

Lecture 11

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2 spin - 1/2 system

$$|t\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle)$$

Probability to measure spin 1 to be $\hbar/2$ along z-axis We define spin 1 to be when the first spin is up and the second is down. Spin 2 is when the first spin is down and the second is up.

$$P_t = \langle t | \left(|\uparrow\rangle \langle \uparrow| \otimes I \right) | t \rangle = \frac{1}{2} \left(\langle \uparrow | \otimes \langle \downarrow | + \langle \downarrow | \otimes \langle \uparrow | \right) \underbrace{\left(|\uparrow\rangle \langle \uparrow| \otimes I \right) \left(|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle \right)}_{|\uparrow\rangle \otimes |\downarrow\rangle}$$

$$P_t = \frac{1}{2} \left(\langle \uparrow \uparrow | \langle \downarrow \downarrow | + \langle \downarrow \uparrow | \langle \uparrow \downarrow | \right) = \frac{1}{2}$$

Total Angular Momentum

The total angular momentum is the sum of orbital and spin angular momentum

$$\vec{J} = \vec{L} + \vec{S}$$

This is actually represented on the form:

$$\vec{J} = \vec{L} \otimes I + I \otimes \vec{S}$$

Commutation

$$\begin{aligned} [\hat{L}, \hat{S}] &= [\hat{L} \otimes I, \hat{S} \otimes I] \\ (L \otimes I) (I \otimes S) - (I \otimes S) (L \otimes I) \\ L \otimes S - L \otimes S &= 0 \end{aligned}$$

Total spin of two spin 1/2 particles

Total spin of two particles spinning up

$$\begin{aligned} \vec{S}^{\text{tot}} &= \vec{S} \otimes I + I \otimes \vec{S} \\ \vec{S}^{\text{tot}} &= S_z \otimes I + I \otimes S_z \\ S_z^{\text{tot}} |\uparrow\rangle \otimes |\uparrow\rangle &= \left(S_z \otimes I + I \otimes S_z \right) |\uparrow\rangle \otimes |\uparrow\rangle = \frac{\hbar}{2} |\uparrow\rangle \otimes |\uparrow\rangle + \frac{\hbar}{2} |\uparrow\rangle \otimes |\uparrow\rangle = \hbar |\uparrow\rangle \otimes |\uparrow\rangle \end{aligned}$$

Total spin of one particle spinning up and one spinning down

$$S_z^{\text{tot}} |\uparrow\rangle \otimes |\downarrow\rangle = \left(S_z \otimes I + I \otimes S_z \right) |\uparrow\rangle \otimes |\downarrow\rangle = \frac{\hbar}{2} |\uparrow\rangle \otimes |\downarrow\rangle + \frac{\hbar}{2} |\uparrow\rangle \otimes |\downarrow\rangle = 0$$

Naturally we see that the total spin is just the sum of the spin of each particle. Its easy then to see that

$$S_z^{\text{tot}} |\downarrow\rangle \otimes |\uparrow\rangle = 0 \quad , \quad S_z^{\text{tot}} |\downarrow\rangle \otimes |\downarrow\rangle = -\hbar |\downarrow\rangle \otimes |\downarrow\rangle$$

Eigenstate of total spin S^{tot} squared

$$\begin{aligned}\vec{S}^{\text{tot}^2} &= \vec{S}^{\text{tot}} \cdot \vec{S}^{\text{tot}} \\ \vec{S}^{\text{tot}^2} &= \left(\vec{S} \otimes I + I \otimes \vec{S} \right) \cdot \left(\vec{S} \otimes I + I \otimes \vec{S} \right) \\ \vec{S}^{\text{tot}^2} &= \vec{S}^2 \otimes I + I \otimes \vec{S}^2 + 2 \left(S_x \otimes S_x + S_y \otimes S_y + S_z \otimes S_z \right)\end{aligned}$$

We must define S_i in terms of the raising and lowering operators.

$$\begin{aligned}S_x \otimes S_x &= \frac{1}{4} \left(S_+ + S_- \right) \otimes \left(S_+ + S_- \right) \\ S_y \otimes S_y &= \frac{1}{4} \left(S_+ S_- \right) \otimes \left(S_+ + S_- \right) \\ S_x \otimes S_x + S_y \otimes S_y &= \frac{1}{2} \left(S_+ \otimes S_- + S_- \otimes S_+ \right)\end{aligned}$$

We can now derive the total:

$$S^{\text{tot}^2} = S^2 \otimes I + I \otimes S^2 + S_+ \otimes S_- + S_- \otimes S_+ + 2S_z \otimes S_z$$

Eigenstate of total spin S^{tot} squared for two particles spinning up

$$\begin{aligned}S^{\text{tot}^2} |\uparrow\rangle \otimes |\uparrow\rangle &= \left(S^2 \otimes I + I \otimes S^2 + S_+ \otimes S_- + S_- \otimes S_+ + 2S_z \otimes S_z \right) |\uparrow\rangle \otimes |\uparrow\rangle \\ S^{\text{tot}^2} |\uparrow\rangle \otimes |\uparrow\rangle &= \hbar^2 \frac{3}{4} |\uparrow\rangle \otimes |\uparrow\rangle + \hbar^2 \frac{3}{4} |\uparrow\rangle \otimes |\uparrow\rangle + \hbar^2 |\text{NULL}\rangle \otimes |\downarrow\rangle + |\text{NULL}\rangle + 2 \left(\frac{\hbar}{2} \right)^2 |\uparrow\rangle \otimes |\uparrow\rangle \\ S^{\text{tot}^2} |\uparrow\rangle \otimes |\uparrow\rangle &= \hbar^2 \left(\frac{3}{4} + \frac{3}{4} + \frac{1}{2} \right) |\uparrow\rangle \otimes |\uparrow\rangle = 2\hbar^2 |\uparrow\rangle \otimes |\uparrow\rangle\end{aligned}$$

Eigenstate of total spin S^{tot} of two particles spinning down

$$S^{\text{tot}^2} |\downarrow\rangle \otimes |\downarrow\rangle = 2\hbar^2 |\downarrow\rangle \otimes |\downarrow\rangle$$

Eigenstate of total spin S^{tot} of one particle spinning up and one spinning down

$$\begin{aligned}S^{\text{tot}^2} |\uparrow\rangle \otimes |\downarrow\rangle &= \left(S^2 \otimes I + I \otimes S^2 + S_+ \otimes S_- + S_- \otimes S_+ + 2S_z \otimes S_z \right) |\uparrow\rangle \otimes |\downarrow\rangle \\ \hbar^2 \frac{3}{4} |\uparrow\rangle \otimes |\downarrow\rangle &+ \hbar^2 \frac{3}{4} |\uparrow\rangle \otimes |\downarrow\rangle + 0 + \hbar^2 |\downarrow\rangle \otimes |\uparrow\rangle + 2 \frac{\hbar}{2} \cdot \left(\frac{\hbar}{2} \right) |\uparrow\rangle \otimes |\downarrow\rangle \\ &\hbar^2 \left(|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle \right)\end{aligned}$$

Eigenstate of total spin S^{tot} of one particle down up and one spinning up

$$\vec{S}^{\text{tot}^2} |\downarrow\rangle \otimes |\uparrow\rangle = \hbar^2 \left(|\downarrow\rangle \otimes |\uparrow\rangle + |\uparrow\rangle \otimes |\downarrow\rangle \right)$$

Conclusion: $|\uparrow\rangle \otimes |\downarrow\rangle$ and $|\downarrow\rangle \otimes |\uparrow\rangle$ are not eigenstates but $|\uparrow\rangle \otimes |\uparrow\rangle$ and $|\downarrow\rangle \otimes |\downarrow\rangle$ are.

Creating Linear Combinations

$$S^{\text{tot}^2} \left(|\uparrow\rangle \otimes |\downarrow\rangle - |\downarrow\rangle \otimes |\uparrow\rangle \right) = 0$$

This is an eigenstate. Also known as the singlet state

$$S^{\text{tot}^2} \left(|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle \right) = 2\hbar^2 \left(|\uparrow\rangle \otimes |\downarrow\rangle + |\downarrow\rangle \otimes |\uparrow\rangle \right)$$

This is one of the three parts of the triplet state. The other being $|\uparrow\rangle \otimes |\uparrow\rangle$ and $|\downarrow\rangle \otimes |\downarrow\rangle$

General Combinations of ang.mom. states

$$|sm\rangle = \sum_{m_1 m_2} C_{m_1 m_2}^{S_1 S_2 S} |s_1 m_1\rangle \otimes |s_2 m_2\rangle$$

where C is the Clebsch-Gordan coefficient.

$$|s_1 m_1\rangle \otimes |s_2 m_2\rangle = \sum_{m} C_{m_1 m_2}^{S_1 S_2 S} |sm\rangle$$

a spin - 3 / 2 with a spin-1 to get a total spin 3/2, m = 3/2

$$|3/2, +3/2\rangle$$

Reading from the table we see the column of 3/2 and +3/2, having values 3/5 square root of 3/5 for $m_1 = +3/2$ and $s_1 = 0$. We get the negative of the square root of 2/5 for $m_1 = +1/2$ and $s_1 = 1$.

$$\left| \underbrace{3/2}_s, + \underbrace{3/2}_m \right\rangle = \sqrt{\frac{3}{5}} \left| \underbrace{3/2}_{s_1}, + \underbrace{3/2}_{m_1} \right\rangle \otimes \left| \underbrace{1}_{s_2}, \underbrace{0}_{m_2} \right\rangle - \sqrt{\frac{2}{5}} \left| \underbrace{3/2}_{s_1}, + \underbrace{1/2}_{m_1} \right\rangle \otimes \left| \underbrace{1}_{s_2}, \underbrace{+1}_{m_2} \right\rangle$$