Lab Assignment 14 Semi-Lagrangian Methods and Strang Splitting

ACS II

Spring 2020

Assigned April 16, 2020 Due April 23, 2020

Semi-Lagrangian Methods

Consider the 1D linear advection equation

$$\frac{\partial c}{\partial t} + \frac{\partial}{\partial x}(u(x,t)c) = 0,$$

where u(x,t) is a given velocity field, and c(x,t) is the concentration of some species. This can be solved through a variety of finite difference/finite element approaches. There also exists an alternative approach to this problem, where one seeks to track not the flow in/out of a control volume, but the path of a specific fluid parcel in time. This gives rise to the semi-Lagrangian approach. The semi-Lagrangian approach is split up into two components for each time-step.

- For each node i at time step k, integrate back along the characteristic curve to time step k-1, identifying the location of the fluid parcel that ends up at node i, x_i^{k-1}
- \bullet Interpolate for the scalar quantity $c(x_i^{k-1},t^k)$

The integration step can be accomplished by any ODE integrator, such as forward Euler or any scheme in the Runge-Kutta method. The interpolation step can be handled by a polynomial interpolant.

Algorithm 1: Algorithm for a single time step of the semi-Lagrangian method.

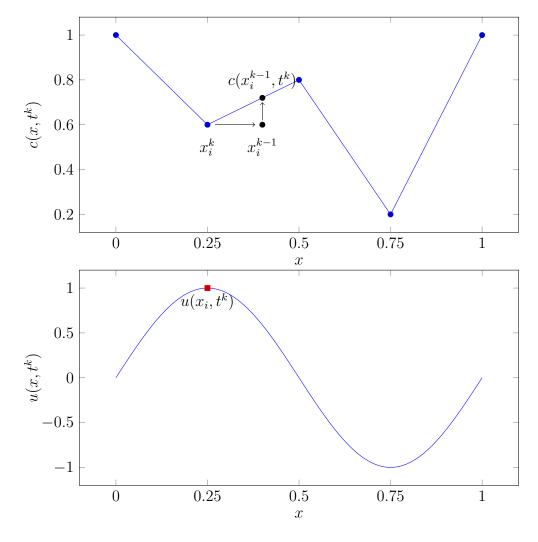


Figure 1: Top: $c(x, t^k)$. Bottom: given $u(x, t^k)$.

Exercise I

For the 1D linear advection problem

$$\frac{\partial c}{\partial t} + u(x,t)\frac{\partial c}{\partial x} = 0, \qquad 0 \le x \le 1, \ 0 \le t \le 8$$

Implement the semi-Lagrangian method for this problem with $c_0(x) = \sin^2(\pi x)$, u(x,t) = 1 and with periodic boundary conditions. Use the forward Euler method with $\Delta t = 0.01$ for the integration step and a linear interpolation scheme for the interpolation step. Estimate the spatial order of accuracy by doing a grid order convergence study, and discuss the observed accuracy of the method.

Operator Splitting

Consider the advection-diffusion equation

$$\frac{\partial c}{\partial t} + \frac{\partial}{\partial x}(u(x,t)c) - \alpha \frac{\partial^2 c}{\partial x^2} = 0.$$

One way to solve this is to use a combination of methods discussed earlier in the course. In particular we will use centered in space-backward in time finite difference solve to solve the diffusion equation $c_t - c_{xx} = 0$. In the above section of this lab you've solved the advection equation $c_t + (u(x,t)c)_x = 0$ using the semi-Lagrangian method. In a seminal 1968 paper, Gilbert Strang introduced the notion of operator splitting, where separate operators in a full PDE are decoupled and solved separately, with there contributions from each being added together. For the advection-diffusion equation, letting $\mathcal{L}_A(c) = (u(x,t)c)_x$ and $\mathcal{L}_D(c) = -\alpha c_{xx}$, the second-order operator splitting method is given as.

- Solve $c_t + \mathcal{L}_A(c)$ for $\Delta t/2$
- Solve $c_t + \mathcal{L}_D(c)$ for Δt
- Solve $c_t + \mathcal{L}_A(c)$ for $\Delta t/2$

Centered in Space-Backwards in Time

To solve the diffusion equation, you will implement a second order centered finite difference solver in space and combine it with first order backwards Euler in time. Recall that on a

standard uniform grid with spacing h, we can discretize c_{xx} as

$$c_{xx}(x_i) \approx \frac{C_{i-1} - 2C_i + C_{i+1}}{h^2}.$$
 (1)

Term c_t can be discretized using a first order Euler approximation,

$$c_t(x_i) \approx \frac{C^{k+1} - C_i^k}{\Delta t}. (2)$$

Combining (2) and using the implicit time level k+1 in (1), gives

$$-\alpha \Delta t C_{i-1}^{k+1} + (2\alpha \Delta t + h^2) C_i^{k+1} - \alpha \Delta t C_{i+1}^{k+1} = h^2 C_i^k.$$
 (3)

Ignoring boundary conditions for the time being, this is a tridiagonal system that must be solved at each time step,

$$\begin{pmatrix} 1 + 2\lambda & -\lambda & & & \\ -\lambda & 1 + 2\lambda & -\lambda & & & \\ & & \ddots & & \\ & & -\lambda & 1 + 2\lambda & -\lambda \\ & & & -\lambda & 1 + 2\lambda \end{pmatrix} \cdot \begin{pmatrix} C_0^{k+1} \\ C_1^{k+1} \\ \vdots \\ C_{N-1}^{k+1} \\ C_N^{k} \end{pmatrix} = \begin{pmatrix} C_0^k \\ C_1^k \\ \vdots \\ C_{N-1}^k \\ C_N^k \end{pmatrix},$$

where $\lambda = \alpha \Delta t/h^2$. This system can be solved efficiently using the Thomas algorithm. If we have periodic boundary conditions however, the system is no longer tridiagonal,

$$\begin{pmatrix}
1+2\lambda & -\lambda & 0 & \dots & -\lambda \\
-\lambda & 1+2\lambda & -\lambda & & & \\
0 & & \ddots & & 0 \\
& & -\lambda & 1+2\lambda & -\lambda \\
-\lambda & 0 & \dots & -\lambda & 1+2\lambda
\end{pmatrix} \cdot \begin{pmatrix}
C_0^{k+1} \\
C_1^{k+1} \\
\vdots \\
C_{N-1}^{k+1} \\
C_N^{k+1} \\
\end{pmatrix} = \begin{pmatrix}
C_0^k \\
C_1^k \\
\vdots \\
C_{N-1}^k \\
C_N^k \\
\end{pmatrix} .$$
(4)

We can no longer apply the Thomas algorithm directly, however we can still solve this system in O(N) operations using the Sherman-Morrison formula:

$$(A + \mathbf{u} \otimes \mathbf{v})^{-1} = A^{-1} - \frac{(A^{-1}\mathbf{u}) \otimes (\mathbf{v}^T A^{-1})}{1 + \mathbf{v}^T A^{-1}\mathbf{u}}.$$
 (5)

We are trying to solve the system $(A + \mathbf{u} \otimes \mathbf{v})\mathbf{x} = \mathbf{b}$. Let $\mathbf{y} = A^{-1}\mathbf{b}$ and $\mathbf{z} = A^{-1}\mathbf{u}$. Then

(5) gives:

$$\mathbf{x} = \left(A^{-1} - \frac{(A^{-1}\mathbf{u}) \otimes (\mathbf{v}^T A^{-1})}{1 + \mathbf{v}^T A^{-1} \mathbf{u}}\right) \mathbf{b}$$

$$= A^{-1}\mathbf{b} - \frac{(A^{-1}\mathbf{u}) \otimes (\mathbf{v}^T A^{-1})}{1 + \mathbf{v}^T A^{-1} \mathbf{u}} \mathbf{b}$$

$$= \mathbf{y} - \frac{\mathbf{z}(\mathbf{v} \cdot \mathbf{y})}{1 + \mathbf{v} \cdot \mathbf{z}}$$

The algorithm to compute \mathbf{x} proceeds as follows:

- 1. solve $A\mathbf{y} = \mathbf{b}$ using Thomas algorithm
- 2. solve $A\mathbf{z} = \mathbf{u}$ using Thomas algorithm

3.
$$\mathbf{x} = \mathbf{y} - \mathbf{z}(\mathbf{v} \cdot \mathbf{y})/(1 + \mathbf{v} \cdot \mathbf{z})$$

It remains to define \mathbf{u} , \mathbf{v} and A. The matrix in (4) can be replaced by $A + \mathbf{u} \otimes \mathbf{v}$, where

$$A = \begin{pmatrix} 2\lambda & -\lambda \\ -\lambda & 1 + 2\lambda & -\lambda \\ & \ddots & \\ & -\lambda & 1 + 2\lambda & -\lambda \\ & & -\lambda & 1 + 2\lambda - \lambda^2 \end{pmatrix}, \quad \mathbf{v} = \mathbf{u} = \begin{pmatrix} 1 \\ 0 \\ \vdots \\ 0 \\ -\lambda \end{pmatrix}.$$

Exercise II

Implement the Strang splitting method using semi-Lagrangian for the advection term and centered space, backwards in time for the diffusion term to solve the advection-diffusion equation

$$\frac{\partial c}{\partial t} + \frac{\partial}{\partial x}(u(x,t)c) - \alpha \frac{\partial^2 c}{\partial x^2} = 0, \qquad 0 \le x \le 1, \ 0 \le t \le 8,$$

with $c_0(x) = \sin^2(\pi x)$, u(x,t) = 1, $\alpha = 0.01$, with periodic boundary conditions. Plot your results at t = 0, 2, 4, 8 and discuss the behavior of your solution.

Submission and Grading

To get credit for this assignment, you must upload the following information to Canvas 11:59pm, April 23, 2020

- your source code files
- Your report as a single file in pdf format, including results from your work and relevant discussion of your observations, results, and conclusions.

This information must be received by 11:59pm, April 23, 2020. Upload the required documents to Canvas. As stated in the course syllabus, late assignment submissions will be subject to a 10% point penalty per 24 hours past the due date at time of submission, to a maximum reduction of 50%, according to the formula:

 $[\mathit{final\ score}] = [\mathit{raw\ score}] \text{ - } \mathit{min}(\text{ 0.5, 0.1 * [\# of\ days\ past\ due}]\text{) * [maximum\ score]}$