

# N-dimensional Laplace Equation

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Consider the N-dimensional Laplace equation

$$\nabla^2 f = \sum_{n=1}^N \partial_n^2 f = 0 \quad (1)$$

where  $f : \mathbb{R}^N \rightarrow \mathbb{R}$  is some scalar field. Suppose  $f$  can be written as a product of scalar valued functions, i.e

$$f(x_i : i = 1, \dots, N) = \prod_{j=1}^N \Phi_j(x_j) \quad (2)$$

so that the derivatives become

$$\partial_n^2 f = \partial_n^2 \Phi_n(x_n) \prod_{\substack{j=1 \\ j \neq n}}^N \Phi_j(x_j) . \quad (3)$$

Hence, the Laplace equation can now be written as

$$\sum_{n=1}^N \partial_n^2 \Phi_n(x_n) \prod_{\substack{j=1 \\ j \neq n}}^N \Phi_j(x_j) = 0 . \quad (4)$$

Dividing both sides of this equation by (2) then yields,

$$\sum_{n=1}^N \frac{\partial_n^2 \Phi_n}{\Phi_n} = 0 . \quad (5)$$

This equation is satisfied iff

$$\frac{\partial_n^2 \Phi_n}{\Phi_n} = -\lambda_n \implies \partial_n^2 \Phi_n + \lambda_n \Phi_n = 0, \quad n < N \quad (6)$$

$$\frac{\partial_N^2 \Phi_N}{\Phi_N} = \sum_{k=1}^{N-1} \lambda_k \implies \partial_N^2 \Phi_N - \Phi_N \sum_{k=1}^{N-1} \lambda_k = 0 \quad (7)$$

The corresponding characteristic equations are thus

$$K_n^2 + \lambda_n = 0 \implies K_n = \pm i \sqrt{\lambda_n} \quad (8)$$

$$K_N^2 - \sum_{k=1}^{N-1} \lambda_k = 0 \implies K_N = \pm \sqrt{\sum_{k=1}^{N-1} \lambda_k} \quad (9)$$

Hence, the solutions are

$$\Phi_n(x_n) = C_{1,n} \cos(x_n \sqrt{\lambda_n}) + C_{2,n} \sin(x_n \sqrt{\lambda_n}) \quad (10)$$

$$\Phi_N(x_N) = D_{1,N} \sinh \left( x_N \sqrt{\sum_{k=1}^{N-1} \lambda_k} \right) + D_{2,N} \cosh \left( x_N \sqrt{\sum_{k=1}^{N-1} \lambda_k} \right) \quad (11)$$

We now introduce some boundary conditions

$$\begin{cases} f(x_n = 0) &= 0 \implies \Phi_n(0) = 0, \forall n \\ f(x_n = L_n) &= 0 \implies \Phi_n(L_n) = 0, \forall n < N \\ f(x_N = \pi) &= \Psi(x_n : n < N) \end{cases} \quad (12)$$

where  $\Psi : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ . The first boundary condition then yield

$$\Phi_n(0) = C_1 = 0 \quad (13)$$

$$\Phi_N(0) = D_2 = 0 \quad (14)$$

The second boundary condition together with  $\Phi_n$  then gives the following Sturm-Liouville boundary value problem

$$\Phi_n(L_n) = C_{2,n} \sin(L_n \sqrt{\lambda_n}) = 0 \quad (15)$$

which is solved by the eigenvalue

$$\sqrt{\lambda_n} = \frac{\alpha_n \pi}{L_n} \quad (16)$$

where  $\alpha_n \in \mathbb{Z}$ . Thus,

$$\Phi_n(x_n) = C_{\alpha_n} \sin\left(\frac{\alpha_n \pi x_n}{L_n}\right), \quad n < N \quad (17)$$

$$\Phi_N(X_N) = D_{\alpha_1, \dots, \alpha_{N-1}} \sinh\left(x_N \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}}\right) \quad (18)$$

To this end, we can write one linearly independent solution to the N-dimensional Laplace equation as

$$f_{\alpha_1, \dots, \alpha_{N-1}} = C_{\alpha_1, \dots, \alpha_{N-1}} \sinh\left(x_N \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}}\right) \prod_{j=1}^{N-1} \sin\left(\frac{\alpha_j \pi x_j}{L_j}\right)$$

and hence, by the superposition principle we can write the full solution as the infinite series

$$f = \sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} C_{\alpha_1, \dots, \alpha_{N-1}} \sinh\left(x_N \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}}\right) \prod_{j=1}^{N-1} \sin\left(\frac{\alpha_j \pi x_j}{L_j}\right)$$

The final boundary condition then leads to the following equation

$$\sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} C_{\alpha_1, \dots, \alpha_{N-1}} \sinh\left(\pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}}\right) \prod_{j=1}^{N-1} \sin\left(\frac{\alpha_j \pi x_j}{L_j}\right) = \Psi(x_n)$$

Without justification, we now multiply both sides of this equation by

$$\prod_{j=1}^{N-1} \sin\left(\frac{\tilde{\alpha}_j \pi x_j}{L_j}\right)$$

to give

$$\sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} C_{\alpha_1, \dots, \alpha_{N-1}} \sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right) \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) = \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right)$$

We can now integrate both sides over the region  $\{0 < x_n < L_n\}$

$$\begin{aligned} & \int_0^{L_1} \dots \int_0^{L_{N-1}} \sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} C_{\alpha_1, \dots, \alpha_{N-1}} \sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right) \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \\ &= \int_0^{L_1} \dots \int_0^{L_{N-1}} \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \end{aligned}$$

moving the integral inside the sum then gives

$$\begin{aligned} & \sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} C_{\alpha_1, \dots, \alpha_{N-1}} \sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \int_0^{L_1} \dots \int_0^{L_{N-1}} \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right) \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \\ &= \int_0^{L_1} \dots \int_0^{L_{N-1}} \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\tilde{\alpha}_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \end{aligned}$$

Notice that due to the orthogonality of the sine function, every term in the series on the LHS will vanish, except for when  $\alpha_j = \tilde{\alpha}_j$ . Thus,

$$\begin{aligned} & C_{\alpha_1, \dots, \alpha_{N-1}} \sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \int_0^{L_1} \dots \int_0^{L_{N-1}} \prod_{j=1}^{N-1} \sin^2 \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \\ &= \int_0^{L_1} \dots \int_0^{L_{N-1}} \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \end{aligned}$$

Hence, we can express the coefficient  $C_{\alpha_1, \dots, \alpha_{N-1}}$  as

$$C_{\alpha_1, \dots, \alpha_{N-1}} = \frac{\int_0^{L_1} \dots \int_0^{L_{N-1}} \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1}}{\sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \int_0^{L_1} \dots \int_0^{L_{N-1}} \prod_{j=1}^{N-1} \sin^2 \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1}}$$

Thus, we can finally write the full solution to the Laplace equation as

$$f(x_1, \dots, x_N) = \sum_{\alpha_1=1}^{\infty} \dots \sum_{\alpha_{N-1}=1}^{\infty} \frac{\int_0^{L_1} \dots \int_0^{L_{N-1}} \Psi(x_n) \prod_{j=1}^{N-1} \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1} \sinh \left( x_n \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \prod_{j=1}^{N-1} \sin \left( \frac{\alpha_j \pi x_j}{L_j} \right)}{\sinh \left( \pi \sqrt{\sum_{k=1}^{N-1} \frac{\alpha_k^2 \pi^2}{L_k^2}} \right) \int_0^{L_1} \dots \int_0^{L_{N-1}} \prod_{j=1}^{N-1} \sin^2 \left( \frac{\alpha_j \pi x_j}{L_j} \right) dx_1 \dots dx_{N-1}}$$