

# Chapt7Fig6b.m

## Fermi-Dirac Distribution at Lower Carrier Density

### Non-Degenerate Semiconductor Regime

#### Semiconductor Physics Documentation

#### Abstract

This document analyzes `Chapt7Fig6b.m`, which plots the Fermi-Dirac distribution function at a carrier density of  $n = 10^{17} \text{ cm}^{-3}$ , one order of magnitude lower than `Chapt7Fig6a.m`. This lower density places the system closer to the non-degenerate (Maxwell-Boltzmann) regime, providing insight into the classical-quantum transition.

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# 1 Theoretical Foundation

## 1.1 Fermi-Dirac Statistics Review

**Theorem 1** (Fermi-Dirac Distribution). *The probability of occupation of a quantum state with energy  $E$  is:*

$$f(E) = \frac{1}{e^{(E-\mu)/k_B T} + 1} \quad (1)$$

where  $\mu$  is the chemical potential (Fermi level) and  $T$  is the temperature.

## 1.2 Degeneracy Parameter

**Definition 1** (Degeneracy Parameter). *The degeneracy of an electron gas is characterized by:*

$$\eta = \frac{\mu}{k_B T} \quad (2)$$

- $\eta \gg 1$ : Degenerate regime (quantum statistics essential)
- $\eta \ll -1$ : Non-degenerate regime (classical approximation valid)
- $|\eta| \lesssim 1$ : Transition regime

## 1.3 Non-Degenerate Limit: Maxwell-Boltzmann Approximation

**Theorem 2** (Maxwell-Boltzmann Limit). *When  $E - \mu \gg k_B T$  for all occupied states, the Fermi-Dirac distribution reduces to:*

$$f(E) \approx e^{-(E-\mu)/k_B T} \quad (3)$$

*Proof.* When  $(E - \mu)/k_B T \gg 1$ , we have  $e^{(E-\mu)/k_B T} \gg 1$ , so:

$$f(E) = \frac{1}{e^{(E-\mu)/k_B T} + 1} \approx \frac{1}{e^{(E-\mu)/k_B T}} = e^{-(E-\mu)/k_B T} \quad (4)$$

□

## 1.4 Chemical Potential Dependence on Carrier Density

**Theorem 3** (Non-Degenerate Carrier Density). *In the non-degenerate limit, the carrier density is:*

$$n = N_c \exp\left(\frac{\mu}{k_B T}\right) \quad (5)$$

where  $N_c = 2 \left(\frac{m^* k_B T}{2\pi\hbar^2}\right)^{3/2}$  is the effective density of states.

**Corollary 1** (Chemical Potential in Non-Degenerate Regime).

$$\mu = k_B T \ln\left(\frac{n}{N_c}\right) \quad (6)$$

When  $n < N_c$ , the chemical potential is negative (below the band edge).

## 1.5 Effective Density of States for GaAs

For GaAs with  $m^* = 0.07m_0$ :

$$N_c = 2 \left( \frac{0.07m_0 k_B T}{2\pi\hbar^2} \right)^{3/2} \quad (7)$$

At  $T = 300$  K:

$$N_c \approx 4.7 \times 10^{17} \text{ cm}^{-3} \quad (8)$$

Since  $n = 10^{17} \text{ cm}^{-3} < N_c$ , the system is near the non-degenerate regime at room temperature.

## 2 Line-by-Line Code Analysis

### 2.1 Key Parameter Difference

```
1 n=1.e17
```

$$n = 1 \times 10^{17} \text{ cm}^{-3} \quad (9)$$

*This is the critical difference from Chapt7Fig6a.m. The carrier density is 10 times lower, pushing the system toward the non-degenerate regime.*

### 2.2 Physical Constants (Identical to Chapt7Fig6a)

```
1 m0=9.109382;
```

$$m_0 = 9.109382 \times 10^{-31} \text{ kg} \quad (10)$$

*Bare electron mass.*

```
1 m1=0.07;
```

$$m^* = 0.07 m_0 \quad (11)$$

*Effective mass for GaAs conduction band electrons.*

```
1 echarge=1.6021764;
```

$$e = 1.6021764 \times 10^{-19} \text{ C} \quad (12)$$

*Elementary charge.*

```
1 hbar=1.05457159;
```

$$\hbar = 1.05457159 \times 10^{-34} \text{ J} \cdot \text{s} \quad (13)$$

*Reduced Planck constant.*

```
1 kB=8.61734e-5;
```

$$k_B = 8.61734 \times 10^{-5} \text{ eV/K} \quad (14)$$

*Boltzmann constant in convenient units.*

```
1 rerr=1.e-3;
```

$$\epsilon_{\text{rel}} = 10^{-3} \quad (15)$$

*Relative error tolerance for chemical potential calculation.*

## 2.3 Temperature Array

```
1 Temperature=[0.1,100,200,300,400];
```

$$T \in \{0.1, 100, 200, 300, 400\} \text{ K} \quad (16)$$

*Same temperature range as Chapt7Fig6a for direct comparison.*

```
1 nTemp=length(Temperature)
```

$$N_{\text{temp}} = 5 \quad (17)$$

*Number of temperature curves to plot.*

## 2.4 Energy Grid Setup

```
1 nplotpoints=300;
2 Emin=0;
3 Emax=200;
4 Estep=(Emax-Emin)/nplotpoints;
```

$$E_k = k \cdot \frac{200 \text{ meV}}{300} \quad \text{for } k = 0, 1, \dots, 300 \quad (18)$$

*Energy discretization from 0 to 200 meV.*

## 2.5 Main Computation Loop

```
1 for j=1:1:nTemp
2     kelvin=Temperature(j);
3     kBT=1000.*kelvin*kB;
4     beta=1./kBT;
```

$$\beta = \frac{1}{k_B T} \quad [\text{meV}^{-1}] \quad (19)$$

*Compute inverse thermal energy for each temperature.*

```
1 mu1=chempot(kelvin,m1,n,rerr);
```

$$\mu = \mu(T, m^* = 0.07m_0, n = 10^{17} \text{ cm}^{-3}) \quad (20)$$

*At lower carrier density,  $\mu$  will be smaller (possibly negative at high  $T$ ).*

## 2.6 Fermi Function Evaluation Loop

```
1 for k = 1:1:nplotpoints+1
2     energy=((k-1)*Estep)+Emin;
3     E(k)=energy;
4     Prob=fermi(beta,energy,mu1);
5     FD(k)=Prob;
6 end;
```

$$f(E_k) = \frac{1}{\exp[\beta(E_k - \mu)] + 1} \quad (21)$$

*Evaluate Fermi-Dirac distribution at each energy point.*

## 2.7 Plotting

```

1      hold on; figure(1); plot(E, log(FD)); axis([0 Emax -10 0]);
2  end
3  xlabel('Energy, E (meV)');
4  ylabel('ln(occupation factor)');
5  ttl=sprintf('Chapt7Fig6b, m*=%4.2f m0, Tmax=%4.1f K, n=%7.2e cm-3',
6  m1, kelvin, n);
  title(ttl);

```

$$\ln f(E) \text{ vs } E \quad (22)$$

*Logarithmic plot showing the exponential tail behavior more prominently.*

## 3 Physical Interpretation and Comparison

### 3.1 Effect of Lower Carrier Density

Comparing  $n = 10^{17} \text{ cm}^{-3}$  (this code) with  $n = 10^{18} \text{ cm}^{-3}$  (Chapt7Fig6a):

Property	$n = 10^{17} \text{ cm}^{-3}$	$n = 10^{18} \text{ cm}^{-3}$
Chemical potential $\mu$	Lower	Higher
Degeneracy $\eta = \mu/k_B T$	Smaller	Larger
Regime	Non-degenerate	Degenerate

Table 1: Comparison of regimes

### 3.2 Chemical Potential Estimates

At  $T = 300 \text{ K}$  with  $N_c \approx 4.7 \times 10^{17} \text{ cm}^{-3}$ :

For  $n = 10^{17} \text{ cm}^{-3}$ :

$$\mu \approx k_B T \ln \left( \frac{10^{17}}{4.7 \times 10^{17}} \right) \approx k_B T \ln(0.21) \approx -1.56 k_B T \approx -40 \text{ meV} \quad (23)$$

For  $n = 10^{18} \text{ cm}^{-3}$ :

$$\mu \approx k_B T \ln \left( \frac{10^{18}}{4.7 \times 10^{17}} \right) \approx k_B T \ln(2.13) \approx 0.76 k_B T \approx 20 \text{ meV} \quad (24)$$

### 3.3 Implications for the Plot

With lower carrier density:

- The chemical potential is lower (possibly negative at high  $T$ )
- The Fermi function drops off more steeply at low energies
- The system more closely follows Maxwell-Boltzmann statistics
- The logarithmic plot shows nearly linear behavior over a wider range

### 3.4 Temperature Dependence

At very low temperature ( $T = 0.1$  K):

$$k_B T \approx 8.6 \times 10^{-6} \text{ eV} = 8.6 \times 10^{-3} \text{ meV} \quad (25)$$

Even at low carrier density, the thermal energy is so small that the distribution still shows sharp step-like behavior, with all electrons in states below the Fermi energy.

At  $T = 400$  K:

$$k_B T \approx 34.5 \text{ meV} \quad (26)$$

The distribution is significantly broadened, and with  $\mu < 0$ , the occupation probability at  $E = 0$  is already less than 0.5.

## 4 Numerical Considerations

### 4.1 The chempot Function

The function `chempot(kelvin, m1, n, rerr)` solves:

$$n = \frac{1}{2\pi^2} \left( \frac{2m^*}{\hbar^2} \right)^{3/2} \int_0^\infty \frac{\sqrt{E}}{e^{(E-\mu)/k_B T} + 1} dE \quad (27)$$

This is typically done using:

1. Initial guess based on non-degenerate approximation:  $\mu_0 = k_B T \ln(n/N_c)$
2. Newton-Raphson iteration to refine
3. Convergence criterion:  $|\mu_{i+1} - \mu_i| < \epsilon_{\text{rel}} \cdot |\mu_i|$

### 4.2 The fermi Function

The function `fermi(beta, E, mu)` computes:

$$f = \frac{1}{e^{\beta(E-\mu)} + 1} \quad (28)$$

Numerical care is needed for large arguments:

- If  $\beta(E - \mu) > 700$ :  $f \approx 0$
- If  $\beta(E - \mu) < -700$ :  $f \approx 1$

## 5 Summary

This code demonstrates the Fermi-Dirac distribution in the non-degenerate regime, where the carrier density is comparable to or less than the effective density of states. The key physical insight is that reducing the carrier density by a factor of 10 shifts the chemical potential significantly lower, bringing the system closer to classical (Maxwell-Boltzmann) behavior while still retaining quantum effects at low temperatures.