

Fabry-Pérot Optical Resonator: Short Cavity / VCSEL Regime

Detailed Analysis of `Chapt9Fig6b.m`

Generated Documentation

Abstract

This document provides comprehensive theoretical foundations and line-by-line code analysis for `Chapt9Fig6b.m`, which calculates the intensity spectrum of a short-cavity Fabry-Pérot resonator ($L = 3 \mu\text{m}$) with high reflectivity mirrors ($R_1 = R_2 = 0.95$). This configuration is representative of vertical-cavity surface-emitting lasers (VCSELs), demonstrating large free spectral range and true single-mode operation capability.

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1 Theoretical Foundation

1.1 Axioms of Short-Cavity Resonators

Axiom 1 (Mode Spacing Scaling). *Free spectral range is inversely proportional to cavity length:*

$$\Delta\nu_{FSR} = \frac{c}{2n_r L} \propto \frac{1}{L} \quad (1)$$

Short cavities have widely spaced modes.

Axiom 2 (Single-Mode Condition). *Single longitudinal mode operation occurs when the gain bandwidth is narrower than the FSR:*

$$\Delta\nu_{gain} < \Delta\nu_{FSR} \quad (2)$$

Axiom 3 (High-Q Requirement). *Short cavities require high reflectivity to maintain sufficient photon lifetime:*

$$\tau_p = \frac{n_r L}{c(1 - \sqrt{R_1 R_2})} \quad (3)$$

1.2 Fundamental Definitions

Definition 1 (VCSEL Cavity). *A vertical-cavity surface-emitting laser has:*

- *Cavity length $L \sim \lambda$ to few λ (typically 1–5 μm)*
- *Distributed Bragg Reflector (DBR) mirrors with $R > 0.99$*
- *Vertical emission perpendicular to substrate*

Definition 2 (Effective Cavity Length). *Including DBR penetration depth:*

$$L_{eff} = L_{active} + 2L_{DBR} \quad (4)$$

where $L_{DBR} \approx \lambda/(4\Delta n)$ is the DBR penetration.

Definition 3 (Mode Number for Short Cavity).

$$m = \frac{2n_r L}{\lambda} \quad (5)$$

For $L = 3 \mu\text{m}$, $n_r = 3.3$, $\lambda = 1.31 \mu\text{m}$: $m \approx 15$.

1.3 Core Theorems

Theorem 1 (High-Finesse Airy Function). *For $R \rightarrow 1$, the Airy function approaches a series of narrow Lorentzians:*

$$I(\nu) \approx I_{max} \sum_m \frac{(\delta\nu/2)^2}{(\nu - \nu_m)^2 + (\delta\nu/2)^2} \quad (6)$$

where $\delta\nu = \Delta\nu_{FSR}/\mathcal{F}$ and $\nu_m = m \cdot \Delta\nu_{FSR}$.

Theorem 2 (Finesse for High Reflectivity). *When $R \rightarrow 1$:*

$$\mathcal{F} = \frac{\pi\sqrt{R}}{1-R} \approx \frac{\pi}{1-R} \quad (7)$$

For $R = 0.9025$ ($R_1 = R_2 = 0.95$): $\mathcal{F} = \pi/0.0975 = 32.2$.

Theorem 3 (Threshold Condition for Short Cavity). *The threshold gain must compensate mirror loss over the short cavity:*

$$\Gamma g_{th}L = \ln \frac{1}{\sqrt{R_1 R_2}} + \alpha_i L \quad (8)$$

For $R = 0.95$ and $L = 3 \mu\text{m}$:

$$\Gamma g_{th} = \frac{\ln(1/0.95)}{3 \times 10^{-4}} + \alpha_i \approx 171 + \alpha_i \text{ cm}^{-1} \quad (9)$$

Theorem 4 (Single-Mode Stability). *Single-mode operation is ensured when:*

$$\Delta\nu_{FSR} > \Delta\nu_{gain} \quad (10)$$

For short cavity with $\Delta\nu_{FSR} = 15 \text{ THz}$ and typical gain bandwidth 30 THz , the nearest mode is significantly detuned from gain peak.

2 Line-by-Line Code Analysis

2.1 Initialization and Parameters

```

1 %Chapt9Fig6b
2 %Fabry-Perot optical resonator
3 clear;
4 clf;
5 err=0.00001;
6 c=3e8;           %speed of light in vacuum
7 nr=3.3;          %effective refractive index
8 lambda0=1310e-9; %center emission wavelength
9 Lc=3e-6;         %Cavity length

```

Short cavity parameters:

$$c = 3 \times 10^8 \text{ m/s} \quad (11)$$

$$n_r = 3.3 \quad (12)$$

$$\lambda_0 = 1310 \text{ nm} \quad (13)$$

$$L_c = 3 \mu\text{m} \quad (14)$$

Cavity length reduced by factor of 100 compared to edge-emitters.

```

1 r1=0.95;          %Reflectivity of mirror1
2 r2=0.95;          %Reflectivity of mirror2
3 r=r1*r2;          %Optical loss per round-trip

```

High-reflectivity symmetric mirrors:

$$R_1 = R_2 = 0.95 \quad (\text{DBR mirrors}) \quad (15)$$

$$R = R_1 R_2 = 0.9025 \quad (16)$$

Typical of multi-layer DBR stacks with 20–30 periods.

2.2 Finesse and Spectral Parameters

```

1 F=pi*r^0.5/(1-r)          %Optical Finesse
2 Fconst=(2*F/pi)^2;
3 Imax=1/((1-r)^2);         %Peak intensity normalized to IO

```

High finesse:

$$\mathcal{F} = \frac{\pi\sqrt{0.9025}}{1 - 0.9025} = \frac{\pi \times 0.95}{0.0975} = \frac{2.98}{0.0975} = 30.6 \quad (17)$$

Coefficient of finesse:

$$F_{const} = \left(\frac{2\mathcal{F}}{\pi} \right)^2 = \frac{4R}{(1-R)^2} = \frac{3.61}{0.0095} = 380 \quad (18)$$

Peak intensity enhancement:

$$I_{max} = \frac{1}{(1-R)^2} = \frac{1}{0.0095} = 105 \quad (19)$$

Enormous enhancement—photon density 100× higher inside cavity.

```

1 f0=c*1e-12/lambda0;          %center frequency in THz
2 deltaf=c*1e-12/(2*Lc*nr)    %FSR in THz
3 deltawavelength=lambda0^2/(2*Lc*nr)  %FSR in m

```

Large free spectral range:

$$f_0 = \frac{3 \times 10^8}{1.31 \times 10^{-6}} = 229 \text{ THz} \quad (20)$$

$$\Delta f_{FSR} = \frac{c}{2n_r L} = \frac{3 \times 10^8}{2 \times 3.3 \times 3 \times 10^{-6}} = 15.2 \text{ THz} \quad (21)$$

Mode spacing in wavelength:

$$\Delta\lambda_{FSR} = \frac{\lambda_0^2}{2n_r L} = \frac{(1.31)^2 \times 10^{-12}}{19.8 \times 10^{-6}} = 87 \text{ nm} \quad (22)$$

Mode linewidth:

$$\delta\nu = \frac{\Delta\nu_{FSR}}{\mathcal{F}} = \frac{15.2 \text{ THz}}{30.6} = 0.50 \text{ THz} = 500 \text{ GHz} \quad (23)$$

2.3 Extended Frequency Scan

```

1 for i=[1:1:10000]
2     Frequency(i)=(f0+i*0.003);
3     Fwavelength(i)=c/Frequency(i);
4     x=sin(pi*Frequency(i)/deltaf);
5     x2=x*x;
6     Intensity(i)=Imax/(1+(Fconst*(x2)));
7 end

```

Key differences from previous codes:

- 10,000 points (vs 1,000)
- Step size 0.003 THz (vs 0.001 THz)
- Total span: 30 THz

This covers approximately two FSR periods to show multiple resonances.

Airy function with high F_{const} :

$$I(\nu) = \frac{105}{1 + 380 \sin^2\left(\frac{\pi\nu}{15.2 \text{ THz}}\right)} \quad (24)$$

At anti-resonance:

$$I_{min} = \frac{105}{1 + 380} = 0.28 \quad (25)$$

Contrast ratio: $I_{max}/I_{min} = 105/0.28 = 375$.

2.4 Plotting

```

1 figure(1);
2 plot(Frequency, Intensity);
3 %axis([f0, f0+1, 0, 5]);
4 hold on;
5 xlabel('Frequency ,\nu(THz)'), ylabel('Intensity');
6 ttl=sprintf('Chapt9Fig6b ,r1=%4.2f ,r2=%4.2f ,nr=%4.2f ,f0=%7.2e THz ,%
Lc=%7.2e m',r1,r2,nr,f0,Lc)
7 title(ttl);
8
9 figure(2);
10 plot(Fwavelength*10^-12, Intensity);
11 xlabel('Wavelength ,\lambda(m)'), ylabel('Intensity');
12 title(ttl);

```

Note: The axis command is commented out to show full range.

3 Comparison: Edge-Emitter vs VCSEL

Parameter	Fig6a1 (Edge, Fresnel)	Fig6a2 (Edge, Coated)	Fig6b (VCSEL)
L_c	300 μm	300 μm	3 μm
R_1, R_2	0.29, 0.29	0.4, 0.8	0.95, 0.95
R	0.082	0.32	0.9025
\mathcal{F}	0.98	2.61	30.6
I_{max}	1.19	2.16	105
$\Delta\nu_{FSR}$	152 GHz	152 GHz	15.2 THz
$\delta\nu$	155 GHz	58 GHz	500 GHz
Mode number m	1511	1511	15

4 Physical Interpretation

4.1 VCSEL Mode Structure

With only $m \approx 15$ longitudinal modes possible:

- Mode spacing 15 THz ≈ 87 nm in wavelength
- Only one mode typically within gain bandwidth
- Inherent single longitudinal mode operation

4.2 Gain-Mode Alignment

The condition for lasing:

$$g(\lambda_m) > g_{th} \quad (26)$$

With 87 nm mode spacing, only the mode nearest gain peak lases.

4.3 Threshold Considerations

Despite high R , threshold gain is high due to short path:

$$g_{th} = \frac{1}{\Gamma L} \ln \frac{1}{R} = \frac{1}{0.03 \times 3 \times 10^{-4}} \ln(1.11) = \frac{0.104}{9 \times 10^{-6}} = 11,600 \text{ cm}^{-1} \quad (27)$$

(assuming $\Gamma = 0.03$ for thin QW active region)

This requires quantum well active regions with high material gain.

4.4 Photon Lifetime

$$\tau_p = \frac{n_r}{c \cdot \alpha_m} = \frac{n_r L}{c \cdot \ln(1/R)} = \frac{3.3 \times 3 \times 10^{-6}}{3 \times 10^8 \times 0.103} = 0.32 \text{ ps} \quad (28)$$

Very short photon lifetime leads to:

- High modulation bandwidth potential
- Reduced relaxation oscillation effects
- Fast switching capability

4.5 Intensity Enhancement

The factor $I_{max} = 105$ means:

- Photon density inside cavity $100 \times$ incident
- Compensates for short gain length
- Essential for achieving lasing in thin active region

5 Numerical Methods

5.1 Resolution Requirements

With $\delta\nu = 500$ GHz and step size 3 GHz:

$$\text{Points per FWHM} = \frac{500}{3} \approx 167 \quad (29)$$

Excellent resolution of resonance peaks.

5.2 Span Coverage

Total scan: $10,000 \times 0.003 = 30$ THz, covering:

$$\frac{30 \text{ THz}}{15.2 \text{ THz}} \approx 2 \text{ FSR periods} \quad (30)$$

Shows two complete resonance peaks.

5.3 Numerical Precision

With $F_{const} = 380$, near resonance:

$$I \approx \frac{105}{1 + 380 \times (\pi\delta\nu/\Delta\nu_{FSR})^2} \quad (31)$$

Standard double precision handles this without issues.

6 Summary

This code simulates a short-cavity, high-finesse Fabry-Pérot resonator representative of VCSELs:

Key parameters:

- Cavity length $L = 3 \mu\text{m}$ ($100\times$ shorter than edge-emitter)
- High reflectivity $R = 0.95$ (DBR mirrors)
- Finesse $\mathcal{F} = 30.6$ ($12\times$ higher than coated edge-emitter)

Resulting characteristics:

- Large FSR = 15 THz (87 nm)—ensures single-mode operation
- Few longitudinal modes ($m \approx 15$)
- High intensity enhancement ($105\times$)—compensates short gain path
- Sharp resonances with contrast ratio 375:1

VCSEL advantages demonstrated:

- Inherent single longitudinal mode selection
- High modulation bandwidth from short τ_p
- Circular beam from vertical geometry (not shown in this 1D model)

This simulation captures the essential physics that makes VCSELs fundamentally different from edge-emitting lasers.