

Chapt8Exercise7b.m

Thomas-Fermi vs RPA Screening at Low Density Breakdown of Thomas-Fermi Approximation

Semiconductor Physics Documentation

Abstract

This document analyzes Chapt8Exercise7b.m, which compares Thomas-Fermi and RPA scattering rates at a low carrier density of $n = 10^{14} \text{ cm}^{-3}$. At this density, the Thomas-Fermi approximation breaks down, revealing the importance of non-local screening effects captured by RPA.

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1 Theoretical Foundation

1.1 Validity Conditions for Thomas-Fermi Screening

Theorem 1 (Thomas-Fermi Validity). *The Thomas-Fermi approximation is valid when:*

$$q \ll k_F \quad \text{and} \quad q_{TF} \sim k_F \quad (1)$$

This requires the electron gas to be sufficiently dense (degenerate).

Definition 1 (Degeneracy Criterion). *An electron gas is degenerate when:*

$$E_F \gg k_B T \quad (2)$$

where the Fermi energy is $E_F = \hbar^2 k_F^2 / (2m^)$.*

1.2 Low-Density Regime

At $n = 10^{14}$ cm $^{-3}$:

$$k_F = (3\pi^2 n)^{1/3} = (3\pi^2 \times 10^{20})^{1/3} \approx 1.43 \times 10^7 \text{ m}^{-1} \quad (3)$$

Compare to $n = 10^{18}$ cm $^{-3}$:

$$k_F = 3.09 \times 10^8 \text{ m}^{-1} \quad (4)$$

The Fermi wavevector decreases by a factor of ~ 20 .

1.3 Consequences for Screening

Theorem 2 (Weak Screening Limit). *At low density:*

$$q_{TF} \propto \sqrt{k_F} \propto n^{1/6} \quad (5)$$

The screening length increases, making screening less effective.

Corollary 1 (RPA Corrections). *When $q \gtrsim k_F$, the RPA dielectric function deviates significantly from Thomas-Fermi:*

$$\varepsilon_{RPA}(q) \neq \varepsilon_{TF}(q) \quad \text{for } q \gtrsim k_F \quad (6)$$

2 Line-by-Line Code Analysis

2.1 Key Parameter: Low Carrier Density

```
1 clear;clf;
2 n=1e14;
3 n=n*1e6;
```

$$n = 10^{14} \text{ cm}^{-3} = 10^{20} \text{ m}^{-3} \quad (7)$$

Critical difference: carrier density is 10^4 times lower than Chapt8Exercise7a.

2.2 Physical Constants (Same as 7a)

```
1 m0=9.109382e-31;
2 echarge=1.6021764e-19;
3 epsilon0=8.85419e-12;
4 hbar=1.05457159e-34;
5 hbar3=hbar^3;
6 m=0.07*m0;
7 epsilonnr0=13.2;
8 epsilon=epsilon0*epsilonnr0;
```

Material parameters for GaAs remain unchanged.

2.3 Fermi Wavevector

```
1 kF=(3*(pi^2)*n)^(1/3)
```

$$k_F = (3\pi^2 \times 10^{20})^{1/3} \approx 1.43 \times 10^7 \text{ m}^{-1} \quad (8)$$

Much smaller than at high density, indicating weak degeneracy.

2.4 Energy Loop

```
1 E=-0.1*echarge;
2 for j=1:1:2
3 E=E+0.2*echarge;
4 k=sqrt(2*m*E)/hbar;
5 k3=k^3;
```

$$E \in \{0.1, 0.3\} \text{ eV}, \quad k = \frac{\sqrt{2m^*E}}{\hbar} \quad (9)$$

Same electron energies as Chapt8Exercise7a for comparison.

2.5 RPA Constants

```
1 rs0=((3/(4*pi*n))^(1/3))*(m*(echarge^2)/(4*pi*epsilon0*(hbar^2)));
2 xi=((32*(pi^2)/9)^(1/3))/(pi^2);
```

$$r_s = \left(\frac{3}{4\pi n} \right)^{1/3} \cdot \frac{m^*e^2}{4\pi\epsilon_0\hbar^2} \quad (10)$$

At low density, r_s is larger, indicating stronger electron-electron correlations.

2.6 Thomas-Fermi Wavevector

```
1 qTF=sqrt(kF*m*echarge^2/(epsilon*(pi^2)*(hbar^2)));
```

$$q_{TF} = \sqrt{\frac{k_F m^* e^2}{\epsilon \pi^2 \hbar^2}} \quad (11)$$

Smaller k_F leads to smaller q_{TF} , weaker screening.

2.7 Angular Discretization

```
1 theta=[pi/1800:pi/1800:pi];
2 q=2*k*sin(theta/2);
3 eta=sin(theta/2);
4 deta=pi*cos(theta/2)./2/180;
5 eta3=eta.^3;
```

Same angular grid as Chapt8Exercise7a.

2.8 RPA Dielectric Function

```
1 x=q/kF;
2 absx=abs(x+2)./abs(x-2);
3 RPAepsilonnr=epsilon0+((rs0./x.^3)*xi.* (x+((1-(x.^2)/4).*log(absx)
)));
```

$$x = q/k_F \quad (\text{now larger values since } k_F \text{ is smaller}) \quad (12)$$

At low density, $x = q/k_F$ can be $\gg 1$, where TF fails.

2.9 Thomas-Fermi Dielectric Function

```
1 TFepsilon=epsilon*(1+qTF^2./q.^2);
```

$$\varepsilon_{\text{TF}} = \varepsilon \left(1 + \frac{q_{\text{TF}}^2}{q^2} \right) \quad (13)$$

With smaller q_{TF} , the screening term is less effective.

2.10 Scattering Rates

```
1 RPArate=2*pi*m/hbar3/k3*n*(echarge^2/4/pi/epsilon0)^2.*deta./
RPAepsilon.^2./eta3;
2 TFrater =2*pi*m/hbar3/k3*n*(echarge^2/4/pi)^2.*deta./TFepsilon.^2./
eta3;
```

$$\frac{1}{\tau} = \frac{2\pi m^*}{\hbar^3 k^3} n \left(\frac{e^2}{4\pi} \right)^2 \frac{d\eta}{\varepsilon(q)^2 \eta^3} \quad (14)$$

Same formula, but n is lower and screening is weaker.

2.11 Plotting with Restricted Axis

```
1 plot(theta*180/pi, TFrater, 'r');
2 axis([0,35,0,9e9]);
3 ...
4 plot(theta*180/pi, RPArate, 'b');
```

Axis restricted to $\theta \in [0, 35]$ to focus on the forward-scattering region where differences are largest.

3 Physical Interpretation

3.1 Why TF Fails at Low Density

The Thomas-Fermi approximation assumes:

1. Local response: Perturbation at \mathbf{r} only affects density at \mathbf{r}
2. Degenerate limit: All states up to E_F are filled
3. Slow spatial variation: $q \ll k_F$

At $n = 10^{14} \text{ cm}^{-3}$:

- Fermi energy: $E_F \approx 0.1 \text{ meV}$ (non-degenerate at room temperature)
- Typical scattering wavevector: $q \sim k \gg k_F$
- TF assumption $q \ll k_F$ is violated

3.2 RPA Captures the Correct Physics

The RPA (Lindhard) function includes:

$$\varepsilon(q) = 1 + \frac{q_{\text{TF}}^2}{q^2} \cdot g(q/k_F) \quad (15)$$

where $g(x)$ is a correction factor that:

- Equals 1 for $x \ll 1$ (recovers TF)
- Decreases for $x \gtrsim 1$ (weakens screening)
- Vanishes for $x \gg 1$ (bare Coulomb restored)

3.3 Expected Results

- TF (red curve): Underestimates scattering at large angles
- RPA (blue curve): Shows stronger scattering due to weaker screening
- Difference is most pronounced at large q (large angles)

4 Comparison with Chapt8Exercise7a

Parameter	Chapt8Exercise7a	Chapt8Exercise7b
Carrier density n	10^{18} cm^{-3}	10^{14} cm^{-3}
Fermi wavevector k_F	$3.09 \times 10^8 \text{ m}^{-1}$	$1.43 \times 10^7 \text{ m}^{-1}$
Screening	Strong	Weak
TF validity	Good	Poor
TF-RPA agreement	Good	Poor

Table 1: Comparison of the two exercises

5 Summary

This code demonstrates the breakdown of the Thomas-Fermi approximation at low carrier densities. When the electron gas is non-degenerate and the typical scattering wavevector exceeds the Fermi wavevector, the local Thomas-Fermi approximation fails, and the full RPA treatment becomes necessary. The restricted angular range in the plot highlights the region where differences between TF and RPA are most pronounced.