

Quantum Many-Body Simulations of Double Dot System

Alocias Mariadason

Institute of Physics

Table of contents

1. Introduction
2. Methods
3. Implementation
4. Summary and Conclusion

Introduction

Quantum-Dot

- Small semiconductor nanostructures
- 2-10 nanometers with 10—50 particles

Quantum-Dot Model

- Schrödinger equation
 - $\mathcal{H}|\psi\rangle = E|\psi\rangle$

Quantum-Dot Model

- Schrödinger equation

- $\mathcal{H}|\psi\rangle = E|\psi\rangle$

- Hamiltonian

- $$\mathcal{H} = -\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i < j} f(\mathbf{r}_i, \mathbf{r}_j) - \frac{1}{2} \sum_k \frac{\nabla_k^2}{M_k} + \sum_{k < l} g(\mathbf{R}_k, \mathbf{R}_l) + V(\mathbf{R}, \mathbf{r})$$

Quantum-Dot Model

- Schrödinger equation
 - $\mathcal{H}|\psi\rangle = E|\psi\rangle$
- Hamiltonian
 - $\mathcal{H} = -\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i<j} f(\mathbf{r}_i, \mathbf{r}_j) - \frac{1}{2} \sum_k \frac{\nabla_k^2}{M_k} + \sum_{k<l} g(\mathbf{R}_k, \mathbf{R}_l) + V(\mathbf{R}, \mathbf{r})$
- Born-Oppenheimer Approximation
 - Ignore Nuclei
 - $\sum_k \frac{\nabla_k^2}{M_k}$ gone
 - $\sum_{k<l} g(\mathbf{R}_k, \mathbf{R}_l)$ constant

Quantum-Dot Model

- Schrödinger equation
 - $\mathcal{H}|\psi\rangle = E|\psi\rangle$
- Hamiltonian
 - $\mathcal{H} = -\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i<j} f(\mathbf{r}_i, \mathbf{r}_j) - \frac{1}{2} \sum_k \frac{\nabla_k^2}{M_k} + \sum_{k<l} g(\mathbf{R}_k, \mathbf{R}_l) + V(\mathbf{R}, \mathbf{r})$
- Born-Oppenheimer Approximation
 - Ignore Nuclei
 - $\sum_k \frac{\nabla_k^2}{M_k}$ gone
 - $\sum_{k<l} g(\mathbf{R}_k, \mathbf{R}_l)$ constant
 - $\mathcal{H} = -\frac{1}{2} \sum_i \nabla_i^2 + \sum_{i<j} f(\mathbf{r}_i, \mathbf{r}_j) + V(\mathbf{R}, \mathbf{r})$

Quantum-Dot Model

- Interaction

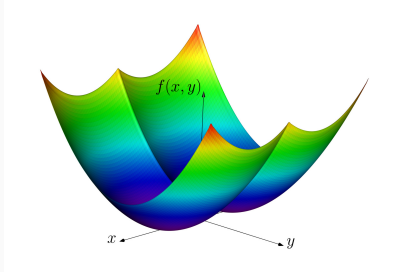
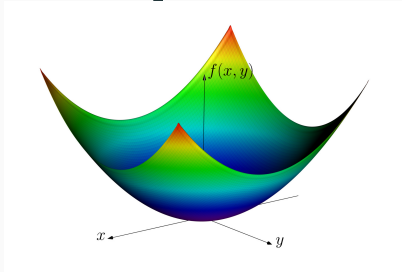
- $f(\mathbf{r}_i, \mathbf{r}_j) = \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$

Quantum-Dot Model

- Interaction: Coulomb repulsion
 - $f(\mathbf{r}_i, \mathbf{r}_j) = \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$
- Confinement: Harmonic Oscillator¹, Double-Well²

$$V(\mathbf{r}) = \frac{1}{2} \omega m r^2$$

$$V(\mathbf{R}, \mathbf{r}) = \frac{1}{2} m \omega^2 (r^2 - \delta R|x| + R^2)$$



¹S. Kvaal. “Harmonic Oscillator Eigenfunction Expansions, Quantum dots, and Effective Interactions”. In: *Phys. Rev. B* 80 (4 2009), p. 045321.

²M. J. A. Schuetz et al. “Nuclear Spin Dynamics in Double Quantum Dots: Multistability, Dynamical Polarization, Criticality, and Entanglement”. In: *Phys. Rev. B* 89 (19 2014), p. 195310.

#NucleiHaveFeelingsTo

Methods

Hartree-Fock Variational Monte-Carlo

Methods: Variational Principle

$$E_0 \leq \frac{\langle \Psi | \mathcal{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle}$$

Methods: Slater Determinant and Energy Functional

- Pauli Principle

Methods: Slater Determinant and Energy Functional

- Pauli Principle
- Slater Determinant
 - $\Psi_T^{\text{AS}} = \frac{1}{\sqrt{N!}} \sum_P (-1)^P \mathcal{P}_P \prod_i \psi_i$
 - $\Psi_T^{\text{S}} = \sqrt{\frac{\prod_{i=1}^N n_i!}{N!}} \sum_P \mathcal{P}_P \prod_i \psi_i$

Methods: Slater Determinant and Energy Functional

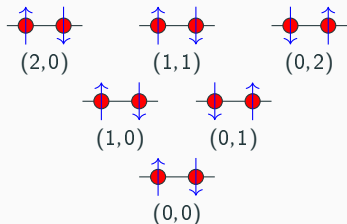
- Pauli Principle
- Slater Determinant
 - $\Psi_T^{\text{AS}} = \frac{1}{\sqrt{N!}} \sum_P (-1)^P \mathcal{P}_P \prod_i \psi_i$
 - $\Psi_T^{\text{S}} = \sqrt{\frac{\prod_{i=1}^N n_i!}{N!}} \sum_P \mathcal{P}_P \prod_i \psi_i$
- $E[\Psi] = \frac{\langle \Psi | \mathcal{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} = \sum_p \langle p | \mathcal{H}_0 | p \rangle + \frac{1}{2} \sum_{p,q} [\langle pq | f_{12} | pq \rangle \pm \langle pq | f_{12} | qp \rangle]$
- $\mathcal{H}_0 = -\frac{1}{2} \sum_i \nabla_i^2 + V(r)$

Methods: Hartree-Fock Method

- Assumptions
 - The Born-Oppenheimer approximation holds.
 - All relativistic effects are negligible.
 - The wavefunction can be described by a single *Slater determinant*.
 - The Mean Field Approximation holds.

Methods: Hartree-Fock Method

- Assumptions
 - The Born-Oppenheimer approximation holds.
 - All relativistic effects are negligible.
 - The wavefunction can be described by a single *Slater determinant*.
 - The Mean Field Approximation holds.



Methods: Hartree-Fock Method

- Constrained minimization

Methods: Hartree-Fock Method

- Constrained minimization
 - Spin orthogonality: $\langle \psi_i | \psi_j \rangle = \delta_{ij}$

Methods: Hartree-Fock Method

- Constrained minimization
 - Spin orthogonality: $\langle \psi_i | \psi_j \rangle = \delta_{ij}$
 - Lagrange Multiplier method

Methods: Hartree-Fock Method

- Constrained minimization
 - Spin orthogonality: $\langle \psi_i | \psi_j \rangle = \delta_{ij}$
 - Lagrange Multiplier method
 - Fock-operator: $\mathcal{F} \equiv \mathcal{H}_0 + \mathcal{J} \pm \mathcal{K}$
 - $\mathcal{J} \equiv \sum_k \langle \psi_k^* | f_{12} | \psi_k \rangle = \int \psi_k^*(\mathbf{r}) f_{12} \psi_k(\mathbf{r}) d\mathbf{r}$
 - $\mathcal{K} \equiv \sum_k \langle \psi_k^* | f_{12} | \psi \rangle = \int \psi_k^*(\mathbf{r}) f_{12} \psi(\mathbf{r}) d\mathbf{r}$
 - $\mathcal{F} | \psi \rangle = \boldsymbol{\varepsilon} | \psi \rangle, \boldsymbol{\varepsilon} = (\varepsilon_0, \dots, \varepsilon_N)$

Methods: Hartree-Fock Method

- Constrained minimization
 - Spin orthogonality: $\langle \psi_i | \psi_j \rangle = \delta_{ij}$
 - Lagrange Multiplier method
 - Fock-operator: $\mathcal{F} \equiv \mathcal{H}_0 + \mathcal{J} \pm \mathcal{K}$
 - $\mathcal{J} \equiv \sum_k \langle \psi_k^* | f_{12} | \psi_k \rangle = \int \psi_k^*(\mathbf{r}) f_{12} \psi_k(\mathbf{r}) d\mathbf{r}$
 - $\mathcal{K} \equiv \sum_k \langle \psi_k^* | f_{12} | \psi \rangle = \int \psi_k^*(\mathbf{r}) f_{12} \psi(\mathbf{r}) d\mathbf{r}$
 - $\mathcal{F} | \psi \rangle = \boldsymbol{\varepsilon} | \psi \rangle, \boldsymbol{\varepsilon} = (\varepsilon_0, \dots, \varepsilon_N)$
 - $N+1$ equations to be solved.

Methods: Hartree-Fock Method

- Integrate out spin
- Pair spins as: $\{\psi_{2I-1}, \psi_{2I}\} = \{\phi_I(\mathbf{r})\alpha(s), \phi_I(\mathbf{r})\beta(s)\}$

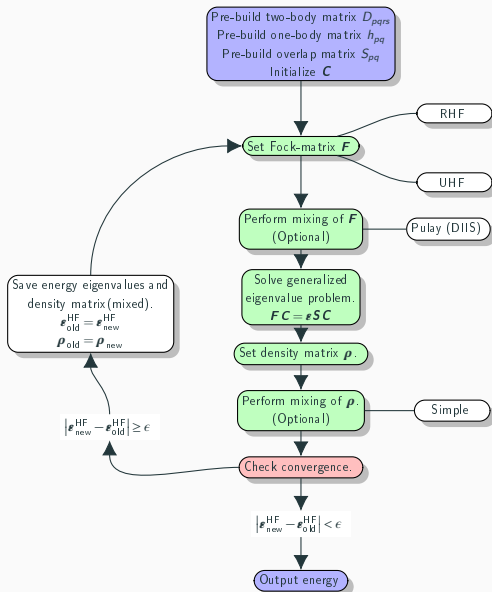
Methods: Hartree-Fock Method

- Integrate out spin
- Pair spins as: $\{\psi_{2I-1}, \psi_{2I}\} = \{\phi_I(\mathbf{r})\alpha(s), \phi_I(\mathbf{r})\beta(s)\}$
- Expand: $\phi_i(\mathbf{r}) = \sum_{p=1}^L C_{pi} \chi_p(\mathbf{r})$

Methods: Hartree-Fock Method

- Integrate out spin
- Pair spins as: $\{\psi_{2I-1}, \psi_{2I}\} = \{\phi_I(\mathbf{r})\alpha(s), \phi_I(\mathbf{r})\beta(s)\}$
- Expand: $\phi_i(\mathbf{r}) = \sum_{p=1}^L C_{pi} \chi_p(\mathbf{r})$
- Roothan-Hall: $\mathbf{F}\mathbf{C}_i = \epsilon \mathbf{S}\mathbf{C}_i$
 - $F_{pq} = h_{pq} + \sum_{pq} \rho_{pq} (2D_{prqs} \pm D_{prsq})$
 - $h_{pq} \equiv \langle p | h | q \rangle$
 - $\rho_{pq} \equiv \sum_{i=1}^{\frac{N}{2}} C_{pi} C_{qi}^*$
 - $D_{pqrs} \equiv \langle pq | f_{12} | rs \rangle$
 - $S_{pq} \equiv \langle p | q \rangle$
- Poople-Nesbet: $\mathbf{F}^+ \mathbf{C}^+ = \epsilon \mathbf{S} \mathbf{C}^+, \mathbf{F}^- \mathbf{C}^- = \epsilon^- \mathbf{S} \mathbf{C}^-$
 - $F_{pq}^{\pm} = h_{pq} + \sum_{k_{\pm}} \sum_{rs} C_{rk_{\pm}}^{\pm\dagger} C_{sk_{\pm}}^{\pm\dagger} [D_{prqs} - D_{prsq}] + \sum_{k_{\mp}} \sum_{rs} C_{rk_{\mp}}^{\mp\dagger} C_{sk_{\mp}}^{\mp\dagger} D_{prqs}$

Methods: Hartree-Fock Method



Implementation

Summary and Conclusion

Questions?

Questions

Questions?