

# note : A simple derivation of Lindblad equation

Orkesh Nurbolat

September 28, 2020

S(principal) and B(bath) are 2 system in  $\mathcal{H}_S$  and  $\mathcal{H}_B$  (L:41)

$$\hat{H}(t) = \hat{H}_s \otimes \mathbf{1}_B + \mathbf{1}_S \otimes \hat{H}_B + \alpha \hat{H}_{SB} \quad (1)$$

for the sake of simplicity : (L:50)

$$\hat{H}(t) = \hat{H}_S + \hat{H}_B + \alpha \hat{H}_{SB} \quad (2)$$

there is : (L:55)

$$\frac{d\hat{\rho}_{SB}}{dt} = -\frac{i}{\hbar} [\hat{H}_S + \hat{H}_B + \alpha \hat{H}_{SB}, \hat{\rho}_{SB}] \quad (3)$$

Hamiltonian is timed as : (L:64)

$$\hat{H}(t) = e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \hat{H}_{SB} e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \quad (4)$$

while the : (L:74)

$$\begin{aligned} \hat{\rho}(t) &= e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \hat{\rho}_{SB} e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \\ \hat{\rho}_{SB}(t) &= e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \hat{\rho} e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \end{aligned} \quad (5)$$

**note** that on the rear end of an equation , central elements are not timed , they are the things at a certain constant time .

(L:91)

$$\frac{d\hat{\rho}(t)}{dt} = -\frac{i}{\hbar} [\hat{H}(t), \hat{\rho}(t)] \quad (6)$$

noting that : (L:98)

$$\hat{\rho}_S(t) = \text{Tr}_B [\hat{\rho}_{SB}(t)] \quad (7)$$

and then we have that (L:106)

$$\hat{\rho}(t) = \hat{\rho}(0) - \frac{i}{\hbar} \alpha \int_0^t [\hat{H}(t'), \hat{\rho}(t')] dt' \quad (8)$$

(L:114)

$$\frac{d\hat{\rho}(t)}{dt} = -\frac{i}{\hbar} [\hat{H}(t), \hat{\rho}(0)] - \frac{\alpha^2}{\hbar^2} \left[ \hat{H}(t), \int_0^t [\hat{H}(t'), \hat{\rho}(t')] dt' \right] \quad (9)$$

considering the relationship between the SB and the S we know that : being aware that : so here what we care about is that :

(L:136)

$$\begin{aligned} & e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \left( [\hat{H}, \hat{\rho}] \right) e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \\ &= e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \left( \hat{H}\hat{\rho} \right) e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} - e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \left( \hat{\rho}\hat{H} \right) e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \end{aligned} \quad (10)$$

and in this context what I know is that : (L:151)

$$\begin{aligned} & e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \left( \hat{H}\hat{\rho} \right) e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \\ &= \left( e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \hat{H} e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \right) \left( e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} \hat{\rho} e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} \right) \end{aligned} \quad (11)$$

by the (11) , we know that this can be done so by this we can say that (L:171)

$$e^{-i/\hbar(\hat{H}_S + \hat{H}_B)t} [\hat{H}, \hat{\rho}] e^{i/\hbar(\hat{H}_S + \hat{H}_B)t} = [\hat{H}(t), \hat{\rho}(t)] \quad (12)$$

note that here the first expression starts from a moment there is no time on it , the time is given by the evaluation we can have the **Born Approximation** applied to (9)

now we like to cut the bath part out of the picture , so we are going to do an trace over it : according to the author :  $\hat{H}$  is defined by the  $\hat{H}_S, \hat{H}_B$  and the interaction term as shown as (2) , so there is a way to define  $\hat{H}_S, \hat{H}_B$  without losing generality , so that (L:215)

$$\text{Tr}_{B \rightarrow \frac{k\alpha}{\hbar}} [\hat{H}(t), \hat{\rho}(0)] [=] 0 \quad (13)$$

if that is true then we could have that : (L:224)

$$\frac{d\hat{\rho}_S(t)}{dt} = -\frac{\alpha^2}{\hbar^2} \text{Tr}_B \left[ \hat{H}(t), \int_0^t [\hat{H}(t'), \hat{\rho}(t')] dt' \right] \quad (14)$$

integrating 14 from  $t$  to the  $t'$  yields: (L:239)

$$\hat{\rho}_S(t') - \hat{\rho}_S(t) = -\frac{\alpha^2}{\hbar^2} \int_t^{t'} \text{Tr}_B \left[ \hat{H}(n'), \int_0^{n'} [\hat{H}(t''), \hat{\rho}(t'')] dt'' \right] dn' \quad (15)$$

or integrating  $t'$  to the  $t$  yields: (L:255)

$$\hat{\rho}_S(t) - \hat{\rho}_S(t') = -\frac{\alpha^2}{\hbar^2} \int_{t'}^t \text{Tr}_B \left[ \hat{H}(n'), \int_0^{n'} [\hat{H}(t''), \hat{\rho}(t'')] dt'' \right] dn' \quad (16)$$

difference of  $\hat{\rho}(t)$  and  $\hat{\rho}(t')$  is of magnitude of  $\alpha^2$  so primary time equation could be written in the following ways without violating **Born approximation** (L:273)

$$\frac{d\hat{\rho}_S(t)}{dt} = -\frac{\alpha^2}{\hbar^2} \text{Tr}_B \left[ \hat{H}(t), \int_0^t [\hat{H}(t'), \hat{\rho}(t)] dt' \right] \quad (17)$$

lets say Bath and System are in full interaction(  $\alpha = 1$  ), and this is an approximation , so there is : (L:286)

$$\frac{d\hat{\rho}_S(t)}{dt} = -\frac{1}{\hbar^2} \text{Tr}_B \left[ \hat{H}(t), \int_0^t [\hat{H}(t'), \hat{\rho}(t)] dt' \right] \quad (18)$$

writing that, (ignoring Bath Hamiltonian time as well) (L:299)

$$\begin{aligned} \hat{\rho}(t) &= \hat{\rho}_S(t) \otimes \hat{\rho}_B(t) \\ &= \hat{\rho}_S(t) \hat{\rho}_B(t) \end{aligned} \quad (19)$$

and take the time to be infinitely long , then : (L:304)

$$\frac{d\hat{\rho}_S(t)}{dt} = -\frac{1}{\hbar^2} \text{Tr}_B \left[ \hat{H}(t), \int_0^\infty [\hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B] dt' \right] \quad (20)$$

## 1 LINDBLAD EQUATION

### 1.1 Hamiltonians and Operators

so there is : (L:320)

$$\hat{H}_{SB} = \hbar \left( \hat{S} \hat{B}^\dagger + \hat{S}^\dagger \hat{B} \right) \quad (21)$$

of course with multiple expression of  $\hat{S}$  it could be expressed as : so : (L:331)

$$\hat{H}_{SB} = \hbar \sum \left( \hat{L} \hat{B}^\dagger + \hat{L}^\dagger \hat{B} \right) \quad (22)$$

and then this  $\hat{L}$  is system and Bath related , how does it have in the system is not interested (L:341)

$$\hat{S}(t) = \hat{S} \quad (23)$$

is not affected by the interaction-picture frame , defining the bath is a bath of bosons : (L:346)

$$\hat{H}_B = \hbar \sum_k \omega_k \hat{a}_k^\dagger \hat{a}_k \quad (24)$$

so the bath have  $k$  modes , and  $\hat{a}^\dagger, \hat{a}$  creates or annihilates them but the operator (L:354)

$$\hat{B} = \sum_k g_k^* \hat{a}_k \quad (25)$$

is not  $\hat{H}_B$  and here  $g_k$  is said to be complex coefficients represing coupling constants then , in the interaction picture,(of course here  $\hat{H}_S + \hat{H}_B$  is not used but only that the  $\hat{H}_B$  is used ) (L:366)

$$\hat{B}(t) = e^{i/\hbar \hat{H}_B t} \hat{B} e^{-i/\hbar \hat{H}_B t} \quad (26)$$

in (26) , taylor slicing it gives the exp parts eventually fill my hand with : (L:375)

$$\begin{aligned} \hat{B}(t) &= e^{i/\hbar \hat{H}_B t} \hat{B} e^{-i/\hbar \hat{H}_B t} \\ &= \sum_k e^{i \sum_{k'} \mathbb{I}(\omega_{k'} \hat{a}_{k'}^\dagger \hat{a}_{k'}) t} g_k^* \hat{a}_k e^{-i \sum_{k''} \mathbb{I}(\omega_{k''} \hat{a}_{k''}^\dagger \hat{a}_{k''}) t} \end{aligned} \quad (27)$$

for the boson we know that : (L:392)

$$\begin{aligned} [\hat{a}_k, \hat{a}_{k'}] &= 0 \\ [\hat{a}_k, \hat{a}_{k'}^\dagger] &= \delta_{k,k'} \\ [\hat{a}_{k'}^\dagger, \hat{a}_k] &= -\delta_{k,k'} \\ [\hat{a}_{k'}^\dagger \hat{a}_{k'}, \hat{a}_k] &= [\hat{a}_{k'}^\dagger, \hat{a}_k] \hat{a}_{k'} + \hat{a}_{k'}^\dagger [\hat{a}_{k'}, \hat{a}_k] = -\delta_{k,k'} \hat{a}_{k'} \end{aligned} \quad (28)$$

however there is : (L:418)

$$\begin{aligned} [\hat{a}_n^\dagger \hat{a}_n \hat{a}_m^\dagger \hat{a}_m, \hat{a}_k] &= [\hat{a}_n^\dagger \hat{a}_n, \hat{a}_k] \hat{a}_m^\dagger \hat{a}_m + \hat{a}_n^\dagger \hat{a}_n [\hat{a}_m^\dagger \hat{a}_m, \hat{a}_k] \\ &= \delta_{n,k} \hat{a}_k \hat{a}_m^\dagger \hat{a}_m + \delta_{m,k} \hat{a}_n^\dagger \hat{a}_n \hat{a}_k \end{aligned} \quad (29)$$

taking it another step further , we have : we also remind ourself that here we could really just take that things, and ofcourse this S is an arbitrary, prod operation will do it in a sorted fashion given partial ordered Set, each element of it is impicitly corresponded to a Natural number (L:436)

$$\begin{aligned} \left[ \prod_{z \in S} \hat{a}_z^\dagger \hat{a}_z, \hat{a}_k \right] &= \sum_{z \in S} \delta_{z,k} \left( \prod_{m \in S \wedge m \subset z} \hat{a}_m^\dagger \hat{a}_m \right) \hat{a}_k \left( \prod_{n \in S \wedge n \supset z} \hat{a}_n^\dagger \hat{a}_n \right) \\ &= \sum_{z \in S} \delta_{z,k} \hat{a}_k \left( \prod_{m \in S \wedge m \subset z} (\hat{a}_m^\dagger \hat{a}_m - \delta_{m,k} \hat{a}_k) \right) \left( \prod_{n \in S \wedge n \supset z} \hat{a}_n^\dagger \hat{a}_n \right) \end{aligned} \quad (30)$$

switching the perspective gives : (L:466)

$$\begin{aligned}
\left[ \hat{H}_B, \hat{a}_k \right] &= -\hbar\omega_k \hat{a}_k \\
\left[ \hat{H}_B^2, \hat{a}_k \right] &= -\hbar\omega_k \left( \hat{a}_k \hat{H}_B + \hat{H}_B \hat{a}_k \right) \\
\left[ \hat{H}_B^n, \hat{a}_k \right] &= -\hbar\omega_k \sum_i^n \left( \prod_{m=0}^{i-1} \hat{H}_B \right) \hat{a}_k \left( \prod_{l=i+1}^n \hat{H}_B \right) \\
&= -\hbar\omega_k \hat{a}_k \sum_i^n \left( \prod_{m=0}^{i-1} \hat{H}_B - \hbar\omega_k \hat{a}_k \right) \left( \prod_{l=i+1}^n \hat{H}_B \right)
\end{aligned} \tag{31}$$

now : (L:499)

$$\begin{aligned}
\left[ \mathbb{e}^{\mathbb{i}\hat{H}_B t/\hbar}, \hat{a}_k \right] &= \sum_n \frac{\left[ (\mathbb{i}\hat{H}_B t/\hbar)^n, \hat{a}_k \right]}{n!} \\
&= \sum_n \frac{(\mathbb{i}t/\hbar)^n \left[ \hat{H}_B^n, \hat{a}_k \right]}{n!} \\
&= \sum_n \frac{(\mathbb{i}t/\hbar)^n \left( -\hbar\omega_k \sum_i^n \hat{H}_B^{i-1} \hat{a}_k \hat{H}_B^{n-i} \right)}{n!}
\end{aligned} \tag{32}$$

and (L:536)

$$\begin{aligned}
\left[ \mathbb{e}^{\mathbb{i}\hat{H}_B t/\hbar}, \hat{a}_k \right] \mathbb{e}^{-\mathbb{i}\hat{H}_B t/\hbar} &= \sum_n \sum_m \frac{(-1)^m (\mathbb{i}t/\hbar)^{n+m} \left( -\hbar\omega_k \sum_i^n \hat{H}_B^{i-1} \hat{a}_k \hat{H}_B^{n-i} \right) \hat{H}_B^m}{n!m!} \\
&= \sum_n \sum_m \frac{(-1)^m (\mathbb{i}t/\hbar)^{n+m} \left( -\hbar\omega_k \hat{a}_k \sum_i^n (\hat{H}_B^{i-1} - \hbar\omega_k \hat{a}_k) \hat{H}_B^{n-i} \right) \hat{H}_B^m}{n!m!}
\end{aligned} \tag{33}$$

for its nth compond ,it has alot of (30) ok , lets just use the CBH formula here which states that: (L:581)

$$\begin{aligned}
\mu(t) &= \mathbb{e}^{\lambda t} \mu \mathbb{e}^{-\lambda t} \\
\frac{\mathbf{d}\mu(t)}{\mathbf{d}t} &= \frac{\mathbf{d}\mathbb{e}^{\lambda t}}{\mathbf{d}t} \mu \mathbb{e}^{-\lambda t} + \mathbb{e}^{\lambda t} \mu \frac{\mathbf{d}\mathbb{e}^{-\lambda t}}{\mathbf{d}t} \\
&= \mathbb{e}^{\lambda t} (\lambda \mu - \mu \lambda) \mathbb{e}^{-\lambda t} \\
&= \mathbb{e}^{\lambda t} [\lambda, \mu] \mathbb{e}^{-\lambda t}
\end{aligned} \tag{34}$$

expanding at (L:591)

$$\frac{\mathbf{d}^2 \mu(t)}{\partial \mathbf{d}t^2} = \mathbb{e}^{\lambda t} [\lambda [\lambda, \mu]] \mathbb{e}^{-\lambda t} \tag{35}$$

now expanding the  $\mu(t)$  from the time 0 gives (L:602)

$$\mu(t) = \mu + [\lambda, \mu]t + \frac{1}{2!}[\lambda, [\lambda, \mu]]t^2 + \dots \quad (36)$$

well now , we know that if (L:607)

$$\begin{aligned} \lambda &= \mathbb{H}\hat{H}/\hbar \\ \mu &= \hat{a}_k \end{aligned} \quad (37)$$

so (L:612)

$$\begin{aligned} [\lambda, \mu] &= -\mathbb{H}\omega_k \hat{a}_k \\ [\lambda, [\lambda, \mu]] &= (-\mathbb{H}\omega_k)^2 \hat{a}_k \end{aligned} \quad (38)$$

so that : (L:617)

$$\begin{aligned} e^{\mathbb{H}\hat{H}t/\hbar} \hat{a}_k e^{-\mathbb{H}\hat{H}t/\hbar} &= \hat{a}_k \left( \sum_n \frac{(-\mathbb{H}\omega_k t)^n}{n!} \right) \\ &= \hat{a}_k e^{-\mathbb{H}\omega_k t} \end{aligned} \quad (39)$$

so then we had (L:630)

$$\hat{B}(t) = \sum_k g_k^* \hat{a}_k e^{-\mathbb{H}\omega_k t} \quad (40)$$

now with that (L:637)

$$\hat{H}_{SB} = \hbar \left( \hat{S} \hat{B}^\dagger + \hat{S}^\dagger \hat{B} \right) \quad (41)$$

this  $\hat{B}$  is a resemble of Jaynes-Cummings model , which is like : (L:644)

$$\hat{H} = \hbar \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right) + \hbar \left( \hat{b}^\dagger \hat{b} + \frac{1}{2} \right) + g \left( \hat{a}^\dagger \hat{b} + \hat{b}^\dagger \hat{a} \right) \quad (42)$$

and then lokoing back to the denity matrix (20) now we want to study that the comutation part : (L:652)

$$\left[ \hat{H}(t), \int_0^t \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \mathbf{d}t' \right] = \int_0^t \left[ \hat{H}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] \mathbf{d}t' \quad (43)$$

the part that cared about is : by the way , is there :  $\hat{\rho}_S(t) \hat{\rho}_B = \hat{\rho}_B \hat{\rho}_S(t)$  ? well I dont think so , becaue  $\hat{\rho}_S(t) \hat{\rho}_B = \hat{\rho}_S(t) \otimes \hat{\rho}_B$   $\hat{\rho}_S(t) = \hat{\rho}_S(t) \otimes \mathbf{1}_B$  and  $\hat{\rho}_B = \mathbf{1}_S \otimes \hat{\rho}_B$  is what it means (L:675)

$$\begin{aligned} \left[ \hat{H}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] &= \hbar \left[ \hat{S} \hat{B}^\dagger + \hat{S}^\dagger, \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] \\ &= \hbar \left[ \hat{S} \hat{B}^\dagger, \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] + \hbar \left[ \hat{S}^\dagger \hat{B}, \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] \end{aligned} \quad (44)$$

lets say for the part that (L:699)

$$\begin{aligned}
& \left[ \hat{S}\hat{B}^\dagger(t), \left[ \hat{H}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] \right] \\
&= \hbar \left[ \hat{S}\hat{B}^\dagger(t), \left[ \hat{S}\hat{B}^\dagger(t') + \hat{S}^\dagger\hat{B}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] \right] \\
&= \hbar \left( \hat{S}\hat{B}^\dagger(t) \left[ \hat{S}\hat{B}^\dagger(t') + \hat{S}^\dagger\hat{B}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] - \left[ \hat{S}\hat{B}^\dagger(t') + \hat{S}^\dagger\hat{B}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] \hat{S}\hat{B}^\dagger(t) \right) \\
&= \hbar \left( \hat{S}\hat{B}^\dagger(t)\hat{S}\hat{B}^\dagger(t')\hat{\rho}_S(t)\hat{\rho}_B + \hat{S}\hat{B}^\dagger(t)\hat{S}^\dagger\hat{B}(t')\hat{\rho}_S(t)\hat{\rho}_B - \hat{S}\hat{B}^\dagger(t)\hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t') - \hat{S}\hat{B}^\dagger(t)\hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t') \right) \\
&- \hbar \left( \hat{S}\hat{B}^\dagger(t')\hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t) + \hat{S}^\dagger\hat{B}(t')\hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t) - \hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t')\hat{S}\hat{B}^\dagger(t) - \hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t')\hat{S}\hat{B}^\dagger(t) \right) \\
&\quad (45)
\end{aligned}$$

grouping similiar terms here gives: so there is  $[\hat{\rho}_b, \hat{S}] = 0$ ,  $[\hat{\rho}_S(t), \hat{S}] \neq 0$ ,  $[\hat{\rho}_b, \hat{B}] \neq 0$ ,  $[\hat{\rho}_S(t), \hat{B}] = 0$  we can then have (L:756)

$$\begin{aligned}
& \left[ \hat{S}\hat{B}^\dagger(t), \left[ \hat{H}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] \right] \\
&= \hbar \left( \hat{S}\hat{S}\hat{\rho}_S(t)\hat{B}^\dagger(t)\hat{B}^\dagger(t')\hat{\rho}_B + \hat{S}\hat{S}^\dagger\hat{\rho}_S(t)\hat{B}^\dagger(t)\hat{B}(t')\hat{\rho}_B - \hat{S}\hat{\rho}_S(t)\hat{S}\hat{B}^\dagger(t)\hat{\rho}_B\hat{B}^\dagger(t') - \hat{S}\hat{\rho}_S(t)\hat{S}^\dagger\hat{B}^\dagger(t)\hat{\rho}_B\hat{B}(t') \right) \\
&- \hbar \left( \hat{S}\hat{\rho}_S(t)\hat{S}\hat{B}^\dagger(t')\hat{\rho}_B\hat{B}^\dagger(t) + \hat{S}^\dagger\hat{\rho}_S(t)\hat{S}\hat{B}(t')\hat{\rho}_B\hat{B}^\dagger(t) - \hat{\rho}_S(t)\hat{S}\hat{S}\hat{\rho}_B\hat{B}^\dagger(t')\hat{B}^\dagger(t) - \hat{\rho}_S(t)\hat{S}^\dagger\hat{S}\hat{\rho}_B\hat{B}(t')\hat{B}^\dagger(t) \right) \\
&\quad (46)
\end{aligned}$$

and on the other wing we have some similiar : (L:783)

$$\begin{aligned}
& \left[ \hat{S}^\dagger\hat{B}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t)\hat{\rho}_B \right] \right] \\
&= \hbar \left( \hat{S}^\dagger\hat{B}(t)\hat{S}\hat{B}^\dagger(t')\hat{\rho}_S(t)\hat{\rho}_B + \hat{S}^\dagger\hat{B}(t)\hat{S}^\dagger\hat{B}(t')\hat{\rho}_S(t)\hat{\rho}_B - \hat{S}^\dagger\hat{B}(t)\hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t') - \hat{S}^\dagger\hat{B}(t)\hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t') \right) \\
&- \hbar \left( \hat{S}\hat{B}^\dagger(t')\hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t) + \hat{S}^\dagger\hat{B}(t')\hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t) - \hat{\rho}_S(t)\hat{\rho}_B\hat{S}\hat{B}^\dagger(t')\hat{S}^\dagger\hat{B}(t) - \hat{\rho}_S(t)\hat{\rho}_B\hat{S}^\dagger\hat{B}(t')\hat{S}^\dagger\hat{B}(t) \right) \\
&\quad (47)
\end{aligned}$$

re arranging it to : (L:808)

$$\begin{aligned}
&= \hbar \left( \hat{S}^\dagger\hat{S}\hat{\rho}_S(t)\hat{B}(t)\hat{B}^\dagger(t')\hat{\rho}_B + \hat{S}^\dagger\hat{S}^\dagger\hat{\rho}_S(t)\hat{B}(t)\hat{B}(t')\hat{\rho}_B - \hat{S}^\dagger\hat{\rho}_S(t)\hat{S}\hat{B}(t)\hat{\rho}_B\hat{B}^\dagger(t') - \hat{S}^\dagger\hat{\rho}_S(t)\hat{S}^\dagger\hat{B}(t)\hat{\rho}_B\hat{B}(t') \right) \\
&- \hbar \left( \hat{S}\hat{\rho}_S(t)\hat{S}^\dagger\hat{B}^\dagger(t')\hat{\rho}_B\hat{B}(t) + \hat{S}^\dagger\hat{\rho}_S(t)\hat{S}^\dagger\hat{B}(t')\hat{\rho}_B\hat{B}(t) - \hat{\rho}_S(t)\hat{S}\hat{S}^\dagger\hat{\rho}_B\hat{B}^\dagger(t')\hat{B}(t) - \hat{\rho}_S(t)\hat{S}^\dagger\hat{S}^\dagger\hat{\rho}_B\hat{B}(t')\hat{B}(t) \right) \\
&\quad (48)
\end{aligned}$$

## 1.2 the particle trace

assume there is no particle in the bath : (L:827)

$$\text{Tr}_B \left[ \hat{B}(t)\hat{B}(t')\hat{\rho}_B \right] = \text{Tr}_B \left[ \hat{B}^\dagger(t)\hat{B}^\dagger(t')\hat{\rho}_B \right] = 0 \quad , \quad \forall t, t' \quad (49)$$

with this only selective are left over, like (L:839)

$$\begin{aligned}
\text{Tr}_B \left[ \hat{S}B(t)^\dagger, \left[ \hat{H}(t'), \hat{\rho}_B \hat{\rho}_S(t) \right] \right] &= \hbar \hat{S} \hat{S}^\dagger \hat{\rho}_S(t) \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] \\
&\quad - \hbar \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{\rho}_B \hat{B}(t') \right] \\
&\quad - \hbar \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \text{Tr}_B \left[ \hat{B}(t') \hat{\rho}_B \hat{B}^\dagger(t) \right] \\
&\quad + \hbar \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} \text{Tr}_B \left[ \hat{\rho}_B \hat{B}(t') \hat{B}^\dagger(t) \right]
\end{aligned} \tag{50}$$

and on the other side there is : (L:854)

$$\begin{aligned}
\text{Tr}_B \left[ \hat{S}^\dagger \hat{B}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] &= \hbar \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \\
&\quad - \hbar \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \text{Tr}_B \left[ \hat{B}(t) \hat{\rho}_B \hat{B}^\dagger(t') \right] \\
&\quad - \hbar \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{\rho}_B \hat{B}(t) \right] \\
&\quad + \hbar \hat{\rho}_S(t) \hat{S} \hat{S}^\dagger \text{Tr}_B \left[ \hat{\rho}_B \hat{B}^\dagger(t') \hat{B}(t) \right]
\end{aligned} \tag{51}$$

now it says that there is ciclic properties of trace : which paissically means that : (L:874)

$$\text{Tr} \left[ \hat{A} \hat{B} \hat{C} \hat{D} \right] = \text{Tr} \left[ \hat{B} \hat{C} \hat{D} \hat{A} \right] = \text{Tr} \left[ \hat{C} \hat{D} \hat{A} \hat{B} \right] = \text{Tr} \left[ \hat{D} \hat{A} \hat{B} \hat{C} \right] \tag{52}$$

rearangement shows : (L:881)

$$\begin{aligned}
\text{Tr}_B \left[ \hat{S}B(t)^\dagger, \left[ \hat{H}(t'), \hat{\rho}_B \hat{\rho}_S(t) \right] \right] &= \hbar \hat{S} \hat{S}^\dagger \hat{\rho}_S(t) \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] - \hbar \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] \\
&\quad + \hbar \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right] - \hbar \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right] \\
&= \hbar \left( \hat{S} \hat{S}^\dagger \hat{\rho}_S(t) - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] \\
&\quad + \hbar \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right]
\end{aligned} \tag{53}$$

and (L:899)

$$\begin{aligned}
\text{Tr}_B \left[ \hat{S}^\dagger \hat{B}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] &= \hbar \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] - \hbar \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \\
&\quad + \hbar \hat{\rho}_S(t) \hat{S} \hat{S}^\dagger \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] - \hbar \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] \\
&= \hbar \left( \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \\
&\quad + \hbar \left( \hat{\rho}_S(t) \hat{S} \hat{S}^\dagger - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right]
\end{aligned} \tag{54}$$



since we like to see how this works with the system hamiltonian , we would like to continue jobs at (44) and then what it give us is that : (L:927)

$$\begin{aligned} \text{Tr}_B \left[ \hat{H}(t), \left[ \hat{H}(t'), \hat{\rho}_S(t) \hat{\rho}_B \right] \right] &= \hbar^2 \left( \hat{S} \hat{S}^\dagger \hat{\rho}_S(t) - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] \\ &+ \hbar^2 \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right] \\ &+ \hbar^2 \left( \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \\ &+ \hbar^2 \left( \hat{\rho}_S(t) \hat{S} \hat{S}^\dagger - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] \end{aligned} \quad (55)$$

so that is it

### 1.3 The expansion of the integrand of the master equation

now we look back to the integration mentioned at the(20) for convenience , we could have that : (L:951)

$$\begin{aligned} F(t) &= \int_0^t \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \mathbf{d}t' \\ G(t) &= \int_0^t \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] \mathbf{d}t' \end{aligned} \quad (56)$$

and if we take the conjugate , it some how just exchange the  $t'$  and  $t$  I think the density matrix has its trace elements summed to 1 but conjugation of the operator  $\hat{B}$  which is not promissed to be hermitian will look like : (L:967)

$$\begin{aligned} \hat{B}^*(t) &= \sum_k g_k \hat{a}_k^* e^{i\omega_k t} \\ (\hat{B}^\dagger(t))^* &= \sum_k g_k^* (\hat{a}_k^\dagger)^* e^{-i\omega_k t} \end{aligned} \quad (57)$$

we know that for the system (L:980)

$$\begin{aligned} \hat{B}(t) \hat{B}^\dagger(t') &= \sum_k \sum_{k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} \hat{a}_k \hat{a}_{k'}^\dagger \\ \left( \hat{B}(t) \hat{B}^\dagger(t') \right)^* &= \sum_k \sum_{k'} g_k g_{k'}^* e^{i(\omega_k t - \omega_{k'} t')} (\hat{a}_k \hat{a}_{k'}^\dagger)^* \\ &= \sum_{k'} \sum_k g_k^* g_{k'} e^{-i(\omega_{k'} t' - \omega_k t)} (\hat{a}_k)^* (\hat{a}_{k'}^\dagger)^* \end{aligned} \quad (58)$$

so I didn't get this . in the perspective of trace : (L:1004)

$$\begin{aligned}
\text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] &= \sum_n \langle n | \sum_k \sum_{k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} \hat{a}_k \hat{a}_{k'}^\dagger \sum_i p_i |i\rangle \langle i| |n\rangle \\
&= \sum_n \langle n | \sum_k \sum_{k'} g_k^* g_{k'} p_n e^{-i(\omega_k t - \omega_{k'} t')} \hat{a}_k \hat{a}_{k'}^\dagger |n\rangle \\
&= \sum_n \langle n | \sum_k g_k^* g_k p_n e^{-i(\omega_k t - \omega_k t')} \hat{a}_k \hat{a}_k^\dagger |n\rangle
\end{aligned} \tag{59}$$

this expression is kind of all real (except the exponential part) if there the ,  $\hat{\rho}_B = \sum_i |i\rangle \langle i| p_i$  is true , if so turning it around would be like (this might not be true , need to recheck ): (L:1038)

$$\begin{aligned}
\left( \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \right)^* &= \sum_n \langle n | \sum_k g_k^* g_k p_n e^{-i(\omega_{k'} t' - \omega_k t)} \hat{a}_k \hat{a}_{k'}^\dagger |n\rangle \\
&= \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right]
\end{aligned} \tag{60}$$

and this gives me a little acceptability to the expression : (L:1057)

$$\begin{aligned}
F(t)^* &= \int_0^t \left( \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] \right)^* dt' \\
&= \int_0^t \left( \text{Tr}_B \left[ \hat{B}(t') \hat{B}^\dagger(t) \hat{\rho}_B \right] \right) dt' \\
G(t)^* &= \int_0^t \left( \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] \right)^* dt' \\
&= \int_0^t \text{Tr}_B \left[ \hat{B}^\dagger(t) \hat{B}(t') \hat{\rho}_B \right] dt'
\end{aligned} \tag{61}$$

so we have that : (L:1079)

$$\begin{aligned}
\frac{d\hat{\rho}_S(t)}{dt} &= - \left( \hat{S} \hat{S}^\dagger \hat{\rho}_S(t) - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) G(t)^* \\
&\quad - \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) F(t) \\
&\quad^* - \left( \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) F(t) \\
&\quad - \left( \hat{\rho}_S(t) \hat{S} \hat{S}^\dagger - \hat{S}^\dagger \hat{\rho}_S(t) \hat{S} \right) G(t)
\end{aligned} \tag{62}$$

#### 1.4 The bath specification

the initial vacuum state of a bath is like : (L:1098)

$$\hat{\rho}_B = (|0\rangle\langle 0| \dots) \otimes (|0\rangle\langle 0| \dots) \tag{63}$$

so that : (L:1109)

$$\begin{aligned}\text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] &= \text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') (|0\rangle|0\rangle\dots\dots) \otimes (\langle 0|\langle 0|\dots\dots) \right] \\ \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] &= \text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) (|0\rangle|0\rangle\dots\dots) \otimes (\langle 0|\langle 0|\dots\dots) \right]\end{aligned}\quad (64)$$

calling some bath states as  $\{|b\rangle\}$  then we can say : (L:1141)

$$\begin{aligned}\text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] &= \sum_b \langle b | \hat{B}(t) \hat{B}^\dagger(t') (|0\rangle|0\rangle\dots\dots) \otimes (\langle 0|\langle 0|\dots\dots) | b \rangle \\ &= (\langle 0|\langle 0|\dots) \sum_b |b\rangle \langle b| \hat{B}(t) \hat{B}^\dagger(t') (|0\rangle|0\rangle\dots) \\ &= (\langle 0|\langle 0|\dots) \hat{B}(t) \hat{B}^\dagger(t') (|0\rangle|0\rangle\dots)\end{aligned}\quad (65)$$

as well as : (L:1176)

$$\text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] = (\langle 0|\langle 0|\dots) \hat{B}^\dagger(t') \hat{B}(t) (|0\rangle|0\rangle\dots) \quad (66)$$

so lets shoot them up : (L:1190)

$$\begin{aligned}\text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] &= (\langle 0|\langle 0|\dots) \sum_{k,k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} \hat{a}_k \hat{a}_{k'}^\dagger (|0\rangle|0\rangle\dots) \\ &= \sum_{k,k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} (\langle 0|\langle 0|\dots) \hat{a}_k \hat{a}_{k'}^\dagger (|0\rangle|0\rangle\dots)\end{aligned}\quad (67)$$

(L:1216)

$$\text{Tr}_B \left[ \hat{B}^\dagger(t') \hat{B}(t) \hat{\rho}_B \right] = \sum_{k',k} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} (\langle 0|\langle 0|\dots) \hat{a}_{k'}^\dagger \hat{a}_k (|0\rangle|0\rangle\dots) \quad (68)$$

with  $\hat{a}_k \hat{a}_{k'}^\dagger = \hat{a}_{k'}^\dagger \hat{a}_k + \delta_{k,k'}$  (L:1236)

$$\begin{aligned}\text{Tr}_B \left[ \hat{B}(t) \hat{B}^\dagger(t') \hat{\rho}_B \right] &= \sum_{k,k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} \delta_{k,k'} + \sum_{k,k'} g_k^* g_{k'} e^{-i(\omega_k t - \omega_{k'} t')} (\langle 0|\langle 0|\dots) \hat{a}_{k'}^\dagger \hat{a}_k (|0\rangle|0\rangle\dots) \\ &= \sum_k |g_k|^2 e^{-i\omega_k(t-t')}\end{aligned}\quad (69)$$

so it gives the things to the  $F(t)$  and we have : (L:1263)

$$\begin{aligned}F(t) &= \sum_k |g_k|^2 \int_0^t e^{-i\omega_k(t-t')} dt' \\ G(t) &= 0\end{aligned}\quad (70)$$

## 1.5 Transition to the continuum

now say (L:1278)

$$J(\omega) = \sum_l |g_l|^2 \delta(\omega - \omega_l) \quad (71)$$

redistribute the k index to the  $J(\omega)$  into the function  $F(t)$  gives (L:1288)

$$F(t) = \int_0^\infty d\omega J(\omega) \int_0^t dt' e^{-i\omega(t-t')} \quad (72)$$

now let's use  $\tau = t - t'$  so  $d\tau = -dt'$  then it gives (L:1301)

$$\int_0^t dt' = - \int_t^0 d\tau = \int_0^t d\tau \quad (73)$$

now it writes that (L:1308)

$$F(t) = \int_0^\infty d\omega J(\omega) \int_0^t d\tau e^{-i\omega\tau} \quad (74)$$

## 1.6 The Markov approximation

so targeting the good old equation (20) we know that we like to have the  $t$  to be infinite on the (74), so (L:1320)

$$\begin{aligned} \int_0^\infty d\tau e^{-i\omega\tau} &= \lim_{\eta \rightarrow 0^+} \int_0^\infty d\tau e^{-i\omega\tau - \eta\tau} \\ &= \lim_{\eta \rightarrow 0^+} \frac{1}{\eta + i\omega} \\ &= \lim_{\eta \rightarrow 0^+} \frac{\eta - i\omega}{\eta^2 + \omega^2} \\ &= \lim_{\eta \rightarrow 0^+} \frac{\eta}{\eta^2 + \omega^2} - \lim_{\eta \rightarrow 0^+} \frac{i\omega}{\eta^2 + \omega^2} \\ &= \pi\delta(\omega) - i\mathcal{P} \frac{1}{\omega} \end{aligned} \quad (75)$$

this  $\mathcal{P}$  stands for Cauchy principal part, then we have: (L:1342)

$$F = \pi \int_0^\infty d\omega J(\omega) \delta(\omega) - i\mathcal{P} \int_0^\infty d\omega J(\omega) \frac{1}{\omega} \quad (76)$$

## 1.7 Final form

(L:1352)

$$\begin{aligned}
F &= \frac{\gamma + \mathbb{i}\varepsilon}{2} \\
\gamma &\equiv 2\pi \int_0^\infty \mathbf{d}\omega J(\omega) \delta(\omega) \\
\varepsilon &\equiv -2P \int_0^\infty \mathbf{d}\omega J(\omega) \frac{1}{\omega}
\end{aligned} \tag{77}$$

with  $G = 0$  in mind we have that equation (62) evolves into (L:1360)

$$\begin{aligned}
\frac{\mathbf{d}\hat{\rho}_S(t)}{\mathbf{d}t} &= - \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \frac{\gamma - \mathbb{i}\varepsilon}{2} \\
&\quad - \left( \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \frac{\gamma + \mathbb{i}\varepsilon}{2} \\
&= -\frac{\gamma}{2} \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger + \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right) \\
&\quad + \frac{\mathbb{i}\varepsilon}{2} \left( \hat{\rho}_S(t) \hat{S}^\dagger \hat{S} - \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger - \hat{S}^\dagger \hat{S} \hat{\rho}_S(t) + \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger \right)
\end{aligned} \tag{78}$$

some how it says that  $\varepsilon = 0$  which can be achived by a good choice of denisty state choosing , for example , we can extend the lower limit of integration to  $-\infty$  noticing that : and eventually : (L:1386)

$$\frac{\mathbf{d}\hat{\rho}_S(t)}{\mathbf{d}t} = \gamma \left( \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger - \frac{1}{2} \{ \hat{S}^\dagger \hat{S}, \hat{\rho}_S(t) \} \right) \tag{79}$$

so back in the original (L:1396)

$$\hat{\rho}_S(t) = \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \hat{\rho}_S \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} \tag{80}$$

and derivationg it will be looking like (L:1406)

$$\begin{aligned}
\frac{\mathbf{d}\hat{\rho}_S(t)}{\mathbf{d}t} &= \frac{\mathbb{i}}{\hbar} \hat{H}_S \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \hat{H}_S \hat{\rho}_S \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} + \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \frac{\mathbf{d}\hat{\rho}_S}{\mathbf{d}t} \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} - \frac{\mathbb{i}}{\hbar} \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \hat{\rho}_S \hat{H}_S \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} \\
&= \frac{\mathbb{i}}{\hbar} \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \left[ \hat{H}_S, \hat{\rho}_S \right] \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} + \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \frac{\mathbf{d}\hat{\rho}_S}{\mathbf{d}t} \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar}
\end{aligned} \tag{81}$$

also , there is that : (L:1434)

$$\left( \hat{S} \hat{\rho}_S(t) \hat{S}^\dagger - \frac{1}{2} \{ \hat{S}^\dagger \hat{S}, \hat{\rho}_S(t) \} \right) = \mathbb{e}^{\mathbb{i}\hat{H}_S t/\hbar} \left( \hat{S} \hat{\rho}_S \hat{S}^\dagger - \frac{1}{2} \{ \hat{S}^\dagger \hat{S}, \hat{\rho}_S \} \right) \mathbb{e}^{-\mathbb{i}\hat{H}_S t/\hbar} \tag{82}$$

so connecting this with the explicit timed (79) and the no explicit timed (82) we have : (L:1450)

$$\gamma \left( \hat{S} \hat{\rho}_S \hat{S}^\dagger - \frac{1}{2} \{ \hat{S}^\dagger \hat{S}, \hat{\rho}_S \} \right) = \frac{\mathbb{I}}{\hbar} [\hat{H}_S, \hat{\rho}_S] + \frac{\mathbf{d}\hat{\rho}_S}{\mathbf{d}t} \quad (83)$$

and arranging it gives : (L:1462)

$$\frac{\mathbf{d}\hat{\rho}_S}{\mathbf{d}t} = -\frac{\mathbb{I}}{\hbar} [\hat{H}_S, \hat{\rho}_S] + \gamma \left( \hat{S} \hat{\rho}_S \hat{S}^\dagger - \frac{1}{2} \{ \hat{S}^\dagger \hat{S}, \hat{\rho}_S \} \right) \quad (84)$$

## 2 conclusion

so just with several  $\hat{L}$  instead of  $\hat{S}$  we can have (L:1479)

$$\frac{\mathbf{d}\hat{\rho}_S}{\mathbf{d}t} = -\frac{\mathbb{I}}{\hbar} [\hat{H}_S, \hat{\rho}_S] + \gamma \sum_j \left( \hat{L}_j \hat{\rho}_S \hat{L}_j^\dagger - \frac{1}{2} \{ \hat{L}_j^\dagger \hat{L}_j, \hat{\rho}_S \} \right) \quad (85)$$