

Differential Equations  
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# A Differential Look at the Watt's Governor

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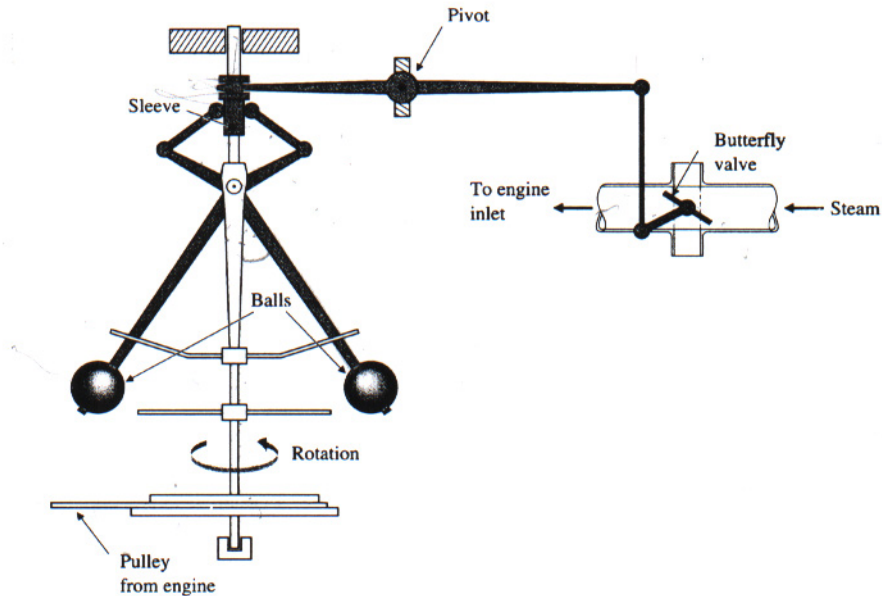


# Introduction

Invented by James Watt in the late 1700s, a governor is an automated speed control that ushered in the industrial revolution.

- Mathematical model.
- Bifurcation.
- Damping.

# Watt's Governor



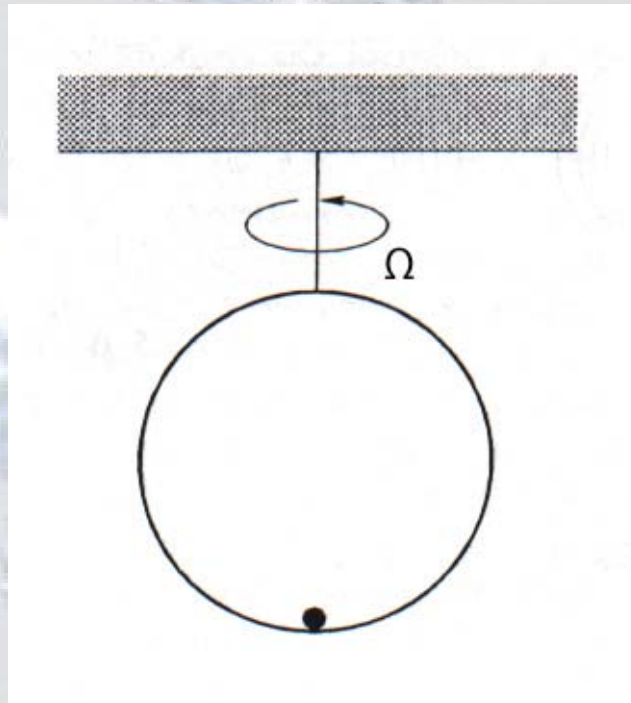
The Watt's governor controlling a steam engine.



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# A Simplified Version

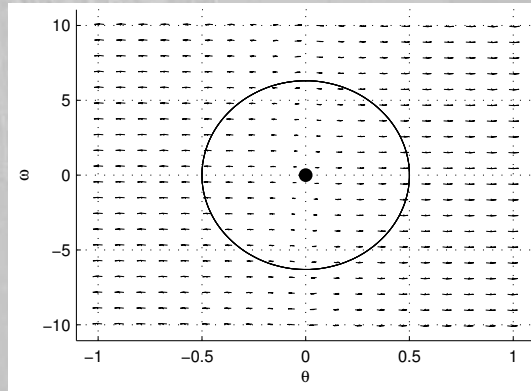


Ball-bearing in a rotating hoop.

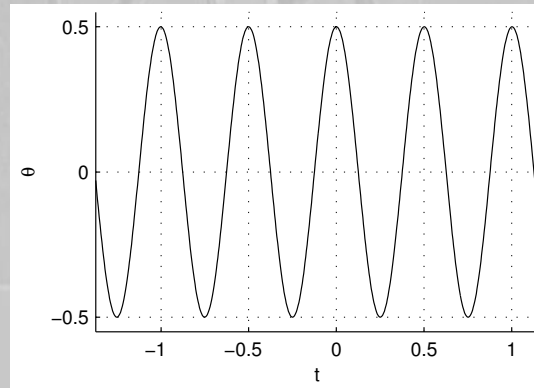


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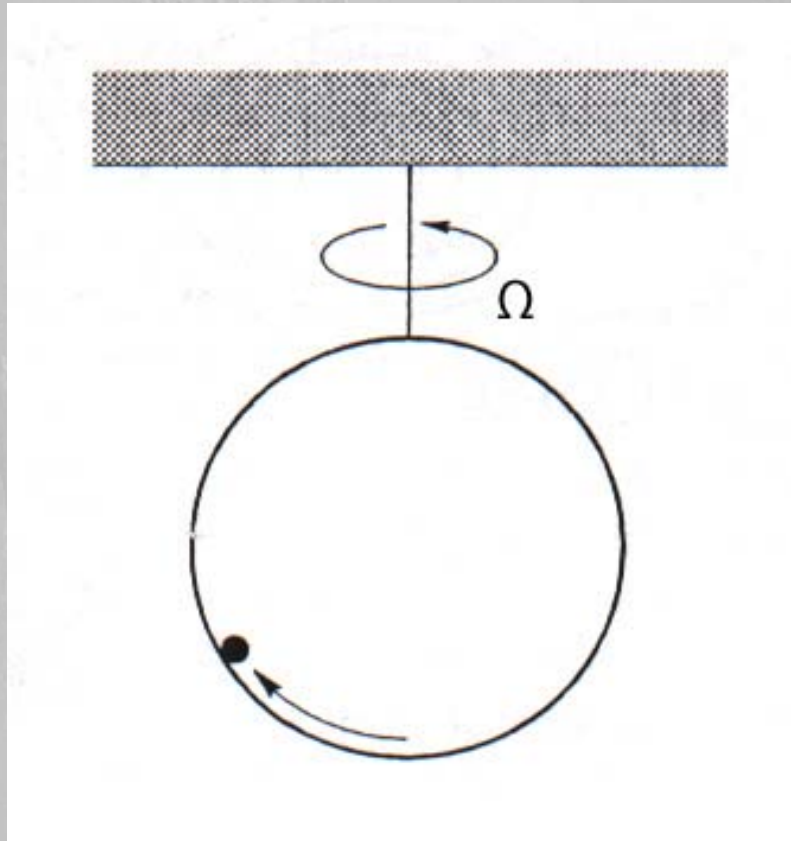


Phase plane for  $\Omega = 1$  rad/sec.



$\theta$  vs.  $t$  for  $\Omega = 1$  rad/sec.



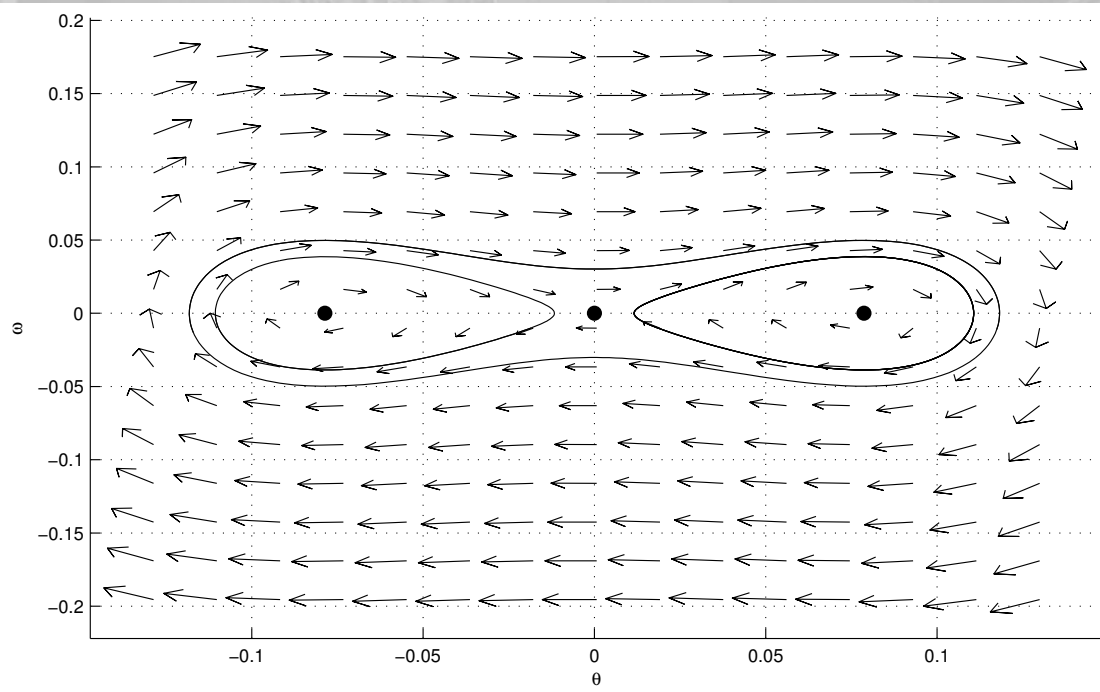


For  $\Omega > 12 \text{ rad/sec}$  the ball moves towards a new equilibrium point.





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$\theta$  vs.  $t$  for  $\Omega = 13$  rad/sec.



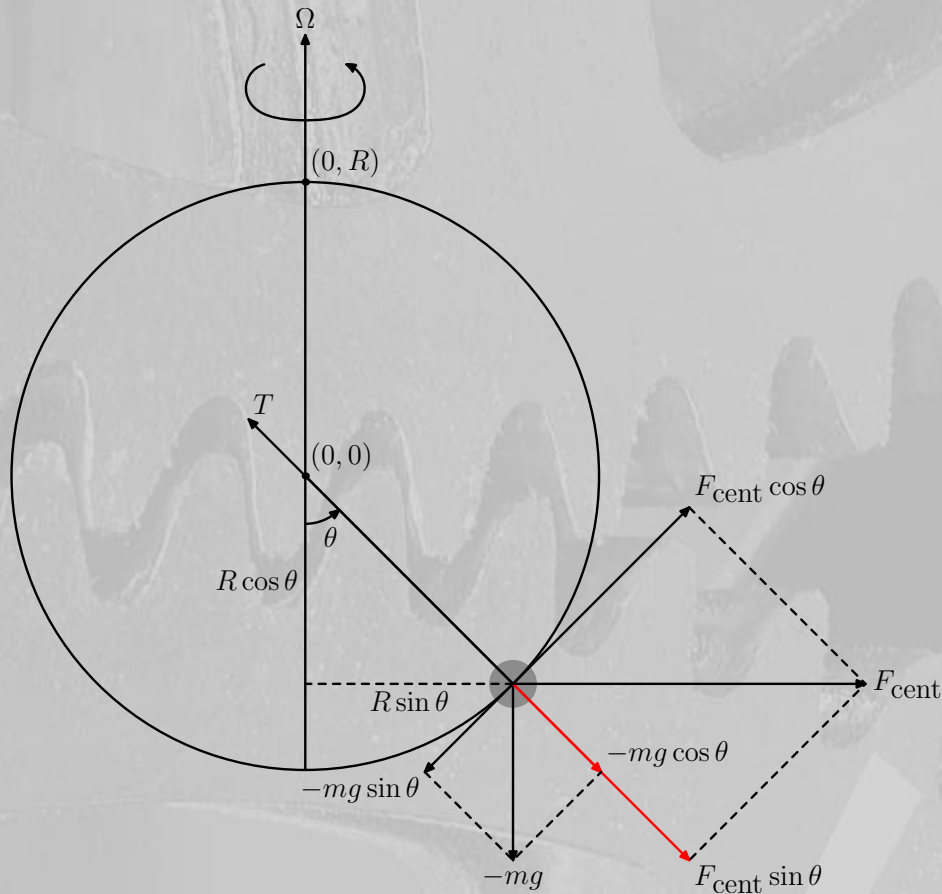


# Identifying the Forces

- Identify the forces that always balance.
- Identify the forces that do not always balance.
- Sum the forces to derive the equations.

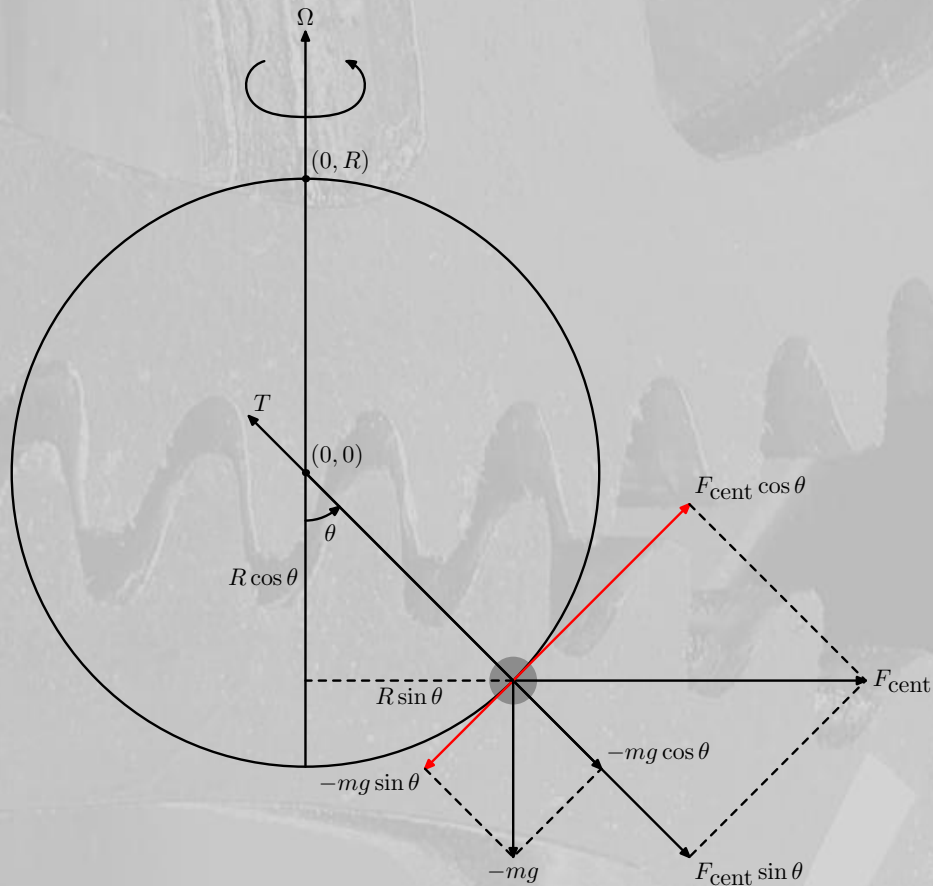






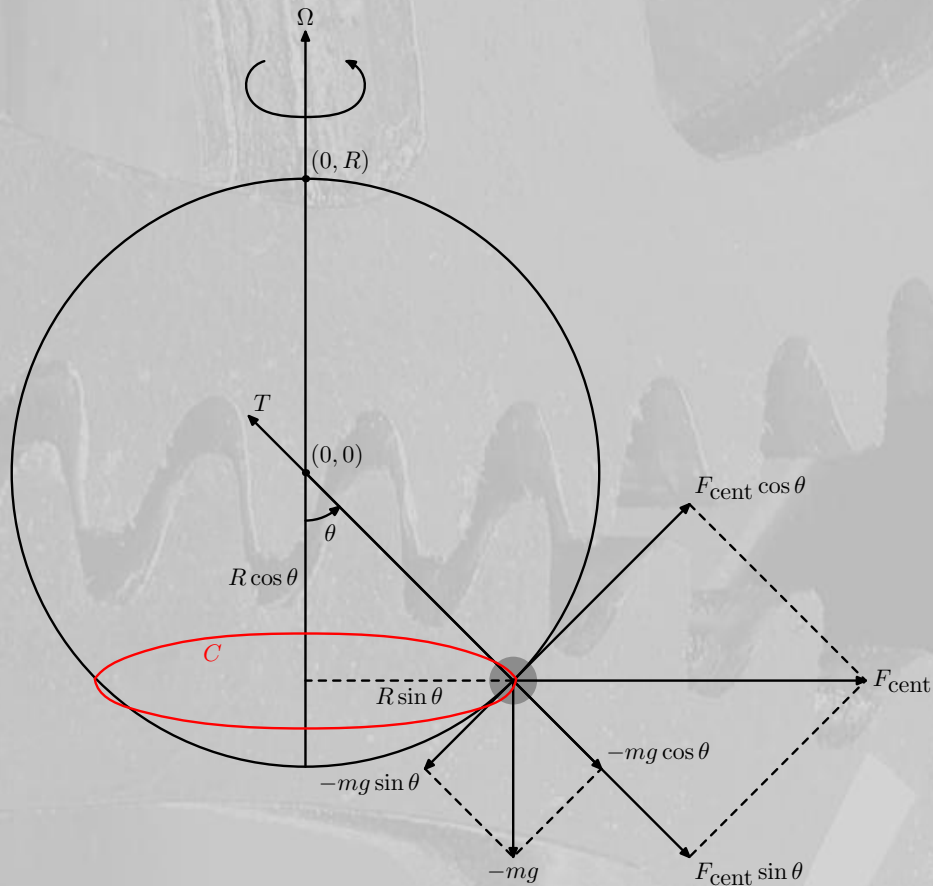
Forces opposing the normal force.





Tangential forces in the vertical plane.





The horizontal path of the ball.





## Finding $F_{\text{cent}}$

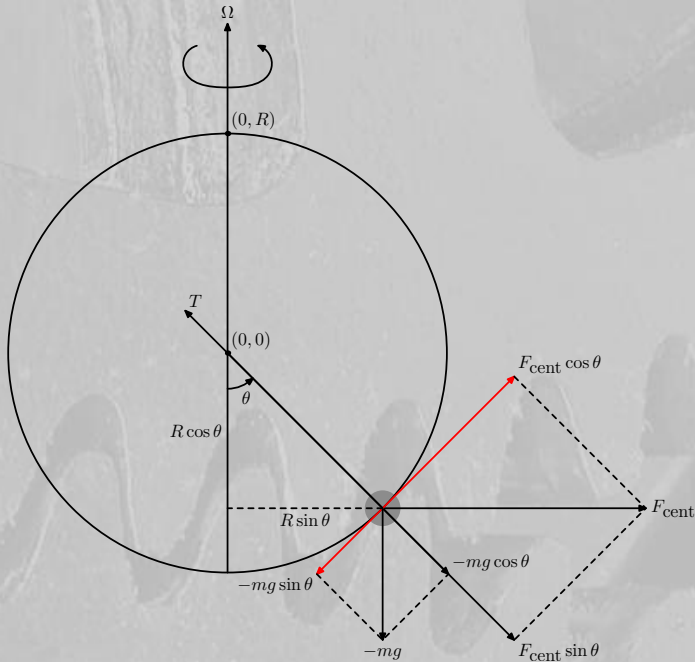
Recall the kinematic identities, and our values.

$$\begin{cases} v_{\text{lin}} = r v_{\text{ang}} \\ a_r = \frac{v_{\text{lin}}^2}{r} \\ F_{\text{cent}} = m a_r \end{cases} \quad \begin{cases} v_{\text{lin}} = (R \sin \theta) \Omega \\ a_r = \frac{[(R \sin \theta) \Omega]^2}{R \sin \theta} = (R \sin \theta) \Omega^2 \\ F_{\text{cent}} = m (R \sin \theta) \Omega^2 \end{cases}$$

In our case,  $\Omega$  is the angular velocity  $v_{\text{ang}}$ , about the center of  $C$  and the radius is  $R \sin \theta$ . The centrifugal force acting on the ball is the mass times  $a_r$ .

$$F_{\text{cent}} = m \Omega^2 R \sin \theta.$$





$$\begin{aligned}ma_T &= F_{\text{cent}} \cos \theta - mg \sin \theta \\mR\ddot{\theta} &= m\Omega^2 R \sin \theta \cos \theta - mg \sin \theta \\mR\ddot{\theta} &= m\Omega^2 R \sin \theta \cos \theta - mg \sin \theta \\\ddot{\theta} &= \Omega^2 \sin \theta \cos \theta - \frac{g}{R} \sin \theta\end{aligned}\tag{1}$$





$$\ddot{\theta} = \Omega^2 \sin \theta \cos \theta - \frac{g}{R} \sin \theta$$

In order to use this equation we must first transpose it into two first order equations.

$$\begin{cases} \theta = \theta \\ \omega = \dot{\theta} \end{cases} \quad \begin{cases} \dot{\theta} = \omega \\ \dot{\omega} = \ddot{\theta} = \Omega^2 \sin \theta \cos \theta - \frac{g}{R} \sin \theta \end{cases}$$

An equilibrium angle means that the forces are balanced and the acceleration is zero.



Set the right side equal to zero.

$$\ddot{\theta} = 0$$

$$\Omega^2 \sin \theta \cos \theta - \frac{g}{R} \sin \theta = 0$$

$$\sin \theta (\Omega^2 \cos \theta - \frac{g}{R}) = 0$$

Therefore,

$$\sin \theta = 0 \quad \text{or} \quad \Omega^2 \cos \theta - \frac{g}{R} = 0.$$

When  $\sin \theta = 0$ ,  $\theta = 0$  or  $\pi$ . To find other equilibrium angles we set the other factor equal to zero.





$$\begin{aligned}\Omega^2 \cos \theta - \frac{g}{R} &= 0 \\ \cos \theta &= \frac{g/R}{\Omega^2}\end{aligned}\quad (2)$$

Cosine is never greater than 1 so we seek  $\Omega$ s that make the right side less than or equal to 1.

$$\begin{aligned}\frac{g}{R\Omega_0^2} &\leq 1. \\ \frac{g}{R} &\leq \Omega_0^2 \\ \sqrt{\frac{g}{R}} &\leq \Omega_0\end{aligned}\quad (3)$$





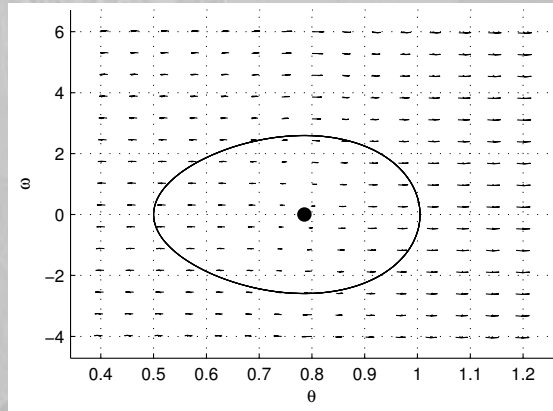
In our case the  $\Omega$  where bifurcation occurs is,

$$\begin{aligned}\sqrt{\frac{9.8}{.06}} &\leq \Omega_0 \\ 12.78 &\leq \Omega_0.\end{aligned}\tag{4}$$

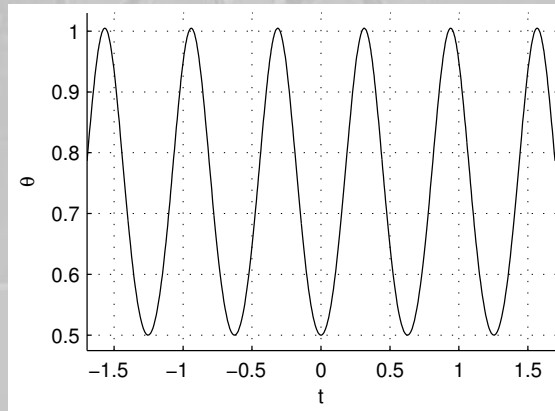
Now we find the  $\Omega$  that produces  $\theta = \pi/4$ .

$$\begin{aligned}\cos \theta &= \frac{g/R}{\Omega^2} \\ \cos \frac{\pi}{4} &= \frac{9.8/.06}{\Omega^2} \\ \sqrt{\frac{9.8}{.06 \cos \frac{\pi}{4}}} &= \Omega \\ 15.2 &= \Omega.\end{aligned}\tag{5}$$





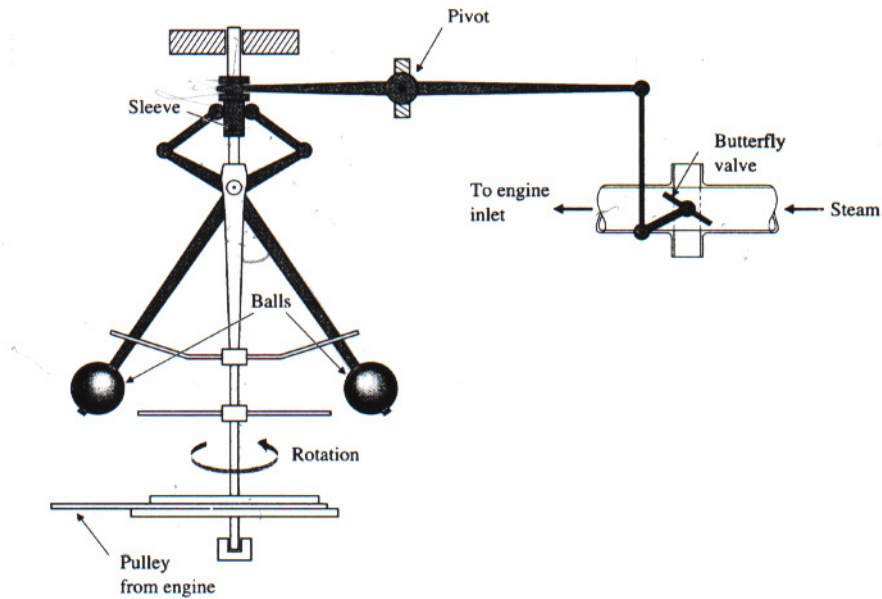
For  $\Omega = 15.2$  rad/sec.



$\theta$  vs.  $t$  for  $\Omega = 15.2$  rad/sec.



# Watt's Governor

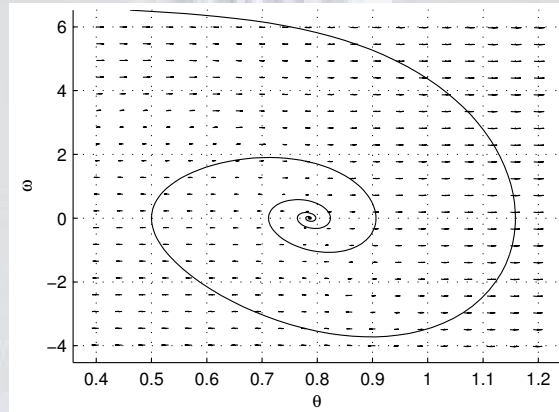


Now we have a governor that will maintain the desired angle but oscillates perpetually. How can we improve this performance?

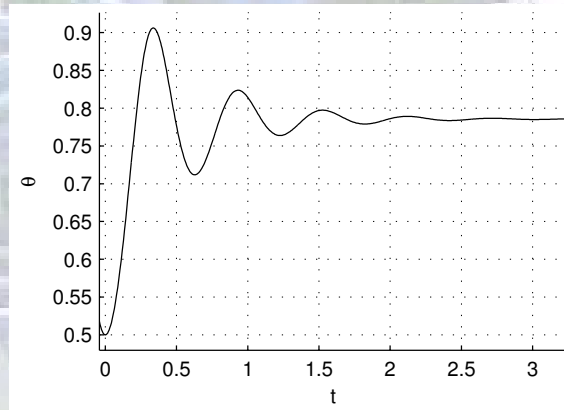
$$\ddot{\theta} = \Omega^2 \sin \theta \cos \theta - \frac{g}{R} m g \sin \theta - k \frac{\dot{\theta}}{m} \quad (6)$$

The damping term is proportional to the angular velocity (in the vertical plane) and is divided by the mass.



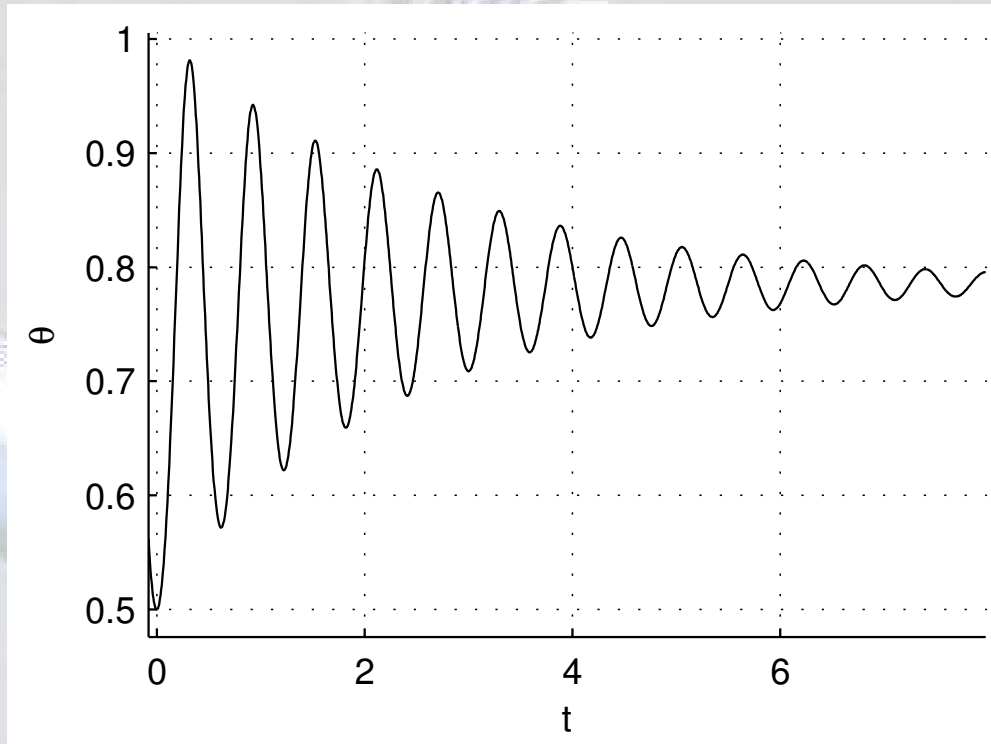


$\Omega = 15.2$  rad/sec with damping term.



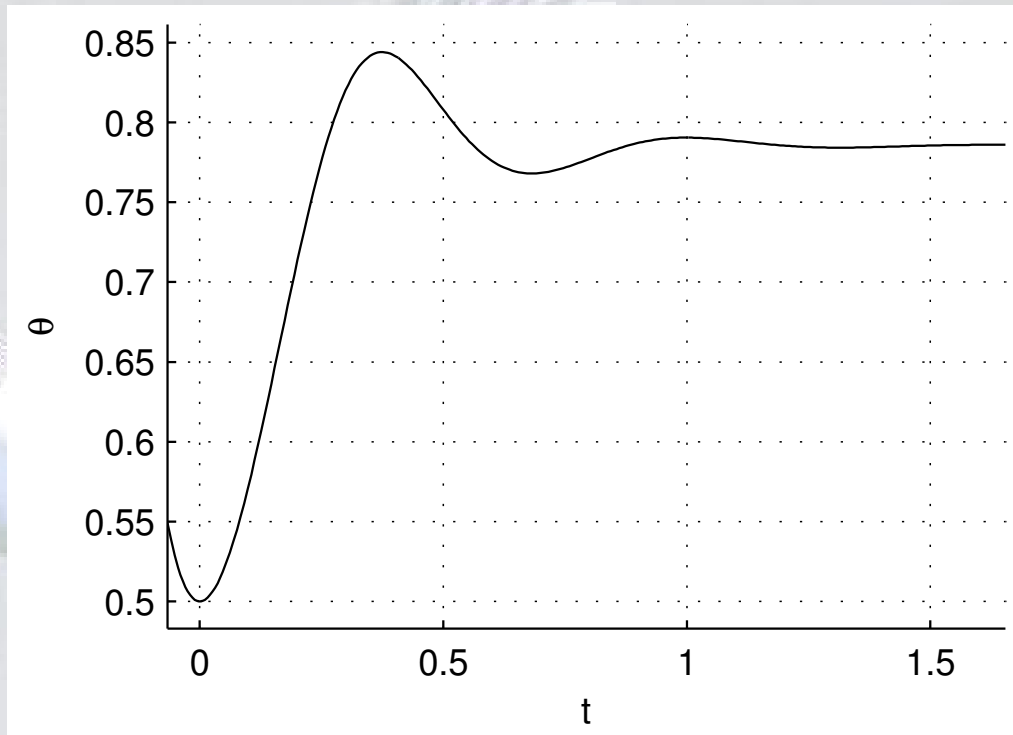
$\theta$  vs.  $t$  for  $\Omega = 15.2$  and damping term.





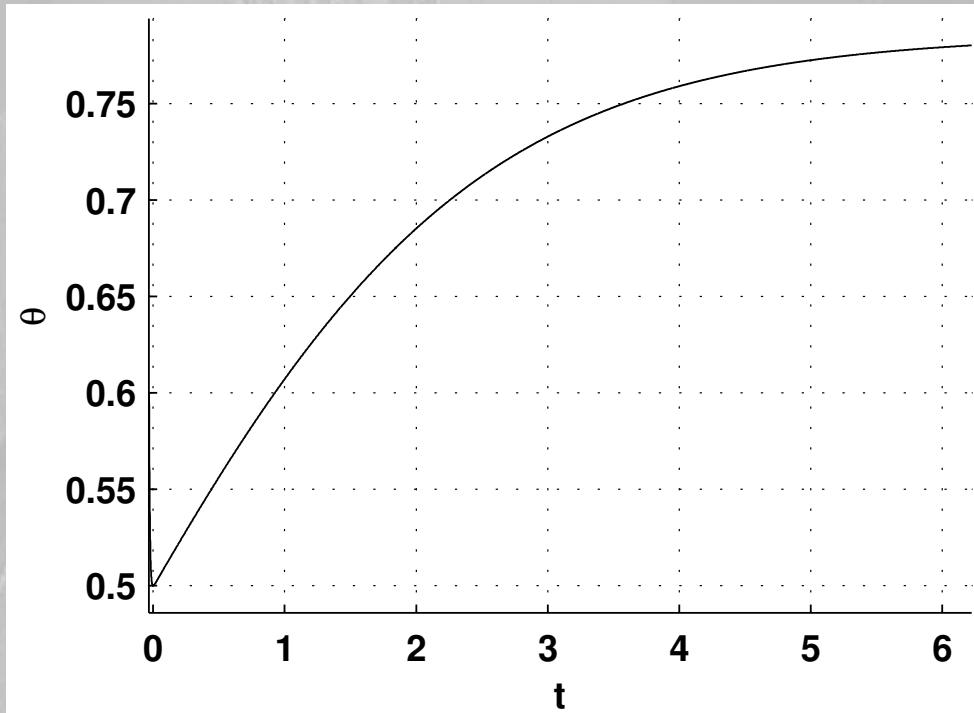
$\theta$  vs.  $t$  for  $\Omega = 15.2$ , damping term, and  $m = 50g$ .





$\theta$  vs.  $t$  for  $\Omega = 15.2$ , damping term, and  $m = 5g$ .





$\theta$  vs.  $t$  for  $\Omega = 15.2$ , damping term, and  $m = .25g$ .

As you can see this also would not be a governor of optimum design. When designing a governor one would have to experiment with the parameters and would undoubtedly be somewhere between  $5g$  and  $1/4g$ .







# Putting It All Together

We have a governor design that will maintain the desired

- Changing  $R$  only effects the where the critical  $\Omega$ s occur but not the oscillatory behavior.
- Increasing the mass reduces the effects of damping, reducing mass increases the effects of damping.
- Changing the damping term has an inverse effect as changing the mass.

