 | 2.42238(4){10} | 7.95247(4){42} | 16.61419(4){44} | 27.54453(3){50} |

$$\psi_p = \sum_l C_{lp}^{\text{HF}} \sum_k C_{kl}^{\text{DW}} \psi_k^{\text{HO}} (\sqrt{\omega} r) J_{\text{Pad\'e}}$$

| ω | N | | | |
|----------|----------------|----------------|-----------------|-----------------|
| | 2 | 4 | 6 | 8 |
| 1.0 | 2.36618(4){10} | 7.90232(4){42} | 16.55609(4){44} | 27.58524(4){50} |

$$\psi_p = \sum_l C_{lp}^{\text{HF}} \sum_k C_{kl}^{\text{DW}} \psi_k^{\text{HO}} (\sqrt{\alpha \omega} r) J_{\text{Pad\'e}}$$

Results: Double-Well Variational Monte-Carlo

| ω | N | | |
|----------|----------------|----------------|-----------------|
| | 2 | 4 | 8 |
| 1.0 | 3.25118(4){11} | 9.17489(4){17} | 28.49671(4){26} |

$$\psi_p = \sum_l C_{lp}^{\text{HF}} \sum_k C_{kl}^{\text{DW}} \psi_k^{\text{HO}} (\sqrt{\omega} r) J_{\text{Pad\'e}}$$

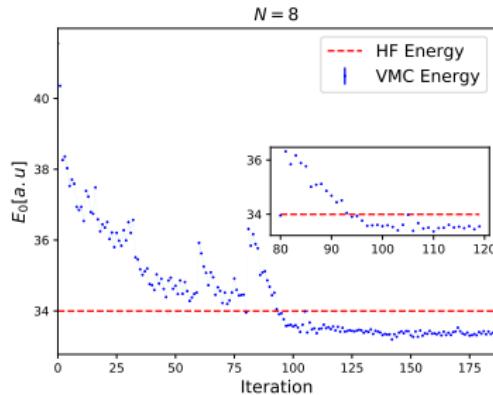
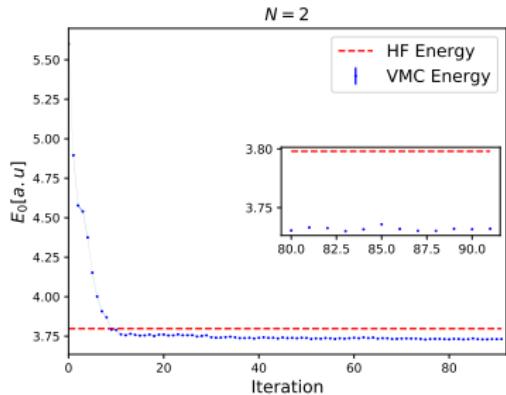
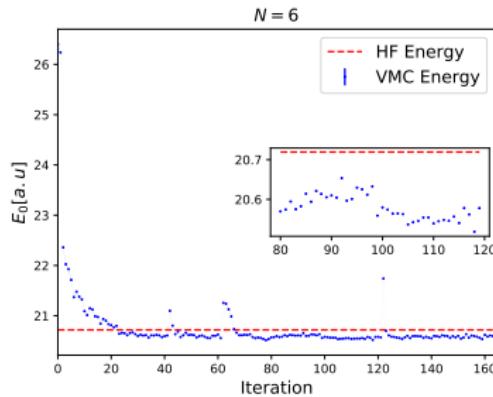
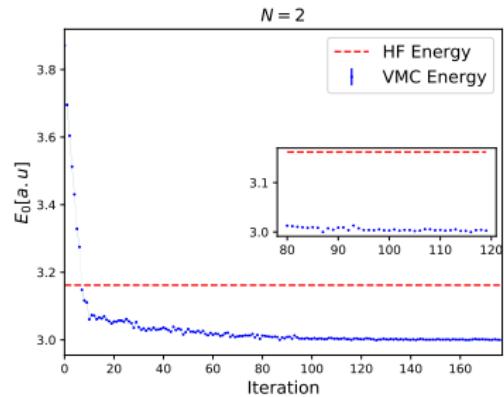
| ω | N | | |
|----------|----------------|----------------|-----------------|
| | 2 | 4 | 8 |
| 1.0 | 3.22226(4){11} | 9.17013(4){17} | 28.62826(4){26} |

$$\psi_p = \sum_l C_{lp}^{\text{HF}} \sum_k C_{kl}^{\text{DW}} \psi_k^{\text{HO}} (\sqrt{\alpha\omega} r) J_{\text{Pad\'e}}$$

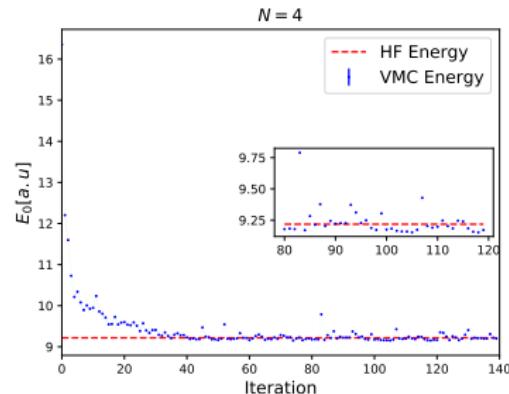
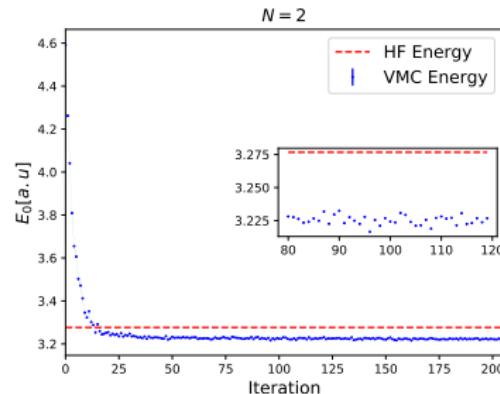
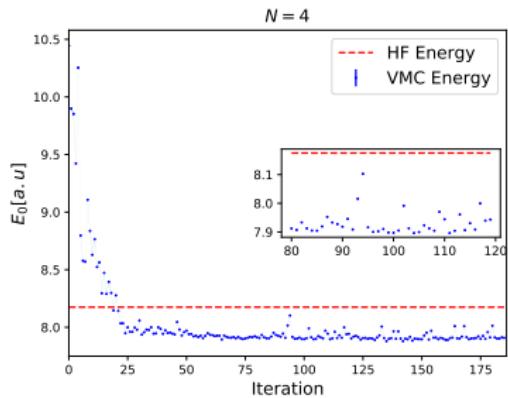
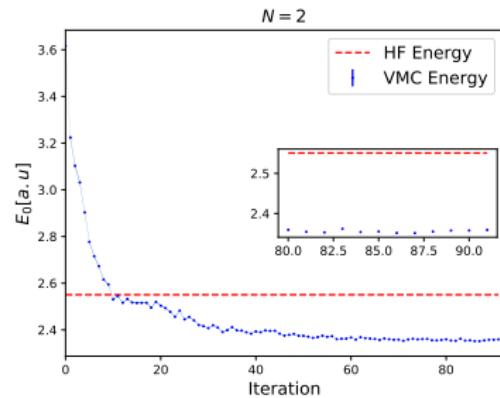
Results: NQS-Jastrow

$$J_{\text{NQS}} = e^{-\sum_{i=1}^N \frac{(r_i - a_i)^2}{2\sigma^2}} \prod_j^M \left(1 + e^{b_j + \sum_{i=1}^N \sum_{d=1}^D \frac{x_i^{(d)} w_{i+d,j}}{\sigma^2}} \right)$$

Results: NQS-Jastrow Harmonic Oscillator



Results: NQS-Jastrow Double-Well



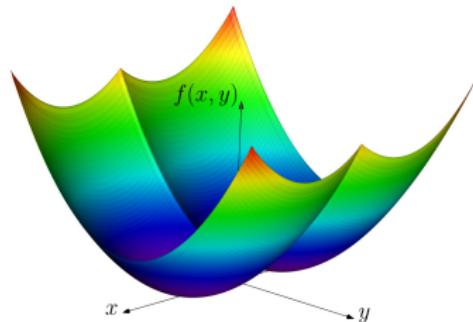
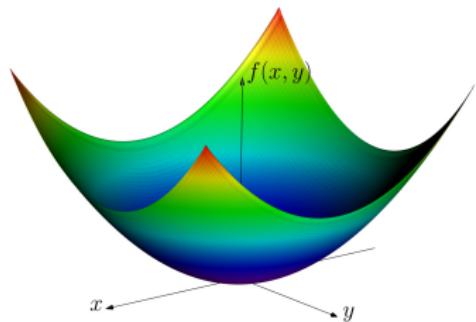
Summary and Conclusion

Summary

- Schrödinger equation: $\mathcal{H}|\psi\rangle = E|\psi\rangle$, $\mathcal{H} = -\sum_i \frac{\nabla_i^2}{2} + f(\mathbf{r}) + V(\mathbf{R}, \mathbf{r})$
- Interaction: $f(\mathbf{r}) = \sum_{i < j} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$
- Confinement: Harmonic Oscillator, Double-Well

$$V(\mathbf{r}) = \frac{1}{2} m \omega^2 r^2$$

$$V(\mathbf{r}) = \frac{1}{2} m \omega^2 (r^2 - \delta R |x| + R^2)$$



**Hartree-Fock
Variational Monte-Carlo**

Summary

$$\begin{aligned}\left\langle \psi_i^{\text{HO}} \middle| \psi_j^{\text{HO}} \right\rangle &= N_i \delta_{ij} \\ \left\langle \psi_i^{\text{HO}} \middle| h^{\text{HO}} \right\rangle &= N_i \varepsilon_i^{\text{HO}} \delta_{ij}\end{aligned}$$

$$\left\langle \psi_i^{\text{HO}} \psi_j^{\text{HO}} \middle| \frac{1}{r_{12}} \right\rangle = \frac{aN_{ijkl}}{\sqrt{2\omega}} \sum_{tuvw}^{ijkl} H_{tuvw}^{ijkl} \sum_{pq}^{t+v, u+w} E_p^{tv} E_q^{uw} (-1)^q \xi_{p+q}(\frac{\omega}{2}, \mathbf{0})$$

$$\begin{aligned}E_t^{i_d+1} &= \frac{1}{2\omega} E_{t-1}^i & \xi_{t_d+1}^n &= t_d \xi_{t_d-1}^{n+1} \\ E_0^0 &= K_{AB} & \xi_0^n &= (-b)^n \zeta_n(0)\end{aligned}$$

$$\zeta_n^{\text{2D}}(x) = \int_{-1}^1 \frac{u^{2n}}{\sqrt{1-u^2}} e^{-u^2 x} du \quad \zeta_n^{\text{3D}}(x) = \int_{-1}^1 u^{2n} e^{-u^2 x} du$$

$$b = \begin{cases} \frac{\omega}{2}, & \text{2D} \\ \omega, & \text{3D} \end{cases}$$

Summary

- Perturbation of harmonic oscillator: $U^{\text{DW}}(r) = V^{\text{HO}}(r) + V_n^{\text{DW}}(r)$
- Expand in HO-functions: $\left| \psi_p^{\text{DW}} \right\rangle = \sum_l C_{lp}^{\text{DW}} \left| \psi_l^{\text{HO}} \right\rangle$
- Eigenvalue equation: $H^{\text{DW}} C^{\text{DW}} = \epsilon^{\text{DW}} C^{\text{DW}}$
 - $H_{ij}^{\text{DW}} = \epsilon_i^{\text{HO}} \delta_{ij} + \langle \psi_i^{\text{HO}} \left| V_n^{\text{DW}} \right| \psi_j^{\text{HO}} \rangle$
- Integral-Elements

$$\langle \psi_p^{\text{DW}} \left| \psi_q^{\text{DW}} \right\rangle = \delta_{pq}$$

$$\langle \psi_p^{\text{DW}} \left| h^{\text{DW}} \right| \psi_q^{\text{DW}} \rangle = \epsilon_p^{\text{DW}} \delta_{pq}$$

$$\left\langle \psi_p^{\text{DW}} \psi_q^{\text{DW}} \left| \frac{1}{r_{12}} \right| \psi_r^{\text{DW}} \psi_s^{\text{DW}} \right\rangle = \sum_{tuvw}^{ijkl} C_{tp}^{\text{DW}} C_{uq}^{\text{DW}} C_{vr}^{\text{DW}} C_{ws}^{\text{DW}} \left\langle \psi_t^{\text{HO}} \psi_u^{\text{HO}} \left| \frac{1}{r_{12}} \right| \psi_v^{\text{HO}} \psi_w^{\text{HO}} \right\rangle$$

Conclusion

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- Stable minimization with NQS-Jastrow

Further Work

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- Rewrite with sparse matrices in Hartree-Fock

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- Extend to non-isotropic gaussians

End

Questions?

Methods: Hartree-Fock

- Assumptions
 - The Born-Oppenheimer approximation holds.
 - All relativistic effects are negligible.
 - The wavefunction can be described by a single *Slater determinant*.
 - The Mean Field Approximation holds.

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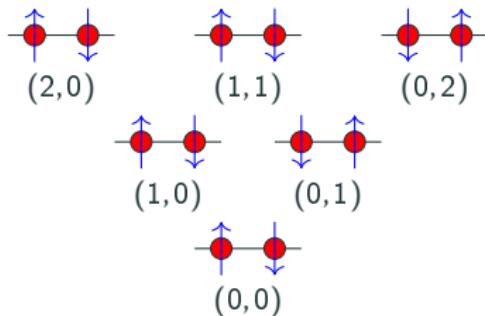
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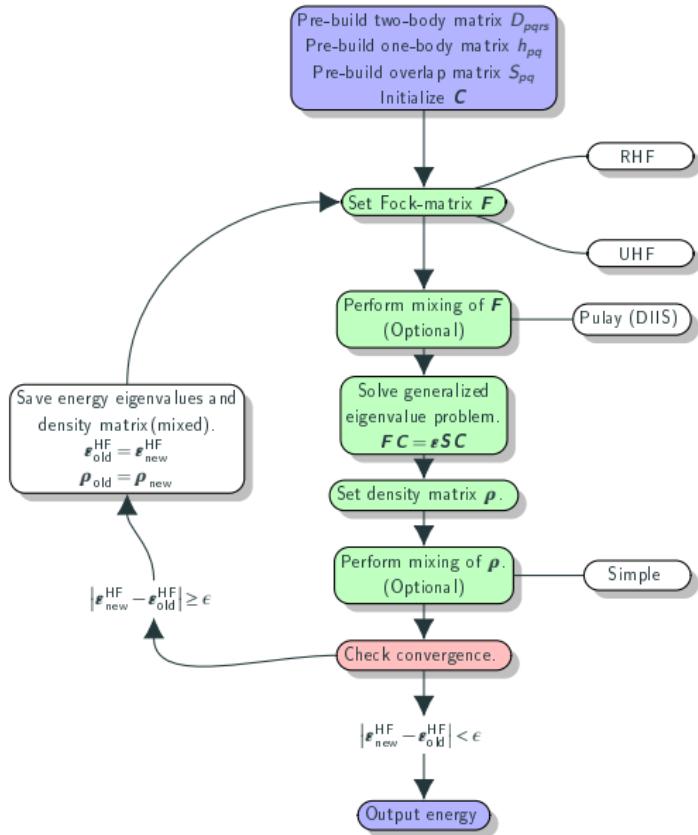
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