

1 Vector Calculus

1.1 Preliminary ideas and some revision of vectors

1.1.1 The Einstein summation convention

In any product of terms, if we have a repeated suffix, then that quantity is considered to be summed over (from 1 to 3, since we will usually be working in three dimensions). For example

$$a_i x_i \text{ is shorthand for } \sum_{i=1}^3 a_i x_i. \quad \equiv \quad a_j x_j$$

1.1.2 The Kronecker delta

This is the quantity δ_{ij} and is defined such that

$$\delta_{ij} = \begin{cases} 1, & i = j; \\ 0, & i \neq j. \end{cases}$$

Example

$$\begin{aligned} \delta_{ij} a_j &= \sum_{j=1}^3 \delta_{ij} a_j = \delta_{i1} a_1 + \delta_{i2} a_2 + \delta_{i3} a_3 \\ &= a_i \end{aligned}$$

Note that the left-hand-side had two different subscripts, while the right-hand-side ends up with only one subscript - this is known as a **contraction**.

1.1.3 The permutation symbol

This is the quantity ε_{ijk} , defined as

$$\varepsilon_{ijk} = \begin{cases} 0, & \text{if any two of } i, j, k \text{ are the same;} \\ 1, & \text{if } i, j, k \text{ is a cyclic permutation of } 1, 2, 3; \\ -1, & \text{if } i, j, k \text{ is an acyclic permutation of } 1, 2, 3. \end{cases}$$

For example

$$\varepsilon_{123} = 1, \quad \varepsilon_{321} = -1, \quad \varepsilon_{133} = 0.$$

We can show, by considering the various cases, that the Kronecker delta and the permutation symbol are connected by the formula

$$\varepsilon_{ijk} \varepsilon_{klm} = \delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}.$$

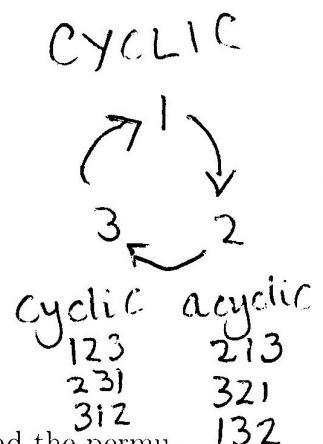
Sum over k

(I will put a proof on blackboard). The quantities δ_{ij} and ε_{ijk} are known as **tensors**.

Exercise: Show this can be rewritten in the alternative form

$$\varepsilon_{ijk} \varepsilon_{ilm} = \delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}.$$

Sum over i



1.1.4 Vector product

Recall that this is the multiplication of two vectors which results in a third vector, perpendicular to the first two. It can be written in the form of a determinant as

$$\mathbf{a} \times \mathbf{b} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}.$$

If $\mathbf{a} \times \mathbf{b} = \mathbf{0}$ then the two vectors are parallel. Recall that $(\mathbf{a} \times \mathbf{b}) = -(\mathbf{b} \times \mathbf{a})$. If we just consider the first component of this vector we can write this as

$$\begin{aligned} a_2 b_3 - a_3 b_2 &= \varepsilon_{123} a_2 b_3 + \varepsilon_{132} a_3 b_2 \\ &= \varepsilon_{ijk} a_j b_k \quad (\text{sum over } j \& k) \end{aligned}$$

since $\varepsilon_{123} = 1$, $\varepsilon_{132} = -1$, and $\varepsilon_{1ij} = 0$ for all other i and j . In general we can write the i th component of $\mathbf{a} \times \mathbf{b}$ as

$$[\mathbf{a} \times \mathbf{b}]_i = \varepsilon_{ijk} a_j b_k$$

1.1.5 Scalar product

This is defined as

$$\begin{aligned} \mathbf{a} \cdot \mathbf{b} &= a_1 b_1 + a_2 b_2 + a_3 b_3 \\ &= a_i b_i \end{aligned}$$

using the summation convention. Recall that if $\mathbf{a} \cdot \mathbf{b} = 0$ then the vectors \mathbf{a} and \mathbf{b} are orthogonal.

1.1.6 Triple scalar product

This is the quantity

$$\begin{aligned} \mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) &= a_i [\underline{\mathbf{b}} \times \underline{\mathbf{c}}]_i = a_i \varepsilon_{ijk} b_j c_k \\ &= \varepsilon_{ijk} a_i b_j c_k \quad \text{sum over } i, j \& k \end{aligned}$$

If this quantity is zero then the vectors $\mathbf{a}, \mathbf{b}, \mathbf{c}$ are coplanar. A useful property of the triple scalar product is that the dot and cross can be swapped without changing the answer, provided the order of the vectors remains unchanged, i.e.

$$\begin{aligned} a_i [\underline{\mathbf{b}} \times \underline{\mathbf{c}}]_i &\xrightarrow{\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}} a \cdot (\mathbf{b} \times \mathbf{c}) \\ &= \varepsilon_{ijk} a_i b_j c_k = (\varepsilon_{kij} a_i b_j) c_k = [\underline{\mathbf{a}} \times \underline{\mathbf{b}}]_k c_k \end{aligned}$$

1.1.7 Triple vector product

This is defined as

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}).$$

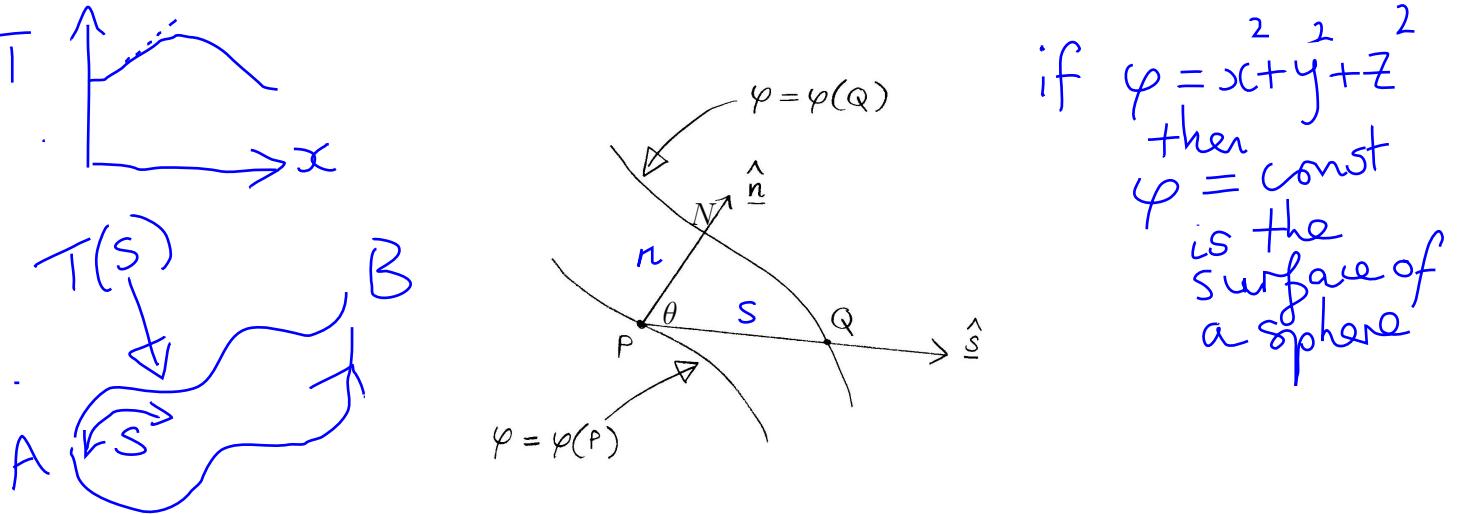
Since $\mathbf{b} \times \mathbf{c}$ is a vector normal to the plane of \mathbf{b} and \mathbf{c} , and $\mathbf{a} \times (\mathbf{b} \times \mathbf{c})$ is normal to $\mathbf{b} \times \mathbf{c}$, it follows that the triple vector product must lie in the plane of \mathbf{b} and \mathbf{c} . In component notation

$$\begin{aligned} [\mathbf{a} \times (\mathbf{b} \times \mathbf{c})]_i &= \varepsilon_{ijk} a_j [\underline{\mathbf{b}} \times \underline{\mathbf{c}}]_k \\ &= \varepsilon_{ijk} a_j \varepsilon_{klm} b_l c_m \\ &= (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) a_j b_l c_m \\ &= a_j b_i c_j - a_j b_j c_i \\ &= (\underline{\mathbf{a}} \cdot \underline{\mathbf{c}}) \mathbf{b}_i - (\underline{\mathbf{a}} \cdot \underline{\mathbf{b}}) \mathbf{c}_i \end{aligned}$$

and so we conclude that

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot \mathbf{c}) \mathbf{b} - (\mathbf{a} \cdot \mathbf{b}) \mathbf{c},$$

which confirms explicitly that the triple vector product indeed lies in the plane of \mathbf{b} and \mathbf{c} .

Figure 1: The surface $\phi = \text{constant}$ through two neighbouring points.

1.2 Gradient

Let ϕ be a differentiable scalar function of position in three dimensions. If P is a general point, ϕ will depend on the position of P , so we may write $\phi = \phi(P)$. The position of P is defined by reference to a coordinate system e.g. if we consider Cartesian coordinates, then P depends on (x, y, z) and hence $\phi = \phi(x, y, z)$, while if we consider cylindrical polar coordinates (r, θ, z) then $\phi = \phi(r, \theta, z)$.

The equation $\phi = \text{constant}$ defines a surface in three dimensions. Varying the constant, we can define a family of surfaces called ‘level surfaces’ or ‘equi- ϕ surfaces’. For example, if ϕ represents pressure, then $\phi = \text{constant}$ defines a family of surfaces over which the pressure is constant. The surface through a **specific point** P is $\phi = \phi(P)$. Let Q be a neighbouring point. (See figure 1). The equation of the level surface through Q is $\phi = \phi(Q)$. We draw the normal to $\phi = \phi(P)$ at P . Suppose that it intersects $\phi = \phi(Q)$ at the point N . Since N is on $\phi = \phi(Q)$ we have $\phi(N) = \phi(Q)$. Let s denote the length along PQ and let n denote the length along PN . Introduce unit vectors \hat{s} and \hat{n} in those directions. We define $\partial\phi/\partial s$ to be the **directional derivative** of ϕ in the direction \hat{s} :

$$\begin{aligned}
 \frac{\partial\phi}{\partial s} &= \lim_{PQ \rightarrow 0} \frac{(\phi(Q) - \phi(P))}{PQ} \\
 &= \lim_{Q \rightarrow P} \frac{(\phi(N) - \phi(P))}{PN} \cdot \frac{PN}{PQ} \\
 &= \lim_{N \rightarrow P} \frac{(\phi(N) - \phi(P))}{PN} \lim_{Q \rightarrow P} \left(\frac{PN}{PQ} \right) \\
 &= \frac{\partial\phi}{\partial n} \cos\theta = \frac{\partial\phi}{\partial n} (\hat{n} \cdot \hat{s})
 \end{aligned}$$

Since $\cos\theta \leq 1$, the maximum directional derivative at P occurs along the normal to $\phi = \phi(P)$ at P .

The vector $\hat{\mathbf{n}} \partial\phi/\partial n$ is called the **gradient** of ϕ at P . We write it as $\text{grad } \phi$ or $\nabla\phi$. The operator grad or ∇ is known as the **vector gradient operator**. We have

$$\frac{\partial\phi}{\partial s} = \hat{\mathbf{s}} \cdot \nabla\phi.$$

1.2.1 Cartesian components of $\nabla\phi$

If $\nabla\phi = A_1\mathbf{i} + A_2\mathbf{j} + A_3\mathbf{k}$ then $\mathbf{i} \cdot \nabla\phi = A_1$. But, by definition, $\mathbf{i} \cdot \nabla\phi = \partial\phi/\partial x$. Hence $A_1 = \partial\phi/\partial x$. Similarly we find $A_2 = \partial\phi/\partial y, A_3 = \partial\phi/\partial z$ and so we have the result:

$$\nabla\phi = \frac{\partial\phi}{\partial x} \hat{\mathbf{i}} + \frac{\partial\phi}{\partial y} \hat{\mathbf{j}} + \frac{\partial\phi}{\partial z} \hat{\mathbf{k}}$$

Example

If $\phi = axy^2 + byz + cx^3z^2$, where a, b, c are constants, find $\nabla\phi$. Also find the directional derivative of ϕ at the point $(1, 4, 2)$ in the direction towards the point $(2, 0, -1)$.

$$\nabla\phi = \hat{i}(ay^2 + 3cx^2z^2) + \hat{j}(2axy + bz) + \hat{k}(by + 2cx^3z)$$

$$P = (1, 4, 2)$$

$$(\nabla\phi)_P = \hat{i}(16a + 12c) + \hat{j}(8a + 2b) + \hat{k}(4b + 4c)$$

$$(1, 4, 2) \quad \underline{s} = (2, 0, -1) - (1, 4, 2) = (1, -4, -3)$$


$$\hat{s} = (\hat{i} - 4\hat{j} - 3\hat{k}) / \sqrt{1^2 + 4^2 + 3^2}$$

$$= (\hat{i} - 4\hat{j} - 3\hat{k}) / \sqrt{26}$$

$$\text{Directional derivative} = (\nabla\phi \cdot \hat{s})_P$$

$$= ((16a + 12c) - 4(8a + 2b) - 3(4b + 4c)) / \sqrt{26}$$

$$= (-16a - 20b) / \sqrt{26}$$

$$\begin{aligned}\hat{\theta} \cdot \hat{i} &= -\sin\theta \\ \hat{\theta} \cdot \hat{j} &= \cos\theta\end{aligned}$$

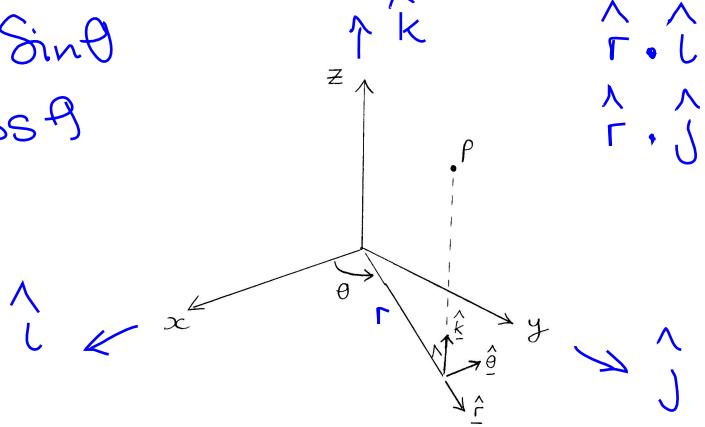


Figure 2: Sketch showing a point P represented by Cartesian coordinates (x, y, z) and cylindrical polar coordinates (r, θ, z) . $x = r \cos\theta$ $y = r \sin\theta$

1.2.2 Cylindrical polar components of $\nabla\phi$

The set-up is as shown in figure 2. We write $\nabla\phi = A_1\hat{r} + A_2\hat{\theta} + A_3\hat{k}$. Then it follows that

$$\begin{aligned}A_1 &= \hat{r} \cdot \nabla\phi \\ &= \hat{i} \cdot \left(\frac{\partial\phi}{\partial x} \hat{i} + \frac{\partial\phi}{\partial y} \hat{j} + \frac{\partial\phi}{\partial z} \hat{k} \right) \\ &= \cos\theta \frac{\partial\phi}{\partial x} + \sin\theta \frac{\partial\phi}{\partial y} + \text{zero} \quad (\hat{i} \cdot \hat{k} = 0) \\ &= \frac{\partial x}{\partial r} \frac{\partial\phi}{\partial x} + \frac{\partial y}{\partial r} \frac{\partial\phi}{\partial y} \\ &= \frac{1}{r} \frac{\partial\phi}{\partial r}\end{aligned}$$

Similarly, we find

$$\begin{aligned}A_2 &= \hat{\theta} \cdot \nabla\phi \\ &= \hat{\theta} \cdot \left(\frac{\partial\phi}{\partial x} \hat{i} + \frac{\partial\phi}{\partial y} \hat{j} + \frac{\partial\phi}{\partial z} \hat{k} \right) \\ &= -\sin\theta \frac{\partial\phi}{\partial x} + \cos\theta \frac{\partial\phi}{\partial y} + \text{zero} \quad (\hat{\theta} \cdot \hat{k} = 0) \\ &= \frac{1}{r} \frac{\partial x}{\partial \theta} \frac{\partial\phi}{\partial x} + \frac{1}{r} \frac{\partial y}{\partial \theta} \frac{\partial\phi}{\partial y} \\ &= \frac{1}{r} \frac{\partial\phi}{\partial \theta}\end{aligned}$$

and $A_3 = \hat{k} \cdot \nabla\phi = \partial\phi/\partial z$. Hence

$$\nabla\phi = \hat{r} \frac{\partial\phi}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial\phi}{\partial \theta} + \hat{k} \frac{\partial\phi}{\partial z}.$$



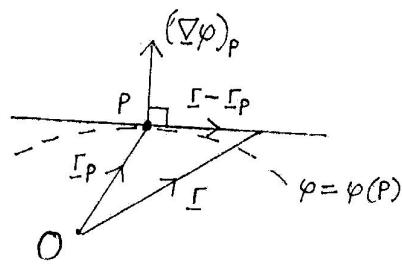


Figure 3: The tangent plane to a surface.

1.2.3 Equation of a tangent plane to $\phi = \phi(P)$

We have that $(\nabla \phi)_P$ is normal to $\phi = \phi(P)$ at P . The equation of the tangent plane is therefore

$$(\mathbf{r} - \mathbf{r}_P) \cdot (\nabla \phi)_P = 0,$$

i.e.

$$\left(\frac{\partial \phi}{\partial x} \right)_P (x - x_P) + \left(\frac{\partial \phi}{\partial y} \right)_P (y - y_P) + \left(\frac{\partial \phi}{\partial z} \right)_P (z - z_P) = 0.$$

↑
 constants

Example

Find the tangent plane to the surface

$$z = e^{-(x^2+y^2)^{1/2}}$$

at the point $x = -1, y = 0$.

N.B. choice of φ
not unique

e.g. $\varphi = z e^{-(x^2+y^2)^{1/2}} = 1$
on S

$$\text{Let } \varphi = z - e^{-(x^2+y^2)^{1/2}}$$

$= 0$ on the
surface

$$\text{Then } \frac{\partial \varphi}{\partial x} = \frac{x}{e^{-(x^2+y^2)^{1/2}}}$$

$$\& \frac{\partial \varphi}{\partial y} = \frac{y}{e^{-(x^2+y^2)^{1/2}}}$$

$$\& \frac{\partial \varphi}{\partial z} = 1$$

P is the
point
 $x = -1 = x_p$
 $y = 0 = y_p$
 $\& z = e^{-1} = z_p$

$$\Rightarrow \left. \frac{\partial \varphi}{\partial x} \right|_P = -e^{-1}; \left. \frac{\partial \varphi}{\partial y} \right|_P = 0; \left. \frac{\partial \varphi}{\partial z} \right|_P = 1.$$

So eqn of tgt plane

$$-e^{-1}(x - (-1)) + 0 + (1)(z - e^{-1}) = 0$$

$$\Rightarrow z = e^{-1}(2+x)$$



1.3 Divergence and Curl

In this section we will assume that \mathbf{A} is a vector function of position in three dimensions, with continuous first partial derivatives.

Since ∇ is a vector operator, we can define formally a scalar product $\nabla \cdot \mathbf{A}$. This is called the **divergence** of the vector \mathbf{A} . We can also define the vector product $\nabla \times \mathbf{A}$, which is called the **curl** of \mathbf{A} . So to summarize we have

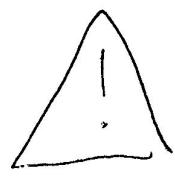
$$\operatorname{div} \mathbf{A} = \nabla \cdot \mathbf{A}, \quad \operatorname{curl} \mathbf{A} = \nabla \times \mathbf{A}.$$

1.3.1 Cartesian form

$$\begin{aligned}\operatorname{div} \mathbf{A} &= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot (A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}) \\ &= \frac{\partial A_1}{\partial x} + \frac{\partial A_2}{\partial y} + \frac{\partial A_3}{\partial z} \quad (\text{SCALAR}) \\ \operatorname{curl} \mathbf{A} &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_1 & A_2 & A_3 \end{vmatrix} \\ (\text{VECTOR}) &= \hat{i} \left(\frac{\partial A_3}{\partial y} - \frac{\partial A_2}{\partial z} \right) - \hat{j} \left(\frac{\partial A_3}{\partial x} - \frac{\partial A_1}{\partial z} \right) + \hat{k} \left(\frac{\partial A_2}{\partial x} - \frac{\partial A_1}{\partial y} \right)\end{aligned}$$

Note that these simple forms for div and curl arise because $\hat{i}, \hat{j}, \hat{k}$ are constant vectors: this is not so in other coordinate systems.

N.B. Note well!



$$\underline{\mathbf{A}} \cdot \underline{\nabla} \neq \underline{\nabla} \cdot \underline{\mathbf{A}} = \operatorname{div} \underline{\mathbf{A}}$$

(scalar)

$$\begin{aligned}(A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}) \cdot \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \\ = A_1 \frac{\partial}{\partial x} + A_2 \frac{\partial}{\partial y} + A_3 \frac{\partial}{\partial z}\end{aligned}$$

~~Scalar~~ operator

Similarly:

$$\underline{\mathbf{A}} \times \underline{\nabla} \neq \underline{\nabla} \times \underline{\mathbf{A}}$$

↑
vector operator

or
 $-\underline{\nabla} \times \underline{\mathbf{A}}$

Examples

(a) If

$$\mathbf{A} = (y^2 \cos x + z^3)\mathbf{i} + (2y \sin x - 4)\mathbf{j} + (3xz^2 + 2)\mathbf{k}$$

find $\operatorname{div} \mathbf{A}$ and $\operatorname{curl} \mathbf{A}$.(b) Find $\operatorname{div} \mathbf{u}$ and $\operatorname{curl} \mathbf{u}$ when (i) $\mathbf{u} = \mathbf{r}$; (ii) $\mathbf{u} = \boldsymbol{\omega} \times \mathbf{r}$, where $\mathbf{r} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k}$, and $\boldsymbol{\omega} = \Omega \mathbf{k}$ with Ω constant.

$$\begin{aligned}\operatorname{div} \mathbf{A} &= \frac{\partial}{\partial x}(y^2 \cos x + z^3) + \frac{\partial}{\partial y}(2y \sin x - 4) + \frac{\partial}{\partial z}(3xz^2 + 2) \\ &= -y^2 \sin x + 2 \sin x + 6xz \\ \operatorname{curl} \mathbf{A} &= \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y^2 \cos x + z^3 & 2y \sin x - 4 & 3xz^2 + 2 \end{vmatrix} = \hat{\mathbf{i}}(0) \\ &\quad - \hat{\mathbf{j}}(3z^2 - 3z^2) \\ &\quad + \hat{\mathbf{k}}(2y \cos x - 2y \cos x) \\ &= \underline{\Omega} \text{ (vector)}\end{aligned}$$

We say
that
 $\underline{\mathbf{A}}$ is IRRATIONAL.

$$\begin{aligned}(\text{b}) (\text{i}) \quad \underline{\mathbf{u}} &= x\hat{\mathbf{i}} + y\hat{\mathbf{j}} + z\hat{\mathbf{k}} \\ \operatorname{div} \underline{\mathbf{u}} &= 3 \quad (\text{dimension of space}) \\ \operatorname{curl} \underline{\mathbf{u}} &= \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ x & y & z \end{vmatrix} = \underline{\Omega}\end{aligned}$$

$$\begin{aligned}(\text{ii}) \quad \underline{\mathbf{u}} &= \underline{\boldsymbol{\omega}} \times \underline{\mathbf{r}} \quad \underline{\boldsymbol{\omega}} = \underline{\Omega} \hat{\mathbf{k}} \\ &= \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ 0 & 0 & \underline{\Omega} \\ x & y & z \end{vmatrix} = -\underline{\Omega} y \hat{\mathbf{i}} + \underline{\Omega} x \hat{\mathbf{j}}\end{aligned}$$

$\operatorname{div} \underline{\mathbf{u}} = 0$ (SOLENOIDAL FIELD)

$$\operatorname{curl} \underline{\mathbf{u}} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -\underline{\Omega} y & \underline{\Omega} x & 0 \end{vmatrix} = \hat{\mathbf{k}}(-\underline{\Omega} - (-\underline{\Omega})) = 2\underline{\Omega} \hat{\mathbf{k}}$$

$\operatorname{curl} \underline{\mathbf{u}}$ is related to
ROTATION

1.4 Operations with the gradient operator

1.4.1 Important sum and product formulae

Note that ∇ is a linear operator, and so:

- (i) $\nabla(\phi_1 + \phi_2) = \nabla\phi_1 + \nabla\phi_2,$
- (ii) $\operatorname{div}(\mathbf{A} + \mathbf{B}) = \operatorname{div}\mathbf{A} + \operatorname{div}\mathbf{B},$
- (iii) $\operatorname{curl}(\mathbf{A} + \mathbf{B}) = \operatorname{curl}\mathbf{A} + \operatorname{curl}\mathbf{B}.$

The proofs of these results follow immediately from the definition of ∇ .

Other key results are:

- (iv) $\nabla(\phi\psi) = \phi\nabla\psi + \psi\nabla\phi,$
- (v) $\operatorname{div}(\phi\mathbf{A}) = \phi\operatorname{div}\mathbf{A} + \nabla\phi \cdot \mathbf{A}.$

$$\begin{aligned}\nabla\phi \cdot \underline{\mathbf{A}} &\equiv \underline{\mathbf{A}} \cdot \nabla\phi \\ &= (\underline{\mathbf{A}} \cdot \nabla)\phi\end{aligned}$$

Proof of (v)

$$\begin{aligned}\operatorname{div}(\phi\mathbf{A}) &= \left(\frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z} \right) \cdot (\phi A_1 \hat{i} + \phi A_2 \hat{j} + \phi A_3 \hat{k}) \\ &= \frac{\partial}{\partial x}(\phi A_1) + \frac{\partial}{\partial y}(\phi A_2) + \frac{\partial}{\partial z}(\phi A_3) \\ &= \underbrace{\phi \left[\frac{\partial A_1}{\partial x} + \frac{\partial A_2}{\partial y} + \frac{\partial A_3}{\partial z} \right]}_{\phi \operatorname{div} \underline{\mathbf{A}}} + \underbrace{A_1 \frac{\partial \phi}{\partial x} + A_2 \frac{\partial \phi}{\partial y} + A_3 \frac{\partial \phi}{\partial z}}_{\nabla\phi \cdot \underline{\mathbf{A}}}\end{aligned}$$

In writing out these proofs it is easier to use the **summation convention** that we introduced earlier. Rather than write (x, y, z) for Cartesian components, we write (x_1, x_2, x_3) and in place of $(\mathbf{i}, \mathbf{j}, \mathbf{k})$ we write $(\hat{\mathbf{e}}_1, \hat{\mathbf{e}}_2, \hat{\mathbf{e}}_3)$. Then we saw earlier that

$$\begin{aligned}\mathbf{A} \cdot \mathbf{B} &= A_i B_i, \\ \mathbf{A} \times \mathbf{B} &= \epsilon_{ijk} \hat{\mathbf{e}}_i A_j B_k\end{aligned}$$

$$[\underline{\mathbf{A}} \times \underline{\mathbf{B}}]_i = \epsilon_{ijk} A_j B_k$$

Also recall the useful result that

$$\epsilon_{ijk} \epsilon_{klm} = \delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}.$$

Thus, under the summation convention:

$$\begin{aligned}\operatorname{div}\mathbf{A} &= \frac{\partial A_i}{\partial x_i} \\ [\nabla\phi]_i &= \frac{\partial \phi}{\partial x_i} \\ [\operatorname{curl} \mathbf{A}]_i &= \epsilon_{ijk} \frac{\partial}{\partial x_j} A_k\end{aligned}\quad \left(\nabla\phi = \hat{\mathbf{e}}_i \frac{\partial \phi}{\partial x_i} \right)$$

where $[\]_i$ indicates the i th component. Using this approach, the proof of (v) takes the form

$$\begin{aligned}\operatorname{div}(\phi\mathbf{A}) &= \frac{\partial}{\partial x_i}(\phi A_i) = \phi \frac{\partial A_i}{\partial x_i} + A_i \frac{\partial \phi}{\partial x_i} \\ &= \phi \operatorname{div} \underline{\mathbf{A}} + (\underline{\mathbf{A}} \cdot \nabla)\phi\end{aligned}$$

Other important results are:

- (vi) $\operatorname{curl}(\phi \mathbf{A}) = \phi \operatorname{curl} \mathbf{A} + \nabla \phi \times \mathbf{A}$,
- (vii) $\operatorname{div}(\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot \operatorname{curl} \mathbf{A} - \mathbf{A} \cdot \operatorname{curl} \mathbf{B}$,
- (viii) $\operatorname{curl}(\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla) \mathbf{A} - \mathbf{B} \operatorname{div} \mathbf{A} - (\mathbf{A} \cdot \nabla) \mathbf{B} + \mathbf{A} \operatorname{div} \mathbf{B}$,
- (ix) $\nabla(\mathbf{A} \cdot \mathbf{B}) = (\mathbf{B} \cdot \nabla) \mathbf{A} + (\mathbf{A} \cdot \nabla) \mathbf{B} + \mathbf{B} \times \operatorname{curl} \mathbf{A} + \mathbf{A} \times \operatorname{curl} \mathbf{B}$.

Example

Prove relation (ix) above. If we work on the RHS we can write

$$\begin{aligned}
 & \underline{\mathbf{B}} \cdot \underline{\nabla} [(\mathbf{B} \cdot \nabla) \mathbf{A} + (\mathbf{A} \cdot \nabla) \mathbf{B} + \mathbf{B} \times \operatorname{curl} \mathbf{A} + \mathbf{A} \times \operatorname{curl} \mathbf{B}]_i \\
 &= \underbrace{B_j \frac{\partial}{\partial x_j} A_i}_{= B_j \frac{\partial A_i}{\partial x_j}} + A_j \frac{\partial}{\partial x_j} B_i + \epsilon_{ijk} B_j (\operatorname{curl} \underline{\mathbf{A}})_k + \epsilon_{ijk} A_j (\operatorname{curl} \underline{\mathbf{B}})_k \\
 &= \dots + \epsilon_{ijk} \left\{ B_j \epsilon_{klm} \frac{\partial A_m}{\partial x_l} + A_j \epsilon_{klm} \frac{\partial B_m}{\partial x_l} \right\} \\
 &= \dots + (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) \left(B_j \frac{\partial A_m}{\partial x_l} + A_j \frac{\partial B_m}{\partial x_l} \right) \\
 &= \cancel{B_j \frac{\partial A_i}{\partial x_j}} + \cancel{A_j \frac{\partial B_i}{\partial x_j}} + B_j \left(\cancel{\frac{\partial A_j}{\partial x_i}} - \cancel{\frac{\partial A_i}{\partial x_j}} \right) + A_j \left(\cancel{\frac{\partial B_j}{\partial x_i}} - \cancel{\frac{\partial B_i}{\partial x_j}} \right) \\
 &= B_j \frac{\partial A_j}{\partial x_i} + A_j \frac{\partial B_j}{\partial x_i} \\
 &= \cancel{\frac{\partial}{\partial x_i}} (A_j B_j) \\
 &= \frac{\partial}{\partial x_i} (\underline{\mathbf{A}} \cdot \underline{\mathbf{B}}) = [\underline{\nabla} (\underline{\mathbf{A}} \cdot \underline{\mathbf{B}})]_i
 \end{aligned}$$

as required.

Note: In the following sections we will assume that our scalar and vector functions possess continuous second derivatives.

1.4.2 The divergence of a gradient: the Laplacian

Consider the operation

$$\begin{aligned}\operatorname{div}(\nabla\phi) &= \left(\mathbf{i}\frac{\partial}{\partial x} + \mathbf{j}\frac{\partial}{\partial y} + \mathbf{k}\frac{\partial}{\partial z}\right) \cdot \left(\frac{\partial\phi}{\partial x}\mathbf{i} + \frac{\partial\phi}{\partial y}\mathbf{j} + \frac{\partial\phi}{\partial z}\mathbf{k}\right) \\ &= \frac{\partial^2\phi}{\partial x^2} + \frac{\partial^2\phi}{\partial y^2} + \frac{\partial^2\phi}{\partial z^2} \\ &\equiv \nabla^2\phi = \frac{\partial^2\phi}{\partial x_i\partial x_i} = \frac{\partial^2\phi}{\partial x_i^2}\end{aligned}$$

repeated subscript

This is to be read as 'del squared ϕ ' or the **Laplacian** of ϕ . The operator ∇^2 is known as the Laplacian operator. We also define the Laplacian of a vector as

$$\nabla^2\mathbf{A} \equiv \frac{\partial^2\mathbf{A}}{\partial x^2} + \frac{\partial^2\mathbf{A}}{\partial y^2} + \frac{\partial^2\mathbf{A}}{\partial z^2}$$

in Cartesian coordinates, and the equation $\nabla^2\phi = 0$ is known as **Laplace's equation**.

Example

If $\phi = x^2 + y^2$, find $\nabla^2\phi$.

$$\nabla^2\phi = f(x, y, z)$$

given POISSON'S EQUATION

$$\frac{\partial^2\phi}{\partial x^2} = 2 = \frac{\partial^2\phi}{\partial y^2}; \quad \frac{\partial^2\phi}{\partial z^2} = 0$$

$$\Rightarrow \nabla^2\phi = 4 \quad \text{so } x^2+y^2 \text{ is a soln of Poisson's eqn}$$

A180

$$\phi = x^2 - y^2 \Rightarrow \nabla^2\phi = 0 \Rightarrow x^2 - y^2 \text{ is a soln of Laplace's eqn}$$

1.4.3 The curl of a gradient

Consider the operation

$$\begin{aligned}
 [\text{curl}(\nabla\phi)]_i &= \varepsilon_{ijk} \frac{\partial}{\partial x_j} (\nabla\phi)_k = \varepsilon_{ijk} \frac{\partial}{\partial x_j} \left(\frac{\partial \phi}{\partial x_k} \right) \\
 &\equiv \frac{1}{2} \varepsilon_{ijk} \frac{\partial}{\partial x_j} \left(\frac{\partial \phi}{\partial x_k} \right) + \frac{1}{2} \varepsilon_{ikj} \frac{\partial}{\partial x_k} \left(\frac{\partial \phi}{\partial x_j} \right) \quad (\varepsilon_{ijk} = -\varepsilon_{ikj}) \\
 &= \frac{1}{2} \varepsilon_{ijk} \frac{\partial}{\partial x_j} \left(\frac{\partial \phi}{\partial x_k} \right) - \frac{1}{2} \varepsilon_{ijk} \frac{\partial}{\partial x_k} \left(\frac{\partial \phi}{\partial x_j} \right) \\
 &= \frac{1}{2} \varepsilon_{ijk} \left\{ \frac{\partial}{\partial x_j} \left(\frac{\partial \phi}{\partial x_k} \right) - \frac{\partial}{\partial x_k} \left(\frac{\partial \phi}{\partial x_j} \right) \right\} = 0 \\
 &\text{(This result can also be established by using tensor notation).} \quad (\text{partial derivatives commute})
 \end{aligned}$$

Example

Consider $\phi = axy^2 + byz + cx^3z^2$ and show explicitly that $\text{curl } \nabla\phi = 0$.

$$\begin{aligned}
 \nabla\phi &= \hat{i}(ay^2 + 3cx^2z^2) + \hat{j}(2axy + bz) \\
 &\quad + \hat{k}(by + 2cx^3z) \\
 \text{curl } (\nabla\phi) &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ ay^2 + 3cx^2z^2 & 2axy + bz & by + 2cx^3z \end{vmatrix} \\
 &= \hat{i}(b - b) - \hat{j}(6cx^2z - 6cx^2z) \\
 &\quad + \hat{k}(2ay - 2ay) \\
 &= 0 \quad \text{as expected.}
 \end{aligned}$$

1.4.4 The divergence of a curl

This is also always zero, as can be seen from the following argument:

$$\begin{aligned}
 \text{div}(\text{curl } \underline{A}) &= \frac{\partial}{\partial x_i} (\text{curl } \underline{A})_i = \varepsilon_{ijk} \frac{\partial}{\partial x_i} \left(\frac{\partial A_k}{\partial x_j} \right) \\
 &\equiv \frac{1}{2} \varepsilon_{ijk} \frac{\partial}{\partial x_i} \left(\frac{\partial A_k}{\partial x_j} \right) + \frac{1}{2} \varepsilon_{jik} \frac{\partial}{\partial x_j} \left(\frac{\partial A_k}{\partial x_i} \right) \\
 &= \frac{1}{2} \varepsilon_{ijk} \left\{ \frac{\partial}{\partial x_i} \left(\frac{\partial A_k}{\partial x_j} \right) - \frac{\partial}{\partial x_j} \left(\frac{\partial A_k}{\partial x_i} \right) \right\} \\
 &= 0 \quad (\text{since partial derivatives commute})
 \end{aligned}$$

Example

Verify that $\text{div}(\text{curl } \underline{A}) = 0$ for the quantity $\underline{A} = y e^x \mathbf{i} + (x^2 + z) \mathbf{j} + y^3 \cos(zx) \mathbf{k}$.

$$\begin{aligned}
 \text{curl } \underline{A} &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y e^x & x^2 + z & y^3 \cos(zx) \end{vmatrix} \\
 &= \hat{i} (3y^2 \cos(zx) - 1) \\
 &\quad - \hat{j} (-y^3 z \sin(zx)) + \hat{k} (2x - e^x)
 \end{aligned}$$

Then

$$\begin{aligned}
 \text{div}(\text{curl } \underline{A}) &= -3y^2 z \sin(zx) + 3y^2 z \sin(zx) + 0 \\
 &= 0 \quad \text{as required.}
 \end{aligned}$$

1.4.5 The curl of a curl

This is the vector quantity

$$\text{curl}(\text{curl } \mathbf{A}).$$

Using tensor notation and the summation convention we can show that

$$\text{curl}(\text{curl } \mathbf{A}) = \nabla(\text{div } \mathbf{A}) - \nabla^2 \mathbf{A}.$$

$$\begin{aligned}\text{Proof } [\text{curl}(\text{curl } \underline{\mathbf{A}})]_i &= \varepsilon_{ijk} \frac{\partial}{\partial x_j} (\text{curl } \underline{\mathbf{A}})_k \\ &= \varepsilon_{ijk} \frac{\partial}{\partial x_j} (\varepsilon_{klm} \frac{\partial A_m}{\partial x_l}) \\ &= (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) \left(\frac{\partial^2 A_m}{\partial x_j \partial x_l} \right) \\ &= \frac{\partial^2 A_j}{\partial x_j \partial x_i} - \frac{\partial^2 A_i}{\partial x_j \partial x_j} \\ &= \frac{\partial}{\partial x_i} \left(\frac{\partial A_j}{\partial x_j} \right) - \frac{\partial^2 A_i}{\partial x_j^2} \\ &= [\nabla(\text{div } \underline{\mathbf{A}})]_i - [\nabla^2 \underline{\mathbf{A}}]_i\end{aligned}$$

Exercise

Calculate $\text{curl}(\text{curl } \mathbf{A})$, $\nabla(\text{div } \mathbf{A})$ and $\nabla^2 \mathbf{A}$ for $\mathbf{A} = y e^x \mathbf{i} + (x^2 + z) \mathbf{j} + y^3 \cos(zx) \mathbf{k}$.

Answers:

$$\begin{aligned}\text{curl}(\text{curl } \underline{\mathbf{A}}) &= (-y^3 \sin zx - y^3 zx \cos zx) \hat{i} \\ &\quad - \hat{j} (2e^x + 3y^2 x \sin zx) \\ &\quad + \hat{k} (y^3 z^2 \cos zx - 6y \cos zx) \\ \nabla(\text{div } \underline{\mathbf{A}}) &= \hat{i} (ye^x - y^3 \sin zx - xy^3 z \cos zx) \\ &\quad + \hat{j} (e^x - 3xy^2 \sin zx) + \hat{k} (-x^2 y^3 \cos zx) \\ \nabla^2 \underline{\mathbf{A}} &= \hat{i} (ye^x) + 2\hat{j} + \hat{k} (-y^3 z^2 \cos zx + 6y \cos zx) \\ &\quad - y^3 x^2 \cos zx\end{aligned}$$

1.4.6 Scalar and vector fields

If, at each point of a region V of space, a scalar function ϕ is defined, we say that ϕ is a **scalar field** over the region V . Similarly, if a vector function \mathbf{A} is also defined at all points of V , then \mathbf{A} is a vector field over V . If $\text{curl } \mathbf{A} = \mathbf{0}$ we say that A is an **irrotational** vector field. If $\text{div } \mathbf{A} = 0$ we say \mathbf{A} is a **solenoidal** vector field. An obvious example of a vector field is the position vector \mathbf{r} of a point in space. In three dimensions:

$$\mathbf{r} = x\mathbf{i} + y\mathbf{j} + z\mathbf{k},$$

$$\begin{aligned}\text{div } \mathbf{r} &= 3 \\ \text{curl } \mathbf{r} &= \left| \begin{array}{ccc} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ x & y & z \end{array} \right| = 0\end{aligned}$$

$$|\mathbf{r}| = r = (x^2 + y^2 + z^2)^{1/2}$$

$$\nabla r = \nabla(x^2 + y^2 + z^2)^{1/2}$$

$$\begin{aligned}&= \left(\hat{\mathbf{i}} \frac{\partial}{\partial x} + \hat{\mathbf{j}} \frac{\partial}{\partial y} + \hat{\mathbf{k}} \frac{\partial}{\partial z} \right) (x^2 + y^2 + z^2)^{1/2} \\ &= (\hat{x}\mathbf{i} + \hat{y}\mathbf{j} + \hat{z}\mathbf{k}) (x^2 + y^2 + z^2)^{-1/2} \\ &= \underline{\Gamma} / r \\ &= \underline{\Gamma}\end{aligned}$$

Example

Find

$$\nabla^2(1/r)$$

$$(r \neq 0)$$

This is $\nabla \cdot (\nabla(1/r))$

First calculate

$$\begin{aligned} \nabla(1/r) &= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) (x^2 + y^2 + z^2)^{-1/2} \\ &= -(\hat{x} + \hat{y} + \hat{z})(x^2 + y^2 + z^2)^{-3/2} \end{aligned}$$

$= -\frac{\underline{r}}{r^3}$

Now

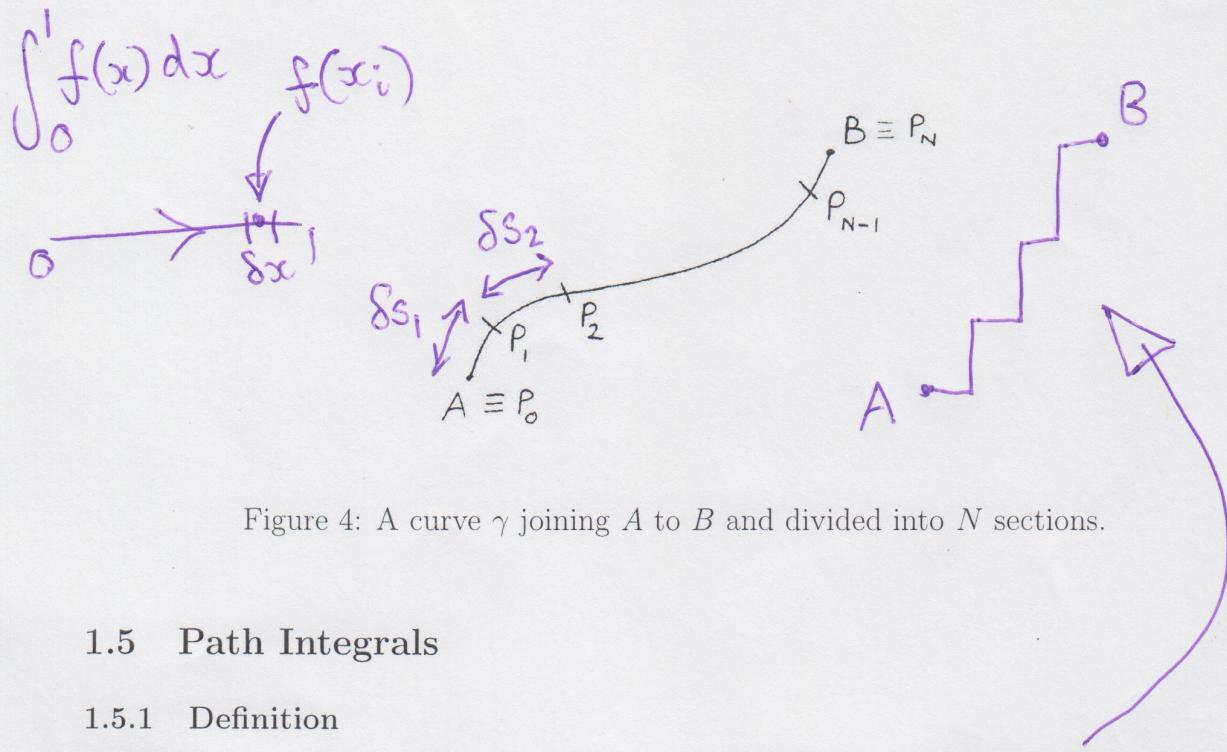
$$\begin{aligned} \nabla \cdot (\nabla(1/r)) &= - \left\{ \frac{\partial}{\partial x} \left(\frac{x}{(x^2 + y^2 + z^2)^{3/2}} \right) + \frac{\partial}{\partial y} \left(\frac{y}{(x^2 + y^2 + z^2)^{3/2}} \right) \right. \\ &\quad \left. + \frac{\partial}{\partial z} \left(\frac{z}{(x^2 + y^2 + z^2)^{3/2}} \right) \right\} \\ &= -\frac{3}{(x^2 + y^2 + z^2)^{3/2}} - \left\{ x \left(-\frac{3}{2} \right) (2x) + y \left(-\frac{3}{2} \right) (2y) \right. \\ &\quad \left. + z \left(-\frac{3}{2} \right) (2z) \right\} \\ &= -\frac{3}{(x^2 + y^2 + z^2)^{3/2}} + \frac{3(x^2 + y^2 + z^2)}{(x^2 + y^2 + z^2)^{5/2}} = 0 \end{aligned}$$

$\therefore \nabla^2(1/r) = 0 \quad (r \neq 0) \Rightarrow \frac{1}{r}$ satisfies Laplace eqn.

In fact $\frac{1}{|\underline{r} - \underline{r}_0|}$ is also a soln. i.e. $\nabla^2 \left(\frac{1}{|\underline{r} - \underline{r}_0|} \right) = 0$

\underline{r}_0 const vector

IMPORTANT FOR
DEVELOPMENT OF PDE THEORY
& GREEN'S FUNCTION (Partial Diff. Eqns)

Figure 4: A curve γ joining A to B and divided into N sections.

1.5 Path Integrals

1.5.1 Definition

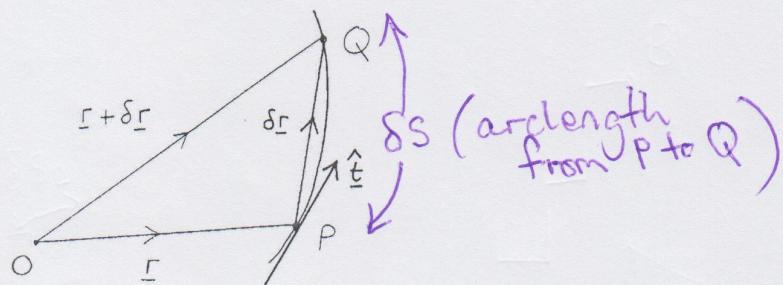
Consider a curve γ (not necessarily in the plane, and not necessarily smooth) joining the points A and B . (See figure 4). Suppose that the curve is divided into N sections: $AP_1, P_1P_2, \dots, P_{N-1}B$. Let $AP_1 = \delta s_1, P_1P_2 = \delta s_2, \dots, P_{N-1}B = \delta s_N$. Next, suppose a function f is defined along this curve γ . We compute the sum

$$f_1\delta s_1 + f_2\delta s_2 + \dots + f_N\delta s_N,$$

where $f_n = f(P_n)$. On increasing N indefinitely, while letting the maximum $\delta s_n \rightarrow 0$, the resulting limit of the sum, if it exists, is called the **path integral** of f along γ , and we write:

$$\int_{\gamma} f ds = \lim_{\substack{N \rightarrow \infty \\ \max(\delta s_n) \rightarrow 0}} \sum_{n=1}^N f_n \delta s_n$$

The function f may be a scalar or a vector.

Figure 5: Diagram showing the tangent vector at a point P .

1.5.2 Path element

See figure 5. Let δs represent the arc PQ and suppose that the vector $\overrightarrow{PQ} = \delta \mathbf{r}$. We define the tangent vector

$$\hat{\mathbf{t}} = \frac{d\mathbf{r}}{ds} = \lim_{\delta s \rightarrow 0} \frac{\delta \mathbf{r}}{\delta s},$$

and the path element

$$d\mathbf{r} = \hat{\mathbf{t}} ds.$$

In Cartesian $d\gamma = dx\hat{i} + dy\hat{j} + dz\hat{k}$

Note that $\hat{\mathbf{t}}$ has length unity because $|\delta \mathbf{r}| \rightarrow \delta s$ as $\delta s \rightarrow 0$. We can then define the quantity

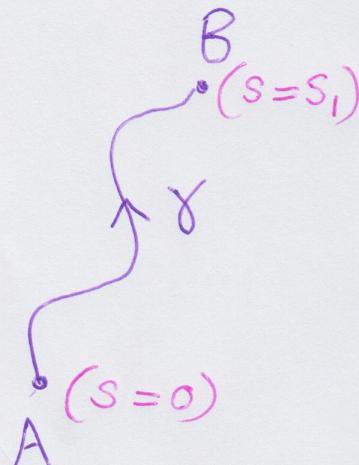
$$\int_{\gamma} \mathbf{F} \cdot d\mathbf{r} = \int_{\gamma} (\mathbf{F} \cdot \hat{\mathbf{t}}) ds$$

1.5.3 Conservative forces

Consider the special case where we have a vector \mathbf{F} of the form

$$\mathbf{F} = \nabla\phi$$

with ϕ a differentiable scalar function. Consider the integral (with γ defined as in figure 3):

$$\begin{aligned}
 \int_{\gamma} \mathbf{F} \cdot d\mathbf{r} &= \int_{\gamma} (\nabla\phi \cdot \hat{\mathbf{t}}) ds \\
 &= \int_{\gamma} \frac{\partial\phi}{\partial x_i} \hat{\mathbf{e}}_i \cdot \frac{dx}{ds} ds \\
 &= \int_{\gamma} \frac{\partial\phi}{\partial x_i} \hat{\mathbf{e}}_i \cdot \frac{dx_j}{ds} \hat{\mathbf{e}}_j ds \\
 &= \int_{\gamma} \frac{\partial\phi}{\partial x_i} \frac{dx_i}{ds} ds \\
 &= \int_{\gamma} \frac{d\phi}{ds} ds \quad (\text{Chain rule}). \\
 &= \int_0^{s_1} \frac{d\phi}{ds} ds \\
 &= [\phi]_A^B \\
 &= \phi(B) - \phi(A)
 \end{aligned}$$


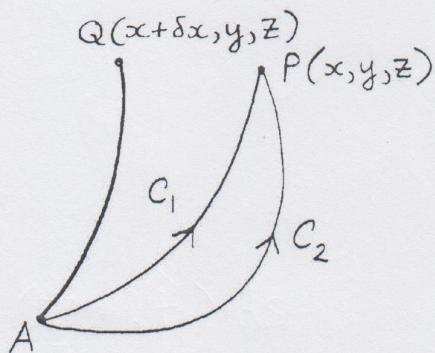
We note that the result is independent of the path γ joining A to B . In particular, if γ is a closed curve (i.e. $B \equiv A$), then we have $\oint_{\gamma} \mathbf{F} \cdot d\mathbf{r} = 0$, where we put a circle on the integral to denote the path is closed. We sometimes refer to such an integral as the circulation of \mathbf{F} around γ .

closed curve
 $\oint_{\gamma} \nabla\phi \cdot d\mathbf{r} = 0$

 Closed curve
 $B \equiv A$

CIRCULATION
of \mathbf{F} around closed curve γ

$\oint_{\gamma} \mathbf{F} \cdot d\mathbf{r}$ (\mathbf{F} not necessarily equal to $\nabla\phi$)

Figure 6: Two curves joining A to P . Q is a neighbouring point.

If a vector field \mathbf{F} has the property that $\oint_{\gamma} \mathbf{F} \cdot d\mathbf{r} = 0$ for **any** closed curve γ , we say that \mathbf{F} is a **conservative field**. Thus, if $\mathbf{F} = \nabla\phi$, then \mathbf{F} is conservative. Conversely, if \mathbf{F} is conservative we can always find a differentiable scalar function ϕ such that $\mathbf{F} = \nabla\phi$. The function ϕ is called the **potential** of the field \mathbf{F} .

Proof of this last part

See figure 6. Let $\mathbf{F} = F_1\mathbf{i} + F_2\mathbf{j} + F_3\mathbf{k}$. Since we know that \mathbf{F} is conservative it must be the case that $\int_A^P \mathbf{F} \cdot d\mathbf{r}$ is independent of the path from A to P and hence

$$\int_{C_1} \mathbf{F} \cdot d\mathbf{r} = \int_{C_2} \mathbf{F} \cdot d\mathbf{r},$$

where C_1 and C_2 are any two curves drawn from A to P . Suppose that the point A is fixed. Then

$$\begin{aligned} \int_A^P \mathbf{F} \cdot d\mathbf{r} &= G(P), \text{ say} \\ &= G(x, y, z) \end{aligned}$$

Let Q be the point $(x + \delta x, y, z)$ and let P be the point (x, y, z) . Consider the quantity

$$\begin{aligned} G(x + \delta x, y, z) - G(x, y, z) &\equiv \int_A^Q \mathbf{F} \cdot d\mathbf{r} - \int_A^P \mathbf{F} \cdot d\mathbf{r} \\ &= \int_P^Q \mathbf{F} \cdot d\mathbf{r} \\ \left(-\int_A^P = \int_P^A \right) & \quad (\text{dy} = dz = 0) \end{aligned}$$

But we can choose the path from P to Q so that only x varies, in which case $d\mathbf{r} = \mathbf{i} dx$. Thus

$$G(x + \delta x, y, z) - G(x, y, z) = \int_x^{x+\delta x} F_1 \, dx$$

and hence

$$\begin{aligned}
 & \text{if } g(\delta x) = \int_x^{x+\delta x} F_1(u) du \quad \frac{\partial G}{\partial x} = \lim_{\delta x \rightarrow 0} \left[G(x+\delta x, y, z) - G(x, y, z) \right] / \delta x \\
 & \text{then} \quad g'(\delta x) = F_1(x+\delta x) \quad = \lim_{\delta x \rightarrow 0} \left(\int_x^{x+\delta x} F_1 dx \right) / \delta x \quad (\text{L'Hopital}) \\
 & \Rightarrow g'(0) = F_1(x) \quad = F_1(x, y, z)
 \end{aligned}$$

Similarly we can show that

$$F_2 = \frac{\partial G}{\partial y}, \quad F_3 = \frac{\partial G}{\partial z}$$

Thus, if \mathbf{F} is conservative then a scalar function (G in this case) can be found such that

$$\mathbf{F} = \nabla G \Rightarrow \text{Curl } \underline{F} = 0$$

Example

For the vector field

(Can check
 $\text{Curl } \underline{F} = 0$)

$$\mathbf{F} = (3x^2 + yz)\mathbf{i} + (6y^2 + xz)\mathbf{j} + (12z^2 + xy)\mathbf{k}$$

find a scalar function $\phi(x, y, z)$ such that $\mathbf{F} = \nabla \phi$. Hence calculate $\int_A^B \mathbf{F} \cdot d\underline{r}$ where $A = (0, 0, 0)$ and $B = (1, 1, 1)$.

$$\begin{aligned}
 & \underline{F} = \nabla \varphi \\
 & \Rightarrow \frac{\partial \varphi}{\partial x} = F_1 = 3x^2 + yz \Rightarrow \varphi = x^3 + xyz + f(y, z) \\
 & \Rightarrow \frac{\partial \varphi}{\partial y} = xz + \cancel{\frac{\partial f}{\partial y}} = F_2 = 6y^2 + xz \Rightarrow f = 2y^3 + g(z) \\
 & \Rightarrow \varphi = x^3 + xyz + 2y^3 + g(z) \\
 & \Rightarrow \cancel{\frac{\partial \varphi}{\partial z}} = xy + g'(z) = F_3 = 12z^2 + xy \\
 & \Rightarrow g(z) = 4z^3 + C
 \end{aligned}$$

$$\Rightarrow \varphi = x^3 + xyz + 2y^3 + 4z^3 + C$$

$$\begin{aligned}
 & \text{Then } \int_A^B \mathbf{F} \cdot d\underline{r} = [\varphi]_{(0,0,0)}^{(1,1,1)} = (8+C) - C \\
 & \qquad \qquad \qquad = 8 //
 \end{aligned}$$

1.5.4 Practical evaluation of path integrals

Suppose we wish to evaluate

$$I = \int_{\gamma} \mathbf{F} \cdot d\mathbf{r}$$

explicitly, where \mathbf{F} is a known function of (x, y, z) and γ is some known curve joining the points $A(x_0, y_0, z_0)$ and $B(x_1, y_1, z_1)$.

Along γ we can write

$$x = x(t), y = y(t), z = z(t) \quad (t_0 \leq t \leq t_1)$$

Here, t is a parameter that takes us along γ with $x(t_0) = x_0, x(t_1) = x_1$ and similarly for y and z . Then we can write

$$d\mathbf{r} = \left(\frac{dx}{dt} \hat{i} + \frac{dy}{dt} \hat{j} + \frac{dz}{dt} \hat{k} \right) dt$$

and hence, with $\mathbf{F} = F_1(t)\mathbf{i} + F_2(t)\mathbf{j} + F_3(t)\mathbf{k}$:

$$I = \int_{\gamma} \mathbf{F} \cdot d\mathbf{r} = \int_{t_0}^{t_1} \left(F_1 \frac{dx}{dt} + F_2 \frac{dy}{dt} + F_3 \frac{dz}{dt} \right) dt$$

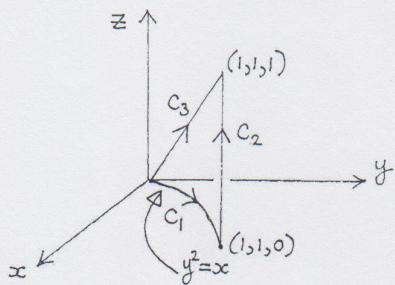


Figure 7: The integration path for this example.

Example (see figure 7)

Evaluate

$$\int_{\gamma} \mathbf{F} \cdot d\mathbf{r} \text{ with } \mathbf{F} = yz\mathbf{i} + xy\mathbf{j} + xz\mathbf{k}$$

when γ joins $(0, 0, 0)$ to $(1, 1, 1)$ along

- (i) $C_1 + C_2$ with C_1 the curve $x = y^2, z = 0$ from $(0, 0, 0)$ to $(1, 1, 0)$ and C_2 is the straight line joining $(1, 1, 0)$ to $(1, 1, 1)$;
- (ii) C_3 is the straight line joining $(0, 0, 0)$ to $(1, 1, 1)$.

(i) On C_1 : $z = dz = 0$ Along C_1 : $\mathbf{F} = 0\hat{i} + t^3\hat{j} + 0\hat{k}$
 $x = y^2$
 $y = t$ ($0 \leq t \leq 1$) $\int_{C_1} \mathbf{F} \cdot d\mathbf{r} = \int_0^1 t^3 \frac{dy}{dt} dt = \frac{1}{4}$
 $x = t^2$

On C_2 : $x = y = 1 \Rightarrow dx = dy = 0$

z goes from 0 to 1 let $z = t$ ($0 \leq t \leq 1$)

Along C_2 : $\mathbf{F} = z\hat{i} + \hat{j} + z\hat{k} = t\hat{i} + \hat{j} + t\hat{k}$
 $\int_{C_2} \mathbf{F} \cdot d\mathbf{r} = \int_0^1 t \underbrace{\frac{dx}{dt}}_{0} + \underbrace{\frac{dy}{dt}}_{0} + \underbrace{t \frac{dz}{dt}}_1 dt = \int_0^1 t dt = \frac{1}{2}$

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_{C_1} \mathbf{F} \cdot d\mathbf{r} + \int_{C_2} \mathbf{F} \cdot d\mathbf{r} = \frac{1}{4} + \frac{1}{2} = \frac{3}{4} //$$

(ii) On C_3 : $x = y = z = t$ ($0 \leq t \leq 1$)

$$\mathbf{F} = t^2\hat{i} + t^2\hat{j} + t^2\hat{k} \quad \int_{C_3} \mathbf{F} \cdot d\mathbf{r} = \int_0^1 3t^2 dt = 1 //$$

Answers not the same $\Rightarrow \mathbf{F}$ is not conservative Can check
 $\text{curl } \mathbf{F} \neq 0$

1.6 Surface integrals

1.6.1 Definition

To define a surface integral of $f = f(P)$ over a surface S , we divide S into elements of area $\delta S_1, \delta S_2, \dots, \delta S_N$. Let f_1, f_2, \dots, f_N be the values of f at typical points P_1, P_2, \dots, P_N of $\delta S_1, \delta S_2, \dots, \delta S_N$ respectively. We calculate the quantity

$$\sum_{n=1}^N f_n \delta S_n.$$

We now let $N \rightarrow \infty$, $\max \delta S_n \rightarrow 0$. The resulting limit, if it exists, is called the **surface integral** of f over S , and we write it as

$$\int_S f dS = \lim_{\substack{N \rightarrow \infty \\ \max(\delta S_n) \rightarrow 0}} \sum_{n=1}^N f_n \delta S_n$$

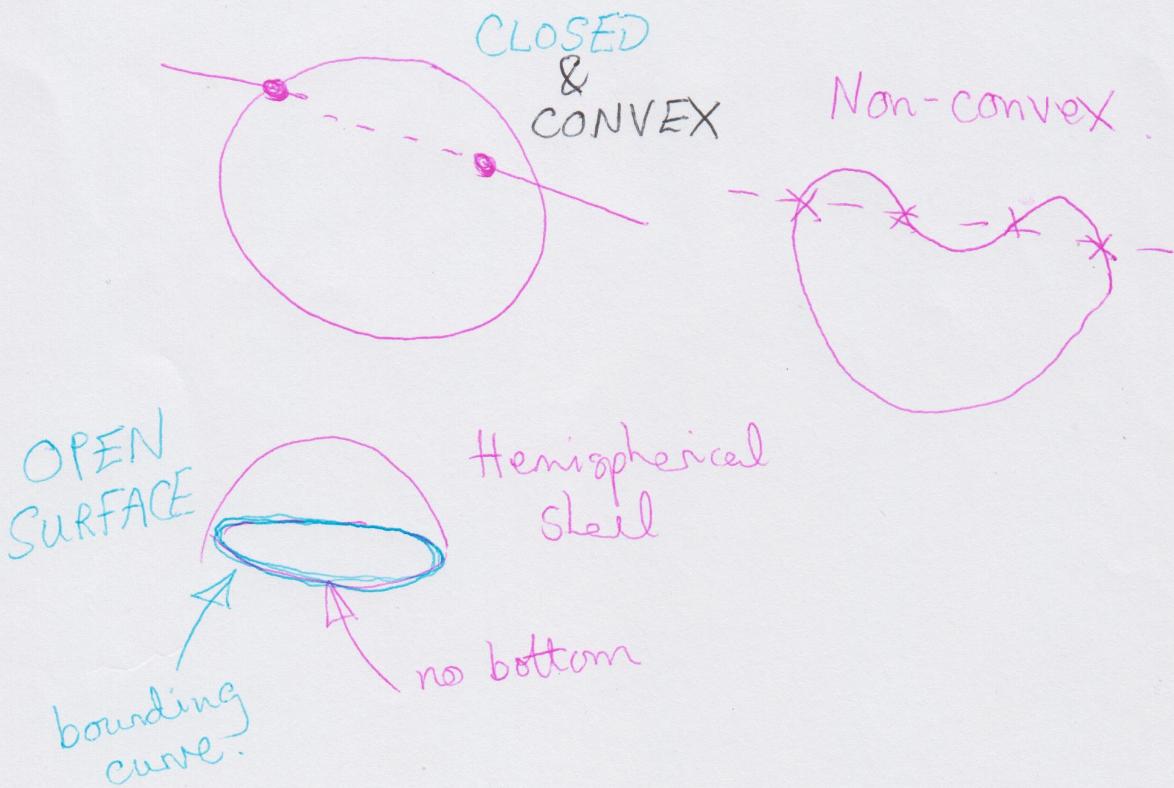
As with the line integral, the function f may be a vector or a scalar.

1.6.2 Types of surfaces

Closed surface: this divides three-dimensional space into two non-connected regions - an interior region and an exterior region;

Convex surface: this is a surface which is crossed by a straight line at most twice;

Open surface: this does not divide space into two non-connected regions - it has a rim which can be represented by a closed curve. (A closed surface can be thought of as the sum of two open surfaces).



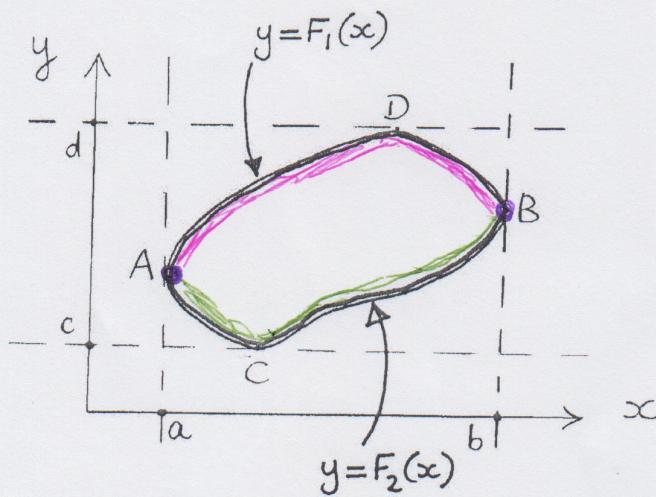


Figure 8: Diagram to illustrate the evaluation of surface integrals.

1.6.3 Evaluation of surface integrals for plane surfaces in the $x - y$ plane

An areal element dS is an ‘infinitesimally small’ element of area of a surface. Even for curved surfaces it can be thought of as approximately plane. The **vector areal element** $d\mathbf{S}$ is the vector $\hat{\mathbf{n}} dS$ where $\hat{\mathbf{n}}$ is the unit vector normal to dS . For plane surfaces dS can be expressed in Cartesian coordinates (x, y) since we may choose the surface to lie in the plane $z = 0$. Thus we can write $dS = dx dy$. (See figure 8). $\hat{\mathbf{n}} = \hat{\mathbf{k}}$ (in 2D)

Let the rectangle $x = a, b$ and $y = c, d$ circumscribe S . We will assume for simplicity that S is convex. (If it isn’t then we split S up into convex sub-regions). Let the equation of the boundary of S be denoted by

$$y = \begin{cases} F_1(x) & \text{upper half } ADB \\ F_2(x) & \text{lower half } ACB \end{cases}$$

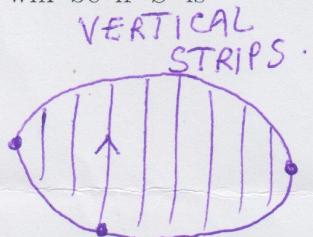
(n.b. we need to ensure these are single-valued functions, which they will be if S is convex). Then

$$\begin{aligned} \text{Area of } S &= \int_{x=a}^{x=b} \int_{y=F_2(x)}^{y=F_1(x)} 1 \cdot dy \quad dx \\ &= \int_a^b (F_1(x) - F_2(x)) dx. \end{aligned}$$

If $f(x, y)$ is any function of position:

$$\int_S f dS = \int_{x=a}^{x=b} \left\{ \int_{y=F_2(x)}^{y=F_1(x)} f(x, y) dy \right\} dx$$

e.g. f could represent density (mass/unit area)



In some situations it may be more convenient to do the x -integration first. If we want to do this we need to write the boundaries in terms of functions of y instead of x . In this case let the boundary be described by

$$x = \begin{cases} G_1(y) & \text{right half } CBD \\ G_2(y) & \text{left half } CAD \end{cases}$$

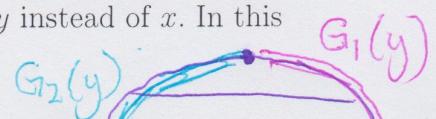
Then

$$\begin{aligned} \text{Area of } S &= \int_{y=c}^{y=d} \left\{ \int_{x=G_2(y)}^{x=G_1(y)} f(x, y) dx \right\} dy \\ &= \int_c^d (G_1(y) - G_2(y)) dy \end{aligned}$$

Horizontal strips

and

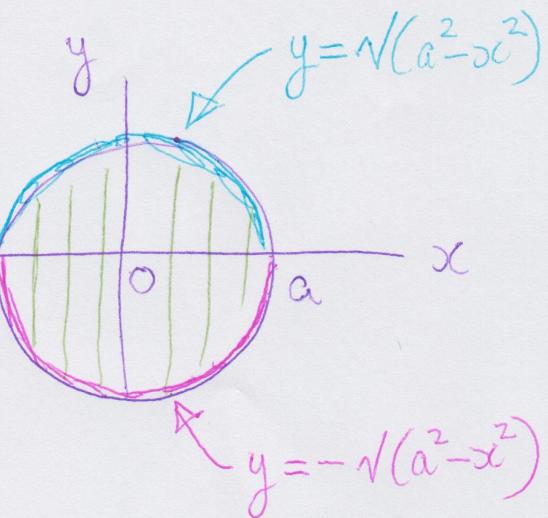
$$\int_S f dS = \int_{y=c}^{y=d} \left\{ \int_{x=G_2(y)}^{x=G_1(y)} f(x, y) dx \right\} dy.$$



1.6.4 Example

Find the area of the circle $x^2 + y^2 = a^2$.

$$\begin{aligned}
 A &= \int_{-a}^a \left\{ \int_{y=-\sqrt{a^2-x^2}}^{y=+\sqrt{a^2-x^2}} dy \right\} dx \\
 &= \int_{-a}^a [y]_{-\sqrt{a^2-x^2}}^{\sqrt{a^2-x^2}} dx \\
 &= 2 \int_{-a}^a \sqrt{a^2-x^2} dx \\
 &= \dots = \pi a^2 //
 \end{aligned}$$



Subst
 $x = a \sin u$

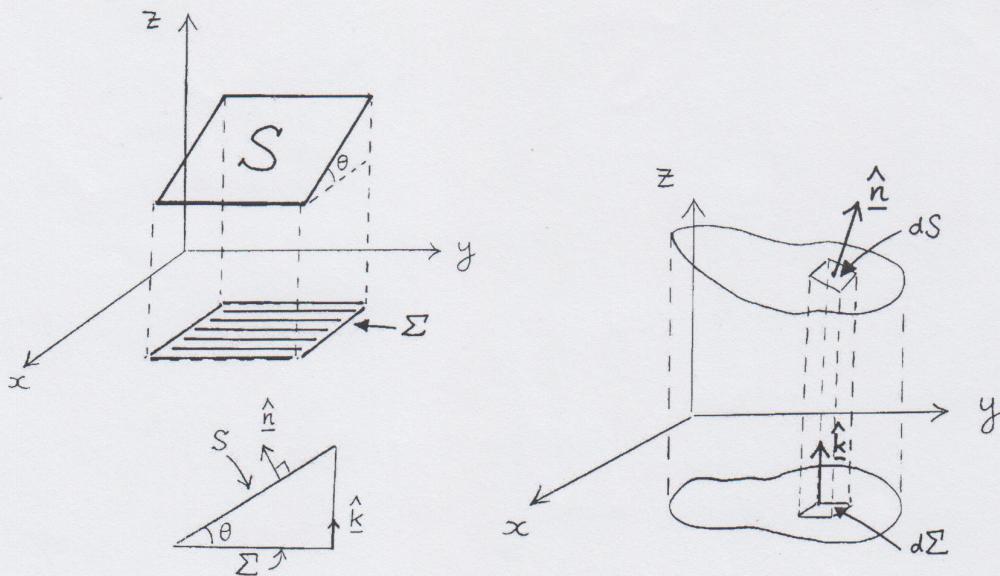


Figure 9: Left: The projection of a plane area S onto the $x - y$ plane. Right: The projection of a curved surface S onto the $x - y$ plane.

1.6.5 Projection of an area onto a plane

Consider first a plane area S (left hand diagram in figure 9). Suppose Σ is the projected area in the $x - y$ plane. Then $\Sigma = S \cos \theta$, where $\cos \theta = |\hat{n} \cdot \hat{k}|$.

Now consider a curved surface. (Right hand diagram in figure 9). If we consider an areal element dS then this will be effectively plane, and so

$$dS = d\Sigma / |\hat{n} \cdot \hat{k}|$$

could
depend
on position

1.6.6 The projection theorem

Let P denote a general point of a surface S which at no point is orthogonal to the direction \mathbf{k} . Then:

$$\int_S f(P) dS = \int_{\Sigma} f(P) \frac{dx dy}{|\hat{\mathbf{n}} \cdot \mathbf{k}|}, \text{ depends on } P$$

where Σ is the projection of S onto the plane $z = 0$, and $\hat{\mathbf{n}}$ is normal to S .

Contains
the
curvature

Proof

$$\begin{aligned} \int_S f(P) dS &= \lim_{N \rightarrow \infty} \sum_{r=1}^N f(P_r) \delta S_r \\ (\text{or } \delta \Sigma_r \rightarrow 0) &\Rightarrow \delta S_r \rightarrow 0 \quad \sum_{r=1}^N f(P_r) \left\{ \frac{\delta \Sigma_r}{|\hat{\mathbf{n}}_r \cdot \mathbf{k}|} + \varepsilon_r \right\} \end{aligned}$$

where $\varepsilon_r \rightarrow 0$ as $\delta S_r \rightarrow 0$. (Here $\hat{\mathbf{n}}_r$ is the unit vector normal to S at P_r and $\delta \Sigma_r$ is the projection of δS_r onto the plane $z = 0$. It therefore follows that

$$\int_S f(P) dS = \int_{\Sigma} f(P) \frac{d\Sigma}{|\hat{\mathbf{n}} \cdot \mathbf{k}|} \quad \text{Related to Jacobian (see later)}$$

as required. Note that $f(P)$ is evaluated at $P(x, y, z)$ on S in both integrals.

If, for example, the equation of S is $z = \phi(x, y)$ then the theorem gives

$$\int_S f(x, y, z) dS = \int_{\Sigma} f(x, y, \phi(x, y)) \frac{dx dy}{|\hat{\mathbf{n}} \cdot \mathbf{k}|}$$

Alternatively, we may choose to project the surface onto $x = 0$ or $y = 0$ to give:

$$\int_S f(P) dS = \int_{\Sigma_x} f(P) \frac{dy dz}{|\hat{\mathbf{n}} \cdot \mathbf{i}|} = \int_{\Sigma_y} f(P) \frac{dx dz}{|\hat{\mathbf{n}} \cdot \mathbf{j}|}$$

where Σ_x is the projection of S onto $x = 0$ and Σ_y is the projection of S onto $y = 0$.

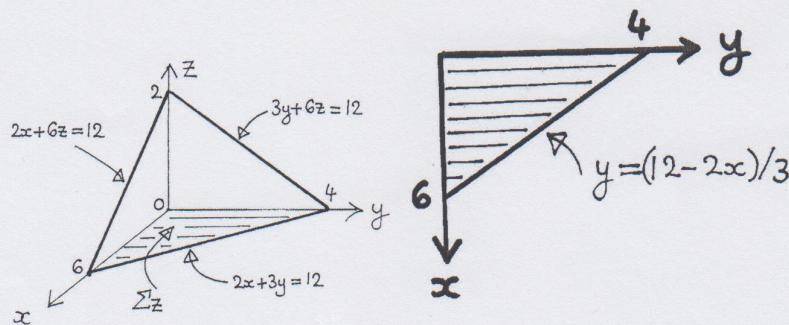


Figure 10: Left: The plane $2x + 3y + 6z = 12$ and its projection onto the $x - y$ plane. Right: The projected region Σ_z viewed from above.

Example of using the projection theorem

[As an exercise try projecting onto $x=0$ or $y=0$]

Evaluate

$$\int_S (y + 2z - 2) dS$$

where S is the part of the plane $2x + 3y + 6z = 12$ in the first octant ($x, y, z \geq 0$), by projecting onto the plane $z = 0$.

Normal to plane is $\nabla(2x+3y+6z) = 2\hat{i} + 3\hat{j} + 6\hat{k}$
 $\Rightarrow \hat{n} = \pm(2\hat{i} + 3\hat{j} + 6\hat{k})/\sqrt{2^2 + 3^2 + 6^2} = \pm(2\hat{i} + 3\hat{j} + 6\hat{k})/7$
 $\Rightarrow |\hat{n}| = 6/7$.

Projecting onto $z=0$

$$\int_S (y + 2z - 2) dS = \int_{\Sigma_z} (y + 2z - 2) \frac{dxdy}{(6/7)}$$

↑ Don't set $z=0$
Here $z = (12 - 2x - 3y)/6$
 $= 2 - \frac{1}{3}x - \frac{1}{2}y$

$$= \frac{7}{6} \int_{\Sigma_z} \left(2 - \frac{2}{3}x\right) dx dy.$$

$$= \frac{7}{6} \int_{x=0}^6 \int_{y=0}^{(12-2x)/3} \left(2 - \frac{2}{3}x\right) dy dx$$

$$= \frac{7}{6} \int_0^6 \left(2 - \frac{2}{3}x\right) \frac{(12-2x)}{3} dx$$

$$= \dots = \frac{14}{27} \int_0^6 (x^2 - 9x + 18) dx = \dots = \frac{28}{3}$$

=

1.7 Volume Integrals

1.7.1 Definition

Consider a volume τ and split it up into N subregions $\delta\tau_1, \delta\tau_2, \dots, \delta\tau_N$. Let P_1, P_2, \dots, P_N be typical points of $\delta\tau_1, \delta\tau_2, \dots, \delta\tau_N$.

Consider the sum

$$\sum_{i=1}^N f(P_i) \delta\tau_i$$

Now let $N \rightarrow \infty, \max \delta\tau_i \rightarrow 0$. If this sum tends to a limit we call it the volume integral of f over τ and write this as

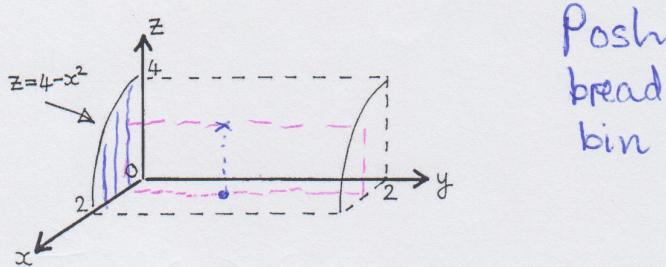
$$\int_{\tau} f d\tau.$$

The function f may be a vector or a scalar.

1.7.2 Volume element

In Cartesian coordinates the volume element

$$d\tau = dx dy dz.$$

Figure 11: The volume τ for the example.

Example

Evaluate

$$\int_{\tau} (2x + y) d\tau$$

e.g.
density

when τ is the volume enclosed by the parabolic cylinder $z = 4 - x^2$ and the planes $x = y = z = 0$ and $y = 2$.

$$I = \int_{x=0}^{x=2} \int_{y=0}^{y=2} \int_{z=0}^{z=4-x^2} (2x+y) dz dy dx$$

$$= \int_0^2 \int_0^2 (2x+y)(4-x^2) dy dx$$

$$= \int_0^2 \int_0^2 (8x - 2x^3 + 4y - x^2y) dy dx$$

$$= \int_0^2 \left[8xy - 2x^3y + 2y^2 - \frac{x^2}{2}y^2 \right]_{y=0}^{y=2} dx$$

$$= \int_0^2 (16x - 4x^3 + 8 - 2x^2) dx$$

$$= \dots = 80/3 //$$

$$\Rightarrow = \int_0^2 (4-x^2) \left[\frac{2xy + y^2}{2} \right]_{y=0}^2 dx$$

$$= \int_0^2 (4-x^2)(4x+2) dx = \dots = 80/3$$

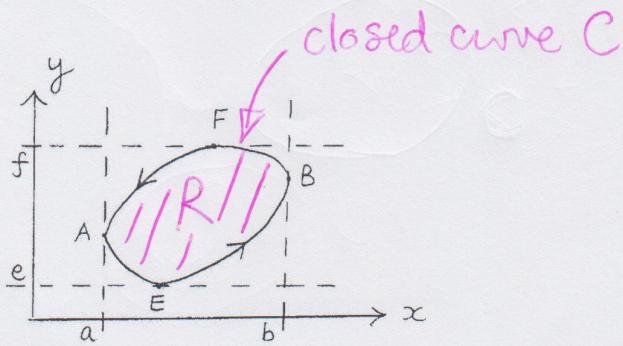


Figure 12: Diagram for proof of Green's theorem.

1.8 Results relating line, surface and volume integrals

1.8.1 Green's theorem in the plane

Suppose R is a closed plane region bounded by a simple plane closed convex curve in the $x - y$ plane. Let L, M be continuous functions of x, y having continuous derivatives throughout R . Then:

$$\oint_C (L \, dx + M \, dy) = \int_R \left(\frac{\partial M}{\partial x} - \frac{\partial L}{\partial y} \right) \, dx \, dy,$$

where C is the boundary of R described in the counter-clockwise (positive) sense.

Proof. We draw a rectangle formed by the tangent lines $x = a, b$ and $y = e, f$ (figure 12). This rectangle circumscribes C . Let $x = X_1(y)$, $x = X_2(y)$ be the equations of EAF and EBF respectively. We then can write

$$\begin{aligned}
 \int_R \frac{\partial M}{\partial x} \, dx \, dy &= \int_e^f \left\{ \int_{X_1(y)}^{X_2(y)} \left(\frac{\partial M}{\partial x} \right) dx \right\} dy \\
 &= \int_e^f M(X_2(y), y) - M(X_1(y), y) \, dy \\
 &= \int_e^f M(X_2(y), y) \, dy + \int_f^e M(X_1(y), y) \, dy \\
 &\equiv \oint_C M \, dy
 \end{aligned}$$

Hori_z strips

Now, let the equations of AEB and AFB be $y = Y_1(x)$, $y = Y_2(x)$. Then

$$\begin{aligned}
 \int_R \frac{\partial L}{\partial y} \, dx \, dy &= \int_a^b \left\{ \int_{Y_1(x)}^{Y_2(x)} \left(\frac{\partial L}{\partial y} \right) dy \right\} dx \\
 &= \int_a^b L(x, Y_2(x)) - L(x, Y_1(x)) \, dx \\
 &= - \left\{ \int_a^b L(x, Y_1(x)) \, dx + \int_b^a L(x, Y_2(x)) \, dx \right\} \\
 &= - \oint_C L \, dx
 \end{aligned}$$

vert. strips

Hence result.

1.8.2 Vector forms of Green's Theorem

(i) (2D Stokes Theorem). Let $\mathbf{F} = L\mathbf{i} + M\mathbf{j}$, and $d\mathbf{r} = dx\mathbf{i} + dy\mathbf{j}$. Then

$$\operatorname{curl} \mathbf{F} = \left(\frac{\partial M}{\partial x} - \frac{\partial L}{\partial y} \right) \mathbf{k}.$$

Over the region R we can write $dx dy = dS$. Thus using Green's theorem:

$$\begin{aligned} \oint_C \mathbf{F} \cdot d\mathbf{r} &= \int_R \mathbf{k} \cdot \operatorname{curl} \mathbf{F} dS \\ &= \int_R \operatorname{curl} \mathbf{F} \cdot d\mathbf{S}. \end{aligned}$$

$d\mathbf{S} = \hat{\mathbf{k}} dS$

in
2D

↑
dirn
of normal

This result can be generalized to three dimensions (see **Stokes theorem** later).

(ii) (Divergence theorem in 2D). This time let $\mathbf{F} = M\mathbf{i} - L\mathbf{j}$. Then

$$\operatorname{div} \mathbf{F} = \frac{\partial M}{\partial x} - \frac{\partial L}{\partial y}$$

and so Green's theorem can be rewritten as

$$\int_R \operatorname{div} \mathbf{F} dx dy = \oint_C F_1 dy - F_2 dx.$$

$F_1 \uparrow + F_2 \uparrow$
 $F_1 = M$
 $F_2 = -L$



Now it can be shown (exercise) that

$$\hat{\mathbf{n}} ds = (dy \mathbf{i} - dx \mathbf{j})$$

where s is arclength along C , and $\hat{\mathbf{n}}$ is the unit normal to C . Therefore we can rewrite Green's theorem as

$$\int_R \operatorname{div} \mathbf{F} dx dy = \oint_C \mathbf{F} \cdot \hat{\mathbf{n}} ds.$$

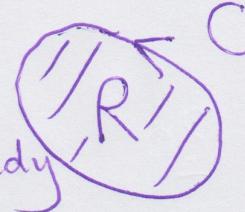
This result also turns out to be true in three dimensions, where it is known as the **Divergence Theorem**.

Example

Show that the area enclosed by a simple closed curve with boundary C can be expressed as

$$\frac{1}{2} \oint_C x dy - y dx.$$

Use this result to calculate the area of an ellipse.

$$G-T \quad \oint_C L dx + M dy = \int_R \left(\frac{\partial M}{\partial x} - \frac{\partial L}{\partial y} \right) dx dy$$


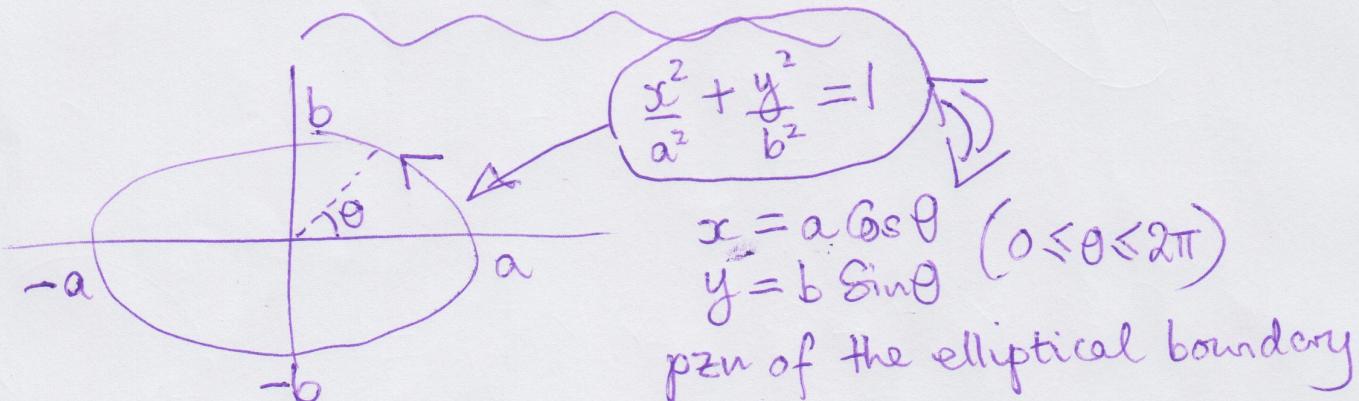
$$\text{Choose } L = -y$$

$$M = x$$

$$G-T \Rightarrow \oint_C -y dx + x dy = \int_R (1+1) dx dy$$

= twice the area

$$\therefore \text{area of } R = \frac{1}{2} \oint_C x dy - y dx \quad \text{of } R$$



$$xdy - ydx = [(a \cos \theta)(b \cos \theta) - (b \sin \theta)(-a \sin \theta)] d\theta$$

$$= ab d\theta$$

$$\therefore \oint_C x dy - y dx = \int_0^{2\pi} ab d\theta = 2\pi ab$$

= twice the area of R

$$\therefore \text{area of ellipse} = \pi ab \quad (\text{as expected})$$

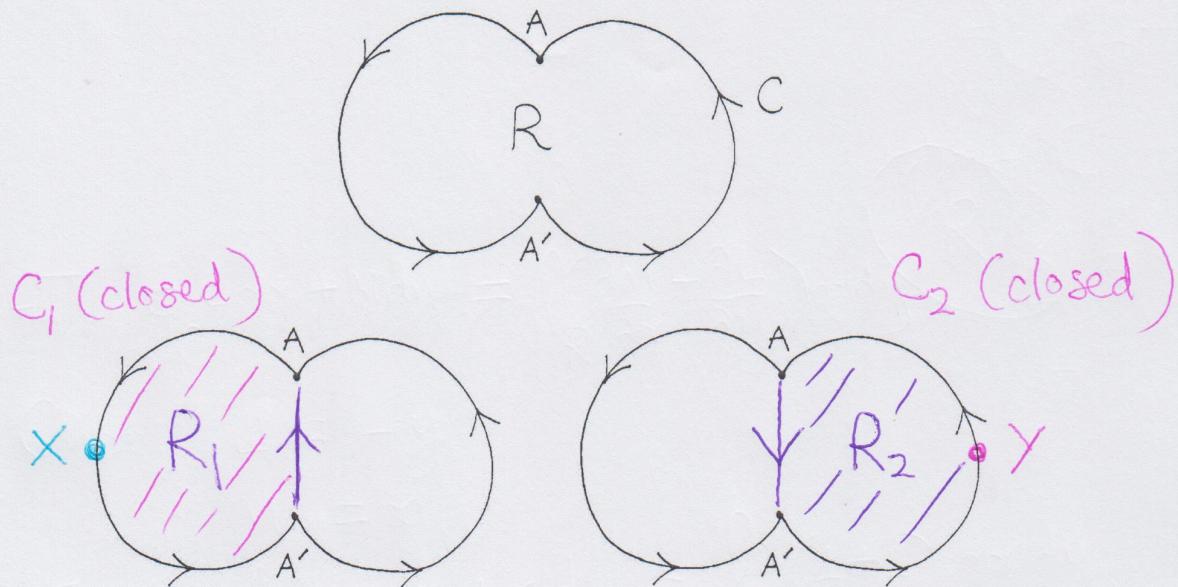


Figure 13: A non-convex boundary.

1.8.3 Extensions of Green's theorem in the plane

Green's theorem is true for more complicated geometries than that assumed in the proof given above. e.g. if C is not convex, but has the shape given in figure 13. We can join the points A, A' so as to form 2 (or more) simple convex closed curves C_1, C_2 enclosing R_1, R_2 where $R_1 + R_2 = R$. Then:

$$\oint_{C_1} \mathbf{F} \cdot d\mathbf{r} + \oint_{C_2} \mathbf{F} \cdot d\mathbf{r} = \int_{R_1} (\operatorname{curl} \mathbf{F}) \cdot d\mathbf{s} + \int_{R_2} (\operatorname{curl} \mathbf{F}) \cdot d\mathbf{s}$$

$$= \int_R (\operatorname{curl} \mathbf{F}) \cdot d\mathbf{s}$$

Now

$$\oint_{C_1} = \int_{AXA'} + \int_{A'A}^A$$

$$\oint_{C_2} = \int_{A'YA} + \int_A^{A'}$$

$$\left[\int_{A'}^A = - \int_A^{A'} \right]$$

and so

$$\oint_C \mathbf{F} \cdot d\mathbf{r} = \oint_{C_1} \mathbf{F} \cdot d\mathbf{r} + \oint_{C_2} \mathbf{F} \cdot d\mathbf{r} = \int_R \operatorname{curl} \mathbf{F} \cdot d\mathbf{s}$$

We see therefore that the theorem still holds.

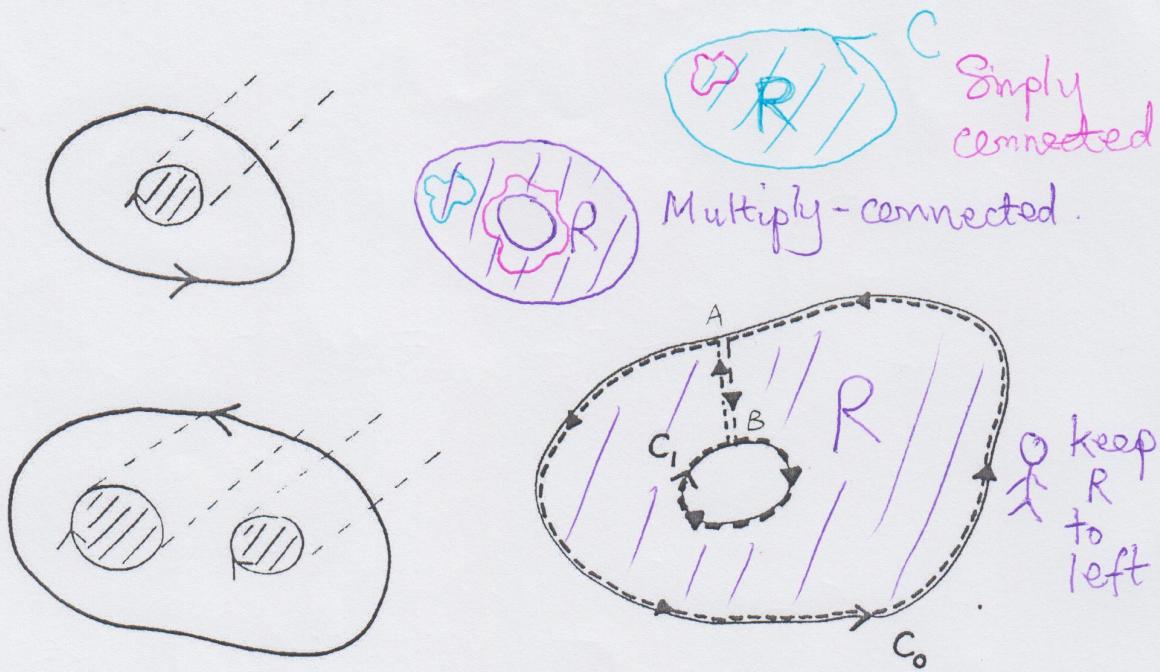


Figure 14: Left: Examples of doubly- and triply-connected regions. Right: Green's theorem in a multiply-connected region.

1.8.4 Green's theorem in multiply-connected regions

A region R is said to be **simply-connected** if any closed curve drawn in R can be shrunk to a point without leaving R . If we restrict ourselves to two dimensions then any region with a hole in it is not simply-connected (left-hand picture in figure 14). A region which is not simply-connected is said to be **multiply-connected**.

If R is multiply-connected, Green's theorem is still true provided C is now interpreted as the entire (outer and inner) boundary, with C described so that the region R is always on the left (right hand picture in figure 14).

For example if we have a doubly-connected region, we describe the outer boundary C_0 in an anti-clockwise fashion and the inner boundary C_1 clockwise. We can then join the point A on C_0 to the point B on C_1 by the line AB . This line then divides R in such a way that it is a simply connected region bounded by the closed curve $C_0 + AB + C_1 + BA$. Then, by Green's theorem:

$$\int_R \operatorname{curl} \mathbf{F} \cdot d\mathbf{S} = (\oint_{C_0} \mathbf{F} \cdot d\mathbf{r} + \oint_A^B \mathbf{F} \cdot d\mathbf{r} + \oint_{C_1} \mathbf{F} \cdot d\mathbf{r} + \oint_B^A \mathbf{F} \cdot d\mathbf{r}) (\mathbf{F} \cdot d\Gamma)$$

and therefore it follows that

$$\int_R \operatorname{curl} \mathbf{F} \cdot d\mathbf{S} = \oint_{C_0} \mathbf{F} \cdot d\mathbf{r} + \oint_{C_1} \mathbf{F} \cdot d\mathbf{r}$$

outer boundary
(anti-clockwise) inner boundary
(clockwise)

where $C = C_0 + C_1$.

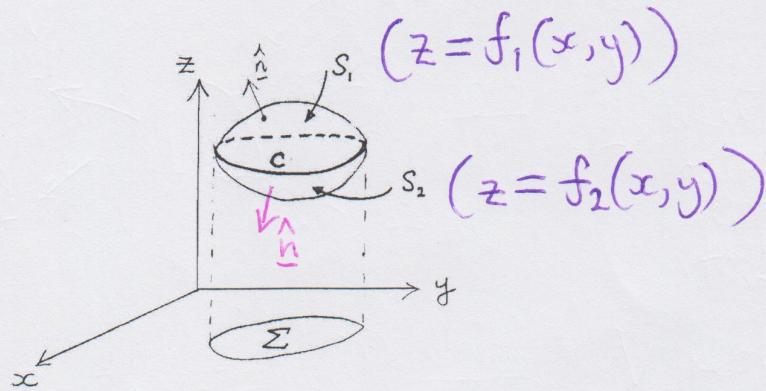


Figure 15: Diagram for the proof of the divergence theorem.

1.8.5 Flux

If S is a surface then the flux of \mathbf{A} across S is defined as

$$\int_S \mathbf{A} \cdot \hat{\mathbf{n}} dS.$$

If S is a closed surface then, by convention, we always draw the unit normal $\hat{\mathbf{n}}$ out of S .

1.8.6 The divergence theorem

If τ is the volume enclosed by a closed surface S with unit outward normal $\hat{\mathbf{n}}$ and \mathbf{A} is a vector field with continuous derivatives throughout τ , then:

denotes
closed
surface


$$\oint_S \mathbf{A} \cdot \hat{\mathbf{n}} dS = \int_{\tau} \operatorname{div} \mathbf{A} d\tau.$$

Proof

We will assume that S is convex and that τ is simply connected, with no interior boundaries. Let $\mathbf{A} = (A_1, A_2, A_3)$ and $\hat{\mathbf{n}} = (l, m, n)$. We have to prove that

$$\int_S (lA_1 + mA_2 + nA_3) dS = \int_{\tau} \left(\frac{\partial A_1}{\partial x} + \frac{\partial A_2}{\partial y} + \frac{\partial A_3}{\partial z} \right) dx dy dz$$

Project S onto the plane $z = 0$ (figure 15). The cylinder with normal cross-section Σ and generators parallel to the z -axis circumscribes S and it touches S along the curve C which divides S into two open surfaces, S_1 (upper) and S_2 (lower). Both S_1 and S_2 have projection Σ in the plane $z = 0$. Suppose the equations of S_1 and S_2 are $z = f_1(x, y)$ and $z = f_2(x, y)$ respectively. Then:

$$\begin{aligned}
 \int_{\tau} \frac{\partial A_3}{\partial z} dx dy dz &= \int_C \frac{\partial A_3}{\partial z} dz dx dy \\
 &= \int_{\Sigma} [A_3(x, y, f_1(x, y)) - A_3(x, y, f_2(x, y))] dx dy
 \end{aligned}$$

$$\hat{n} = l\hat{i} + m\hat{j} + n\hat{k}$$

Now, using the projection theorem:

$$\begin{aligned}\int_{S_1} n A_3 dS &= \sum n A_3(x, y, f_1(x, y)) \frac{dxdy}{|\hat{n} \cdot \hat{k}|} \\ &= \sum A_3(x, y, f_1(x, y)) dxdy.\end{aligned}$$

Similarly:

$$\begin{aligned}\int_{S_2} n A_3 dS &= \sum n A_3(x, y, f_2(x, y)) \frac{dxdy}{|\hat{n} \cdot \hat{k}|} \\ &= - \sum A_3(x, y, f_2(x, y)) dxdy\end{aligned}$$

Thus:

$$\oint_S n A_3 dS = \sum [A_3(x, y, f_1(x, y)) - A_3(x, y, f_2(x, y))] dxdy$$

and therefore

$$\int_{\tau} \frac{\partial A_3}{\partial z} d\tau = \oint_S n A_3 dS$$

Similarly, by projecting onto the planes $x = 0$ and $y = 0$:

$$\int_{\tau} \frac{\partial A_1}{\partial x} d\tau = \oint_S l A_1 dS$$

and

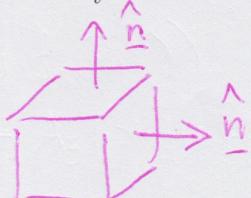
$$\int_{\tau} \frac{\partial A_2}{\partial y} d\tau = \oint_S m A_2 dS$$

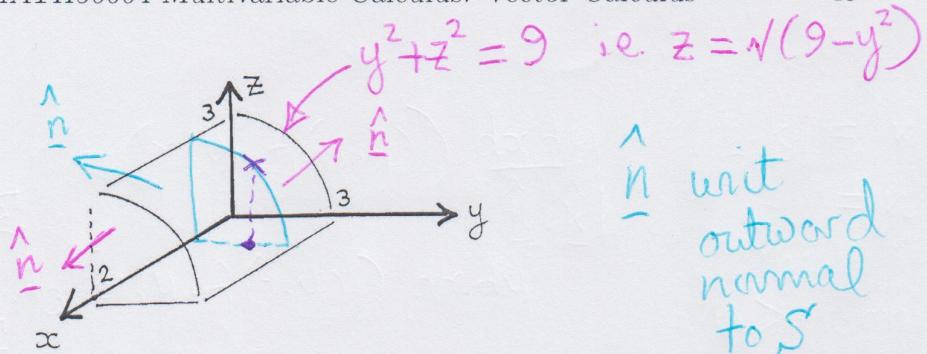
and hence

$$\oint_S \mathbf{A} \cdot \hat{n} dS = \int_C \operatorname{div} \underline{A} dC$$

as required.

Note that the surface S need not necessarily be smooth - it could be, for example, a cube or a tetrahedron.



Figure 16: The surface S in the example. S has five sides.

∇ region
enclosed
by S .

Example

Evaluate

$$\int_S \mathbf{A} \cdot \hat{\mathbf{n}} dS \text{ if } \mathbf{A} = 2x^2y \mathbf{i} - y^2 \mathbf{j} + 4xz^2 \mathbf{k},$$

and S is the surface of the region in the first octant bounded by $y^2 + z^2 = 9$, $x = 2$ and $x = y = z = 0$.

$$\begin{aligned}
 \text{Div thm } \oint_S \mathbf{A} \cdot \hat{\mathbf{n}} dS &= \iiint_V \operatorname{div} \mathbf{A} dV \\
 &= \iiint_V (4xy - 2y + 8xz) dV \\
 &= \int_{x=0}^2 \int_{y=0}^{y=3} \int_{z=0}^{z=\sqrt{9-y^2}} (4xy - 2y + 8xz) dz dy dx \\
 &= \int_{x=0}^2 \int_{y=0}^3 \left[4xyz - 2yz + 4xz^2 \right]_{z=0}^{z=\sqrt{9-y^2}} dy dx \\
 &= \int_0^3 \int_0^2 4xy\sqrt{9-y^2} - 2y\sqrt{9-y^2} + 4x(9-y^2) dx dy \\
 &= \int_0^3 \left[2x^2y\sqrt{9-y^2} - 2xy\sqrt{9-y^2} + 2x^2(9-y^2) \right]_{x=0}^{x=2} dy \\
 &= \int_0^3 4y\sqrt{9-y^2} + 8(9-y^2) dy = \dots = \underline{\underline{180}}
 \end{aligned}$$

\nwarrow subst.
 $u = y^2$