

1 Killing-Yano tensors

An anti-symmetric rank-2 tensor $\omega_{\mu\nu} = -\omega_{\nu\mu}$ is called a **Killing-Yano tensor** if it satisfies

$$\nabla_\alpha \omega_{\beta\gamma} + \nabla_\beta \omega_{\alpha\gamma} = 0.$$

- a.** Show that for any Killing-Yano tensor $\omega_{\alpha\beta}$, the symmetric tensor $\omega_{\alpha\gamma}\omega_{\beta}{}^{\gamma}$ is a Killing tensor. (In some sense Killing-Yano tensors can be thought of as the “square root” of a Killing tensor.)
- b.** Show that for any Killing-Yano tensor $\omega_{\mu\nu}$ and geodesic $x^{\mu}(s)$, the (co-)vector $V_{\mu} = \omega_{\mu\alpha} \frac{dx^{\alpha}}{ds}$ is parallel transported only the geodesic, i.e.

$$\frac{dx^\alpha}{ds} \nabla_\alpha V^\mu = 0.$$

Note: Kerr also has a Killing-Yano tensor.

$$a) \quad \omega_s^{\sigma} = g^{\sigma\epsilon} \omega_{\epsilon s} \quad \rightarrow \quad \nabla_a \omega_s^{\sigma} + \nabla_s \omega_a^{\sigma} = g^{\sigma\epsilon} \nabla_a \omega_{\epsilon s} + g^{\sigma\epsilon} \nabla_s \omega_{\epsilon a} \\ = g^{\sigma\epsilon} (\nabla_a \omega_{\epsilon s} + \nabla_s \omega_{\epsilon a}) = 0$$

Thus $\nabla_a (w_{bc} w^c) + \nabla_b (w_{ac} w^c) + \nabla_c (w_{ab} w^c) + \nabla_c (w_{bc} w^a)$
 $+ \nabla_d (w_{ad} w^c) + \nabla_d (w_{cd} w^a)$

$$\begin{aligned} & \therefore \cancel{\omega_f^6} (\cancel{\nabla_a} L_{\phi 6} + \nabla_6 \omega_a) + (\cancel{\omega_{\phi 6}} \cancel{\nabla_a} \omega_f^6 + \omega_{a6} \cancel{\nabla_6} \omega_f^6) \\ & + \omega_{\phi 6} (\cancel{\nabla_a} \omega_{\phi 6}) + \omega_a^6 (\cancel{\nabla_6} \omega_{\phi 6}) + \omega_{f6} (\cancel{\nabla_a} \omega_{\phi 6}^6 + \nabla_6 \omega_a^6) \\ & + \cancel{\omega_{\phi 6}} \cancel{\nabla_f} L_{a6} + \omega_{a6} \cancel{\nabla_f} \omega_{\phi 6} + \cancel{\omega_a^6} \cancel{\nabla_f} \omega_{\phi 6} + \omega_{\phi 6} \cancel{\nabla_f} \omega_a^6 = 0 \end{aligned}$$

e) $N_D^P \frac{dx^q}{ds} \nabla_a V^{\mu} = 0 \quad V^{\mu} = \omega_{\mu q} \frac{dx^q}{ds}$

$$\nabla_\sigma V^\mu = \nabla_\sigma \left(w_{\mu\alpha} \frac{dx^\alpha}{ds} \right) = \left(\nabla_\sigma w_{\mu\alpha} \right) \frac{dx^\alpha}{ds} + w_{\mu\alpha} \left(\nabla_\sigma \frac{dx^\alpha}{ds} \right)$$

$$\frac{dx^\sigma}{ds} \nabla_\sigma V^\mu = \frac{dx^\sigma}{ds} \left(\nabla_\sigma w_{\mu\alpha} \right) \frac{dx^\alpha}{ds} + w_{\mu\alpha} \frac{dx^\sigma}{ds} \left(\nabla_\sigma \frac{dx^\alpha}{ds} \right) = \frac{dx^\sigma}{ds} \partial_\sigma \frac{dx^\alpha}{ds} + \Gamma_{\sigma\beta}^\alpha \frac{dx^\sigma}{ds} \frac{dx^\beta}{ds} = 0$$

The only thing that remains is $\frac{dx^\sigma}{ds} (\nabla_\sigma \omega_{\mu\nu}) \frac{dx^\mu}{ds}$

$$\nabla_\sigma \omega_{\mu\nu} = -\nabla_\mu \omega_{\sigma\nu} = \nabla_\mu \omega_{\nu\sigma} = -\nabla_\nu \omega_{\mu\sigma}$$

$$(\nabla_\sigma \omega_{\mu\nu}) \frac{dx^\mu}{ds} \frac{dx^\nu}{ds} = -\nabla_\sigma \omega_{\nu\mu} \frac{dx^\mu}{ds} \frac{dx^\nu}{ds} \stackrel{\mu \leftrightarrow \nu}{=} -\nabla_\sigma \omega_{\nu\mu} \frac{dx^\nu}{ds} \frac{dx^\mu}{ds} = 0$$

2 Vortical solutions revisited

In class we introduced “vortical” geodesic solutions as solutions for which $z = \cos \theta$ oscillates between $0 \leq z_1 \leq z_2 \leq 1$. We argued that these solutions exist only when

$$\mathcal{E}^2 - \mu > 0 \quad \mathcal{E}^2 > \mu \quad (1)$$

$$\mathcal{L}^2 \leq a^2(\mathcal{E}^2 - \mu) \quad \mathcal{E}^2 = \mu \quad (2)$$

$$-(|\mathcal{L}| - |a|\sqrt{\mathcal{E}^2 - \mu})^2 \leq \mathcal{Q} \leq 0 \quad \mathcal{E} = \pm \sqrt{\mu} \quad (3)$$

Let's examine these solutions a little further.

a. Define

$$\hat{Q} = \frac{Q}{a^2(E^2 + \mu)}$$

$$\hat{L}^2 = \frac{L^2}{a^2(E^2 + \mu)}$$

$$u = \cos^2 \theta$$

$$\hat{P}_\theta = -\frac{P_\theta}{a^2(E^2 + \mu)}$$

Show that with this notation we get

$$\hat{P}_\theta(u) = u^2 + (\hat{Q} + \hat{L}^2 - 1)u - \hat{Q}.$$

b. Find the zeroes of $\hat{P}_\theta(u)$.

c. Under what conditions are both roots u_1 and u_2 real lie in the range $0 < u_1 < u_2 < 1$? Can you recover the mentioned conditions for vortical solutions.

We now turn the radial part of vortical solution.

d. Show that for vortical solutions

$$P_r \geq 2M(Q + (L - aE)^2 + \mu r^2).$$

e. Show that for vortical null geodesics ($\mu = 0$) that $P_r > 0$. What does this mean for their radial solutions? Sketch a vortical null trajectory in the maximally extended Penrose diagram for Kerr.

f. Show that for vortical timelike geodesics ($\mu > 0$) we always have that $P_r > 0$ outside r_- .

$$u^2 + (\hat{Q} + \hat{L}^2 - 1)u - \hat{Q} = 0$$

$$\Delta = \sqrt{(\hat{Q} + \hat{L}^2 - 1)^2 + 4\hat{Q}} \quad u_{1,2} = \frac{-(\hat{Q} + \hat{L}^2 - 1) \pm \sqrt{(\hat{Q} + \hat{L}^2 - 1)^2 + 4\hat{Q}}}{2}$$

$$\sqrt{(\hat{Q} + \hat{L}^2 - 1)^2 + 4\hat{Q}} = \hat{Q}^2 + 2(\hat{L}^2 + 1)\hat{Q} + (\hat{L}^2 - 1)^2$$

$$(Q + 1)(L + 1) \quad (Q + 1)(L - 1) \rightarrow \begin{matrix} A+B = (\hat{L}-1) + (\hat{L}+1) \\ AB = (\hat{L}+1)(\hat{L}-1)^2 \end{matrix}$$

Real solutions when $(Q+L^2-1)^2+4Q > 0$

$$(Q+L^2-1)^2 > -4Q \rightarrow Q < 0$$

$$0 < \frac{-(Q+L^2-1) \pm \sqrt{(Q+L^2-1)^2+4Q}}{2} < 1$$

$$-(Q+L^2-1) - \sqrt{(Q+L^2-1)^2+4Q} > 0$$

$$\sqrt{(Q+L^2-1)^2+4Q} > -(Q+L^2-1)$$