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HOMEWORK 2: MONOPOLES

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1 Dirac

- (a) We are interested in the relation between the global properties of a manifold M and the structure of differential forms taking values on its cotangent bundle T^*M at each point of M .

Poincaré's lemma on $M = \mathbb{R}$: Let ω be a p -form ($p \in \{0, 1\}$) constructed from the cotangent space T^*M of M . Then $d\omega = 0$ (ω is closed) implies $\omega = d\lambda$ (ω is closed) where λ is a $(p-1)$ -form (0-form).

Proof: On \mathbb{R} , we can use the identity map as a global coordinate chart. The induced basis on 1-forms is $\{dx\}$ (a smooth frame field) and any 1-forms can be written as $\omega = gdx$ with $g \in C^\infty(\mathbb{R})$. Suppose now that ω is closed: we have $0 = d\omega = \partial_x g dx \wedge dx = 0, \forall g \in C^\infty(\mathbb{R})$ (ω being a 1-form is not restrictive, but would be for \mathbb{R}^n with $n > 1$). Then we take the 0-form $\lambda = G$ where G is any primitive of g ($G(x)$ exists because g is smooth) and apply an exterior derivative to get $d\lambda = gdx$. Because there are no $(0-1)$ -forms there is no need to check the lemma for 0-forms.

Counterexample: Consider the circle smooth manifold $S^1 \subset \mathbb{R}^2$ (embedded as $\{x^2 + y^2 = 1 | (x, y) \in \mathbb{R}^2\}$ for simplicity). It takes at least two charts to cover this manifold and, although on individual charts all closed 1-forms are exact (charts make the manifold look like \mathbb{R} locally), this property is lost globally. Choose the chart map $\theta = \arctan_2$ sending points (x, y) on the circle to their angle with the x axis excluding the point $(1, 0)$ so that the domain is open. With this chart we have the coordinate induced one form frame field $d\theta$ which we use to construct the closed form $\omega = d\theta$. On $(0, 2\pi)$, this form is exact since we have a 0-form $\lambda = F \in C^\infty((0, 2\pi))$ such that $\omega = d\lambda = \partial_\theta F d\theta = d\theta$ forcing $F = \theta + c, c \in \mathbb{R}$ since F has to be a primitive of 1 in the variable θ . The function F is smooth on the chart, but can never be extended to a smooth function over S^1 globally. Indeed, 0 and 2π being identified, a continuous function on S^1 should be constant at the excluded point $(0, 1)$ and this would require $\lim_{\theta \rightarrow 0^+} (\theta + c) = \lim_{\theta \rightarrow 2\pi^-} (\theta + c)$ which is impossible. Therefore there is a closed form on S^1 that is not exact.

- (b) Let $F^{(2)}$ be a 2-form on the 2-sphere S^2 . Suppose $F^{(2)}$ is globally exact implying there is a 1-form ω such that $F^{(2)} = d\omega$. Then we can use Stokes theorem in combination with the fact S^2 has no boundary to write $g = \frac{1}{4\pi} \in F^{(2)} = \frac{1}{4\pi} \int_{\partial S^2} d\omega = 0$.

- (c) Now working in Minkowski space $\mathbb{R}^{1,3}$ in the coordinate chart (t, r, θ, ϕ) built from spherical coordinates on \mathbb{R}^3 , we have the 2-form $F^{(4)} = Q \sin(\theta) d\theta \wedge d\phi$ with $Q \in \mathbb{R}$. We want to determine if $F^{(4)}$ satisfies Maxwell's equations $dF^{(4)} = 0, \quad d \star F^{(4)} = 0$. We have $dF^{(4)} = Q \cos(\theta) d\theta \wedge d\theta \wedge d\phi = 0$. To evaluate the Hodge dual of $F^{(4)}$

- (d)
(e)
(f)
(g)
(h)
(i)

2 Taub-NUT, or the gravitomagnetic monopole

- (a)
(b)
(c)
(d)