

# 1. Install(for Linux )

## 1.1. requirement

```
1 | see  
  | https://www.intel.cn/content/www/cn/zh/developer/tools/oneapi/toolkits.html#base-kit  
2 | Intel® oneAPI Math Kernel Library  
3 | Intel® Distribution for Python*  
4 | Intel® C++ Compiler Classic
```

Suggested to directly install two integrated packages:

1. Intel oneAPI Base Toolkit
2. Intel® oneAPI HPC Toolkit

The current version of oneAPI has removed Python from the above integrated packages, so a separate download of the corresponding Python interpreter,

Intel® Distribution for Python\*

is required.

Please add the necessary configurations to the bashrc file:

```
1 | echo "source /opt/intel/oneapi/setvars.sh" >> ~/.bashrc
```

## 1.2. compile

Unzip the compressed file `FQHE_SPHERE.zip`, enter the extracted directory, ensure that the `icc` and `icpc` compilation commands are available, and verify that the MKL (Math Kernel Library) path is correct.

```
1 | hu@gugugu:~$ echo ${MKLROOT}  
2 | /opt/intel/oneapi/mkl/2022.0.2
```

Execute

```
1 | bash compile.sh
```

## 2. Usage

```
1 | hu@gugugu:~$ which python3  
2 | /opt/intel/oneapi/intelpython/latest/bin/python3
```

## 2.1. ED

The following script directly calculates the energy spectrum of the critical point when  $N = 12$ ,  $Z_2 = 1$ ,  $PH = 1$ , and  $L_z = 0$ .

```
1  #!/opt/intel/oneapi/intelpython/latest/bin/python3
2  # -*- coding:utf-8 -*-
3
4
5  import numpy as np
6  import sys
7  sys.path.append(r"/home/hu/study/Code/FQHE_Package/") # the path to package
8  import FQHE_SPHERE as fqhe
9
10
11  n = 12;
12  Z2 = 1;
13  PH = 1;
14  LZ = 0;
15
16  inter = fqhe.interaction()
17  inter.setpseudo2body(inten=[0.475,0.1])
18  intra = fqhe.interaction()
19  intra.setpseudo2body(inten=[0.0,0.0])
20
21  kk = fqhe.fermion.QHFM(n,n,num_thread=8)
22  kk.setL2ab(a=0.0,b=0.0)
23
24  sol = fqhe.fermion.dl_solver(kk, intra,intra,inter, "H2L2",qn=
    {"Lz":Lz,"Z_2":Z2,"PH":PH}, t=0.316)
25  eigvsL2_p = sol.eigsv(k=5,tol=10**(-6))
26  eigvs = sorted( eigvsL2_p )
27
28
29  ge = eigvsL2_p[0].get("energy")#-1.86886840
30  gap1 = eigvsL2_p[3].get("energy") - ge#-0.62067682 - ge
31  resu = "<Index>\t<L^2>\tPH\tZ_2\t<H+L^2>\t\tE_n=<H>\t\tEn-E0\t\tscale\n"
32  i = 0
33  for vec in eigvs:
34      resu += "{:2d}\t".format( i )
35      resu += "{:+.2f}\t".format( vec.get("L2") )
36      resu += "{:d}\t".format( int(vec.get("PH")) )
37      resu += "{:d}\t".format( int(vec.get("Z_2")) )
38      resu += "{:+.8f}\t".format( vec.eig )
39      resu += "{:+.8f}\t".format( vec.get("energy") )
40      resu += "{:+.8f}\t".format( vec.get("energy")-ge )
41      resu += "{:+.8f}\n".format(3*(vec.get("energy")-ge)/gap1)
42      i += 1
43  print(resu)
```

output

```

1 matvec(self, vector): 96 times
2 eigsv: Measure: H2L2
3 total 0.7496089935302734 : Hamiltonian( 0.5623340606689453 ) lanczos(
  0.17470121383666992 )
4 <Index> <L^2> PH Z_2 <H+L^2> E_n=<H> En-E0
  scale
5 0 -0.00 +1 +1 -1.86886840 -1.86886840 +0.00000000
  +0.00000000
6 1 -0.00 +1 +1 -1.28306494 -1.28306494 +0.58580346
  +1.40796525
7 2 +2.00 +1 +1 -0.87103274 -0.87103274 +0.99783566
  +2.39827525
8 3 +6.00 +1 +1 -0.62067682 -0.62067682 +1.24819158
  +3.00000000
9 4 +6.00 +1 +1 -0.45961655 -0.45961655 +1.40925185
  +3.38710471

```

Detailed explanation:

### 2.1.1. fqhe.fermion.QHFM

Quantum Hall ferromagnetic(QHFM) class

```

1 class QHMFermionDL(FermionDL):
2     def __init__(self, n, n_o, dtype=np.float64, projectionQ=0, num_thread=8):

```

parameters:

```

1 1. n::Integer
2   the number of electrons
3
4 2. n_o::Integer
5   the number of orbits(**not flux**)
6
7 3. dtype: Deprecated, not editable.
8
9 4. projectionQ: Deprecated, not editable.
10
11 5. num_thread::Integer
12   The number of threads used for the computation.

```

method:

```

1 kk.setL2ab(a=0.0,b=0.0)

```

modify the Hamiltonian from  $H$  to:

$$\begin{aligned}
 &H + a \cdot (L^2 - b)^2 \quad \text{if } b \neq 0 \\
 &H + a \cdot L^2 \quad \text{if } b = 0
 \end{aligned}$$

This modification is employed to select specific values of  $\ell$  (angular momentum quantum number), but in general, convergence becomes difficult.

## 2.1.2. Interaction

Two body interaction  $V_{\sigma\sigma'}(r_{12})$  class

```
1 inter = fqhe.interaction()
2 inter.setpseudo2body(inten=[v0,v1,v2,.....])
```

$V_i$  are Haldane Pseudo-potential coefficients.

## 2.1.3. solver

eigensolver

```
1 class fqhe_d1_solver():
2     def __init__(self, ins, intra1, intra2, inter, itype, qn={"Lz": 0},
3         cdw=np.zeros(2, dtype=np.float64), t=0.0):
4         '''
5         ins: fqhe.fermion.QHFM object used for the computation.
6         intra1, intra2: Intralayer interactions.
7         inter: Interlayer interaction.
8         itype: "twobody", "L2" or "H2L2"
9         "twobody": H = H2 = V_{updown} + V_{upup} + V_{downdown} + t *
10         n^x_{00}
11         "L2": H = L^2
12         "H2L2": H = H2 + a * (L^2 - b)^2 or H = H2 + a * L^2
13         qn: Quantum numbers of the Hilbert space
14         {"Lz": 0}: Lz=0
15         {"Lz": 0, "Z_2": 1}: Lz=0 and Z_2=1
16         {"Lz": 0, "Z_2": 1, "PH": 1}: Lz=0, Z_2=1, and PH=1
17         cdw: (invalid)
18         dtype: (invalid)
19         t: Hopping term
20         '''
```

Explanation:

This is a Python class named `fqhe_d1_solver`. The class constructor `__init__` takes several parameters to initialize the solver object. Here's an explanation of the parameters:

- `ins`: An object of the `fqhe.fermion.QHFM` class is used for the computation.
- `intra1, intra2`: Intralayer interactions.
- `inter`: Interlayer interaction.
- `itype`: Specifies the type of Hamiltonian to be used for the computation. It can take one of three values: "twobody", "L2", or "H2L2". Depending on the value, the Hamiltonian will be set accordingly.
- `qn`: Quantum numbers of the Hilbert space represented as a dictionary. It can contain keys such as "Lz", "Z\_2", and "PH" to specify specific quantum numbers for the Hilbert space.

- `cdw`: (invalid) An unspecified parameter, possibly intended for future use.
- `dtype`: (invalid) An unspecified parameter, possibly intended for future use.
- `t`: Hopping term used in the computation.

Note: The comments inside the `__init__` function provide explanations for each parameter's purpose and usage.

```
1 eigvsL2_p = sol.eigsv(k=5,tol=10**(-6))
2 '''
3     k: the number of states from Lanczos
4     tol: the tolerance of Lanczos
5     '''
```

## 2.2. Measure

```
1 eigvsL2_p = sol.eigsv(k=5,tol=10**(-6))
```

After running the Lanczos algorithm to solve for low-energy states, it will return a list composed of instances of the `FQHE_SPHERE.core.eigenvector.eigenvector` class.

```
1 >> type(eigvsL2_p[0])
2 FQHE_SPHERE.core.eigenvector.eigenvector
```

`FQHE_SPHERE.core.eigenvector.eigenvector` is a class that encapsulates an eigenvector along with various quantum numbers (energy,  $L_z$ ,  $L^2$ ,  $Z_2$ , PH). It allows direct access to these properties for inspection.

```
1 >> eigvsL2_p[0]
2 The eigenvector in sector: {'Lz': 0, 'Z_2': 1, 'PH': 1} with E =
  -1.8688683997894542, dimension = 31681.
3 measured qn: {'energy': -1.8688683997894506, 'L2': -5.064787549425714e-16}
4 '''
5 1. The eigenvector in sector: {'Lz': 0, 'Z_2': 1, 'PH': 1}
6    {'Lz': 0, 'Z_2': 1, 'PH': 1} the quantum number of the Hilbert space
7 2. E = -1.8688683997894542
8    eigenvalue of Hamiltonian  $\langle H^2 + a*(L^2-b)^2 \rangle$  or  $\langle H = H^2 + a*L^2 \rangle$ 
9 3. dimension = 31681
10    the dimension of Hilbert space
11 4. measured qn: {'energy': -1.8688683997894506, 'L2': -5.064787549425714e-16}
12    the measured quantum number
13    energy  $\langle H^2 \rangle$ 
14     $L^2$   $\langle L^2 \rangle$ 
15    '''
```

By using the `.vec` method, you can directly obtain the eigenvector in `numpy` array format.

```

1  >> eigvsL2_p[0].vec # vacuum
2  array([ 0.16680578,  0.0822094 , -0.0822094 , ..., -0.04158522,
3         0.04135224,  0.0822094 ])
4  >> np.dot( eigvsL2_p[0].vec, eigvsL2_p[0].vec)
5  0.9999999999999994
6  >> eigvsL2_p[3].vec # energy stress tensor
7  The eigenvector in sector: {'Lz': 0, 'Z_2': 1, 'PH': 1} with E =
   -0.6206768230208931, dimension = 31681.
8  measured qn: {'energy': -0.6206768230208921, 'L2': 5.999999999999992}
9  >> np.dot( eigvsL2_p[0].vec, eigvsL2_p[3].vec)
10 -4.85722573273506e-17

```

When selecting the parameter `itype="H2L2"` in `fqhe.fermion.d1_solver`, it will automatically compute all quantum numbers and save them in the result class.

Alternatively, you can also use class methods to perform the calculation.

```

1
2  class eigenvector():
3      .
4      .
5      .
6      @property
7      def density(self):
8          return self.fqhe.density(self)
9
10     @property
11     def sigma_x(self):
12         return self.fqhe.sigma_x(self)
13
14     @property
15     def Z2(self):
16         if self.qn.get("Z_2")!=None:
17             return self.qn.get("Z_2")
18         elif self.measured_qn.get("Z_2")!=None:
19             return self.measured_qn.get("Z_2")
20         resu = self.fqhe.Z2(self)
21         self.add_measured_qn({"Z_2":resu})
22         return resu
23
24     @property
25     def PH(self):
26         if self.qn.get("PH")!=None:
27             return self.qn.get("PH")
28         elif self.measured_qn.get("PH")!=None:
29             return self.measured_qn.get("PH")
30         resu = self.fqhe.PH(self)
31         self.add_measured_qn({"PH":resu})
32         return resu

```

The calculation of  $L^2$  can be a bit complicated. To compute  $L^2$  in `fqhe.fermion.dl_solver`, choose the parameter `itype="L2"`, and then use a class method to obtain the sparse matrix representation of the  $L^2$  operator.

```
1 >> kk = fqhe.fermion.QHFM(n,n,num_thread=8)
2 >> sol = fqhe.fermion.dl_solver(kk, intra,intra,inter, "L2",qn=
  {"Lz":Lz,"Z_2":Z2,"PH":PH})
3 >> L2 = sol.Hamiltonian() # get the matrix form of L^2
4 >> T = eigvsL2_p[3].vec
5 >> L2.matvec(T)@T # <T|L^2|T>
6 5.999999999999992
```

## 3. Density operator

### 3.1. $\hat{n}_{l,m}^A (m \neq 0)$

Def

$$\begin{aligned}\hat{n}_{l,m}^A &= \int d\Omega Y_{l,m}^*(\Omega) \hat{n}^A(\Omega) \\ &= (2s+1) \sqrt{\frac{2l+1}{4\pi}} \sum_{m_1} (-1)^{3s+m_1+l} \begin{pmatrix} s & l & s \\ -s & 0 & s \end{pmatrix} \begin{pmatrix} s & l & s \\ -m_1 & m & m_1-m \end{pmatrix} \hat{c}_{m_1}^\dagger A \hat{c}_{m_1-m} \\ &= \sqrt{\frac{(2s+1)^2}{4\pi(2l+1)}} \sum_{m_1} (-1)^{3s+l+m+m_1} \langle s, s; s, -s | s, s; l, 0 \rangle \langle s, m_1-m; s, -m_1 | s, s; l, -m \rangle \hat{c}_{m_1}^\dagger A \hat{c}_{m_1-m} \\ &= \sum_{j_1, j_2} \mathcal{C}_{j_1, j_2} \hat{c}_{j_1}^\dagger A \hat{c}_{j_2} \delta_{j_1, j_2+m}\end{aligned}$$

and coefficient

$$\mathcal{C}_{j_1, j_2} = \sqrt{\frac{(2s+1)^2}{4\pi(2l+1)}} (-1)^{3s+l+m+j_1} \langle s, s; s, -s | l, 0 \rangle \langle s, j_2; s, -j_1 | l, -m \rangle \delta_{j_1, j_2+m}$$

Since the operator  $\hat{n}_{l,m}^A (m \neq 0)$  change the QN, [see next Sec. for more detailed usage.](#)

The recommended standard usage is

```
1 kk = fqhe.fermion.QHFM(n,n,dtype=np.float64, num_thread=8)
2 kk.Hilbertspace(qn={"Lz":0})
3 psi1 = kk.nlmAOp(tt[0].vec, l=4, m=-3, A=np.array([[1.0,2],[3,4]]), qn={"Lz":0})
```

and the definition of this interface

```

1 def nlmAOp(self, vec, l=0, m=0, A=np.array([[1.0,0],[0,-1]]), qn={"Lz":0}):
2     '''
3     vec: input state |a>
4     l, m: parameters in operator(definition)
5     A: 2*2 matrix
6     qn: the quantum number of the input state |a>
7     '''

```

Measure  $L^2$

```

1 print(np.linalg.norm(psi1))
2 psi1 /= np.linalg.norm(psi1) # normalization
3 L2 = kk.L2H(qn={"Lz":-3}) # L^2 m=-3
4 L2.matvec(psi1)@psi1/(np.linalg.norm(psi1)**2) # measur L^2
5 >> 20.0

```

## 4. Pairing Operator

### 4.1. Definition

$$\Delta_{l,m}^A = \sum_{m_1} (-1)^{m_1} \sqrt{2l+1} \begin{pmatrix} s & s & l \\ m_1 & m-m_1 & -m \end{pmatrix} c_{m-m_1} A c_{m_1} = \langle s, m_1; s, m-m_1 | l, m \rangle c_{m-m_1} A c_{m_1},$$

$$(\Delta_{l,m}^A)^\dagger = \sum_{m_1} (-1)^{m_1} \sqrt{2l+1} \begin{pmatrix} s & s & l \\ m_1 & m-m_1 & -m \end{pmatrix} c_{m_1}^\dagger A^\dagger c_{m-m_1}^\dagger = \langle s, m_1; s, m-m_1 | l, m \rangle c_{m_1}^\dagger A^\dagger c_{m-m_1}^\dagger$$

and

$$\begin{aligned}
 (\mathcal{O}_{l,m}^{l_a, l_b})^{AB} &= \sum_{m_a, m_b} (-1)^{m_b} \begin{pmatrix} l_a & l_b & l \\ m_a & -m_b & -m \end{pmatrix} (\Delta_{l_a, m_a}^A)^\dagger (\Delta_{l_b, m_b}^B) \\
 &= \sum_{m_a, m_b} \frac{(-1)^{l_a - l_b + m_b + m}}{\sqrt{2l+1}} \langle l_a, m_a; l_b, -m_b | l, m \rangle (\Delta_{l_a, m_a}^A)^\dagger (\Delta_{l_b, m_b}^B)
 \end{aligned}$$

### 4.2. $\Delta_{l,m}^A$ and $(\Delta_{l,m}^A)^\dagger$

First, the operator  $\Delta_{l,m}^A$  transforms Hilbert space from  $\mathcal{H}(N_e, L_z)$  to  $\mathcal{H}(N_e - 2, L_z - m)$  and  $(\Delta_{l,m}^A)^\dagger$  transforms from  $\mathcal{H}(N_e, L_z)$  to  $\mathcal{H}(N_e + 2, L_z + m)$ . For example, if  $|a\rangle \in \mathcal{H}(N_e, L_z)$  and

$$|b\rangle = \Delta_{l,m}^A |a\rangle$$

we have  $|b\rangle \in \mathcal{H}(N_e - 2, L_z - m)$ .

Second, the fermion(boson) statistics require that the operator  $\Delta_{l,m}^A$  vanish if  $2s + 1 + l$  is odd(even).

Since the change of particle number  $N_e$  and  $L_z$  is not considered in the python-cpp interface, we should carefully call this interface.

The recommended standard usage is



```

1 kk = fqhe.fermion.QHFM(n,n,dtype=np.float64, num_thread=8)
2 kk.Hilbertspace(qn={"Lz":0})
3 psi1 = kk.PairingAlmOp(tt[0].vec, l=2, m=-2, A=np.array([[1.0,0],[0,0]]),
4 daggerQ=0, qn={"Lz":0})
5 print(np.linalg.norm(psi1))

```

where the first two lines are used to set the Hilbert space in cpp code to  $\mathcal{H}(N_e = 12, L_z = 0)$ . The first line is used to **set the particle number** and the second is to set  $L_z$ .

The third line is the interface of operators  $\Delta_{l,m}^A$  and  $(\Delta_{l,m}^A)^\dagger$

```

1 def PairingAlmOp(self, vec, l=0, m=0, A=np.array([[1.0,0],[0,-1]]), daggerQ=0, qn=
2 {"Lz":0}):
3     '''
4     vec: input state |a>
5     l, m: parameters in operator(definition)
6     A: 2*2 matrix
7     daggerQ: 0 for  $\Delta_{l,m}^A$  and 1 for  $(\Delta_{l,m}^A)^\dagger$ 
8     qn: the quantum number of the input state |a>
9     '''

```

and the forth line is used to check the norm of result.

In this example, we set  $N_e = 12, L_z = 0, l = 2, m = -2$ , thus  $2s + 1 + l = 14$  is even and the result state is in Hilbert space  $\mathcal{H}(N_e = 10, L_z = 2)$ .

After calling this function, the code automatically set the Hilbert space in cpp code to  $\mathcal{H}(N_e = 10, L_z = 2)$ .

### Measure $L^2$

```

1 psi1 /= np.linalg.norm(psi1) # normalization
2 L2 = kk.L2H(qn={"Lz":2}) # L^2
3 L2.matvec(psi1)@psi1/(np.linalg.norm(psi1)**2) # measur L^2
4 >> 6.0000000000000036

```

Although the Hilbert space in cpp code has been set to the correct, we still need to tell the program the value of the angular momentum  $L_z = 2$ . (qn={"Lz":2}).

### Example for $(\Delta_{l,m}^A)^\dagger$

```

1 kk = fqhe.fermion.QHFM(n,n,dtype=np.float64, num_thread=8)
2 kk.Hilbertspace(qn={"Lz":0})
3 psi1 = kk.PairingAlmOp(tt[0].vec, l=4, m=3, A=np.array([[1.0,2],[3,4]]),
4 daggerQ=0, qn={"Lz":0})#daggerQ=1
5 print(np.linalg.norm(psi1))
6 psi1 /= np.linalg.norm(psi1) # normalization
7 L2 = kk.L2H(qn={"Lz":3}) # L^2 qn={"Lz":3}
8 L2.matvec(psi1)@psi1/(np.linalg.norm(psi1)**2) # measur L^2
9 >> 19.999999999999996

```

### 4.3. $(\mathcal{O}_{l,m}^{l_a,l_b})^{AB}$

The right pairing operator

$$\begin{aligned}
 (\mathcal{O}_{l,m}^{l_a,l_b})^{AB} &= \sum_{m_a, m_b} (-1)^{m_b} \begin{pmatrix} l_a & l_b & l \\ m_a & -m_b & -m \end{pmatrix} (\Delta_{l_a, m_a}^A)^\dagger (\Delta_{l_b, m_b}^B) \\
 &= \frac{(-1)^{l_a - l_b + m_b + m}}{\sqrt{2l+1}} \sum_{m_a, m_b} \langle l_a, m_a; l_b, -m_b | l_a, l_b; l, m \rangle (\Delta_{l_a, m_a}^A)^\dagger (\Delta_{l_b, m_b}^B) \\
 &= \frac{(-1)^{l_a - l_b + m_b + m}}{\sqrt{2l+1}} \sum_{m_a, m_b, m_1, m_2} \langle l_a, m_a; l_b, -m_b | l_a, l_b; l, m \rangle \langle s, m_1; s, m_a - m_1 | s, s; l_a, m_a \rangle \\
 &\quad \langle s, m_2; s, m_b - m_2 | s, s; l_b, m_b \rangle c_{m_1}^\dagger A^\dagger c_{m_a - m_1}^\dagger c_{m_b - m_2} B c_{m_2} \\
 &= \sum_{j_1, j_2, j_3, j_4} \mathcal{C}_{j_1, j_2, j_3, j_4} c_{j_1}^\dagger A^\dagger c_{j_2}^\dagger c_{j_3} B c_{j_4} \delta_{j_1 + j_2 - j_3 - j_4, m}
 \end{aligned}$$

and let

$$\begin{pmatrix} j_1 & \rightarrow & m_1 & m_1 & \rightarrow & j_1 \\ j_2 & \rightarrow & m_a - m_1 & m_2 & \rightarrow & j_4 \\ j_3 & \rightarrow & m_b - m_2 & m_a & \rightarrow & j_1 + j_2 \\ j_4 & \rightarrow & m_2 & m_b & \rightarrow & j_3 + j_4 \end{pmatrix}$$

we have the coefficient

$$\begin{aligned}
 \mathcal{C}_{j_1, j_2, j_3, j_4} &= \frac{(-1)^{l_a - l_b + m_b + m}}{\sqrt{2l+1}} \langle l_a, j_1 + j_2; l_b, -j_3 - j_4 | l_a, l_b; l, m \rangle \\
 &\quad \langle s, j_1; s, j_2 | s, s; l_a, j_1 + j_2 \rangle \langle s, j_4; s, j_3 | s, s; l_b, j_3 + j_4 \rangle \delta_{j_1 + j_2 - j_3 - j_4, m}
 \end{aligned}$$

```

1 kk = fqhe.fermion.QHFM(n,n,dtype=np.float64, num_thread=8)
2 kk.Hilbertspace(qn={"Lz":0})
3 psil = kk.PairingABlm1albp(tt[0].vec, l=4, m=-2, la=2, lb=2, A=np.array([[1.0,2],
  [3,4]]), B=np.array([[1.0,3],[5,7]]), qn={"Lz":0}) #

```

The third line is the interface of operators

```

1 def PairingABlm1albp(self, vec, l=0, m=0, la=0, lb=0, A=np.array([[1.0,0],
  [0,-1]]), B=np.array([[1.0,0],[0,-1]]), qn={"Lz":0}):
2     '''
3     vec: input state |a>
4     l, m, la, lb: parameters in operator(definition)
5     A, B: 2*2 matrix
6     qn: the quantum number of the input state |a>
7     '''

```

check the  $L^2$  of result

```
1 print(np.linalg.norm(psi1))
2 psi1 /= np.linalg.norm(psi1) # normalization
3 L2 = kk.L2H(qn={"Lz":-2}) # L^2
4 L2.matvec(psi1)@psi1/(np.linalg.norm(psi1)**2) # measur L^2
5 >> 20.0
```