

1.1

let

$(\epsilon_{\mu\nu\rho} l^{(\mu} n^{\nu)}) + (\epsilon_{\mu\nu\rho} m^{(\mu} \bar{m}^{\nu)})$ then $(\epsilon_{\mu\nu\rho} l^{(\mu} n^{\nu)}) - (\epsilon_{\mu\nu\rho} m^{(\mu} \bar{m}^{\nu)}) \neq 0$. But

$$\epsilon_{\mu\nu\rho} (l^{(\mu} n^{\nu)} - m^{(\mu} \bar{m}^{\nu)}) = -\frac{1}{2} \epsilon_{\mu\nu\rho} (-2l^{(\mu} n^{\nu)} + m^{(\mu} \bar{m}^{\nu)}) = -\frac{1}{2} \epsilon_{\mu\nu\rho} 0 = 0$$

$$\text{thus } \epsilon_{\mu\nu\rho} l^{(\mu} n^{\nu)} = \epsilon_{\mu\nu\rho} m^{(\mu} \bar{m}^{\nu)}$$

1.2)

$$\psi_0 = \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} m^{\delta} \rightarrow \epsilon_{\mu\nu\rho} l^{\mu} (m^{\nu} + \epsilon a l^{\nu}) l^{\sigma} (m^{\delta} + \epsilon a l^{\delta})$$

$$= \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} m^{\delta} + \epsilon_{\mu\nu\rho} l^{\mu} (\epsilon a l^{\nu}) l^{\sigma} m^{\delta}$$

$$\epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} \epsilon a l^{\delta} + o(\epsilon^2)$$

$$\epsilon_{\mu\nu\rho} l^{\mu} l^{\nu} = -\epsilon_{\mu\nu\rho} l^{\mu} l^{\nu} = 0$$

We have shown $\psi_0 \rightarrow \psi_0 + o(\epsilon^2)$

$$\psi_1 = \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} n^{\delta} \rightarrow \epsilon_{\mu\nu\rho} l^{\mu} (m^{\nu} + \epsilon a l^{\nu}) l^{\sigma} (n^{\delta} + \epsilon a^* m^{\delta} + \epsilon a \bar{m}^{\delta})$$

$$= \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} n^{\delta} + \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} \epsilon a^* m^{\delta}$$

$$+ \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} \epsilon a \bar{m}^{\delta}$$

$$= \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} n^{\delta} + \epsilon_{\mu\nu\rho} l^{\mu} m^{\nu} l^{\sigma} (a^* m^{\delta} + a \bar{m}^{\delta})$$

$$g^{\mu\nu} = -2 l^{(\mu} n^{\nu)} + 2 m^{(\mu} \bar{m}^{\nu)} \quad (=) \quad g^{\theta\delta} = -2 l^{(\theta} n^{\delta)} + 2 m^{(\theta} \bar{m}^{\delta)}$$

$$g^{\theta\delta} l^\gamma = -2 l^{(\theta} n^{\delta)} l^\gamma + 2 m^{(\theta} \bar{m}^{\delta)} l^\gamma$$

$$g^{\theta\delta} l^\gamma = -l^\theta n^\delta l^\gamma - l^\delta n^\theta l^\gamma + m^\theta l^\delta \bar{m}^\gamma + m^\delta \bar{m}^\theta l^\gamma$$

$$\begin{aligned} \langle \epsilon_{\gamma\delta} l^\gamma g^{\theta\delta} l^\gamma \rangle &= -\langle \epsilon_{\gamma\delta} l^\gamma l^\theta n^\delta l^\gamma \rangle - \langle \epsilon_{\gamma\delta} l^\gamma l^\delta n^\theta l^\gamma \rangle \\ &\quad + \langle \epsilon_{\gamma\delta} l^\gamma m^\theta l^\delta \bar{m}^\gamma \rangle + \langle \epsilon_{\gamma\delta} l^\gamma \bar{m}^\theta l^\delta m^\gamma \rangle \end{aligned}$$

Note that $\langle \epsilon_{\gamma\delta} l^\gamma \bar{m}^\theta l^\delta m^\gamma \rangle = \langle \epsilon_{\gamma\delta} l^\gamma \bar{m}^\theta l^\gamma m^\delta \rangle \xrightarrow{\text{rename } \gamma \leftrightarrow \delta} \langle \epsilon_{\delta\gamma} l^\gamma m^\theta l^\delta \bar{m}^\gamma \rangle$

Note $\langle \epsilon_{\gamma\delta} l^\gamma g^{\theta\delta} l^\gamma \rangle = 0$ ϵ is the traceless part of $R_{\mu\nu\sigma}$

Moreover $\langle \epsilon_{\gamma\delta} \rangle = -\langle \epsilon_{\theta\gamma\delta} \rangle$ but $l^\alpha l^\sigma = l^\sigma l^\alpha$ so this contraction is 0

Finally $\langle \epsilon_{\gamma\delta} l^\gamma n^\theta l^\delta l^\gamma \rangle = -\langle \epsilon_{\gamma\delta} l^\gamma n^\theta l^\gamma l^\delta \rangle \xrightarrow[\text{rename } \gamma \leftrightarrow \delta]{} -\langle \epsilon_{\gamma\delta} l^\gamma n^\theta l^\delta l^\gamma \rangle = 0$

thus $\langle \epsilon_{\gamma\delta} l^\gamma m^\theta l^\delta \bar{m}^\gamma \rangle + \langle \epsilon_{\gamma\delta} l^\gamma \bar{m}^\theta l^\delta m^\gamma \rangle - 2 \langle \epsilon_{\gamma\delta} l^\gamma m^\theta l^\delta \bar{m}^\gamma \rangle = 0$

$$\therefore \psi_1 \rightarrow \psi_1 + \epsilon \alpha^* \psi_0$$

$$\begin{aligned} \psi_2 &= \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \bar{m}^\gamma n^\delta \rangle \rightarrow \langle \epsilon_{\gamma\delta} l^\gamma (m^\theta + \epsilon \alpha l^\theta) (\bar{m}^\gamma + \epsilon \alpha^* l^\gamma) (n^\delta + \epsilon \alpha^* m^\delta + \epsilon \alpha \bar{m}^\delta) \rangle \\ &= \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \bar{m}^\gamma n^\delta \rangle + \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \bar{m}^\gamma \alpha^* \epsilon n^\delta \rangle + \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \bar{m}^\gamma \epsilon \alpha \bar{m}^\delta \rangle \\ &\quad + \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \epsilon \alpha^* l^\gamma n^\delta \rangle \end{aligned}$$

\nearrow Symmetric $\bar{m}^\delta \bar{m}^\gamma$

$$\begin{aligned} \langle \epsilon_{\gamma\delta} l^\gamma m^\theta \bar{m}^\gamma n^\delta \rangle \quad g^{\theta\gamma} &= -2 l^{(\theta} n^{\gamma)} + 2 m^{(\theta} \bar{m}^{\gamma)} \\ \langle \epsilon_{\gamma\delta} l^\gamma g^{\theta\gamma} n^\delta \rangle &= -\langle \epsilon_{\gamma\delta} l^\gamma l^{(\theta} n^{\gamma)} n^\delta \rangle - \langle \epsilon_{\gamma\delta} l^\gamma n^{(\theta} l^{\gamma)} n^\delta \rangle \\ &\quad + \langle \epsilon_{\gamma\delta} l^\gamma m^{(\theta} \bar{m}^{\gamma)} n^\delta \rangle + \langle \epsilon_{\gamma\delta} l^\gamma \bar{m}^{(\theta} m^{\gamma)} n^\delta \rangle \end{aligned}$$

$$\therefore \text{Cov} \int l^a m^b m^{-b} m^d = \text{Cov} \int l^a n^b l^b m^d = \text{Cov} \int l^a m^b l^b n^d$$

$$\stackrel{a \rightarrow b}{=} \text{Cov} \int l^a m^b l^b n^d = \text{Cov} \int l^a m^b l^b n^d$$

$$\therefore \psi_2 \rightarrow \psi_2 + 2a^* \epsilon \psi_1$$

$$\begin{aligned} \psi_3 &= \text{Cov} \int l^a n^b \bar{m}^c n^d \rightarrow \text{Cov} \int l^a (n^b + a^* \epsilon m^b + a \epsilon \bar{m}^b) (\bar{m}^c + a^* \epsilon l^c) (n^d + a^* \epsilon m^d + a \epsilon \bar{m}^d) \\ &= \text{Cov} \int l^a n^b \bar{m}^c n^d + \text{Cov} \int l^a n^b \bar{m}^c a^* \epsilon m^d + \text{Cov} \int l^a n^b \bar{m}^c a \epsilon \bar{m}^d \\ &\quad + \text{Cov} \int l^a n^b a^* \epsilon l^c n^d + \text{Cov} \int l^a a^* \epsilon m^b \bar{m}^c n^d + \text{Cov} \int l^a a \epsilon \bar{m}^b \bar{m}^c n^d \end{aligned}$$

$$\text{Note } \text{Cov} \int l^a n^b l^c n^d = \text{Cov} \int l^a n^b l^c n^d + \underbrace{\text{Cov} \int l^a l^b n^c n^d}_{\text{this is 0 so we can add it with no problem}} = 2 \text{Cov} \int l^a l^b n^c n^d$$

focus on these terms

$$\text{Cov} \int l^a n^b \bar{m}^c n^d + 2 \text{Cov} \int l^a \bar{m}^b m^c n^d + \text{Cov} \int \cancel{l^a \bar{m}^b \bar{m}^c n^d}$$

$$g^a = -l^a n^b - n^a l^b + m^a \bar{m}^b + \bar{m}^a m^b$$

$$\begin{aligned} \text{Cov} \int g^a \bar{m}^b \bar{m}^c &= -\text{Cov} \int l^a \bar{m}^b \bar{m}^c n^d - \text{Cov} \int n^a \bar{m}^b \bar{m}^c l^d + \text{Cov} \int m^a \bar{m}^b \bar{m}^c \bar{m}^d \\ &\quad + \text{Cov} \int \bar{m}^a \bar{m}^b \bar{m}^c m^d \end{aligned}$$

$$\begin{aligned} \text{But } \text{Cov} \int n^a \bar{m}^b \bar{m}^c l^d &= \text{Cov} \int n^a \bar{m}^b \bar{m}^c l^d \\ &= \text{Cov} \int l^a \bar{m}^b \bar{m}^c n^d \end{aligned}$$

$$\therefore \text{Cov} \int l^a \bar{m}^b \bar{m}^c n^d = -\text{Cov} \int l^a \bar{m}^b \bar{m}^c n^d = 0$$

Now

$$\text{Cov} \int l^a n^b \bar{m}^c n^d + 2 \text{Cov} \int l^a \bar{m}^b m^c n^d = 0 \quad \text{Since}$$

$$[a, [b, c]] = 0$$

$$\text{and } (\epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma m^\delta + 2(\epsilon_{\alpha\beta\gamma\delta} l^\alpha \bar{m}^\beta m^\gamma n^\delta =$$

$$(\epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma m^\delta + \epsilon_{\alpha\beta\gamma\delta} l^\alpha \bar{m}^\beta m^\gamma n^\delta + \epsilon_{\alpha\beta\gamma\delta} l^\alpha m^\beta n^\gamma \bar{m}^\delta = \\ = (\epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma m^\delta + \epsilon_{\alpha\beta\gamma\delta} \bar{m}^\beta m^\gamma n^\delta l^\alpha + \epsilon_{\alpha\beta\gamma\delta} l^\alpha \bar{m}^\beta m^\gamma n^\delta = 0$$

So what is left is

$$2(\epsilon_{\alpha\beta\gamma\delta} l^\alpha m^\beta \bar{m}^\gamma n^\delta) a^\dagger \epsilon + (\epsilon_{\alpha\beta\gamma\delta} l^\alpha m^\beta \bar{m}^\gamma n^\delta) a^\dagger \epsilon + \psi_3$$

$$\therefore \psi_3 \rightarrow \psi_3 + 3a^\dagger \epsilon (\epsilon_{\alpha\beta\gamma\delta} l^\alpha m^\beta \bar{m}^\gamma n^\delta) = \psi_3 + 3a^\dagger \epsilon \psi_2$$

$$\psi_n = (\epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma n^\delta)$$

$$\psi_n \rightarrow (\epsilon_{\alpha\beta\gamma\delta} (n^\alpha + a^\dagger \epsilon m^\alpha + a \epsilon \bar{m}^\alpha) (\bar{m}^\beta + \epsilon a^\dagger l^\beta) (n^\gamma + a^\dagger \epsilon m^\gamma + a \epsilon \bar{m}^\gamma) (\bar{m}^\delta + a^\dagger \epsilon l^\delta) \\ = (\epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma n^\delta + \overset{(1)}{\epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma a^\dagger \epsilon l^\delta} + \epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta a^\dagger \epsilon m^\gamma \bar{m}^\delta \quad (3) \\ + \epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta a \epsilon \bar{m}^\gamma \bar{m}^\delta + \epsilon_{\alpha\beta\gamma\delta} n^\alpha \epsilon a^\dagger l^\beta m^\gamma \bar{m}^\delta + \epsilon_{\alpha\beta\gamma\delta} a^\dagger \epsilon m^\alpha \bar{m}^\beta m^\gamma n^\delta \quad (4) \\ + \epsilon_{\alpha\beta\gamma\delta} a \epsilon \bar{m}^\alpha m^\beta n^\gamma \bar{m}^\delta) \\ 0$$

$$(1) \epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma l^\delta = (\delta_{\alpha\beta} n^\alpha l^\delta) \bar{m}^\beta m^\gamma = (\delta_{\alpha\beta} a l^\delta) n^\alpha \bar{m}^\beta m^\gamma = \epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma n^\delta$$

$$(2) \epsilon_{\alpha\beta\gamma\delta} n^\alpha l^\beta m^\gamma \bar{m}^\delta = (\epsilon_{\alpha\beta\gamma\delta} l^\beta n^\alpha m^\gamma \bar{m}^\delta) = \epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma n^\delta$$

$$(3) \epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma a^\dagger \epsilon l^\delta = - \overset{\text{O symmetry}}{\epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma a^\dagger \epsilon l^\delta} + \overset{\text{O symmetry}}{\epsilon_{\alpha\beta\gamma\delta} n^\alpha l^\beta m^\gamma \bar{m}^\delta} + \overset{\text{O symmetry}}{\epsilon_{\alpha\beta\gamma\delta} n^\alpha n^\beta l^\gamma \bar{m}^\delta} \\ = \epsilon_{\alpha\beta\gamma\delta} n^\alpha l^\beta m^\gamma \bar{m}^\delta = (\epsilon_{\alpha\beta\gamma\delta} l^\beta n^\alpha m^\gamma \bar{m}^\delta) = \epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma n^\delta$$

$$(4) \epsilon_{\alpha\beta\gamma\delta} m^\alpha \bar{m}^\beta n^\gamma \bar{m}^\delta = (\delta_{\alpha\beta} m^\alpha \bar{m}^\delta) n^\gamma \bar{m}^\beta = \epsilon_{\alpha\beta\gamma\delta} n^\alpha \bar{m}^\beta m^\gamma \bar{m}^\delta = (3)$$

$$\therefore (3) + (4) = 2 \epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma n^\delta$$

$$\therefore (1) + (2) + (3) + (4) = 4 \underbrace{\epsilon_{\alpha\beta\gamma\delta} l^\alpha n^\beta \bar{m}^\gamma n^\delta}_{\psi_2} \text{ thus, } \psi_n \rightarrow \psi_n + a^\dagger \epsilon 4 \psi_2$$

$$c) \quad \psi_1 = (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}})$$

$$\begin{aligned} \psi_1 &\rightarrow (\epsilon \delta_{\gamma\delta} (l^a + \epsilon b^* m^a + \epsilon b \bar{m}^a) (m^b + \epsilon b n^b) (l^{\bar{c}} + \epsilon b^* m^{\bar{c}} + \epsilon b \bar{m}^{\bar{c}}) (m^{\bar{d}} + \epsilon b n^{\bar{d}})) \\ &= (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} \epsilon b n^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a m^b \epsilon b^* m^{\bar{c}} m^{\bar{d}} \\ &\quad + \epsilon \delta_{\gamma\delta} l^a m^b \epsilon b \bar{m}^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a \epsilon b n^b l^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} \epsilon b^* m^a m^b l^{\bar{c}} m^{\bar{d}} \\ &\quad + \epsilon \delta_{\gamma\delta} \epsilon b \bar{m}^a m^b l^{\bar{c}} m^{\bar{d}}) \end{aligned}$$

(1) (2) (3) (4)

$$(1) \quad (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}}) = \psi_1$$

$$\begin{aligned} (2) \quad (\epsilon \delta_{\gamma\delta} l^a m^b \bar{m}^{\bar{c}} m^{\bar{d}}) &= -(\epsilon \delta_{\gamma\delta} l^a m^b m^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a m^b n^{\bar{c}} m^{\bar{d}} \\ &\quad + \epsilon \delta_{\gamma\delta} l^a n^b l^{\bar{c}} m^{\bar{d}}) = \epsilon \delta_{\gamma\delta} l^{\bar{c}} m^{\bar{d}} l^a m^b = (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}}) \\ &= \psi_1 \end{aligned}$$

$$(3) \quad (\epsilon \delta_{\gamma\delta} l^a n^b l^{\bar{c}} m^{\bar{d}}) = \epsilon \delta_{\gamma\delta} l^{\bar{c}} m^{\bar{d}} l^a n^b = (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}}) = \psi_1$$

$$(4) \quad (\epsilon \delta_{\gamma\delta} \bar{m}^a m^b l^{\bar{c}} m^{\bar{d}}) = \epsilon \delta_{\gamma\delta} l^{\bar{c}} m^{\bar{d}} \bar{m}^a m^b = (\epsilon \delta_{\gamma\delta} l^a m^b \bar{m}^{\bar{c}} m^{\bar{d}}) = (2)$$

$$\therefore (1) + (2) + (3) + (4) = 4\psi_1$$

$$\text{thus } \psi_0 \rightarrow \psi_0 + \epsilon b \psi_1$$

$$\begin{aligned} \psi_1 &= (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}}) \rightarrow (\epsilon \delta_{\gamma\delta} (l^a + \epsilon b^* m^a + \epsilon b \bar{m}^a) (m^b + \epsilon b n^b) (l^{\bar{c}} + \epsilon b^* m^{\bar{c}} + \epsilon b \bar{m}^{\bar{c}}) m^{\bar{d}}) \\ &= (\epsilon \delta_{\gamma\delta} l^a m^b l^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a m^b \epsilon b^* m^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} l^a m^b \epsilon b \bar{m}^{\bar{c}} m^{\bar{d}} + \\ &\quad \epsilon \delta_{\gamma\delta} l^a \epsilon b n^b l^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} \epsilon b^* m^a m^b l^{\bar{c}} m^{\bar{d}} + \epsilon \delta_{\gamma\delta} \epsilon b \bar{m}^a m^b l^{\bar{c}} m^{\bar{d}}) \end{aligned}$$

(1) (2) (3) (4)

Symmetry

$$(1): \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta = \int d^4x m^b \gamma_n^\delta \lambda^a m^c = \int d^4x m^a n^b \lambda^c m^d \quad \text{0 symmetry}$$

$$= - \int d^4x m^a \lambda^b n^c \gamma_m^\delta + \int d^4x m^a \cancel{m^b \bar{m}^c} \gamma_m^\delta + \int d^4x m^a \bar{m}^b \cancel{\gamma_m^\delta} \quad \text{0 symmetry}$$

$$\therefore \int d^4x m^a n^b \lambda^c m^d = - \int d^4x m^a \lambda^b n^c m^d$$

$$\text{but } \int d^4x m^a n^b \lambda^c m^d = \int d^4x n^b m^a m^d \lambda^c \xrightarrow{\text{rename } a \rightarrow d} = \int d^4x n^b n^d m^a \lambda^c$$

$$= \int d^4x m^a \lambda^c n^b m^d$$

$$= \int d^4x m^a \lambda^b n^c m^d$$

$$\therefore \int d^4x m^a \lambda^b n^c m^d = - \int d^4x m^a \lambda^b n^c m^d = 0 \quad \text{thus } (1) = 0$$

$$(2): \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta = \psi_2 \quad \text{0 symmetry}$$

$$(3): \int d^4x \lambda^a n^b \lambda^c \gamma_n^\delta = - \int d^4x \lambda^a \cancel{\lambda^b} n^c \gamma_n^\delta + \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta \quad \psi_2$$

$$+ \int d^4x \lambda^a \bar{m}^b m^c \gamma_n^\delta$$

$$(4): \int d^4x \bar{m}^a m^b \lambda^c \gamma_n^\delta = - \int d^4x m^a \bar{m}^b \lambda^c \gamma_n^\delta + \int d^4x \lambda^a n^b \lambda^c \gamma_n^\delta \quad \text{this is (3)}$$

$$+ \int d^4x \lambda^a \cancel{\lambda^b} n^c \gamma_n^\delta$$

$$= - \int d^4x m^a \bar{m}^b \lambda^c \gamma_n^\delta + \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta \quad \text{0 symmetry}$$

$$+ \int d^4x \lambda^a \bar{m}^b m^c \gamma_n^\delta$$

$$(3) + (4)$$

$$\therefore \int d^4x \bar{m}^a m^b \lambda^c \gamma_n^\delta + 2 \int d^4x \lambda^a \bar{m}^b m^c \gamma_n^\delta = 0 \quad \text{Same as question b}$$

$$\text{thus } \psi_1 \rightarrow \psi_1 + 3 \epsilon b \psi_2$$

$$\psi_2 = \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta = \int d^4x (\lambda^a + b \epsilon^* m^a + b \epsilon \bar{m}^a) (m^b + \epsilon b n^b) (\bar{m}^c + \epsilon b^* n^c) \gamma_n^\delta$$

$$= \int d^4x \lambda^a m^b \bar{m}^c \gamma_n^\delta + \int d^4x \lambda^a m^b \epsilon b^* \gamma_n^\delta + \int d^4x \lambda^a \epsilon b n^b \bar{m}^c \gamma_n^\delta \quad \text{0 symmetry} \quad (1)$$

$$+ \int d^4x b \epsilon^* m^a m^b \bar{m}^c \gamma_n^\delta + \int d^4x b \epsilon \bar{m}^a m^b \bar{m}^c \gamma_n^\delta \quad (2)$$

$$(1): \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \psi_3$$

0 Symmetry

0 Symmetry

$$(2): \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = - \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

$$= \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \psi_3$$

$$(1) + (2) = 2\psi_3$$

$$\therefore \psi_2 \rightarrow \psi_2 + 2b\psi_3$$

$$\psi_3 = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d \rightarrow \int d^4x \delta^4(x) (\bar{\psi}^a + \epsilon b^* \bar{\psi}^a + \epsilon b \bar{\psi}^a) \gamma^\mu (\psi^b + \epsilon b^* \psi^b + \epsilon b \psi^b) \bar{\psi}^c \gamma^\nu \psi^d$$

$$= \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d (1)$$

$$+ \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d (2)$$

$$(1) \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = - \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d \stackrel{\text{rename}}{=} - \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

$$= \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d + \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

0 Symmetry 0 Symmetry

$$\therefore \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = - \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

$$\text{But } \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

$$\text{rename} = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d \therefore \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = - \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = 0$$

$$\therefore (1) = 0$$

$$(2): \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d$$

$$\text{rename} = \int d^4x \delta^4(x) \bar{\psi}^a \gamma^\mu \psi^b \bar{\psi}^c \gamma^\nu \psi^d = \psi_4$$

$$\therefore \psi_3 \rightarrow \psi_3 + \epsilon b \psi_4$$

$$\begin{aligned}
 \psi_1 &= (\not{a} \not{x} \not{p} \not{m} \not{\sigma} \not{n} \not{r} \not{m}) \rightarrow (\not{a} \not{x} \not{p} \not{n} (\not{m} \not{\sigma} + i \not{b} \not{n} \not{\sigma}) \not{n} \not{r} (\not{m} \not{\sigma} + i \not{b} \not{n} \not{\sigma})) \quad \text{symmetry} \\
 &= (\not{a} \not{x} \not{p} \not{n} \not{m} \not{\sigma} \not{n} \not{r} \not{m}) + (\not{a} \not{x} \not{p} \not{n} \not{a} i \not{b} \not{n} \not{\sigma} \not{n} \not{r} \not{m}) + (\not{a} \not{x} \not{p} \not{n} \not{m} \not{\sigma} \not{n} \not{r} i \not{b} \not{n} \not{\sigma}) \\
 &\quad \text{symmetry} \\
 &= (\not{a} \not{x} \not{p} \not{n} \not{m} \not{\sigma} \not{n} \not{r} \not{m}) = \psi_1 + O(\epsilon^2)
 \end{aligned}$$

$$\begin{aligned}
 d) \quad \psi_0 &= (\not{a} \not{x} \not{p} \not{l} \not{m} \not{\sigma} \not{l} \not{x} \not{m}) \rightarrow (\not{a} \not{x} \not{p} (\not{l} \not{a} - \epsilon \not{c} \not{l} \not{a}) (\not{m} \not{\sigma} + i \epsilon \not{d} \not{m} \not{\sigma}) (\not{l} \not{x} - \epsilon \not{c} \not{l} \not{x}) (\not{m} \not{\sigma} + i \epsilon \not{d} \not{m} \not{\sigma})) \\
 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} \not{m}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} i \epsilon \not{d} \not{m} \not{\sigma}) - (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \epsilon \not{c} \not{l} \not{x} \not{m}) \\
 &\quad + (\not{a} \not{x} \not{p} \not{l} \not{a} i \epsilon \not{d} \not{m} \not{\sigma} \not{l} \not{x} \not{m}) - (\not{a} \not{x} \not{p} \epsilon \not{c} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} \not{m}) \\
 &= \psi_0 + i \epsilon \not{d} \psi_0 - \epsilon \not{c} \psi_0 + i \epsilon \not{d} \psi_0 - \epsilon \not{c} \psi_0 = \psi_0 + 2(i \epsilon \not{d} - \epsilon \not{c}) \psi_0
 \end{aligned}$$

$$\begin{aligned}
 \psi_1 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} \not{n}) = (\not{a} \not{x} \not{p} (\not{l} \not{a} - \epsilon \not{c} \not{l} \not{a}) (\not{m} \not{\sigma} + i \epsilon \not{d} \not{m} \not{\sigma}) (\not{l} \not{x} - \epsilon \not{c} \not{l} \not{x}) (\not{n} \not{\sigma} + \epsilon \not{c} \not{n} \not{\sigma})) \\
 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} \not{n}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{l} \not{x} \epsilon \not{c} \not{n} \not{\sigma}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} (-\epsilon \not{c} \not{l} \not{x}) \not{n} \not{\sigma}) \\
 &\quad + (\not{a} \not{x} \not{p} \not{l} \not{a} i \epsilon \not{d} \not{m} \not{\sigma} \not{l} \not{x} \not{n}) + (\not{a} \not{x} \not{p} (-\epsilon \not{c} \not{l} \not{a}) \not{m} \not{\sigma} \not{l} \not{x} \not{n}) \\
 &= \psi_1 + \cancel{\epsilon \not{c} \psi_1} - \cancel{\epsilon \not{c} \psi_1} + i \epsilon \not{d} \psi_1 - \epsilon \not{c} \psi_1 = \psi_1 - \epsilon (\not{c} - i \not{d}) \psi_1
 \end{aligned}$$

$$\begin{aligned}
 \psi_2 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{m} \not{x} \not{n}) = (\not{a} \not{x} \not{p} (\not{l} \not{a} - \epsilon \not{c} \not{l} \not{a}) (\not{m} \not{\sigma} + i \epsilon \not{d} \not{m} \not{\sigma}) (\not{m} \not{\sigma} - i \epsilon \not{d} \not{m} \not{\sigma}) (\not{n} \not{\sigma} + \epsilon \not{c} \not{n} \not{\sigma})) \\
 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{m} \not{x} \not{n}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{m} \not{x} \epsilon \not{c} \not{n} \not{\sigma}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} (-i \epsilon \not{d}) \not{m} \not{x} \not{n}) \\
 &\quad + (\not{a} \not{x} \not{p} \not{l} \not{a} i \epsilon \not{d} \not{m} \not{\sigma} \not{m} \not{x} \not{n}) - \epsilon \not{c} (\not{a} \not{x} \not{p} \not{l} \not{a} \not{m} \not{\sigma} \not{m} \not{x} \not{n}) = \psi_2
 \end{aligned}$$

$$\begin{aligned}
 \psi_3 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{n} \not{\sigma} \not{m} \not{x} \not{n}) = (\not{a} \not{x} \not{p} (\not{l} \not{a} - \epsilon \not{c} \not{l} \not{a}) (\not{n} \not{\sigma} + \epsilon \not{c} \not{n} \not{\sigma}) (\not{m} \not{\sigma} - i \epsilon \not{d} \not{m} \not{\sigma}) (\not{n} \not{\sigma} + \epsilon \not{c} \not{n} \not{\sigma})) \\
 &= (\not{a} \not{x} \not{p} \not{l} \not{a} \not{n} \not{\sigma} \not{m} \not{x} \not{n}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{n} \not{\sigma} \not{m} \not{x} \epsilon \not{c} \not{n} \not{\sigma}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \not{n} \not{\sigma} (-i \epsilon \not{d}) \not{m} \not{x} \not{n}) + (\not{a} \not{x} \not{p} \not{l} \not{a} \epsilon \not{c} \not{n} \not{\sigma} \not{m} \not{x} \not{n}) \\
 &\quad + (\not{a} \not{x} \not{p} (-\epsilon \not{c} \not{l} \not{a}) \not{n} \not{\sigma} \not{m} \not{x} \not{n}) = \psi_3 + \epsilon \not{c} \psi_3 - i \epsilon \not{d} \psi_3 + \epsilon \not{c} \psi_3 - \cancel{\epsilon \not{c} \psi_3} = \psi_3 + \epsilon (\not{c} - i \not{d}) \psi_3
 \end{aligned}$$

$$\begin{aligned}
\psi_n &= (\epsilon \gamma^\delta l^\alpha \bar{m}^\beta n^\gamma m^\delta) = (\epsilon \gamma^\delta (n^\alpha + \epsilon c n^\alpha) (\bar{m}^\beta - i \epsilon d \bar{m}^\beta) (n^\gamma + \epsilon c n^\gamma) (\bar{m}^\delta - i \epsilon d \bar{m}^\delta)) \\
&= (\epsilon \gamma^\delta n^\alpha \bar{m}^\beta n^\gamma m^\delta) - (\epsilon \gamma^\delta n^\alpha \bar{m}^\beta n^\gamma i \epsilon d \bar{m}^\delta) + (\epsilon \gamma^\delta n^\alpha \bar{m}^\beta \epsilon c n^\gamma m^\delta) \\
&+ (\epsilon \gamma^\delta n^\alpha (-i \epsilon d \bar{m}^\beta) n^\gamma m^\delta) + \epsilon c (\epsilon \gamma^\delta n^\alpha \bar{m}^\beta n^\gamma m^\delta) \\
&= \psi_n - i \epsilon d \psi_n + \epsilon c \psi_n - i \epsilon d \psi_n + \epsilon c \psi_n = \psi_n + 2 \epsilon (c - i d) \psi_n
\end{aligned}$$

c)

First of all, as was discussed in class Kerr spacetime has two double principal null vectors and thus, one can always choose both l and n such that $\psi_0 = \psi_1 = \psi_3 = \psi_4 = 0$. Looking at the previous questions, one can see that rotations of ψ_0 combine only the background values (0th order in ϵ) of ψ_1 and ψ_0 into the perturbed value (1st order in ϵ) of ψ_0 . Thus one concludes that, since both the background values of ψ_0 and ψ_1 vanish, ψ_0 remains invariant under infinitesimal rotations.

One sees that rotations of ψ_1 (type II) and ψ_3 (type I) combine background values of ψ_2 , which doesn't vanish, into their perturbed values. Thus, these two cannot be invariant under infinitesimal rotations. Since the background value of ψ_2 is non vanishing, we do not consider it as invariant under infinitesimal rotations, and finally as for ψ_0 , infinitesimal rotations of ψ_4 mix only the background values of ψ_4 and ψ_3 into the perturbed value of ψ_4 , and since both vanish, ψ_4 remains invariant under infinitesimal rotations as well.

f)

Since the ψ fields transform as scalars and the background values of $\psi_0, \psi_1, \psi_3, \psi_4$ vanish, normal gauge transformations do not affect these scalars. Thus, they are gauge invariant.

g) $\psi_2 = \psi_2^{\text{Kerr}} + \delta\psi_2$ In class we saw how do gauge transformations work with the metric. for $g_{\mu\nu} = g'_{\mu\nu} + \epsilon h_{\mu\nu}$ then a gauge transformation

↑
perturbation

acted on $h_{\mu\nu}$ as $h_{\mu\nu} \rightarrow h_{\mu\nu} + \nabla_\mu \xi_\nu + \nabla_\nu \xi_\mu$

Similarly here, the gauge transformation will act on $\delta\psi$ like $\delta\psi \rightarrow \delta\psi + \mathcal{L}_\xi \psi^{\text{Kerr}}$ with $\mathcal{L}_\xi \psi^{\text{Kerr}} = \xi^\mu \partial_\mu \psi^{\text{Kerr}} + \xi^\alpha \omega_\alpha \psi^{\text{Kerr}}$. To have $\psi_2 = \psi_2^{\text{Kerr}}$, $\delta\psi_2$ has to vanish after the gauge transformation

$$\therefore \xi^r \partial_r \psi^{kur} + \xi^\theta \partial_\theta \psi^{kur} = -\nabla \psi \quad \text{assuming } \psi = \psi'(r, \theta)$$

$$\Rightarrow \xi^r \frac{3M}{(r - ia \cos \theta)^n} + \xi^\theta \frac{3ia M \sin \theta}{(r - ia \cos \theta)^n} = -\psi'(r, \theta)$$

$$\Rightarrow \xi^r \frac{3M (r + ia \cos \theta)^n}{(r^2 + a^2 \cos^2 \theta)^n} + \xi^\theta \frac{3ia M \sin \theta (r + ia \cos \theta)^n}{(r^2 + a^2 \cos^2 \theta)^n} = -\psi'(r, \theta)$$

Assuming $\xi^r, \xi^\theta, r, \theta, M, a$ to be real, we can find two equations for the 2 unknowns

ξ^r and ξ^θ By equating $\text{Re}[LHS] = \text{Re}[\psi'(r, \theta)]$ and $\text{Im}[LHS] = \text{Im}[\psi'(r, \theta)]$

See Mathematica for the solution

Problem 2)

$$q) h_+ = \frac{h}{\omega} \left(\frac{GM_z}{c^2} \right)^{5/3} \left(\frac{M_{JW}(t)}{c} \right)^{2/3} \left(\frac{1 + e - e^2 \cos \phi(t)}{2} \right) \cos(\phi(t))$$

$$h_- = \frac{h}{\omega} \left(\frac{GM_z}{c^2} \right) \left(\frac{M_{JW}(t)}{c} \right)^{2/3} \cos \phi(t) \sin \phi(t). \quad \text{See the mathematical notebook for the solution}$$

b) In the previous question we showed that $h(t) = \frac{h}{\omega} \left(\frac{GM_z}{c^2} \right)^{5/3} \left(\frac{M_{JW}(t)}{c} \right)^{2/3} \cos[\phi(t) + \alpha]$

for simplicity I will absorb α in the constant ϕ_0 of $\phi(t)$

$$\text{then we can write } A(t) = \frac{h}{\omega} \left(\frac{GM_z}{c^2} \right)^{5/3} \left(\frac{M_{JW}(t)}{c} \right)^{2/3} \rightarrow h(t) = A(t) \cos(\phi(t))$$

In lecture notes we calculated the Fourier transform of a function which has this form using the stationary phase approx: $\tilde{h} = \int dt A(t) \cos \phi(t) e^{2i\omega t}$

The stationary phase approx, assumes $\frac{d \ln A}{dt} \ll \dot{\phi}$, $\ddot{\phi} \ll (\dot{\phi})^2$. then we rewrite

$\cos \phi = \frac{1}{2} (e^{i\phi} + e^{-i\phi})$ and we keep only the first term since the second term cannot create a stationary point. For the first term we have $\dot{\phi}(t_k) = \omega$. finally we evaluate everything in retarded time since we are far enough to need here redshift

$$\tilde{h} \approx \frac{1}{2} e^{i\omega t} \int dt A(t_k) e^{i\omega t_k - i\phi(t_k)} \quad \text{Taylor expand around } t_k \text{ and use } \int dx e^{-ix^2} = \sqrt{\pi} e^{i\pi/4}$$

$$\text{to get } \tilde{h} = \frac{1}{2} A(t_k) \left(\frac{\omega}{\ddot{\phi}(t_k)} \right)^{1/2} e^{i(\omega t_k - \phi(t_k) - \pi/4 + \omega t_d/c)}$$

$$\text{now we have to relate } t_k \text{ and } t: \phi(t) = -2 \left(\frac{5GM_z}{c^3} \right)^{5/8} (t_c - t)^{3/8} + \phi_0$$

$$\dot{\phi}(t_*) = n\dot{\phi} \rightarrow -\mathcal{A} \left(\frac{5GM_c}{c^3} \right)^{-5/8} \cdot \frac{5}{8} (t_c - t)^{-3/8} = \mathcal{A} n \dot{\phi}$$

$$\rightarrow t_* - t_c = -\frac{5}{256 n \dot{\phi}} \frac{1}{8/3} \left(\frac{GM_c}{c^3} \right)^{-5/3}$$

$$\phi(t_*) = \phi_0 - \frac{c^5}{15 \dot{\phi}^{5/3} G^{5/3} M_c^{5/3} n^{5/3}}$$

the full phase is $\psi = n\dot{\phi}(t_* + \text{def}/c) - \phi(t_*) - n/\eta$ using $t_* = \underbrace{t_* - t_c}_{\text{I know this}} + t_c$

mathematically $\psi = n\dot{\phi}((t_* - t_c) + \frac{\text{def}}{c}) - \phi(t_*) - n/\eta + n\dot{\phi}t_c$

$$\psi = -\phi_0 + \frac{13 c^5}{128 \dot{\phi}^{5/3} G^{5/3} M_c^{5/3} n^{5/3}} - \frac{n}{\eta} - n\dot{\phi} \frac{\text{def}}{c} + n\dot{\phi} t_c$$

$$\psi = \psi_0 + n\dot{\phi} \left(t_c - \frac{\text{def}}{c} \right) + \frac{13 c^5}{128 \dot{\phi}^{5/3} (GM_c n)^{5/3}}$$

we see reference time to t_c
effective distance def
reference phase $\psi_0 = -\phi_0 - n/\eta$

$$\tilde{A}(t) = \frac{1}{2} A(t_*) \left(\frac{n}{\dot{\phi}(t_*)} \right)^{1/2}$$

for the rest of this problem
in $\phi: M_c \rightarrow M_z$

M_c should be M_z ?
or $M_c = \frac{M_z}{1+z}$ but then
we have another parameter z

$$= \frac{2}{\text{def}} \left(\frac{GM_z}{c^2} \right)^{5/3} \left(\frac{n \dot{\phi}_{\text{gw}}(t_*)}{c} \right)^{2/3} \left(\frac{n}{\dot{\phi}(t_*)} \right)^{1/2}$$

during the inspiral phase

$$= \frac{2}{\text{def}} \sqrt{\frac{5}{6}} \left(\frac{GM_z}{c^2} \right)^{5/3} \left(\frac{n \dot{\phi}_{\text{gw}}(t_*)}{c} \right)^{2/3} \frac{c^{5/2}}{4 \dot{\phi}^{1/6} (GM_z)^{5/6} n^{4/3}}$$

\uparrow
used to be M_c

$$\dot{\phi}(t) = \frac{1}{(1+z) n} \left(\frac{5}{256} \frac{1}{t_c - t} \right)^{3/4} \left(\frac{GM_c}{c^3} \right)^{-5/8}$$

$$= \frac{1}{n} \left(\frac{5}{256} \frac{1}{t_c - t} \right)^{3/4} \left(\frac{GM_z}{c^3} \right)^{-5/8}$$

and without $\tilde{A} = \left(\frac{5}{c^2} \right)^{1/4} \frac{(GM_z)^{5/6}}{\text{def} 2\sqrt{3} c^{3/2} n^{2/3}} \dot{\phi}^{-7/4}$

this Fourier expansion is valid until the system reaches its ISCO thus as long as
the emission is dominated by quadrupole radiation $f_{\text{max}} = 2f_{\text{ISCO}} = \frac{1}{3n\sqrt{6}} \frac{c^3}{GM_{\text{tot}}}$

$$1) \quad h = \underbrace{\left(\frac{5}{2} \right)^{1/6} \frac{G^{5/8}}{2\sqrt{3} \, c^{3/2} \rho^{1/2}} M_7^{5/8} f^{-7/6}}_{\sigma_1} \exp \left[i \left(\varphi_0 + \eta f \left(t_c - \frac{t_c f f}{c} \right) \right) + \frac{13 \, c^5 \, f^{-5/3}}{\underbrace{128 (6\pi)^{5/3}}_{\sigma_2} M_t^{5/3}} \right]$$

$$h = \sigma_1 M_7^{5/6} f^{-7/6} \exp \left[i \left(\varphi_0 + \eta f \left(t_c - \frac{t_c}{c} \right) + \sigma_2 f^{-5/3} M_t^{-5/3} \right) \right]$$

See mathematica