专题 常用曲线坐标系下的流体力学基本方程

一.正交曲线坐标系

盾角坐标系 xi xz x3

正交曲线坐村条 9, 9, 9,

对于一个矢量尺=从2+从2+从2+从2+从2

定义其在正交曲线坐标系下的基实为已; = 3次 2+3% 2+3% 3+3% 6 元意, 正交曲线坐标系的基矢量在空间各处大小, 方向均不同, 且一般不为单位矢量.

那么,一个矢量在在正交曲线系下可以表示为

$$\vec{a} = a_1'\vec{e}_1 + a_2'\vec{e}_2 + a_3'\vec{e}_3$$

为了方便,我们一般用沿基文方向的单位矢量来分解一个矢量.

为此,我们定义 ni 为单位基文量, hi 为尺度因子

$$\hat{N}_{i} = \frac{\vec{e}_{i}}{h_{i}} \qquad h_{i} = |\vec{e}_{i}| = \sqrt{\frac{2\chi_{i}}{2q_{i}}^{2} + \frac{2\chi_{i}}{2q_{i}}^{2}} + \frac{2\chi_{i}}{2q_{i}}^{2} + \frac{2\chi_{i}}{2q_{i}}^{2}$$

现在,可以用单位基矢量表示矢量石:

$$\vec{\alpha} = \alpha_1 \hat{n}_1 + \alpha_2 \hat{n}_2 + \alpha_3 \hat{n}_3$$

1.单位基矢的偏导数

我们利用二阶连续偏导数可以交换求导次序的性质,有

$$\frac{\partial \vec{e_1}}{\partial \hat{e_2}} = \frac{\partial \vec{e_2}}{\partial \hat{e_1}} = \frac{\partial^2 \vec{R}}{\partial \hat{e_1} \partial \hat{e_2}}$$

而
$$\begin{cases} \frac{\partial \vec{E}}{\partial q_2} = \frac{\partial}{\partial q_2} (h_1 \hat{n}_1) = h_1 \frac{\partial \hat{n}_1}{\partial q_2} + \frac{\partial h_1}{\partial q_2} \hat{n}_1 \end{cases}$$
 注意到正交曲级的性质指出 $\frac{\partial \hat{n}_1}{\partial q_1} \parallel \hat{n}_1$ 因此两部分分别对应相等

$$\Rightarrow \begin{cases} h_1 \frac{\partial \hat{h}_1}{\partial \varrho_2} = \frac{\partial h_2}{\partial \varrho_1} \hat{n}_2 \\ h_2 \frac{\partial \hat{n}_2}{\partial \varrho_1} = \frac{\partial h_1}{\partial \varrho_2} \hat{n}_1 \end{cases} \Rightarrow \frac{\partial \hat{n}_1}{\partial \varrho_2} = \frac{1}{h_1} \frac{\partial h_2}{\partial \varrho_1} \hat{n}_2 \quad (i \neq j)$$

利用
$$\hat{n}_{i} = \hat{n}_{j} \times \hat{n}_{k}$$
 (i.j.k为 1.2.3的一个顺序排列)有 $\frac{\partial \hat{n}_{i}}{\partial t_{i}} = \frac{\partial}{\partial t_{i}} (\hat{n}_{j} \times \hat{n}_{k}) = \frac{\partial \hat{n}_{j}}{\partial t_{i}} \times \hat{n}_{k} + \hat{n}_{j} \times \frac{\partial \hat{n}_{k}}{\partial t_{i}}$

$$= \frac{1}{h_{j}} \frac{\partial h_{i}}{\partial t_{j}} \hat{n}_{i} \times \hat{n}_{k} + \hat{n}_{j} \times \frac{1}{h_{k}} \frac{\partial h_{i}}{\partial t_{k}} \hat{n}_{i}$$

$$= -\frac{1}{h_{j}} \frac{\partial h_{i}}{\partial t_{j}} \hat{n}_{j} - \frac{1}{h_{k}} \frac{\partial h_{i}}{\partial t_{k}} \hat{n}_{k}$$

至此,我们得到了正交曲线系下单位基文的偏导数:

$$\begin{cases}
\frac{\partial \hat{n}_{i}}{\partial q_{j}} = \frac{1}{h_{i}} \frac{\partial h_{j}}{\partial q_{i}} \hat{n}_{j} & (i \neq j) \\
\frac{\partial \hat{n}_{i}}{\partial q_{i}} = -\frac{1}{h_{j}} \frac{\partial h_{i}}{\partial q_{j}} \hat{n}_{j} - \frac{1}{h_{k}} \frac{\partial h_{i}}{\partial q_{k}} \hat{n}_{k}
\end{cases}$$

2.正交曲线系下的 Nabla算符▽

由此得到正交曲线系下的 Nabla 算符

$$\nabla = \frac{n_1^2}{h_1} \frac{\partial}{\partial q_1} + \frac{n_2^2}{h_2} \frac{\partial}{\partial q_2} + \frac{n_3^2}{h_3} \frac{\partial}{\partial q_3}$$

(1)梯度

$$\nabla a = \frac{\hat{n_1}}{h_1} \frac{\partial \alpha}{\partial q_1} + \frac{\hat{n_2}}{h_2} \frac{\partial \alpha}{\partial q_2} + \frac{\hat{n_3}}{h_3} \frac{\partial \alpha}{\partial q_3}$$

(2)散度

$$\nabla \cdot \overrightarrow{\alpha} = \left(\frac{\overrightarrow{h_1}}{h_1} \frac{\partial}{\partial q_1} + \frac{\overrightarrow{h_2}}{h_2} \frac{\partial}{\partial q_2} + \frac{\overrightarrow{h_3}}{h_3} \frac{\partial}{\partial q_3}\right) \cdot \left(\alpha_1 \overrightarrow{h_1} + \alpha_2 \overrightarrow{h_2} + \alpha_3 \overrightarrow{h_3}\right)$$

$$= \frac{\overrightarrow{h_1}}{h_1} \left(\frac{\partial \alpha_1}{\partial q_1} \overrightarrow{h_1} + \frac{\partial \alpha_2}{\partial q_2} \overrightarrow{h_2} + \frac{\partial \alpha_3}{\partial q_1} \overrightarrow{h_3} + \alpha_1 \frac{\partial \overrightarrow{h_1}}{\partial q_1} + \alpha_2 \frac{\partial \overrightarrow{h_2}}{\partial q_1} + \alpha_3 \frac{\partial \overrightarrow{h_3}}{\partial q_1}\right) + \cdots$$

$$= \frac{\overrightarrow{h_1}}{h_1} \left[\frac{\partial \alpha_1}{\partial q_1} \overrightarrow{h_1} + \alpha_1 \left(-\frac{1}{h_2} \frac{\partial h_1}{\partial q_2} \overrightarrow{h_1} - \frac{1}{h_3} \frac{\partial h_1}{\partial q_3} \overrightarrow{h_3}\right) + \alpha_2 \frac{1}{h_2} \frac{\partial h_1}{\partial q_2} \overrightarrow{h_1} + \alpha_3 \frac{1}{h_3} \frac{\partial h_1}{\partial q_3} \overrightarrow{h_1}\right] + \cdots$$

$$= \frac{\overrightarrow{h_1}}{h_1} \left(\frac{\partial \alpha_1}{\partial q_1} + \frac{\alpha_2}{h_2} \frac{\partial h_1}{\partial q_2} + \frac{\alpha_3}{h_3} \frac{\partial h_1}{\partial q_3}\right) \overrightarrow{h_1} + \cdots$$

$$= \frac{1}{h_{1}} \left(\frac{\partial \alpha_{1}}{\partial q_{1}} + \frac{\alpha_{2}}{h_{2}} \frac{\partial h_{1}}{\partial q_{2}} + \frac{\alpha_{3}}{h_{3}} \frac{\partial h_{1}}{\partial q_{3}} \right)$$

$$+ \frac{1}{h_{2}} \left(\frac{\partial \alpha_{2}}{\partial q_{2}} + \frac{\alpha_{3}}{h_{3}} \frac{\partial h_{2}}{\partial q_{3}} + \frac{\alpha_{1}}{h_{1}} \frac{\partial h_{2}}{\partial q_{1}} \right)$$

$$+ \frac{1}{h_{3}} \left(\frac{\partial q_{3}}{\partial q_{3}} + \frac{\alpha_{1}}{h_{1}} \frac{\partial h_{3}}{\partial q_{1}} + \frac{\alpha_{2}}{h_{2}} \frac{\partial h_{3}}{\partial q_{2}} \right)$$

$$= \frac{1}{h_{1}h_{2}h_{3}} \left[\left(h_{2}h_{3} \frac{\partial \alpha_{1}}{\partial q_{1}} + \alpha_{1}h_{3} \frac{\partial h_{2}}{\partial q_{1}} + \alpha_{1}h_{2} \frac{\partial h_{3}}{\partial q_{1}} \right) + \cdots \right]$$

$$= \frac{1}{h_{1}h_{2}h_{3}} \sum_{i=1}^{3} \frac{\partial}{\partial q_{i}} \left(\frac{h_{1}h_{2}h_{3}}{h_{1}} \alpha_{i} \right)$$

$$\nabla \times \overrightarrow{\alpha} = \left(\frac{\mathring{n}_{1}}{h_{1}}\frac{\partial}{\partial q_{1}} + \frac{\mathring{n}_{2}}{h_{2}}\frac{\partial}{\partial q_{2}} + \frac{\mathring{n}_{3}}{h_{3}}\frac{\partial}{\partial q_{3}}\right) \times \left(a_{1}\mathring{n}_{1} + q_{2}\mathring{n}_{2} + q_{3}\mathring{n}_{3}}\right)$$

$$= \frac{\mathring{n}_{1}}{h_{1}}\frac{\partial}{\partial q_{1}} \times \left(\alpha_{1}\mathring{n}_{1} + \alpha_{2}\mathring{n}_{2} + q_{3}\mathring{n}_{3}\right) + \cdots$$

$$= \frac{\mathring{n}_{1}}{h_{1}} \times \left(\frac{\partial \alpha_{1}}{\partial q_{1}}\mathring{n}_{1} + \frac{\partial q_{2}}{\partial q_{2}}\mathring{n}_{2} + \frac{\partial q_{3}}{\partial q_{1}}\mathring{n}_{3} + q_{1}\frac{\partial \mathring{n}_{1}}{\partial q_{1}} + a_{2}\frac{\partial \mathring{n}_{2}}{\partial q_{1}} + a_{3}\frac{\partial \mathring{n}_{3}}{\partial q_{1}}\right) + \cdots$$

$$= \frac{\mathring{n}_{1}}{h_{1}} \times \left[\frac{\partial \alpha_{2}}{\partial q_{1}}\mathring{n}_{2} + \frac{\partial q_{3}}{\partial q_{1}}\mathring{n}_{2} + \alpha_{1}\left(-\frac{1}{h_{2}}\frac{\partial h_{1}}{\partial q_{2}}\mathring{n}_{2} - \frac{1}{h_{3}}\frac{\partial h_{1}}{\partial q_{3}}\mathring{n}_{3}\right)\right] + \cdots$$

$$= \frac{1}{h_{1}} \left[\left(\frac{\alpha_{1}}{h_{3}}\frac{\partial h_{1}}{\partial q_{3}} - \frac{\partial q_{3}}{\partial q_{1}}\right)\mathring{n}_{2} + \left(\frac{\partial q_{2}}{\partial q_{2}} - \frac{\alpha_{1}}{h_{2}}\frac{\partial h_{1}}{\partial q_{2}}\right)\mathring{n}_{3}\right]$$

$$+ \frac{1}{h_{2}} \left[\left(\frac{\alpha_{3}}{h_{1}}\frac{\partial h_{2}}{\partial q_{1}} - \frac{\partial q_{1}}{\partial q_{2}}\right)\mathring{n}_{3} + \left(\frac{\partial \alpha_{3}}{\partial q_{2}} - \frac{\alpha_{2}}{h_{3}}\frac{\partial h_{2}}{\partial q_{3}}\right)\mathring{n}_{1}\right]$$

$$+ \frac{1}{h_{3}} \left[\left(\frac{\alpha_{3}}{h_{1}}\frac{\partial h_{3}}{\partial q_{1}} - \frac{\partial q_{2}}{\partial q_{3}}\right)\mathring{n}_{1} + \left(\frac{\partial \alpha_{1}}{\partial q_{3}} - \frac{\alpha_{3}}{h_{3}}\frac{\partial h_{3}}{\partial q_{1}}\right)\mathring{n}_{2}\right]$$

$$= \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \hat{n_1} & h_2 \hat{n_2} & h_3 \hat{n_3} \\ \frac{\partial}{\partial q_1} & \frac{\partial}{\partial q_2} & \frac{\partial}{\partial q_3} \\ h_1 q_1 & h_2 q_2 & h_3 q_3 \end{vmatrix}$$

至此,我们得到了正交曲线系下的基本运算关系下面讨论的流体作如下假设: ①均匀 P=cunst ②不受体积力

由于 P= const , 上述方程化为 V·T=0

⑴柱坐标

$$\begin{cases} x = r\cos\theta \\ y = r\sin\theta \\ z = z \end{cases} \Rightarrow \begin{cases} h_r = 1 \\ h_\theta = r \\ h_z = 1 \end{cases}$$

$$\nabla \cdot \overrightarrow{U} = \frac{1}{r} \frac{\partial}{\partial r} (r U_r) + \frac{1}{r} \frac{\partial}{\partial \theta} U_{\theta} + \frac{1}{r} \frac{\partial}{\partial z} (r U_z)$$
$$= \frac{\partial U_r}{\partial r} + \frac{U_r}{r} + \frac{1}{r} \frac{\partial U_{\theta}}{\partial \theta} + \frac{\partial U_z}{\partial z} = 0$$

四球坐标、

$$\begin{cases} X = r \sin \theta \cos \theta \\ Y = r \sin \theta \sin \theta \end{cases} \Rightarrow \begin{cases} hr = 1 \\ h\theta = r \sin \theta \\ h\phi = r \sin \theta \end{cases}$$

$$\nabla \cdot \overrightarrow{U} = \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial r} (r^2 \sin \theta U_r) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} (r \sin \theta U_\theta) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial y} (r U_y)$$

$$= \frac{\partial U_r}{\partial r} + \frac{2U_r}{r} + \frac{1}{r} \frac{\partial U_\theta}{\partial \theta} + \frac{U_\theta \cot \theta}{r} + \frac{1}{r \sin \theta} \frac{\partial U_\theta}{\partial y} = 0$$

三 动量方程

由之前的假设 P= const, 子与,上述方程化为

由本构关系

マ·台=-マ·(ps)+2μマ·芭 其中第二项可以化简 (マ·色); = aj eij = aj = (ai bj + aj aj vi) = = (ai aj bj + aj aj vi) 写成实体形式 , 并利用连续性方程 マ·ゼーの , 得 マ·色 = = [マ(ロ·ゼ) + マゼ] = = 1 マゼ

关于方程左边第二项:

下面我们讨论具体的曲线坐标系下方程的形式:

⑴柱坐标

$$\begin{cases} \chi = r \cos \theta \\ y = r \sin \theta \end{cases} \begin{cases} h_r = 1 \\ h_{\delta} = r \\ h_{z} = 1 \end{cases} \begin{pmatrix} \partial_r \\ \partial_{\theta} \end{pmatrix} \begin{pmatrix} \hat{r} & \hat{\theta} & \hat{z} \end{pmatrix} = \begin{pmatrix} \hat{0} & \hat{0} & \hat{0} \\ \hat{\theta} & -\hat{r} & \hat{0} \\ \hat{0} & \hat{0} & \hat{0} \end{pmatrix}$$

$$\begin{split} (\overrightarrow{U} \cdot \overrightarrow{V}) \overrightarrow{V} &= (\overrightarrow{U}_r \stackrel{\hat{r}}{r} + \overrightarrow{u} \stackrel{\hat{\theta}}{\theta} + \overrightarrow{U}_z \stackrel{\hat{Z}}{Z}) \cdot (\stackrel{\hat{r}}{r} \partial_r + \frac{\stackrel{\hat{\theta}}{\theta}}{r} \partial_\theta + \stackrel{\hat{Z}}{Z} \partial_z) (\overrightarrow{U}_r \stackrel{\hat{r}}{r} + \overrightarrow{U} \stackrel{\hat{\theta}}{\theta} + \overrightarrow{U}_z \stackrel{\hat{Z}}{Z}) \\ &= (\overrightarrow{U}_r \partial_r + \frac{\overrightarrow{U}_\theta}{r} \partial_\theta + \overrightarrow{U}_z \stackrel{\hat{Z}}{Z}) (\overrightarrow{U}_r \stackrel{\hat{r}}{r} + \overrightarrow{U}_\theta \stackrel{\hat{\theta}}{\theta} + \overrightarrow{U}_z \stackrel{\hat{Z}}{Z}) \\ &= \overrightarrow{U}_r \frac{\partial \overrightarrow{U}_r}{\partial r} \stackrel{\hat{r}}{r} + \overrightarrow{U}_r \frac{\partial \overrightarrow{U}_\theta}{\partial r} \stackrel{\hat{\theta}}{\theta} + \overrightarrow{U}_r \frac{\partial \overrightarrow{U}_z}{\partial r} \stackrel{\hat{z}}{Z} \\ &+ \frac{\cancel{U}_\theta}{r} \frac{\partial \overrightarrow{U}_r}{\partial \theta} \stackrel{\hat{r}}{r} + \frac{\cancel{U}_\theta}{r} \frac{\partial \cancel{U}_\theta}{\partial z} \stackrel{\hat{\theta}}{\theta} + \frac{\cancel{U}_\theta}{r} \frac{\partial \cancel{U}_z}{\partial z} \stackrel{\hat{z}}{Z} \\ &+ \frac{\cancel{U}_\theta}{r} \frac{\partial \overrightarrow{U}_r}{\partial z} \stackrel{\hat{r}}{r} + \frac{\cancel{U}_\theta}{r} \frac{\partial \cancel{U}_\theta}{\partial z} \stackrel{\hat{\theta}}{\theta} + \frac{\cancel{U}_z}{r} \frac{\partial \cancel{U}_z}{\partial z} \stackrel{\hat{z}}{Z} \\ &+ \cancel{U}_z \frac{\partial \overrightarrow{U}_r}{\partial z} \stackrel{\hat{r}}{r} + \cancel{U}_z \frac{\partial \cancel{U}_\theta}{\partial z} \stackrel{\hat{\theta}}{\theta} + \cancel{U}_z \frac{\partial \cancel{U}_z}{\partial z} \stackrel{\hat{z}}{Z} \end{aligned}$$

(2) 球坐标

$$\begin{cases} X = r \sin \theta \cos \theta \\ Y = r \sin \theta \sin \theta \end{cases} \begin{cases} hr = 1 \\ h\theta = r \\ h\varphi = r \sin \theta \end{cases} \begin{pmatrix} \hat{r} & \hat{\theta} & \hat{\phi} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 \\ \hat{\theta} & -\hat{r} & 0 \\ \sin \theta & \cos \theta & -\sin \theta & -\cos \theta & -\sin \theta \end{pmatrix}$$

$$Z = r \cos \theta \qquad \begin{pmatrix} hr = 1 \\ h\varphi = r \sin \theta & \partial \phi \end{pmatrix} \begin{pmatrix} \hat{r} & \hat{\theta} & \hat{\phi} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 \\ \hat{\theta} & -\hat{r} & 0 \\ \sin \theta & \cos \theta & -\sin \theta & -\cos \theta & -\sin \theta \end{pmatrix}$$

$$\begin{split} & = (U_r \hat{r} + U_b \hat{\theta} + U_y \hat{q}^2) \cdot (\hat{r} \partial_r + \frac{\hat{\theta}}{r} \partial_b + \frac{\hat{q}}{r \sin \theta} \partial_y) (U_r \hat{r} + U_b \hat{\theta} + U_y \hat{q}^2) \\ & = (U_r \hat{r} + \frac{U_b}{r} \partial_b + \frac{U_y}{r \sin \theta} \partial_y) (U_r \hat{r} + U_b \hat{\theta} + U_y \hat{q}^2) \\ & = U_r \frac{\partial U_r}{\partial r} \hat{r} + U_r \frac{\partial U_b}{\partial r} \hat{\theta} + U_r \frac{\partial U_y}{\partial r} \hat{q}^2 \\ & + \frac{U_b}{r} (\frac{\partial U_r}{\partial \theta} - U_b) \hat{r} + \frac{U_b}{r} (\frac{\partial U_b}{\partial \theta} + U_r) \hat{\theta}^2 + \frac{U_b}{r} \frac{\partial U_r}{\partial \theta} \hat{q}^2 \\ & + \frac{U_b}{r \sin \theta} (\frac{\partial U_r}{\partial y} - \sin \theta U_r) \hat{r} + \frac{U_b}{r \sin \theta} (\frac{\partial U_b}{\partial y} - \cos \theta U_p) \hat{\theta} \\ & + \frac{U_r}{r \sin \theta} (\frac{\partial U_r}{\partial y} - \sin \theta U_r - \cos \theta U_\theta) \hat{q}^2 \\ & = (U_r \frac{\partial U_r}{\partial r} + \frac{U_b}{r} \frac{\partial U_r}{\partial \theta} + \frac{U_r}{r \sin \theta} \frac{\partial U_r}{\partial y} - \frac{U_r^2 + U_r^2}{r}) \hat{r} \\ & + (U_r \frac{\partial U_r}{\partial r} + \frac{U_r}{r} \frac{\partial U_r}{\partial \theta} + \frac{U_r}{r \sin \theta} \frac{\partial U_r}{\partial y} - \frac{U_r^2 + U_r^2}{r}) \hat{\theta} \\ & + (U_r \frac{\partial U_r}{\partial r} + \frac{U_r}{r} \frac{\partial U_r}{\partial \theta} + \frac{U_r}{r \sin \theta} \frac{\partial U_r}{\partial y} - \frac{U_r^2 U_r^2}{r} - \frac{U_r^2 \cot \theta}{r}) \hat{\theta} \\ & = \partial_r^2 + \frac{1}{r} \partial_r + \frac{1}{r^2} \partial_\theta^2 + \frac{U_r}{r \sin \theta} \partial_y^2) (\hat{r} \partial_r + \frac{\hat{\theta}}{r} \partial_\theta + \frac{\hat{\phi}}{r \sin \theta} \partial_q^2) \\ & = \partial_r^2 + \frac{1}{r} \partial_r + \frac{1}{r^2} \partial_\theta^2 + \frac{1}{r^2 \sin^2 \theta} \partial_y^2 + \frac{1}{r} \partial_r + \frac{cot\theta}{r^2} \partial_\theta^2) \\ & = \frac{\partial^2 U_r}{\partial r^2} \hat{r} + \frac{\partial^2 U_r}{\partial r} \hat{\theta} + \frac{\partial^2 U_r}{r \sin \theta} \hat{r} + \frac{1}{r^2 \partial \theta} \hat{r} + \frac{1$$

$$= \left(\frac{\partial^{2}V_{1}}{\partial Y^{2}} + \frac{2}{r}\frac{\partial^{2}V_{1}}{\partial r} + \frac{1}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta^{2}} + \frac{1}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y^{2}} + \frac{\cot\theta}{r^{2}}\frac{\partial^{2}V_{1}}{\partial \theta} - \frac{2U_{1}}{r^{2}} - \frac{2}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta} - \frac{2U_{1}}{r^{2}}\right)$$

$$- \frac{2}{r^{2}\sin\theta}\frac{\partial^{2}V_{2}}{\partial y})\hat{r}$$

$$+ \left(\frac{\partial^{2}V_{2}}{\partial r^{2}} + \frac{2}{r}\frac{\partial^{2}V_{2}}{\partial r} + \frac{1}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta^{2}} + \frac{1}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y} + \frac{\cot\theta}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta} + \frac{2}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta} - \frac{U_{2}}{r^{2}\sin^{2}\theta}\right)$$

$$- \frac{2Ux_{2}\theta}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y})\hat{\theta}$$

$$+ \left(\frac{\partial^{2}U_{2}}{\partial r^{2}} + \frac{2}{r}\frac{\partial^{2}V_{2}}{\partial r} + \frac{1}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta^{2}} + \frac{1}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y^{2}} + \frac{\cot\theta}{r^{2}}\frac{\partial^{2}V_{2}}{\partial \theta} - \frac{U_{2}\theta}{r^{2}\sin^{2}\theta} + \frac{2}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y} + \frac{2}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y} + \frac{2}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y} - \frac{U_{2}\theta}{r^{2}\sin^{2}\theta} + \frac{2}{r^{2}\sin^{2}\theta}\frac{\partial^{2}V_{2}}{\partial y} + \frac{2}{r^{2}\sin^{2}\theta}$$