# FYS5419: Project 1

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# **Imports**

```
In [ ]: ### Regular imports ###
        import warnings
        import numpy as np
        import sympy as sp
        from time import time
        from numba import njit
        from tabulate import tabulate
        import matplotlib.pyplot as plt
        from dataclasses import dataclass
        from numpy import float64, complex128, ndarray
        from typing import Annotated, Literal, Callable
        ### Customizations ###
        np.set printoptions(precision=3, sign=' ')
        plt.rcParams.update({'font.size': 12,
                              'font.family': 'serif',
                              'figure.figsize': (10, 6),
                              'mathtext.fontset': 'cm',
                              'axes.prop cycle': plt.cycler(color=["#3d0dce",
                                                                    "#fflele",
                                                                    "#61ffa6",
                                                                    "#86db0f"
                                                                    "#f711ff"
                                                                    "#64B5CD"]),
                              'lines.linewidth': 3,
                              'legend.fontsize': 10,
        tensor prod = np.kron
        ### Qiskit imports ###
        from qiskit import QuantumCircuit
        from qiskit aer import AerSimulator
        from giskit.visualization import plot histogram
```

# **Types**

```
In []: ### Matplotlib Figure type ###
Figure = Annotated[plt.Figure, "Figure"]

### Indexing types ###
Indexing_2qbit = Annotated[Literal[0, 1,], "Valid 2 qubit indexing"]
Indexing_3qbit = Annotated[Literal[0, 1, 2], "Valid 3 qubit indexing"]

### Vectors ###
Array2D = Annotated[ndarray[2, float], "2D array"]
Array2D_c = Annotated[ndarray[2, complex], "2D complex array"]
```

```
Array4D = Annotated[ndarray[4, float], "4D array"]
Array4D_c = Annotated[ndarray[4, complex], "4D complex array"]

### Matrices ###
Array2x2 = Annotated[ndarray[(ndarray[2, float]), "2x2 matrix"]

Array2x2_c = Annotated[ndarray[(ndarray[2, complex], ndarray[2, complex])], "2x2 complex matrix"]

Array4x4 = Annotated[ndarray[(ndarray[4, float], ndarray[4, float])], "4x4 matrix"]

Array4x4_c = Annotated[ndarray[(ndarray[4, complex], ndarray[4, complex])], "4x4 complex matrix"]
```

# a)

```
In [ ]: @njit
        def qubit_basis() -> tuple[Array2D, Array2D]:
             Creates the qubits standard qubit basis: |0) and |1).
             Returns
             _ _ _ _ _ _
                 q0: Array2D
                     |0\rangle = [1, 0]
                 q1: Array2D
                    |1\rangle = [0, 1]
             q0: Array2D = np.array([1, 0])
             q1: Array2D = np.array([0, 1])
             return q0, q1
        @njit
        def pauli() -> tuple[Array2x2 c, Array2x2 c, Array2x2 c]:
             Creates the Pauli matrices \sigma x, \sigma y, and \sigma z.
             Returns
                 σ x: Array2x2 c
                     Pauli X
                 σ_y: Array2x2_c
                     Pauli Y
                 σ z: Array2x2 c
                     Pauli Z
             \sigma_x: Array2x2_c = np.array([[0, 1], [1, 0]], dtype=complex128)
             \sigma_y: Array2x2_c = np.array([[0, -1j], [1j, 0]], dtype=complex128)
```

```
\sigma_z: Array2x2_c = np.array([[1, 0], [0, -1]], dtype=complex128)
    return σ x, σ y, σ z
def hadamard() -> Array2x2:
    Creates the Hadamard gate.
    Returns
    _ _ _ _ _
        H: Array2x2
             Hadamard gate
    H: Array2x2 = 1/np.sqrt(2) * np.array([[1, 1], [1, -1]])
    return H
def phase() -> Array2x2 c:
    Creates the phase gate.
    Returns
    _ _ _ _ _ _ _
         S: Array2x2 c
              Phase gate
    S = np.array([[1, 0], [0, 1j]])
    return S
@njit
def cnot() -> Array4x4 c:
    Creates the CNOT gate.
    Returns
         CNOT: Array4x4 c
              CNOT gate
    CNOT: Array4x4 c = np.array([[1, 0, 0, 0],
                                       [0, 1, 0, 0],
                                       [0, 0, 0, 1],
                                       [0, 0, 1, 0]], dtype=complex128)
    return CNOT
def create bell states() -> tuple[Array4D, Array4D, Array4D]:
    Creates the Bell states |\Phi^+\rangle, |\Phi^-\rangle, |\Psi^+\rangle, and |\Psi^-\rangle.
    Returns
         Φ 00: Array4x4
              |\Phi^+\rangle = (|00\rangle + |11\rangle)/\sqrt{2}
         Φ_10: Array4x4
             |\Phi^-\rangle = (|\Theta0\rangle - |11\rangle)/\sqrt{2}
         ¥ 01: Array4x4
             |\Psi^+\rangle = (|01\rangle + |10\rangle)/\sqrt{2}
         ¥ 11: Array4x4
```

```
|\Psi^-\rangle = (|01\rangle - |10\rangle)/\sqrt{2}
    q0, q1 = qubit basis()
           Array2x2 = hadamard()
    CNOT: Array4D c = cnot()
    q0 H: Array2D = H @ q0
    q1 H: Array2D = H @ q1
    \Phi_{00}: Array4D = CNOT @ tensor_prod(q0_H, q0) # |\Phi_{00}| = (|00| + |11|)/\sqrt{2}
    \Phi 10: Array4D = CNOT @ tensor prod(q1 H, q0) # |\Phi^-\rangle = (|\theta\theta\rangle - |11\rangle)/\sqrt{2}
    \Psi 01: Array4D = CNOT @ tensor prod(q0 H, q1) # |\Psi^+\rangle = (|\theta 1\rangle + |1\theta\rangle)/\sqrt{2}
    \Psi 11: Array4D = CNOT @ tensor prod(q1 H, q1) # |\Psi^-\rangle = (|\theta 1\rangle - |1\theta\rangle)/\sqrt{2}
    return Φ 00, Φ 10, Ψ 01, Ψ 11
def create system vectors(n qubits: int) -> list[ndarray[float64]]:
    Create the system vectors for a given number of qubits.
    Parameters
     _ _ _ _ _ _ _ _ _ _
    n qubits : int
         The number of qubits in the system.
    Returns
    _ _ _ _ _ _
    system vectors : list[ndarray[int]]
         A list of system vectors, where each vector represents a possible st
    q0, q1 = qubit basis()
    system vectors = []
    for i in range(2**n qubits):
         binary str = f'{i:0{n qubits}b}'
         state = np.array([1])
         for bit in binary str:
             state = tensor_prod(state, q0 if bit == '0' else q1)
         system vectors.append(state)
    return system vectors
def measure qubit(qubit: Indexing 2qbit, bell state: Array4D) -> Literal[0,
    Makes a measurement on a specified cubit in a Bell state
    Parameters
         qubit: int
             Which cubit to measure. In a two-qubit state one could pass eith
         bell state: Array4D
             The state to measure
    Returns
```

```
state: Literal[0, 1]
            What state the qubit is in. Either 0 or 1
   if qubit not in [0, 1]:
        raise ValueError(f"Invalid qubit index. Must be either 0 or 1, not {
   n = len(bell state)
   n \text{ qubits} = int(np.log2(n))
   measurement probabilities = {'0': 0, '1': 0}
   for i in range(n):
        binary str = f'\{i:0\{n \text{ qubits}\}b\}'[\text{qubit}] \# Count in binary, 0-padded
        prob = bell state[i]
        measurement probabilities[binary str] += np.abs(prob)**2
   # Weighted random choice depending on the state
    result = np.random.choice([0, 1], p=[*measurement probabilities.values()
    return result
def measure all qubits(bell state: Array4D) -> Literal['00', '01', '10', '11
   Makes a measurement on all qubits in a Bell state
   Parameters
        bell state: Array4D
           The state to measure
   Returns
        state: Literal['00', '01', '10', '11']
            What state the qubits are in. For a two-qubit system, this would
   measurement probabilities = {'00': 0, '01': 0, '10': 0, '11': 0}
    for i, bin str in enumerate(measurement probabilities.keys()):
        prob = bell state[i]
        measurement probabilities[bin str] += np.abs(prob)**2
   # Weighted random choice depending on the state
    result = np.random.choice(a=[*measurement probabilities.keys()],
                              p=[*measurement probabilities.values()])
    return result
```

# **Playing with Qubits**

- The following explores how the qubit states can be combined to form composite states.
- I also explore how their connection to the bit string representation of the state.
- There seems to be a natural connection between the bit string representation and the order of the qubits in the tensor product. This is used to define the order of the qubits in the tensor product.

```
+----+
          Product
                   Expected
                              Calculated
  Binary |
     00 I
          |0⟩⊗|0⟩
                 | [ 1
                     0 0
                          01 | [1
                                 0
                                   0
                                      01 I
                0 ] |
                     1
                        0
                          0] | [0
          |0⟩⊗|1⟩
                | [0 0 1 0] | [0 0 1 0] |
     10 |
          |1⟩⊗|0⟩
     11 \mid |1\rangle \otimes |1\rangle
```

```
In [ ]: # Checking the states being as expected
          n \text{ qubits} = 3
          v1, v2, v3, v4, v5, v6, v7, v8 = create system vectors(n qubits)
           bin strs = [f'{i:0{n qubits}b}' for i in range(2**n qubits)]
          q0, q1 = qubit basis()
           table = tabulate([[bin_strs[0], |0\rangle \otimes |0\rangle \otimes |0\rangle, tensor_prod(tensor_prod(q0, c
                                  [bin_strs[1], |0\rangle \otimes |0\rangle \otimes |1\rangle, tensor_prod(tensor_prod(q0, c
                                  [bin_strs[2], |0\rangle \otimes |1\rangle \otimes |0\rangle, tensor_prod(tensor_prod(q0, c
                                  [bin\_strs[3], '|0) \otimes |1) \otimes |1)', tensor_prod(tensor_prod(q0, c
                                  [bin strs[4], |1\rangle \otimes |0\rangle \otimes |0\rangle, tensor prod(tensor prod(q1, c
                                  [bin_strs[5], |1\rangle \otimes |0\rangle \otimes |1\rangle, tensor_prod(tensor_prod(q1, c
                                  [bin_strs[6], '|1\rangle \otimes |1\rangle \otimes |0\rangle', tensor_prod(tensor_prod(q1, c
                                  [bin strs[7], |1\rangle \otimes |1\rangle \otimes |1\rangle, tensor prod(tensor prod(q1, c
                                  headers = ['Binary', 'Product', 'Expected', 'Calculated']
                                  tablefmt = 'outline',
                                  colalign = ['right', 'center', 'center']
           print(table)
```

```
Binary | Product |
                                          Expected
                                                                          Calculated
000 \mid |0\rangle \otimes |0\rangle \otimes |0\rangle \mid [1 \ 0 \ 0 \ 0 \ 0 \ 0 \ 0] \mid [1 \ 0 \ 0
0] |
        001 \mid |0\rangle \otimes |0\rangle \otimes |1\rangle \mid [0 \ 1 \ 0 \ 0 \ 0 \ 0
                                                            0] | [ 0 1
0] |
        010 \mid |0\rangle \otimes |1\rangle \otimes |0\rangle \mid [0 \ 0 \ 1 \ 0 \ 0 \ 0
                                                           0] | [ 0 0 1 0 0
0] [
        011 \mid |0\rangle \otimes |1\rangle \otimes |1\rangle \mid [0 \ 0 \ 0 \ 1 \ 0 \ 0]
                                                            0] | [ 0 0
0] [
        100 \mid |1\rangle \otimes |0\rangle \otimes |0\rangle \mid [0 \ 0 \ 0 \ 0 \ 1 \ 0 \ 0
                                                           0] | [0 0 0 0
                                                                                  1
0] |
        101 \mid |1\rangle \otimes |0\rangle \otimes |1\rangle \mid [0 0 0 0 0
                                                0
                                                    1
                                                       0
                                                           0] | [ 0 0
                                                                           0
0] |
        110 \mid |1\rangle \otimes |1\rangle \otimes |0\rangle \mid [0 \ 0 \ 0 \ 0]
                                                0 0
                                                            0] | [ 0 0
0] [
        111 \mid |1\rangle \otimes |1\rangle \otimes |1\rangle \mid [0 0 0 0 0 0 0 1] \mid [0 0 0 0 0 0]
1] |
----+
```

# **Exploring Pauli Matrices**

```
In [ ]: q0, q1 = qubit basis()
          \sigma_x, \sigma_y, \sigma_z = pauli()
          q0 x: Array2D c = \sigma x @ q0
          q0 y: Array2D c = \sigma y @ q0
          q0 z: Array2D c = \sigma z @ q0
          q1 x: Array2D c = \sigma x @ q1
          q1_y: Array2D_c = \sigma_y @ q1
          q1 z: Array2D c = \sigma z @ q1
           print(f'\sigma_x|0) = \{q0_x\}')
           print(f'\sigma y|0) = \{q0 y\}')
          print(f'\sigma z|0) = \{q0 z\}')
          print()
          print(f'\sigma x|1) = \{q1 x\}')
          print(f'\sigma y|1) = \{q1 y\}')
          print(f'\sigma_z|1) = \{q1_z\}')
         \sigma \times |0\rangle = [0.+0.j \ 1.+0.j]
         \sigma y | 0 \rangle = [0.+0.j 0.+1.j]
         \sigma z | 0 \rangle = [1.+0.j 0.+0.j]
         \sigma_x|1\rangle = [1.+0.j 0.+0.j]
         \sigma y|1\rangle = [0.-1.j 0.+0.j]
         \sigma_z|1\rangle = [0.+0.j -1.+0.j]
```

# **Exploring Gates**

```
In []: H: Array2x2 = hadamard()
    S: Array2x2_c = phase()

    q0_H = H @ q0
    q0_S = S @ q0

    q1_H = H @ q1
    q1_S = S @ q1

    print(f'H|0) = {q0_H}')
    print(f'S|0) = {q0_S}')
    print()
    print(f'H|1) = {q1_H}')
    print(f'S|1) = {q1_S}')

H|0) = [ 0.707    0.707]
    S|0) = [ 1.+0.j    0.+0.j]

H|1) = [ 0.707 -0.707]
    S|1) = [ 0.+0.j    0.+1.j]
```

### **Bell States**

Example of a Bell State and how to create it using a circuit. In this case, the Bell State  $|\Phi^+\rangle$  is created using a Hadamard gate and a CNOT gate:

$$|\Phi^+
angle = rac{1}{\sqrt{2}} \Big(|00
angle + |11
angle \Big)$$

image.png

### **Creation by Applying Gates**

#### **Direct Creation**

```
In [ ]: q 00, q 01, q 10, q 11 = create system vectors(2)
          \Phi 00 = 1/np.sqrt(2) * (q 00 + q 11)
          \Phi 10 = 1/np.sqrt(2) * (q_00 - q_11)
          \Psi 01 = 1/np.sqrt(2) * (q 01 + q 10)
          \Psi_11 = 1/np.sqrt(2) * (q_01 - q_10)
          print(f'|\Phi^+) = \{\Phi \ 00\}'\}
          print(f'|\Phi^-) = \{\Phi \ 10\}')
          print(f'|\Psi^+) = \{\Psi \ 01\}')
          print(f'|\Psi^-) = \{\Psi \ 11\}'\}
        |\Phi^{+}\rangle = [0.707 \ 0. \ 0.
                                           0.707]
         |\Phi^-\rangle = [0.707 0.
                                   0. -0.707]
        |\Psi^{+}\rangle = [0. 0.707 0.707 0.]
        |\Psi^{-}\rangle = [0. 0.707 - 0.707 0.
                                                  ]
```

### Acting on the Bell State with Gates

```
In []: H = hadamard()
          CNOT = cnot()
          I 2 = np.eye(2)
          # Applying the Hadamard and CNOT gate to the first Bell state
           \Phi 00 H = tensor prod(H, I 2) @ \Phi 00
           \Phi 00 H CNOT = CNOT @ \Phi 00 H
          print(f'|\Phi^+) =
                                        {Φ 00}')
          print(f'(H \otimes I) | \Phi^+) = \{\Phi \cup \emptyset \mid H\}'\}
          print(f'CNOT(H\otimes I)|\Phi^+) = \{\Phi \ \Theta\Theta \ H \ CNOT\}'\}
                              [ 0.707 0.
         |\Phi^+\rangle =
                                                   0.
                                                             0.7071
         (H\otimes I)|\Phi^{+}\rangle = [0.5 \ 0.5 \ 0.5 \ -0.5]
         CNOT(H \otimes I) | \Phi^+ \rangle = [0.5+0.j 0.5+0.j -0.5+0.j 0.5+0.j]
```

### **Testing by Measuring Bell States**

```
In []: # Making a measurment
q1_results = {0: 0, 1: 0}
q2_results = {0: 0, 1: 0}
n_lambdas = 10_000
for i in range(n_lambdas):
    q1_res = measure_qubit(0, Φ_00)
    q1_results[q1_res] += 1
    q2_res = measure_qubit(1, Φ_00)
    q2_results[q2_res] += 1
```

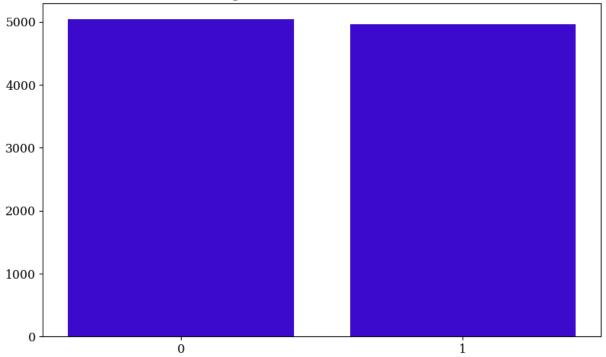
```
In [12]: print(f'Odds of measuring qubit 0 as 0:, {q1_results[0]/n_lambdas: .2%}')
    print(f'Odds of measuring qubit 0 as 1:, {q1_results[1]/n_lambdas: .2%}')

plt.bar(q1_results.keys(), q1_results.values())
    plt.xticks([0, 1])
    plt.title("Number of times qubit 0 has been measured as either 0 or 1")
```

```
plt.savefig('figs/a_q0_measurement.pdf')
plt.show()
```

Odds of measuring qubit 0 as 0:, 50.41% Odds of measuring qubit 0 as 1:, 49.59%

#### Number of times qubit 0 has been measured as either 0 or 1

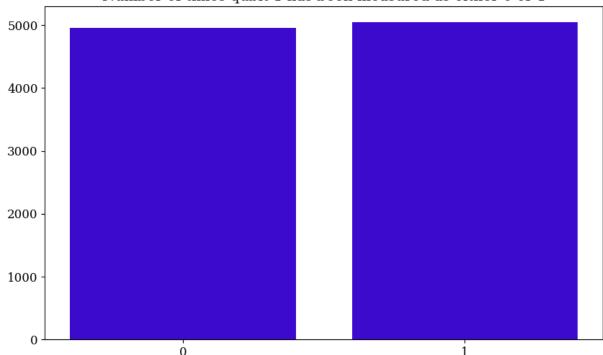


```
In []: print(f'Odds of measuring qubit 0 as 0:, {q2_results[0]/n_lambdas: .2%}')
    print(f'Odds of measuring qubit 0 as 1:, {q2_results[1]/n_lambdas: .2%}')

    plt.bar(q2_results.keys(), q2_results.values())
    plt.xticks([0, 1])
    plt.title("Number of times qubit 1 has been measured as either 0 or 1")
    plt.savefig('figs/a_q1_measurement.pdf')
    plt.show()
```

Odds of measuring qubit 0 as 0:, 49.57% Odds of measuring qubit 0 as 1:, 50.43%

#### Number of times qubit 1 has been measured as either 0 or 1

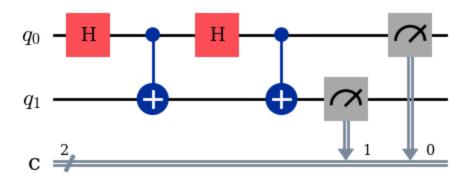


# Circuit and Measurements in Qiskit

```
In []: qc = QuantumCircuit(2, 2)
    qc.h(0)
    qc.cx(0, 1)
    qc.cx(0, 1)

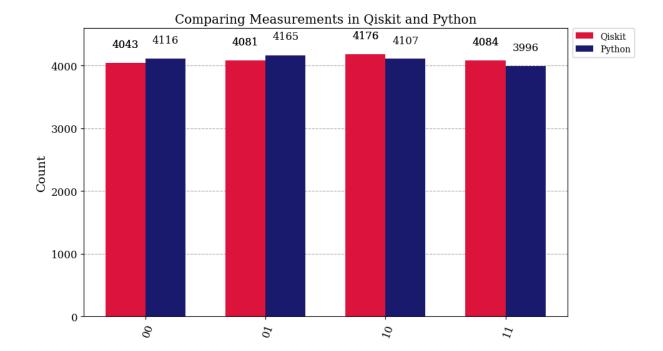
# Measureing the qubits
    qc.measure(1, 1)
    qc.measure(0, 0)
    qc.draw('mpl')
```

Out[]:



```
In [ ]: simulator = AerSimulator()
    n_lambdas = 2**14
    results_qk = simulator.run(qc, shots=n_lambdas).result().get_counts(qc)
```

```
results qk = dict(sorted(results qk.items(), key=lambda x: x[0])) # Sort the
 q00, q01, q10, q11 = results <math>qk.values()
 print('CNOT (H|\Phi^+)) (Qiskit):')
 print('----')
 print(f'Odds of measuring both gubits as 0:
                                                           {q00/n lambdas: .2
 print(f'Odds of measuring qubit 0 as 0, and qubit 1 as 1: {q10/n lambdas: .2
 print(f'Odds of measuring qubit 0 as 1, and qubit 1 as 0: {q10/n lambdas: .2
                                                           {q11/n_lambdas: .2
 print(f'Odds of measuring both gubits as 1
 # Comparing with python implementation
 H = hadamard()
 CNOT = cnot()
 Φ 00 = create bell states()[0]
 state = CNOT @ tensor prod(H, I 2) @ \Phi 00
 results py = \{f'\{i:02b\}': 0 \text{ for } i \text{ in } range(4)\}
 for i in range(n lambdas):
     res = measure all qubits(state)
     results py[res] += 1
 q00, q01, q10, q11 = results py.values()
 print()
 print('CNOT (H|\Phi^+)) (Python):')
 print('----')
 print(f'Odds of measuring both qubits as 0:
                                                           {q00/n lambdas: .2
 print(f'Odds of measuring qubit 0 as 0, and qubit 1 as 1: {q01/n lambdas: .2
 print(f'Odds of measuring qubit 0 as 1, and qubit 1 as 0: {q10/n lambdas: .2
 print(f'Odds of measuring both qubits as 1:
                                                         {q11/n lambdas: .2
 %matplotlib inline
 fig = plot histogram([results qk, results py], title="Comparing Measurements")
                                          legend=['Qiskit', 'Python'],
                                          color=['crimson', 'midnightblue'] )
 fig.savefig('figs/a giskit vs python.pdf')
 fig.savefig('selected results/a qiskit vs python.pdf')
CNOT (H|\Phi^+)) (Qiskit):
Odds of measuring both gubits as 0:
                                                   24.68%
Odds of measuring qubit 0 as 0, and qubit 1 as 1: 25.49%
Odds of measuring qubit 0 as 1, and qubit 1 as 0: 25.49%
Odds of measuring both gubits as 1
                                                   24.93%
CNOT (H|\Phi^+\rangle) (Python):
-----
Odds of measuring both qubits as 0:
                                                   25.12%
Odds of measuring qubit 0 as 0, and qubit 1 as 1: 25.42\%
Odds of measuring qubit 0 as 1, and qubit 1 as 0: 25.07%
Odds of measuring both qubits as 1:
                                                   24.39%
```



### Conclusion

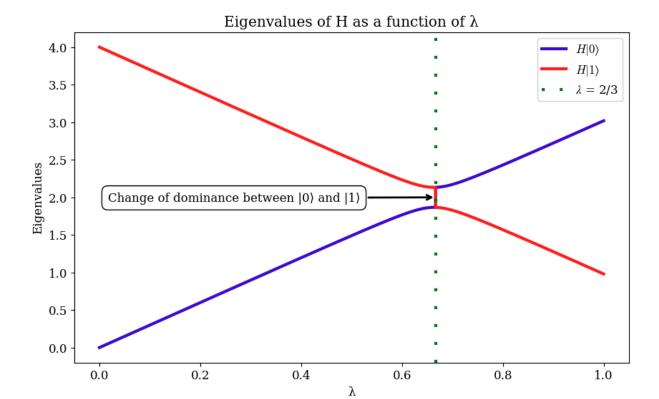
- As we can see, the qubit states can be combined to form composite states using the tensor product of  $|0\rangle$  and  $|1\rangle$ .
- We looked at the bell state  $|\Phi^+\rangle$ , in which both qubits are either in the state  $|0\rangle$ , represented by  $|00\rangle$ , or  $|1\rangle$  represented by  $|11\rangle$ .
- Applying a Hadamard and CNOT gate to  $|\Phi^+\rangle$ , we get the expected results of approximately a 25% chance of measuring either  $|00\rangle$ ,  $|01\rangle$ ,  $|10\rangle$  or  $|11\rangle$ .

# b)

```
E1 = 0
E2 = 4
E = (E1 + E2) / 2
\Omega = (E1 - E2) / 2
\sigma x, \sigma y, \sigma z = pauli()
I_2 = np.eye(2, dtype=complex128)
H0 = E*I_2 + \Omega*\sigma_z
V11 = 3
V22 = -V11
V12 = V21 = 0.2
c = (V11 + V22) / 2
\omega z = (V11 - V22) / 2
\omega \times = V12
HI = c*I 2 + \omega z*\sigma z + \omega x*\sigma x
H = H0 + \lambda *HI
return H
```

### **Finding Eigenvalues**

```
In [ ]: n lambdas = 10001
        lambdas = np.linspace(0, 1, n lambdas)
        energy eigvals = [np.linalg.eigvals(Hamiltonian 1qbit(<math>\lambda)).real for \lambda in lamb
         plt.plot(lambdas, energy eigvals)
        x = 2/3
        y = (energy eigvals[int(2/3*n lambdas)][0] + energy eigvals[int(2/3*n lambdas)][0]
         plt.annotate(text=r"Change of dominance between |0) and |1)",
                      xy=(x, y),
                      xytext=(x-.65, y-.05),
                      arrowprops=dict(arrowstyle='->', lw=2),
                      bbox=dict(facecolor='white',
                      edgecolor='black',
                      boxstyle='round,pad=0.5')),
         plt.xlabel('\lambda')
         plt.axvline(x=2/3, color='green', linestyle=(0, (1, 5)))
         plt.ylabel('Eigenvalues')
         plt.title('Eigenvalues of H as a function of \lambda')
         plt.legend([r'$H|0)$', r'$H|1)$', r'$\lambda$ = 2/3'], fontsize='medium')
         plt.savefig('figs/b eigenvalues.pdf')
         plt.savefig('selected_results/b eigenvalues.pdf')
         plt.show()
```



### Conclusion

• As the interacting term  $\lambda$  increases, naturally, the  $|1\rangle$  component of the eigenvectors begins to dominate the eigenvalues as it represents the interaction terms.

# c)

```
X = np.array([[0, 1], [1, 0]], dtype=complex128)
    return np.cos(\theta*0.5) * I 2-1j * np.sin(\theta*0.5) * X
@njit
def Ry(operation := Array2x2 c:
    Rotation around the y-axisChange of groundstate bewteen energy eigenstat
    Parameters
    _ _ _ _ _ _ _ _ _
        φ: float
            The angle (radians) to rotate by
    Returns
    _ _ _ _ _ _
        Ry: Array2x2 c
            The complex 2x2 rotation matrix
    I 2 = np.eye(2)
    Y = np.array([[0, -1j], [1j, 0]])
    return np.cos(\phi*0.5) * I 2-1j * np.sin(\phi*0.5) * Y
@njit
def Energy lqbit(\theta: float, \phi: float, \lambda: float) -> float:
    Calculates the energy eigenvalues of the Hamiltonian for a 1-qubit syste
    Parameters
    _ _ _ _ _ _ _ _ _ _
        θ: float
            The angle (radians) to rotate by around the x-axis
        φ: float
            The angle (radians) to rotate by around the y-axis
        λ: float
            The interaction strength
    Returns
    _ _ _ _ _ _
        E: float
            The energy eigenvalue (real)
    q0, q1 = qubit basis()
    q0: Array2D c = q0.astype(complex128) # Promoted to complex for njit
    basis = q0 \# |0\rangle: Will be rotated anyways so no need to use |1\rangle
    rotated basis = Rx(\theta) @ Ry(\phi) @ basis
    E = rotated_basis.conj().T @ Hamiltonian_lqbit(λ) @ rotated_basis
    assert abs(E.imag) < le-14, f'Energy is complex. Something went wrong: E</pre>
    return E. real
@njit
def VQE lqbit(n iterations: int, n: float, \lambda: float = 0) -> float:
    Variational Quantum Eigensolver using Gradient Descent to find the minim
```

```
Parameters
         N iterations: int
             Number of iterations
         η: float
             Learning rate
         λ: float
             Interaction strength of the Hamiltonian
    Returns
         Energy(\theta, \phi, \lambda): float
             The lowest energy eigenvalue found in n iterations (real)
    \pi = np.pi
    \theta = 2 \pi n \operatorname{random.rand}()
    \phi = 2*\pi*np.random.rand()
    for _ in range(n_iterations):
         \Delta E_{\Delta \theta} = (Energy_1qbit(\theta+\pi/2, \phi, \lambda) - Energy_1qbit(\theta-\pi/2, \phi, \lambda)) / 2
         \Delta E_\Delta \phi = (Energy\_1qbit(\theta, \phi+\pi/2, \lambda) - Energy\_1qbit(\theta, \phi-\pi/2, \lambda)) / 2
         \theta = \eta * \Delta E \Delta \theta
         φ -= η * ΔΕ Δφ
    E = Energy lqbit(θ, φ, λ)
    return E
@dataclass
class Benchmark Results():
    Dataclass for the benchmark results
    Parameters
     _____
         eig vals vge: ndarray
             The energy eigenvalues found by the VQE algorithm
         eig vals numpy: ndarray
             The energy eigenvalues found by numpy
         lambdas: ndarray
             Interaction strengths
         alg name: str
             The name of the VQE algorithm
         switch point: float (optional) = None
             The interaction strength where the groundstate changes
    eig vals vqe: ndarray
    eig vals numpy: ndarray
    lambdas: ndarray
    alg name: str
    switch point: float = None
def benchmark VQE(VQE alg: Callable, Hamiltonian: Callable, n iterations: ir
    Benchmarks a VQE algorithm for a given number of iterations and learning
    Parameters
```

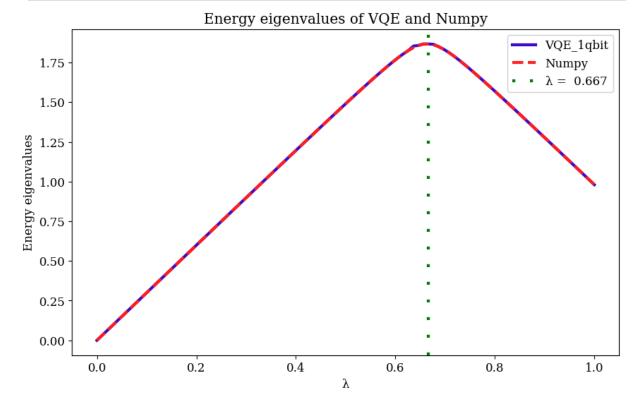
```
VQE alg: Callable
            The VQE algorithm to benchmark
        Hamiltonian: Callable
            The Hamiltonian to pass to numpy
        n iterations: int
            Number of iterations
        η: float
            Learning rate
        lambdas: ndarray
            Interaction strengths to benchmark
    Returns
        result: Benchmark Results
            The benchmark results stored in a dataclass
    start = time()
    eig_vals_vqe = np.array([VQE_alg(n_iterations, \eta, \lambda) for \lambda in lambdas])
    vqe time cold = time() - start
    start = time()
    eig_vals_vqe = np.array([VQE_alg(n_iterations, \eta, \lambda) for \lambda in lambdas])
    vge time warm = time() - start
    start = time()
    eig vals numpy = [np.min(np.linalg.eigvals(Hamiltonian(<math>\lambda))).real for \lambda i
    numpy time = time() - start
    name = VQE alg. name
    table = [["Method", "Time (s)", "Times Slower than Numpy"],
             [f"{name} (cold)", f"{vqe_time_cold:.2g}", f"{vqe_time_cold/num
             [f"{name} (warm)", f"{vqe time warm:.2g}", f"{vqe time warm/num
             ["Numpy", f"{numpy time:.2g}", "X"]
    print(tabulate(table, headers="firstrow", tablefmt="grid"))
    result = Benchmark Results(eig vals vqe, eig vals numpy, lambdas, name)
    return result
def plot relative error(result: Benchmark Results) -> Figure:
    Plots a comparison between the VQE algorithm and numpy
    Parameters
        result: Benchmark Results
            The benchmark results stored in a named tuple
    eig vals vqe = result.eig vals vqe
    eig vals numpy = result.eig vals numpy
    lambdas = result.lambdas
    switch_point = result.switch point
```

```
alg name = result.alg name
    warnings.filterwarnings("ignore", category=RuntimeWarning) # Ignore Runt
    relative error = (eig vals vqe - eig vals numpy) / eig vals numpy
    warnings.filterwarnings("default", category=RuntimeWarning) # Turn warni
    plt.plot(lambdas, relative error, label=f'Relative error between {alg ne
    plt.axvline(x=switch_point, color='green', linestyle=(0, (1, 5)), label=
    width = 1/15
    plt.axvspan(switch point - width, switch point + width , color='red', al
    zeros = np.where(relative error == 0)[0]
    if len(zeros) > 0:
        plt.scatter(lambdas[zeros], relative_error[zeros], color='black', la
    plt.xlabel('λ')
    plt.ylabel('Relative error')
    percentage formatter = lambda x, : f'\{x*100:1g\}\%'
    plt.gca().yaxis.set major formatter(plt.FuncFormatter(percentage formatt
    plt.legend()
    return plt.gcf()
def plot eigvals(result: Benchmark Results) -> Figure:
    Plots a comparison between the VQE algorithm and numpy
    Parameters
        result: Benchmark Results
            The benchmark results stored in a named tuple
    eig_vals_vqe = result.eig_vals_vqe
    eig vals numpy = result.eig vals numpy
    lambdas = result.lambdas
    switch point = result.switch point
    alg name = result.alg name
    plt.plot(lambdas, eig_vals_vqe, label=f'{alg_name}')
    plt.plot(lambdas, eig vals numpy, label='Numpy', linestyle='--')
    plt.axvline(x=switch point, color='green', linestyle=(0, (1, 5)), label=
    plt.xlabel('λ')
    plt.ylabel('Energy eigenvalues')
    plt.legend(fontsize = 'medium')
    return plt.gcf()
```

# Comparing VQE and Exact Eigenvalues

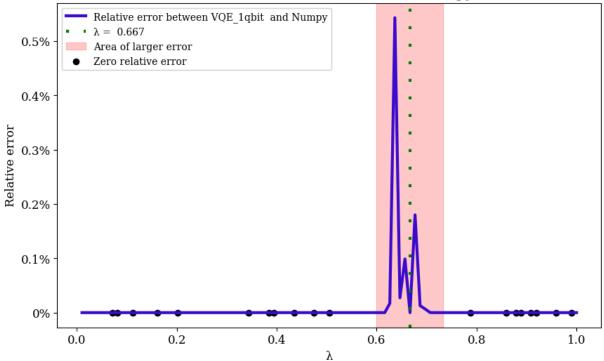
```
In []: n iterations = 100
      \eta = 1/2
      n lambdas = 100
      lambdas = np.linspace(0, 1, n lambdas)
      result_1qbit: Benchmark_Results = benchmark_VQE(VQE_1qbit, Hamiltonian_1qbit
      switch point = 2/3
      result lqbit.switch point = switch point
     +----+
                       Time (s) | Times Slower than Numpy
     +----+
      VQE 1qbit (cold) |
                        2.2
                              | 2316.11
                             | 57.98
      VQE 1qbit (warm) |
                        0.054
     Numpy
                        0.00094 | X
```

```
In [20]: fig = plot_eigvals(result_1qbit)
   plt.title('Energy eigenvalues of VQE and Numpy')
   plt.savefig('figs/c_vqe_vs_numpy.pdf')
   plt.show()
```



```
In [ ]: fig = plot_relative_error(result_1qbit)
    plt.title('Relative error between VQE and Numpy')
    plt.savefig('figs/c_relative_error.pdf')
    plt.show()
```

Relative error between VQE and Numpy



### Conclusion

- As we can see, the VQE method is able to approximate the eigenvalues of the Hamiltonian matrix, with a relative error of  $\approx 10^{-6}\%$
- With 128 iterations, the VQE was  $\approx 2000$  times slower on a cold start than the exact eigenvalue method. After jit compilation, it was only  $\approx 80-90$  times slower. How this scales with larger matrices is currently unknown.
- ullet Larger error around the  $\lambda=2/3$  region, might not be a coincidence.

# d)

```
H: Array4x4 c
             The complex 4x4 Hamiltonian matrix
    \epsilon 00, \epsilon 10, \epsilon 01, \epsilon 11 = 0.0, 2.5, 6.5, 7.0
    # np.diag not supported by numba
    H0: Array2x2 = np.array([[\epsilon00, 0, 0, 0],
                                [0, \epsilon 10, 0, 0],
                                [0, 0, ε01, 0],
                                [0, 0, 0, ε11]], dtype=complex128)
    Hx = 2
    Hz = 3
    \sigma_x, \sigma_y, \sigma_z = pauli()
    HI: Array2x2 = tensor prod(Hx*\sigma x, \sigma x) + tensor prod(Hz*\sigma z, \sigma z)
    H: Array2x2 = H0 + \lambda *HI
    return H
def density matrix groundstate(λ: float) -> Array4x4:
    Creates the density matrix of the lowest energy state
    Parameters
    _____
        λ: float
            The interaction strength
    Returns
    _____
        ρ0: Array4x4
             The density matrix of the lowest energy state
    H = Hamiltonian 2qbit(\lambda)
    \psi 0: Array4D = np.linalg.eigh(H)[1][:, 0]
    ρ0: Array4x4 = np.outer(ψ_0, ψ_0.conj())
    return ρ0
def partial trace(qubit: Indexing 2qbit, ρ: Array4x4) -> float:
    Partial trace over qubit
    Parameters
    _ _ _ _ _ _ _ _ _ _
        qubit: int
             Which qubit to trace over. Uses 0-based indexing
         ρ: Array4x4
            The density matrix
    Returns
    _ _ _ _ _ _
        Tr(ρ partial): float
             The partial trace of the density matrix
    q0, q1 = create_system_vectors(1)
    I 2 = np.eye(2)
```

```
if qubit == 0:
        op0 = tensor prod(q0, I 2)
        op1 = tensor prod(q1, I 2)
    elif qubit == 1:
        op0 = tensor prod(I 2, q0)
        op1 = tensor prod(I 2, q1)
    else:
        raise ValueError('qubit must be 0 or 1')
    return op0.conj() @ \rho @ op0.T + op1.conj() @ \rho @ op1.T
def entropy(\lambda: float, \epsilon: float = 1e-12) -> float:
    Calculates the entropy of the reduced density matrix
    Parameters
        λ: float
            The interaction strength
        ε: float (optional)
            Small value to avoid log(0). Default is 1e-12
    Returns
    _ _ _ _ _ _
        S: float
            The entropy of the reduced density matrix
    \rho 0 = density matrix groundstate(\lambda)
    \rho A = partial trace(0, \rho0)
    S = -np.trace(\rho A @ np.log2(\rho A + \epsilon))
    # Sanity check
    assert np.iscomplex(S) == False, f'Entropy is complex. Something went wr
    return S.real
```

### **Analytical Solution**

#### **Using SymPy**

```
In []: for i, (eigval, eigvec) in enumerate(zip(eigvals, eigvecs)):
    print(f'Eigenvalue {i}:')
    display(sp.simplify(eigval))
    print(f'Eigenvector {i}:')
    display(sp.simplify(eigvec), )
    print('--'*24)
    print()
```

Eigenvalue 0:

$$-Hz\lambda + rac{arepsilon_{01}}{2} + rac{arepsilon_{10}}{2} - rac{\sqrt{4Hx^2\lambda^2 + arepsilon_{01}^2 - 2arepsilon_{01}arepsilon_{10} + arepsilon_{10}^2}}{2}$$

Eigenvector 0:

$$\begin{bmatrix} 0 \\ \frac{-\varepsilon_{01}+\varepsilon_{10}-\sqrt{4Hx^2\lambda^2+\varepsilon_{01}^2-2\varepsilon_{01}\varepsilon_{10}+\varepsilon_{10}^2}}{2Hx\lambda} \\ 1 \\ 0 \end{bmatrix}$$

-----

Eigenvalue 1:

$$-Hz\lambda+rac{arepsilon_{01}}{2}+rac{arepsilon_{10}}{2}+rac{\sqrt{4Hx^2\lambda^2+arepsilon_{01}^2-2arepsilon_{01}arepsilon_{10}+arepsilon_{10}^2}}{2}$$

Eigenvector 1:

$$\left[\begin{array}{c}0\\\frac{-\varepsilon_{01}+\varepsilon_{10}+\sqrt{4Hx^2\lambda^2+\varepsilon_{01}^2-2\varepsilon_{01}\varepsilon_{10}+\varepsilon_{10}^2}}{2Hx\lambda}\\1\\0\end{array}\right]$$

-----

Eigenvalue 2:

$$Hz\lambda + rac{arepsilon_{00}}{2} + rac{arepsilon_{11}}{2} - rac{\sqrt{4HxHz\lambda^2 + arepsilon_{00}^2 - 2arepsilon_{00}arepsilon_{11} + arepsilon_{11}^2}}{2}$$

Eigenvector 2:

$$\begin{bmatrix} \frac{\varepsilon_{00}-\varepsilon_{11}-\sqrt{4HxHz\lambda^2+\varepsilon_{00}^2-2\varepsilon_{00}\varepsilon_{11}+\varepsilon_{11}^2}}{2Hx\lambda} \\ 0 \\ 0 \\ 1 \end{bmatrix}$$

-----

Eigenvalue 3:

$$Hz\lambda + rac{arepsilon_{00}}{2} + rac{arepsilon_{11}}{2} + rac{\sqrt{4HxHz\lambda^2 + arepsilon_{00}^2 - 2arepsilon_{00}arepsilon_{11} + arepsilon_{11}^2}}{2}$$

Eigenvector 3:

```
\left[\begin{array}{c} \frac{\varepsilon_{00}-\varepsilon_{11}+\sqrt{4HxHz\lambda^2+\varepsilon_{00}^2-2\varepsilon_{00}\varepsilon_{11}+\varepsilon_{11}^2}}{2Hx\lambda} \\ 0 \\ 0 \\ 1 \end{array}\right]
```

\_\_\_\_\_

#### Using the Results from Above

```
In [ ]: def analytical energy eigenstates(λ: float) -> tuple[Array4D, Array4D, Array
              Calculates the energy eigenstates of the Hamiltonian using the analytica
              Parameters
                   λ: float
                        The interaction strength
              Returns
               _ _ _ _ _ _
                   eigenstates: tuple[Array4x4, Array4x4, Array4x4, Array4x4]
                        The energy eigenstates. The order is the same as the energy eige
              Hx = 2
              Hz = 3
              \epsilon 00, \epsilon 10, \epsilon 01, \epsilon 11 = 0, 2.5, 6.5, 7
              q00, q01, q10, q11 = create system vectors(2)
              sqrt1 = np.sqrt(4*(\lambda*Hx)**2 + \epsilon11**2 - 2*\epsilon11*\epsilon00 + \epsilon00**2)
              sqrt2 = np.sqrt(4*(\lambda*Hx)**2 + \epsilon10**2 - 2*\epsilon10*\epsilon01 + \epsilon01**2)
              warnings.filterwarnings('ignore', category=RuntimeWarning) # Suppress Ru
              v00: Array4D c = q00 * -(\epsilon 11 - \epsilon 00 + sqrt1)/(2*\lambda*Hx) + \
                                          q01 * 0 +
                                          q10 * 0 +
                                          q11 * 1
              v01: Array4D c = q00 * -(\epsilon 11 - \epsilon 00 - sqrt1)/(2*\lambda*Hx) + \
                                          q01 * 0 +
                                          q10 * 0 +
                                          q11 * 1
              v10: Array4D c = q00 * 0 +
                                          q01 * -(-\epsilon 10 + \epsilon 01 + sqrt2)/(2*\lambda*Hx) + \
                                          q10 * 1 +
```

```
q11 * 0
    v11: Array4D c = q00 * 0 +
                                q01 * -(-\epsilon 10 + \epsilon 01 - sqrt2)/(2*\lambda*Hx) + \
                                q10 * 1 +
                                q11 * 0
    warnings.filterwarnings('default', category=RuntimeWarning) # Reset warr
    # Normalizing the vectors
    v00: Array4D = v00 / np.linalg.norm(v00)
    v01: Array4D = v01 / np.linalg.norm(v01)
    v10: Array4D = v10 / np.linalg.norm(v10)
    v11: Array4D = v11 / np.linalg.norm(v11)
    return v00, v01, v10, v11
def analytical energy eigenvalues(\lambda: float) -> tuple[float, float, float, fl
    Calculates the energy eigenvalues of the Hamiltonian using the analytical
    Parameters
     _ _ _ _ _ _ _ _ _ _
         λ: float
              The interaction strength
    Returns
         eigenvalues: tuple[float, float, float, float]
              The energy eigenvalues. The order is the same as the energy eige
    1.1.1
    Hx = 2
    Hz = 3
    \epsilon 00, \epsilon 10, \epsilon 01, \epsilon 11 = 0, 2.5, 6.5, 7
    sqrt1 = np.sqrt(\epsilon 00**2 - 2*\epsilon 00*\epsilon 11 + 4*(\lambda*Hx)**2 + \epsilon 11**2)
    sqrt2 = np.sqrt(\epsilon 01**2 - 2*\epsilon 01*\epsilon 10 + 4*(\lambda*Hx)**2 + \epsilon 10**2)
    eigval_00 = 1/2 * (-sqrt1 + \epsilon 00 + 2*\lambda*Hz + \epsilon 11)
    eigval 01 = 1/2 * ( sqrt1 + \epsilon 00 + 2*\lambda*Hz + \epsilon 11)
    eigval_10 = 1/2 * (-sqrt2 + \epsilon 01 - 2*\lambda*Hz + \epsilon 10)
    eigval 11 = 1/2 * ( sqrt2 + \epsilon01 - 2*\lambda*Hz + \epsilon10)
    return eigval 00, eigval 01, eigval 10, eigval 11
def entropy state(state: Literal[0, 1, 2, 3], \lambda: float, \epsilon: float = 1e-12) ->
    Calculates the entropy for a given energy eigenstate
    Parameters
         state: int
              The state to find the entropy of. Uses 0-based indexing
         λ: float
```

```
The interaction strengths
ε: float (optional)
Small value to avoid log(0). Default is le-12

Returns

S: float
The entropy of the reduced density matrix for the original groun

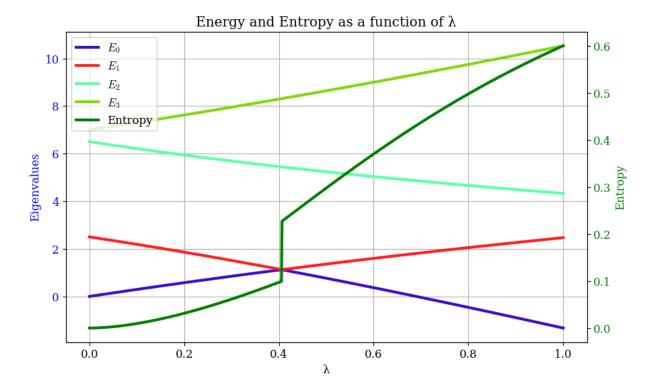
### Array4D = analytical_energy_eigenstates(λ)[state]
p: Array4x4 = np.outer(ψ, ψ.conj())

p_A = partial_trace(0, ρ)
S = -np.trace(ρ_A @ np.log2(ρ_A + ε))

### Sanity check
assert np.iscomplex(S) == False, f'Entropy is complex. Something went wr return S.real
```

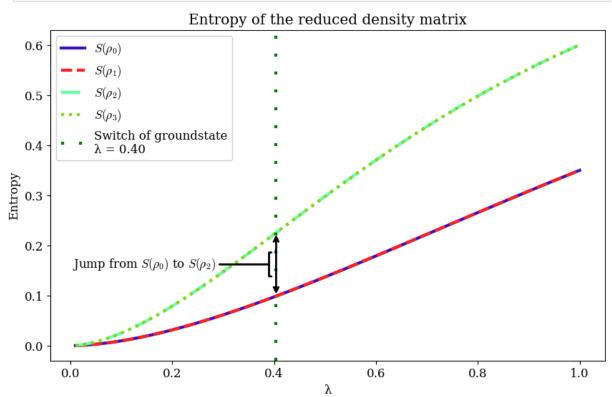
# **Plotting Energy and Entropy**

```
In []: n lambdas = 1000
        lambdas = np.linspace(0, 1, n lambdas)
        energy eigvals = [np.linalg.eigvalsh(Hamiltonian 2qbit(<math>\lambda)) for \lambda in lambdas
        entropies = [entropy(\lambda) for \lambda in lambdas]
        plt.plot(lambdas, energy_eigvals, label=['$E_0$', '$E_1$', '$E_2$', '$E_3$']
        lines energy, labels energy = plt.qca().get legend handles labels()
        plt.xlabel('\lambda')
        plt.ylabel('Eigenvalues', color='blue')
        plt.yticks(color='blue')
        plt.grid()
        plt.twinx()
        plt.plot(lambdas, entropies, color='green', label='Entropy')
        lines entropy, labels entropy = plt.gca().get legend handles labels()
        plt.ylabel('Entropy', color='green')
        plt.yticks(color='green')
        plt.title('Energy and Entropy as a function of \lambda')
        plt.legend(lines energy + lines entropy, labels energy + labels entropy, for
        plt.savefig('figs/d energy entropy.pdf')
        plt.savefig('selected results/d energy entropy.pdf')
        plt.show()
```



```
In [ ]: n lambdas = 100
        lambdas = np.linspace(0, 1, n lambdas)
        entropy 0 = [entropy state(0, \lambda) for \lambda in lambdas]
        entropy 1 = [entropy state(1, \lambda) for \lambda in lambdas]
        entropy 2 = [entropy state(2, \lambda) for \lambda in lambdas]
        entropy 3 = [entropy state(3, \lambda) for \lambda in lambdas]
        energy eigenvalues (\lambda) for \lambda in lambdas
        E0 = energy eigvals[:, 0]
        E1 = energy eigvals[:, 2] # Energy is not sorted the same way as the entropy
        switch point idx = np.argmin(np.abs(E0 - E1))
        gap = entropy 2[switch point idx] - entropy 0[switch point idx]
        plt.plot(lambdas, entropy 0, label=r'$(\rho \ 0)$', linewidth = 3, linestyle='-'
        plt.plot(lambdas, entropy 1, label=r' $S(\rho 1)$', linewidth = 3, linestyle='--
        \label{lambdas} \verb|plt.plot(lambdas, entropy_2, label=r'$S(\rho_2)$', linewidth = 3, linestyle='-.
        plt.plot(lambdas, entropy 3, label=r' $S(\rho 3)$', linewidth = 3, linestyle=':'
        plt.axvline(x=lambdas[switch point idx], color='green', linestyle=(0, (1, 5)
        plt.ylabel('Entropy')
        plt.xlabel('λ')
        plt.title('Entropy of the reduced density matrix')
        plt.legend(fontsize='medium')
        plt.annotate(text="",
                      xy=(lambdas[switch point idx], entropy 0[switch point idx]),
                      xytext=(lambdas[switch_point_idx], entropy_0[switch_point_idx]
                      arrowprops=dict(arrowstyle='<->', lw=2)
        plt.annotate(text="Jump from S(\rho 0) to S(\rho 2)",
                         xy=(lambdas[switch point idx]-.01, entropy 0[switch point id
```

```
xytext=(lambdas[switch_point_idx] - 0.4, entropy_0[switch_point_rowprops=dict(arrowstyle='-[', lw=2)
)
plt.savefig('figs/d_entropy.pdf')
plt.show()
```



### Conclusion

- As the interaction strength  $\lambda$  increases, we see a switch in which the lowest energy state switches from  $|\psi_{00}\rangle$  with corresponding eigenvalue  $E_0$ , to  $|\psi_{01}\rangle$  with corresponding eigenvalue  $E_1$ .
- When calculating entropy, we always look at the lowest energy state. After the switch from  $|\psi_{00}\rangle$  to  $|\psi_{01}\rangle$ , we switch what state we calculate the entropy from. As these two states have different entropies, the entropy makes a sudden jump at the point of the switch.

# e)

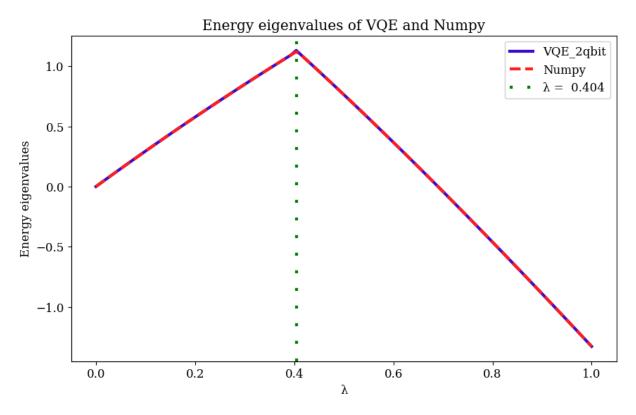
```
In []: @njit def Energy_2qbit(\theta1: float, \theta2: float, \theta3: float, \theta4: float, \lambda: float) -> fl Calculates the energy eigenvalues of the Hamiltonian for a 2-qubit system.
```

```
Parameters
     _ _ _ _ _ _ _ _ _ _
         θ1: float
             The angle (radians) to rotate by around the x-axis of qubit 0
         θ2: float
             The angle (radians) to rotate by around the y-axis of qubit 0
         θ3: float
             The angle (radians) to rotate by around the x-axis of qubit 1
         θ4: float
             The angle (radians) to rotate by around the y-axis of qubit 1
         λ: float
             The interaction strength
    Returns
     _ _ _ _ _ _
         E: float
              The energy eigenvalue (real part)
    R1:
           Array2x2 c = Rx(\theta 1) @ Ry(\theta 2)
           Array2x2 c = Rx(\theta 3) @ Ry(\theta 4)
    CNOT: Array4x4 c = cnot()
    basis = np.array([1, 0, 0, 0], dtype=complex128) # <math>|00\rangle: Will be rotated
    rotated basis: Array4x4 c = CNOT @ tensor prod(R1, R2) @ basis
    E = rotated basis.conj().T @ Hamiltonian_2qbit(\lambda) @ rotated_basis
    assert abs(E.imag) < le-14 , f'Energy is complex. Something went wrong:</pre>
    return E. real
@njit
def VQE 2gbit(N: int, \eta: float, \lambda: float = 0) -> float:
    Variational Quantum Eigensolver using Gradient Descent to find the minim
    Parameters
     _ _ _ _ _ _ _ _ _
         N: int
             Number of iterations
         η: float
             Learning rate
         λ: float
              Interaction strength of the Hamiltonian
    Returns
    _____
         Energy(\theta1, \theta2, \theta3, \theta4, \lambda): float
             The energy eigenvalue (real part)
    1.1.1
    \pi = np.pi
    \theta 1 = 2 \pi n \operatorname{random.rand}()
    \theta 2 = 2 \pi n \operatorname{random.rand}()
    \theta 3 = 2 \pi n \operatorname{random.rand}()
    \theta 4 = 2 \pi n \operatorname{random.rand}()
    # More compact to rename function
```

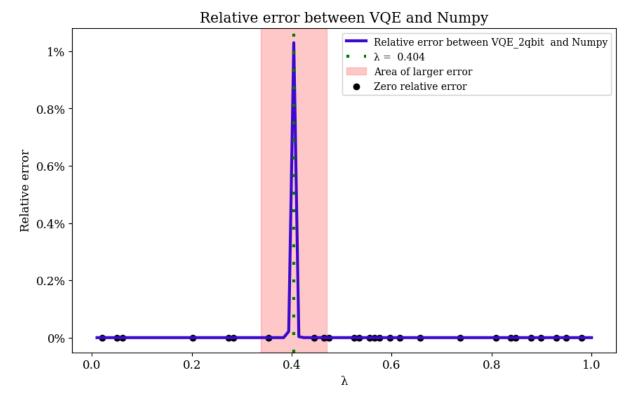
```
E: Callable = Energy_2qbit for _ in range(N):  \Delta E_{\Delta} \theta 1 = (E(\theta 1 + \pi/2, \, \theta 2, \, \theta 3, \, \theta 4, \, \lambda) - E(\theta 1 - \pi/2, \, \theta 2, \, \theta 3, \, \theta 4, \, \lambda)) / 2   \Delta E_{\Delta} \theta 2 = (E(\theta 1, \, \theta 2 + \pi/2, \, \theta 3, \, \theta 4, \, \lambda) - E(\theta 1, \, \theta 2 - \pi/2, \, \theta 3, \, \theta 4, \, \lambda)) / 2   \Delta E_{\Delta} \theta 3 = (E(\theta 1, \, \theta 2, \, \theta 3 + \pi/2, \, \theta 4, \, \lambda) - E(\theta 1, \, \theta 2, \, \theta 3 - \pi/2, \, \theta 4, \, \lambda)) / 2   \Delta E_{\Delta} \theta 4 = (E(\theta 1, \, \theta 2, \, \theta 3, \, \theta 4 + \pi/2, \, \lambda) - E(\theta 1, \, \theta 2, \, \theta 3, \, \theta 4 - \pi/2, \, \lambda)) / 2   \theta 1 -= \eta * \Delta E_{\Delta} \theta 1   \theta 2 -= \eta * \Delta E_{\Delta} \theta 2   \theta 3 -= \eta * \Delta E_{\Delta} \theta 3   \theta 4 -= \eta * \Delta E_{\Delta} \theta 4   \# \textit{Using the actual function name}   E = Energy_2qbit(\theta 1, \, \theta 2, \, \theta 3, \, \theta 4, \, \lambda)   \text{return E}
```

# **Comparing VQE and Exact Eigenvalues**

```
In []: n iterations = 400
      \eta = 1/4
      n lambdas = 100
      lambdas = np.linspace(0, 1, n lambdas)
      result 2qbit: Benchmark Results = benchmark VQE(VQE 2qbit, Hamiltonian 2qbit
      energy eigenvalues (\lambda) for \lambda in lambdas
      E0 = energy eigvals[:, 0]
      E1 = energy_eigvals[:, 2] # Energy is not sorted the same way as the entropy
      switch point idx = np.argmin(np.abs(E0 - E1))
      switch point = lambdas[switch point idx]
      result 2qbit.switch point = switch point
     +----+
              | Time (s) | Times Slower than Numpy |
     | Method
     | VQE 2qbit (cold) | 2.9 | 2300.59
     +----+
     | VQE_2qbit (warm) | 1.4 | 1140.17
     +-----+
              | 0.0012 | X
     +-----+
In [30]: fig = plot eigvals(result 2gbit)
      plt.title('Energy eigenvalues of VQE and Numpy')
      plt.savefig('figs/e vqe vs numpy.pdf')
      plt.show()
```



```
In [ ]: fig = plot_relative_error(result_2qbit)
   plt.title('Relative error between VQE and Numpy')
   plt.savefig('figs/e_relative_error.pdf')
   plt.show()
```



# Conclusion

- The VQE method was 500-1000 times slower than using numpy, as opposed to 50-80 times slower in the 1 qubit case.
  - This might not be surprising as one needs to preform the gradient descent which can't be done in parallel
  - The 2 qubit case involves twice as large matrices and vectors.
- The VQE method for the 2 qubit case needed 4 times the number of iterations to get the same relative error as in the 1 qubit case of about 1%
- There was about one order of magnitude difference in the performance of the VQE method between the 1 and 2 qubit case, compared to using numpy.
- The algorithm could be improved by doing more iterations if the interaction strength  $\lambda$  is near the point where the eigenvalues switch.