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## HOMework 2

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# 1 Two real scalars

- (a) We are considering here perturbative results in the quantum field theory of two interacting real massive scalars  $\varphi$ ,  $\Phi$  with respective masses  $m$  and  $M$ . The lagrangian density describing this theory is

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \varphi)^2 - \frac{1}{2}m^2 \varphi^2 + \frac{1}{2}(\partial_\mu \Phi)^2 - \frac{1}{2}M^2 \Phi^2 - \frac{g}{2!1!} \Phi \varphi^2$$

where  $g$  describes the coupling of the fields and is the parameter of our perturbative expansion. The position-space Feynman rules for perturbative computation of the interacting vacuum  $|\Omega\rangle$   $n$ -point functions  $\langle \Omega | T \varphi(x_1) \cdots \Phi(x_k) \cdots \Phi(x_n) | \Omega \rangle$  for this theory are summed up graphically below:

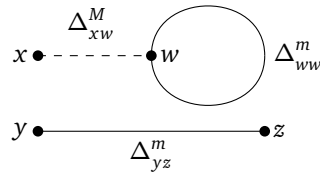


1. Every vertex in a diagram is associated to a four-position variable  $x$ . Its contribution to the symbolic representation of the amplitude is the integral  $-ig \int d^4x$  acting on the propagators from the  $x$  vertex to other vertices.

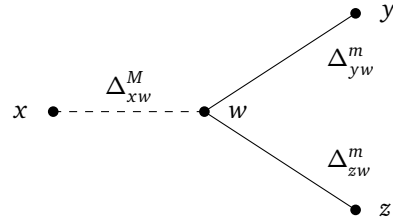
3. Divide the amplitude by the symmetry factor of the diagram  $S = 2/W$  where  $W$  is the number of wick contraction producing the same symbolic diagram expression.  $S$  is computed graphically as the order of the fixed end-points automorphism group of the diagram.

2. Each vertex is the source of two full lines and a dashed line free Klein-Gordon respectively representing a  $\varphi$  Feynman propagator  $\Delta_F^m(x-y)$  and a  $\Phi$  Feynman propagator  $\Delta_F^M(x-y)$  between points  $x$  and  $y$  (vertices or external points  $x_1 \cdots x_n$  of the expanded  $n$ -point function). Each edge of the diagram is symbolically represented as a multiplication by its associated Feynman propagator. In scalar field theory the external points contribute a trivial factor of 1 to the amplitude of the diagram.

- (b) The three-point function  $G(x, y, z) = \langle \Omega | T \Phi(x) \varphi(y) \varphi(z) | \Omega \rangle$  has no  $O(1)$  contributions because the number of contracted fields is odd at this order and no full Wick contractions can be formed: the vacuum is annihilated. At  $O(g)$ , we have the diagrams



$$A = -\frac{ig}{2} \int d^4w \Delta^M(x-w) \Delta^m(w-w) \Delta^m(y-z) \quad B = -ig \int d^4w \Delta^M(x-w) \Delta^m(w-y) \Delta^m(w-z)$$



Where the symmetry factor for the left diagram gains a factor of 2 from exchanging the endpoints of the internal loop. We have the amplitude

$$G(x, y, z) = -ig \int d^4w \Delta^M(x-w) \Delta^m(w-y) \Delta^m(w-z) - \frac{ig}{2} \int d^4w \Delta^M(x-w) \Delta^m(w-w) \Delta^m(y-z) + O(g^2).$$

- (c) i. The fourier transform of the previously computed  $A, B$  contribution to the position-space  $n$ -point function is

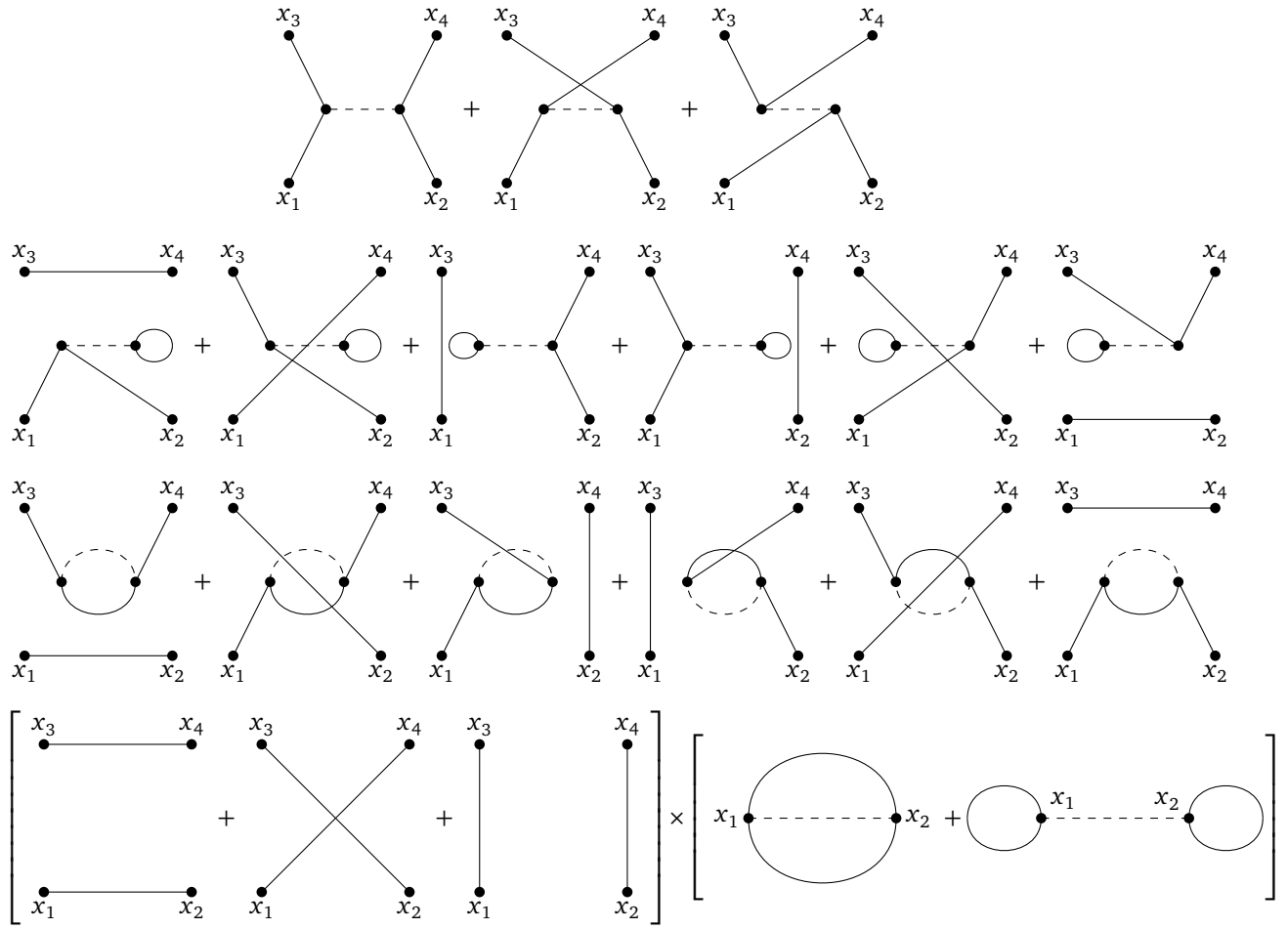
$$\tilde{G}_A(p_1, p_2, p_3) =$$

$$\tilde{G}_B(p_1, p_2, p_3) =$$

ii.

- (d)

(e)



(f)

(g)

(h)

## 2 Acknowledgement

Work in progress, but thanks to Thiago and ChatGPT for help drawing the diagrams