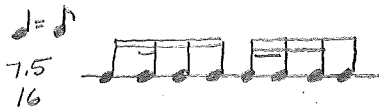


$\text{♩} = 120 \text{ bpm}$



• Musical Notation: Full of redundancies.

Ex.

• All technically expressing the same discretization of time

Music: temporal discretization on a lattice  $\approx 8 \times 10^{-8} \text{ s}$

Attempting to express emotional journey,

A purpose for each moment

Not easily translated into 2D sheet of paper

Comp. relies on performer to elaborate upon instructions

- to dynamically introduce features to the continuum.

and they influence the continuum interpretations by choosing among seemingly redundant notations. Diff. Subliminal instructions

Performers must extrapolate to continuum

Also Occurs in our simulations of Nature on a system with a finite number of degrees of freedom. Today we will be composers of Chiral Symmetry.

- Attempt to Express the low-Energy symmetries of QCD and see that the simulation also dynamically elaborates on our instructions.

As Any good composer, First Be clear on the idea YOU want to Express  
Nature:

outline of moods { 1=16 Determination  
No  
boyish enthusiasm  
productive Discouragement  
0 → a - Hope  
Lüscher - clarity

$$U(N_f)_L \times U(N_f)_R \rightarrow \underbrace{SU(N_f)_L \times SU(N_f)_R \times U(1)_V \times U(1)_A}_{\text{Spontaneous Symm. Breaking}}$$

• goldstone Bosons  
• No parity Partners  $N^*$  1535 MeV

$$\boxed{\{D, \gamma_5\} = 0}$$

But our Story lives with  $U(1)_A$

Whenever have <sup>global</sup> symmetries → Conserved Currents

$$\langle 0 | J_5^\mu(0) J^\mu(x_1) J^\mu(x_2) | 0 \rangle$$

forced to make a decision

$$\partial_\mu J^\mu = 0$$

• conserves fermion #

• couple to photon

$$\frac{\epsilon^{\mu\nu\alpha\beta} k_\mu k_\nu J_\alpha \rightarrow 0}{k^2}$$

• 2 polarizations.

$$\mathcal{L} = \bar{\psi} i \gamma^\mu (\partial_\mu - i e A_\mu) \psi$$

$$\partial_\mu J_5^\mu \sim \epsilon^{\mu\nu\alpha\beta} F_{\mu\nu} F_{\alpha\beta} \neq 0$$

due to instantons + carriers of topological charge

$$\partial_\mu J_5^\mu \sim \Phi_{top} = n_- - n_+ \sim RH$$

Atiyah Singer

LH zero modes of Dirac operator

$$D|4\rangle = 0$$

$$D\gamma_5|4\rangle = 0$$

$$\Rightarrow \gamma_5|4_0\rangle = \pm|4_0\rangle$$

chirality ✓

Obs.  $\pi^0 \rightarrow \gamma\gamma$ , flavor Singlet (anomaly ~~SSB~~)  $m_{\pi^0} \gg m_\pi$

Start w/ Invariant Action

have a measure that is not!

Jump in

$$S = \int d^4x \bar{\psi} \gamma_\mu \partial_\mu \psi \longrightarrow a^d \sum_{x, \mu} \bar{\psi}_x \gamma_\mu \frac{U_\mu(x) \psi_{x+\hat{\mu}} - U_{-\mu}(x) \psi_{x-\hat{\mu}}}{2a}$$

parallel transport gauge

$$U_\mu^\dagger(x, \hat{\mu})$$

prop  $\sim \frac{1}{p^2}$  single pole  $\xrightarrow{LSZ} ?$

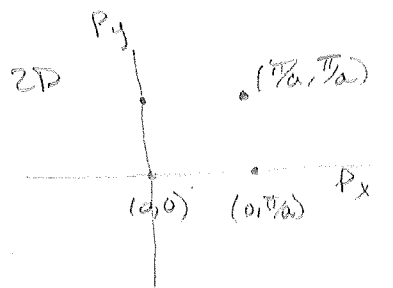
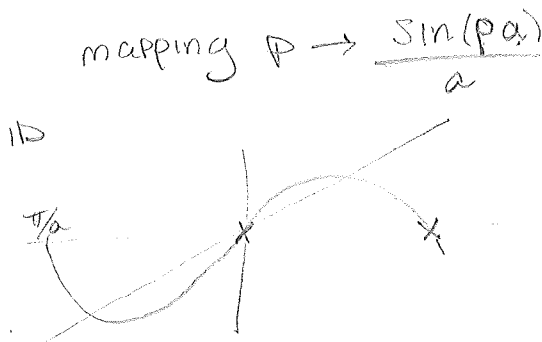
Bilinear Dirac op:

$$D_{x,y} = \sum_\mu \gamma_\mu \frac{1}{2a} (U_\mu(x) \delta_{y, x+\hat{\mu}} - U_{-\mu}(x) \delta_{y, x-\hat{\mu}})$$

$$\tilde{D}^{-1} = \langle \bar{\psi}(-p) \psi(p) \rangle = \left[ i \sum_\mu \gamma_\mu \frac{1}{a} \sin(p_\mu a) \right]^{-1}$$

$$\{ \gamma^\mu, \gamma^\nu \} = 2g^{\mu\nu} \quad \sim \frac{-i \sum_\mu \gamma_\mu \sin(p_\mu a) a}{\sum_\mu \sin^2(p_\mu a)} \xrightarrow{a \rightarrow 0} \frac{-i \not{p}}{p^2}$$

High Momentum Modes? Regulated by lattice. Periodic Brillouin Zone  
 /  $[-\pi/a, \pi/a]$   
 dangerous



Doubles for every dimension!

$$2^4 \text{ or } 16 \text{ in } D=4$$

$$\therefore 16 = 1$$

But Not to Worry. Have a way to comment EFT<sup>s</sup>. Higher Dim. op.

Exactly what Wilson did

back to Action

$$+ a^d \sum_{x, \mu} \frac{1}{2a} \left[ 2\bar{\psi}_x \psi_x - \bar{\psi}_x \psi_{x+\hat{\mu}} - \bar{\psi}_x \psi_{x-\hat{\mu}} \right]$$


$$\sim -\bar{\psi} a \frac{\gamma}{2} \nabla_\mu^\# \nabla_\mu \psi$$

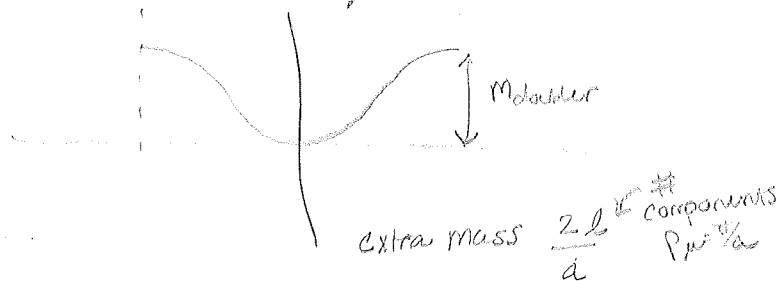
Irrelevant operator

low momentum vanishes : high mom.  $O(\text{cutoff})$  mass term

no  $\gamma_\mu$ , transforms as a mass term (wave# dep)

Denom. of  $\tilde{D}^{-1}$

$$\sum_\mu \sin(p_\mu a)^2$$


$$\sum_\mu \sin(p_\mu a)^2 + \sum_\mu (1 - \cos(p_\mu a))^2$$


extra mass  $\frac{2L}{a}$  components  $p_\mu \frac{\pi}{a}$

Implications:

✓  $a \rightarrow 0$  decouple from theory.

✗ Explicitly broken Chiral Symmetry (w/o masses)

- No protection of  $M_\pi$  from additive mass renormalization

What is keeping the masses so light? Hierarchy.

Rules - think we want - Symmetry of Action  
Not Symmetry of Measure. Continuum  $\rightarrow$   $\chi^{\text{real}}$  Symm.

- Exact Symm. of regulated Action = exact Sym of Quantum theory  
(No Anomalies)

If Anomalous in continuum  $\rightarrow$  broken on lattice

-  $\partial_\mu J^\mu = 0$  1.) Original Action 2.) UV regulator } explicitly violate Symmetry.  
Symm current have non-zero Div.

Nelson Ninomiya (1981)

So. Confusing and Unclear how to move forward

Telling Scientists: 'You can't have what you want.'

Changed Question } how do we express continuum sym. in a regulated theory  
to: what of our continuum qualities must we give up?

1.) Bilinear

$$S = \sum_{x,y} \bar{\psi}(x) M_{xy} \psi(y)$$

2.) Real

$$M = M^T$$

3.) Translation Inv.

$$M_{xy} = D(x-y)$$

4.) Local

$D(p)$  - regular func. of  $p$  in B. zone

5.) Chirally Symmetric

$$\{D, \gamma_5\} = 0$$

6.) Doubler free.

$$I = 1$$

Performer cannot Do it

5

of course, topology.

$$\{D, \gamma^5\} = 0 \Rightarrow D = \sum_{\mu} \gamma_{\mu} d_{\mu}(\varphi)$$

↑  
vector field on torus

$\{\gamma_{\mu}, \gamma^5\} = 0$   
keeps kinetic terms  
from mixing LH and RH

$\Rightarrow$  Zero modes definite chirality  
 $\gamma^5 | \chi_{\pm} \rangle = \pm | \chi_{\pm} \rangle$   
↑  
handedness

Euler Characteristic:  $\chi = 2 - 2H$  Holes: 1

$$\chi = 0$$

Hopf-Poincaré Index theorem

$$\sum \text{Index}(d_{\mu}) = \chi = 0$$

$\Rightarrow$  Zeros of  $d_{\mu}$  come in pairs  
of opp index.

Doublers!

6

Lüscher, Ginsparg and Wilson

$$\psi \rightarrow \psi + i\varepsilon^a \tau^a \gamma_5 \psi \quad \bar{\psi} \rightarrow \bar{\psi} + \bar{\psi} i\varepsilon^a \tau^a \gamma_5$$

generators of  $SU(N_f)$   
 $a \in \{1, \dots, N_f^2 - 1\}$   
Flavor singlet Anal ( $\pi'$ )  
 $a=0 \quad \tau^a = \mathbb{I}$

1982 "A remnant of Chiral Symmetry on the Lattice"

Continuum Action and Spin blocking (Wilson)  
renormalization

$$\{D, \gamma_5\} = a D \gamma_5 D$$

Implications for quark prop ( $D^\dagger$  on both sides)

$$\{D_{xy}^\dagger, \gamma_5\} = a \gamma_5 \delta(x-y)$$

Added A contact term "Loop" correction

Dirac  $\phi$ :

Assume  $\gamma_5$  hermiticity  $\gamma_5 D \gamma_5 = D^\dagger$

$$\det [D - \lambda \mathbb{1}] = \det [D^\dagger - \lambda \mathbb{1}]$$

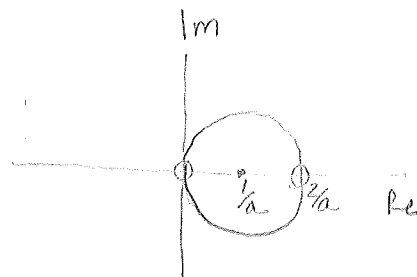
$\lambda$  and  $\lambda^*$  are E. vals : CC pairs

Chirality:  $\gamma_5 M.E.$

$$\lambda (\phi_1, \gamma_5 \phi_2) = \lambda^* (\phi_2, \gamma_5 \phi_1)$$

$$\text{Im } \lambda (\phi_1, \gamma_5 \phi_2) = 0$$

only eigenvalues w/ real eigenvalues can have non-vanishing chirality



G + W could not find solution

- Looking for Local  
nearest neighbor

Spelling it, too restrictive. Next week, solutions

Locality: Couplings in Action decay w/  
distance faster than any signal eg  $e^{-\frac{Kx}{a}}$   
 $K O(1)$  is "local"

So: Currently have (w/ continuum transformations)

Backwards! [ Local Actions + No Doubling  $\Rightarrow$  not a sym of Action  
Fermion Measure: Invariant

Lüscher 1998: Modify transformation

$$\psi \rightarrow \psi + i\varepsilon T^a \gamma_5 (1 - \frac{a}{2} D)$$

$$P_L = \frac{1 - \hat{\gamma}_5}{2} \quad \hat{\gamma}_5 = \gamma_5 (1 - aD)$$

Gauge field Dependent!  
Continuum: locality is a local concept.

Is it a Symmetry?

$$\bar{\psi} D \psi \rightarrow \bar{\psi} \left( 1 + i\varepsilon T^a (1 - \frac{a}{2} D) \gamma_5 \right) D \left( 1 + i\varepsilon T^a \gamma_5 (1 - \frac{a}{2} D) \right) \psi$$

$$\rightarrow \bar{\psi} D \psi + \bar{\psi} i\varepsilon T^a \left[ \gamma_5 D + D \gamma_5 - a D \gamma_5 D \right] \psi + O(\varepsilon^2)$$

vanishes if Sol. to GW!



2

What About the Measure?

Grassman Int  $\rightarrow$  det of trans. Matrix.

$$\text{Det} \left[ \mathbb{1} + i\varepsilon^a T^a \gamma_5 \left( \mathbb{1} - \frac{a}{2} D \right) \right]^2$$

$$\text{Det} X = e^{\overset{\text{LFDs}}{\int} \text{tr} \log X}$$

$$= 1 + 2i\varepsilon^a \text{tr}_{\text{FDS}} \left[ T^a \gamma_5 \left( \mathbb{1} - \frac{a}{2} D \right) \right] + \mathcal{O}(\varepsilon^2) \log(1+x) = x + \dots$$

$$= 1 - \frac{2i\varepsilon^a}{2} \underbrace{\left( \text{tr}_F T^a \right)}_{N_f} \underbrace{\left( \text{tr}_{\text{CD}} \gamma_5 D \right)}_{\sim Q_{\text{top}}} + \mathcal{O}(\varepsilon^2) \neq 1$$

Invariant Action And Anomalous Measure! Not Just Any Anomaly...

$$\frac{a}{2} \text{tr}_{\text{CD}} \gamma_5 D = -\frac{1}{2} \text{tr} \left( \gamma_5 (2 - aD) \right)$$

tr: sum of evals

$$= -\frac{1}{2} \sum_{\lambda} (\phi_{\lambda}, \gamma_5 (2 - aD) \phi_{\lambda})$$

$$= -\frac{1}{2} \sum_{\lambda} (2 - a\lambda) (\phi_{\lambda}, \gamma_5 \phi_{\lambda})$$

$\langle \gamma_5 \rangle$   
 vanishes unless  $\lambda$  real  
 $\lambda = 0, \lambda = 2$

$\sum$  over Zero Mode Chiralities

$$= n_- - n_+ = Q_{\text{top}}$$

Just like In Continuum, topologically-related Axial Anomaly

Message: No Geo thm bypassed IF modify  $\chi^{\text{real}}$  sym.

Next: Site to GWT: purchase and Domain Walls