

The Bonn-Gatchina analysis of reaction with production of the η -meson

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Bonn-Gatchina Partial Wave Analysis



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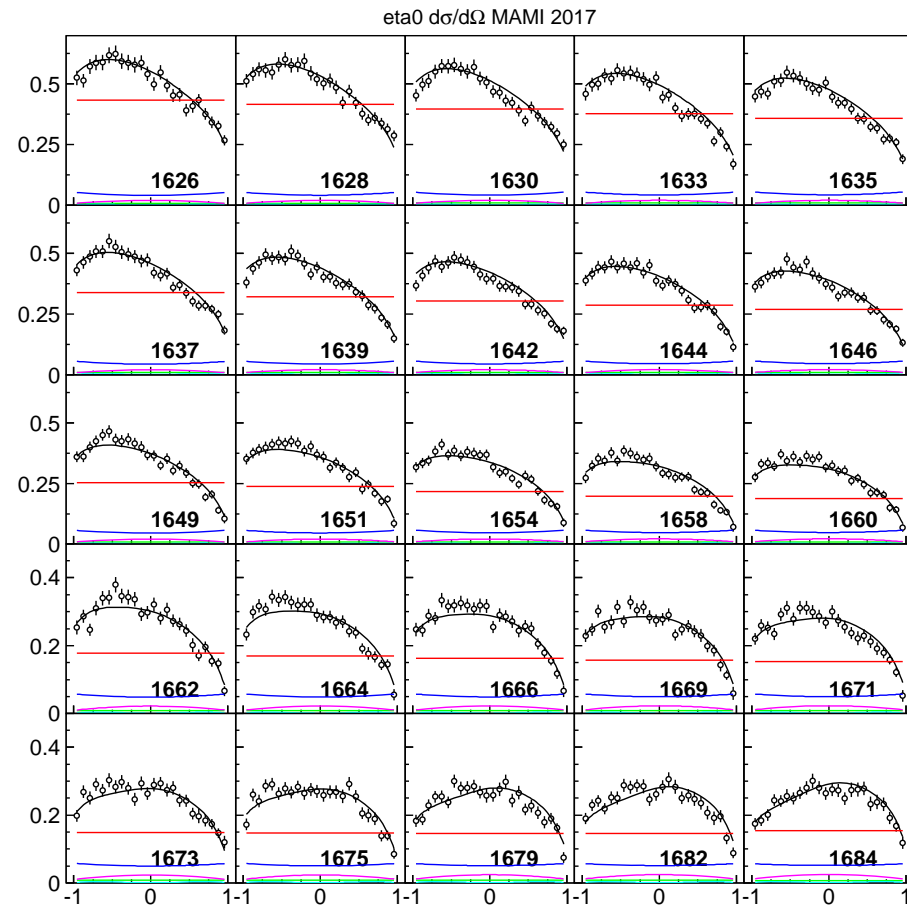
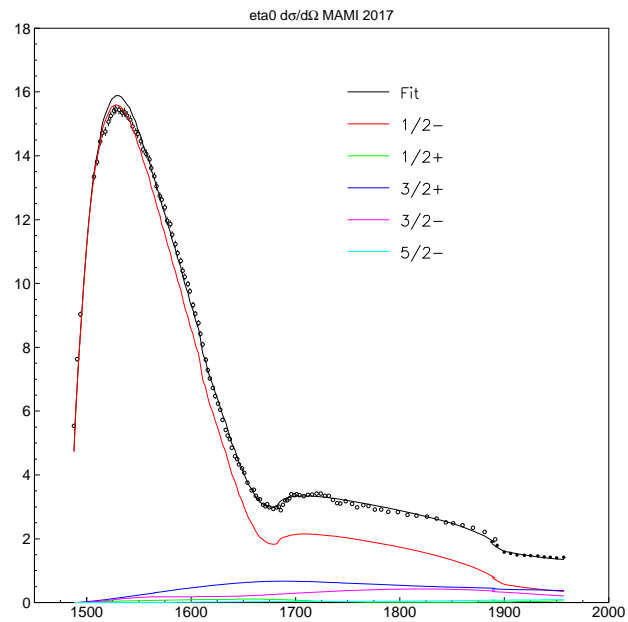
Responsible: Dr. V. Nikonov, E-mail: nikonov@hiskp.uni-bonn.de
Last changes: January 26th, 2010.

Recently included data

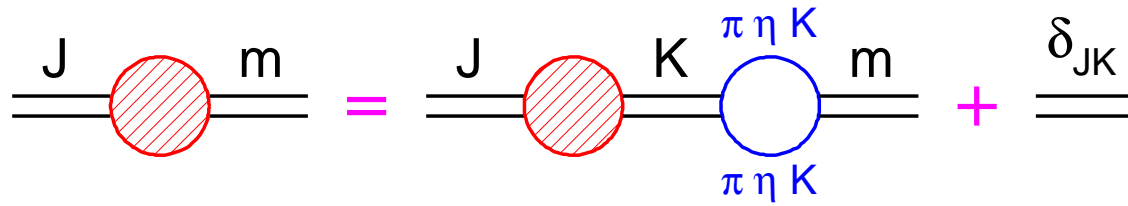
DATA	2011-2016	added in 2016-2018
$\gamma n \rightarrow \Lambda K, \Sigma^- K$		$\frac{d\sigma}{d\Omega}$ (CLAS), E (CLAS)
$\gamma n \rightarrow \pi^- p$	$\frac{d\sigma}{d\Omega}, \Sigma, P$	E, Σ (CLAS)
$\gamma n \rightarrow \eta n$	$\frac{d\sigma}{d\Omega}, \Sigma$	$\frac{d\sigma}{d\Omega}$ (MAMI) $\frac{d\sigma}{d\Omega} (h = \frac{1}{2})$ (CB-ELSA)
$\gamma p \rightarrow \eta p$	$\frac{d\sigma}{d\Omega}, \Sigma$ (GRAAL)	$\frac{d\sigma}{d\Omega}, F, T$ (MAMI) $T, P, H, G,$ (CB-ELSA) E, Σ (CB-ELSA, CLAS)
$\gamma p \rightarrow \eta' p$		$\frac{d\sigma}{d\Omega}, \Sigma$
$\gamma p \rightarrow K^+ \Lambda$	$\frac{d\sigma}{d\Omega}, \Sigma, P, T, C_x, C_z, O_{x'}, O_{z'}$	Σ, P, T, O_x, O_z (CLAS)
$\gamma p \rightarrow K^+ \Sigma^0$	$\frac{d\sigma}{d\Omega}, \Sigma, P, C_x, C_z$	Σ, P, T, O_x, O_z (CLAS)
$\pi^- p \rightarrow \pi^+ \pi^- n$		$d\sigma/d\Omega$ (HADES)
$\pi^- p \rightarrow \pi^- \pi^0 p$		$d\sigma/d\Omega$ (HADES)
$\gamma p \rightarrow \pi^0 \pi^0 p$	$d\sigma/d\Omega, \Sigma, E, I_c, I_s$	T, P, H, F, P_x, P_y (CB-ELSA)
$\gamma p \rightarrow \pi^+ \pi^- p$		$d\sigma/d\Omega, I_c, I_s$ (CLAS)
$\gamma p \rightarrow \omega p$	$d\sigma/d\Omega, \Sigma, \rho_{ij}^k, E, G$ (CB-ELSA)	Σ (CLAS) P,T,F,H (CLAS)
$\gamma p \rightarrow K^* \Lambda$		$d\sigma/d\Omega, \rho_{ij}$

The analysis of the new $\gamma p \rightarrow \eta p$ data.

New MAMI data: a strong cusp effect from the $\eta' p$ channel



N/D based (D-matrix) analysis of the data



$$D_{jm} = D_{jk} \sum_{\alpha} B_{\alpha}^{km}(s) \frac{1}{M_m^2 - s} + \frac{\delta_{jm}}{M_j^2 - s} \quad \hat{D} = \hat{\kappa}(I - \hat{B}\hat{\kappa})^{-1}$$

$$\hat{\kappa} = \text{diag} \left(\frac{1}{M_1^2 - s}, \frac{1}{M_2^2 - s}, \dots, \frac{1}{M_N^2 - s}, R_1, R_2, \dots \right)$$

$$\hat{B}_{ij} = \sum_{\alpha} B_{\alpha}^{ij} = \sum_{\alpha} \int \frac{ds'}{\pi} \frac{g_{\alpha}^{(R)i} \rho_{\alpha}(s', m_{1\alpha}, m_{2\alpha}) g_{\alpha}^{(L)j}}{s' - s - i0}$$

In the present fits we calculate the elements of the B_α^{ij} using one subtraction taken at the channel threshold $M_\alpha = (m_{1\alpha} + m_{2\alpha})$:

$$B_\alpha^{ij}(s) = B_\alpha^{ij}(M_\alpha^2) + (s - M_\alpha^2) \int_{M_\alpha^2}^{\infty} \frac{ds'}{\pi} \frac{g_\alpha^{(R)i} \rho_\alpha(s', m_{1\alpha}, m_{2\alpha}) g_\alpha^{(L)j}}{(s' - s - i0)(s' - M_\alpha^2)}.$$

In this case the expression for elements of the \hat{B} matrix can be rewritten as:

$$B_\alpha^{ij}(s) = g_a^{(R)i} \left(b^\alpha + (s - M_\alpha^2) \int_{M_\alpha^2}^{\infty} \frac{ds'}{\pi} \frac{\rho_\alpha(s', m_{1\alpha}, m_{2\alpha})}{(s' - s - i0)(s' - M_\alpha^2)} \right) g_\beta^{(L)j} = g_a^{(R)i} B_\alpha g_\beta^{(L)j}$$

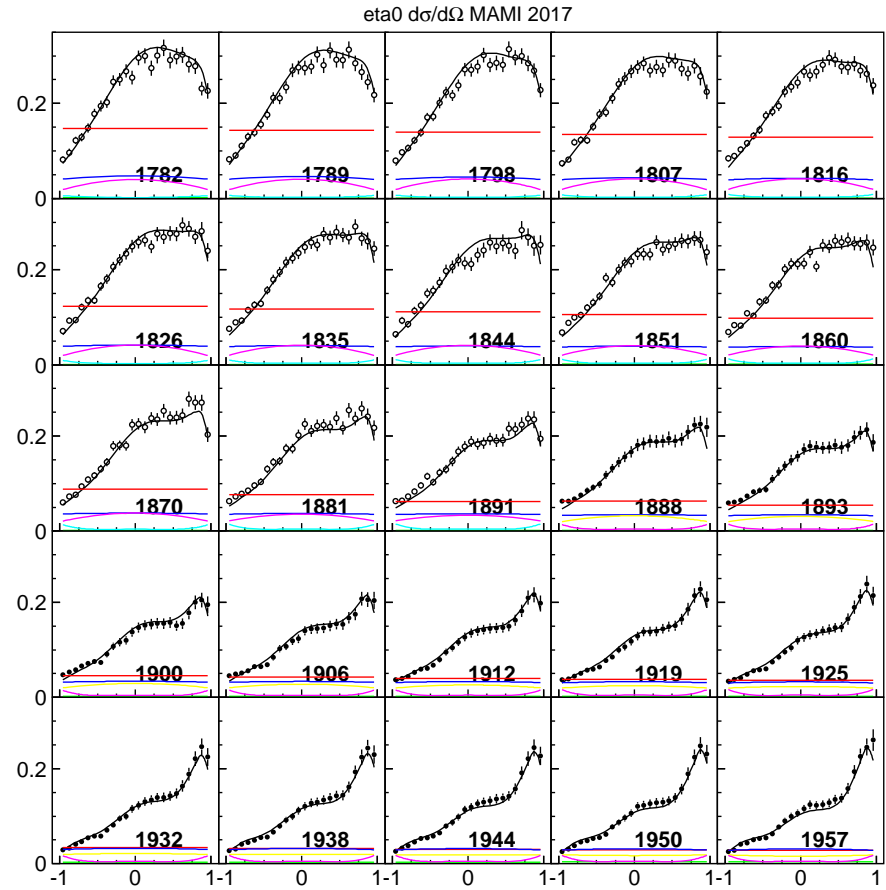
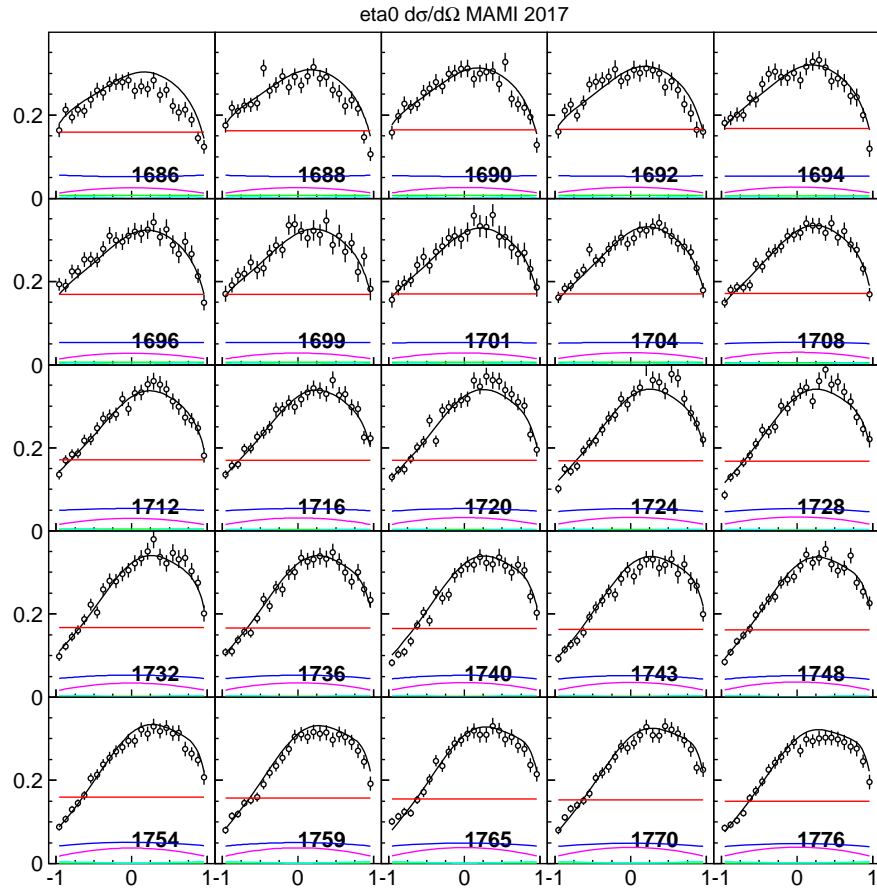
and D-matrix method equivalent to the K-matrix method with loop diagram with real part taken into account:

$$A = \hat{K}(I - \hat{B}\hat{K})^{-1} \quad B_{\alpha\beta} = \delta_{\alpha\beta} B_\alpha$$

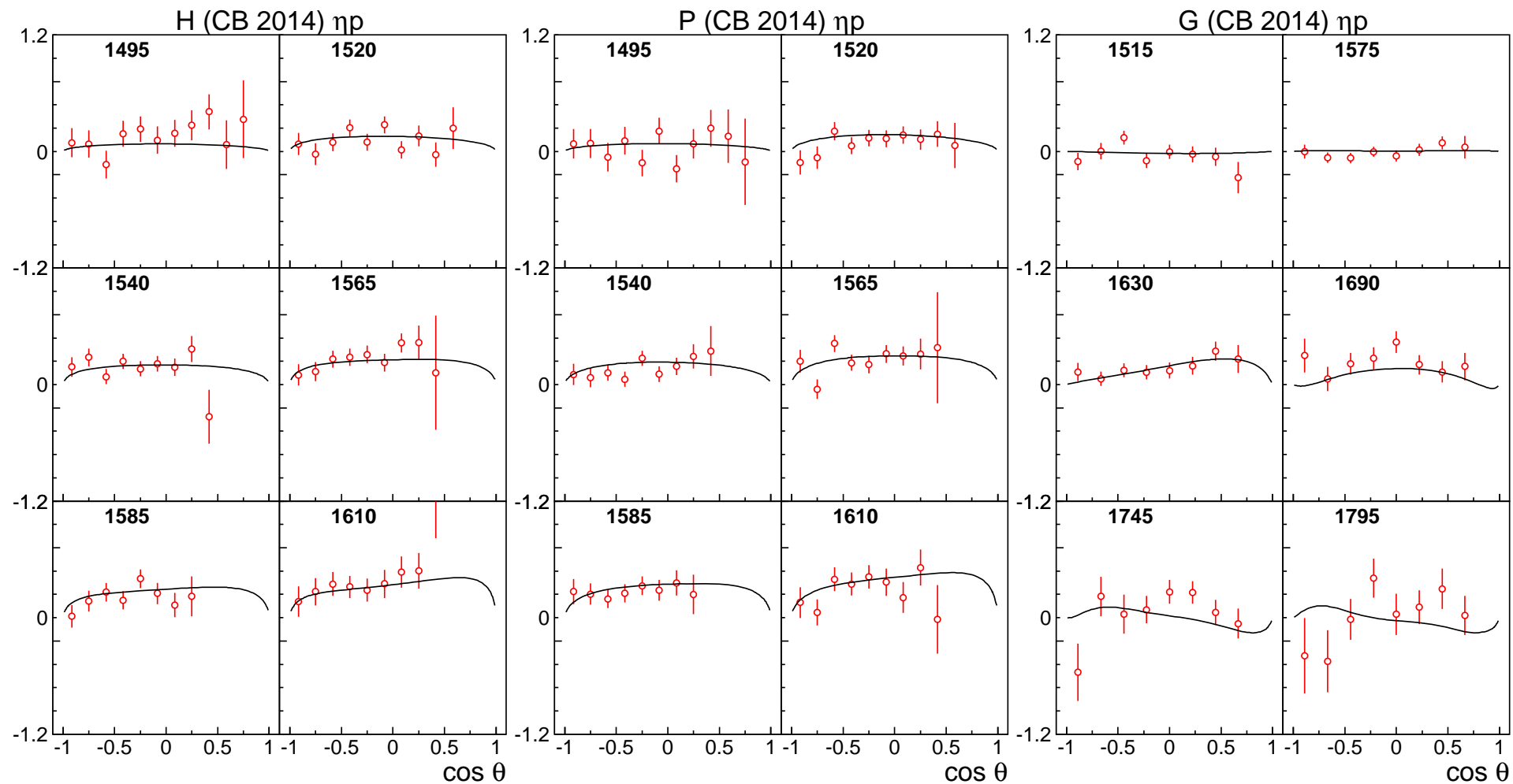
S_{11} -partial wave: πN , ηN , $K\Lambda$, $K\Sigma$, $\Delta(1232)\pi$, ρN (S,D).

	$N(1535)S_{11}$		$N(1650)S_{11}$		$N(1890)S_{11}$	
	K-matrix	D-matrix	K-matrix	D-matrix	K-matrix	D-matrix
M_{pole}	1501 ± 4	1494	1647 ± 6	1651	1900 ± 15	1905
Γ_{pole}	134 ± 11	116	103 ± 8	95	90^{+30}_{-15}	106
Elastic residue	31 ± 4	25	24 ± 3	23	1 ± 1	1.5
Phase	$-(29 \pm 5)^\circ$	-38[°]	$-(75 \pm 12)^\circ$	-62[°]	–	–
Res $\pi N \rightarrow \Delta \pi$	7 ± 4	4	11 ± 3	12	–	–
Phase	$(147 \pm 17)^\circ$	157[°]	$-(30 \pm 20)^\circ$	-40	–	–
$A^{1/2}$ ($\text{GeV}^{-\frac{1}{2}}$)	0.116 ± 0.010	0.107	0.033 ± 0.007	0.029	0.012 ± 0.006	0.010
Phase	$(7 \pm 6)^\circ$	1[°]	$-(9 \pm 15)^\circ$	0[°]	$120 \pm 50^\circ$	150[°]

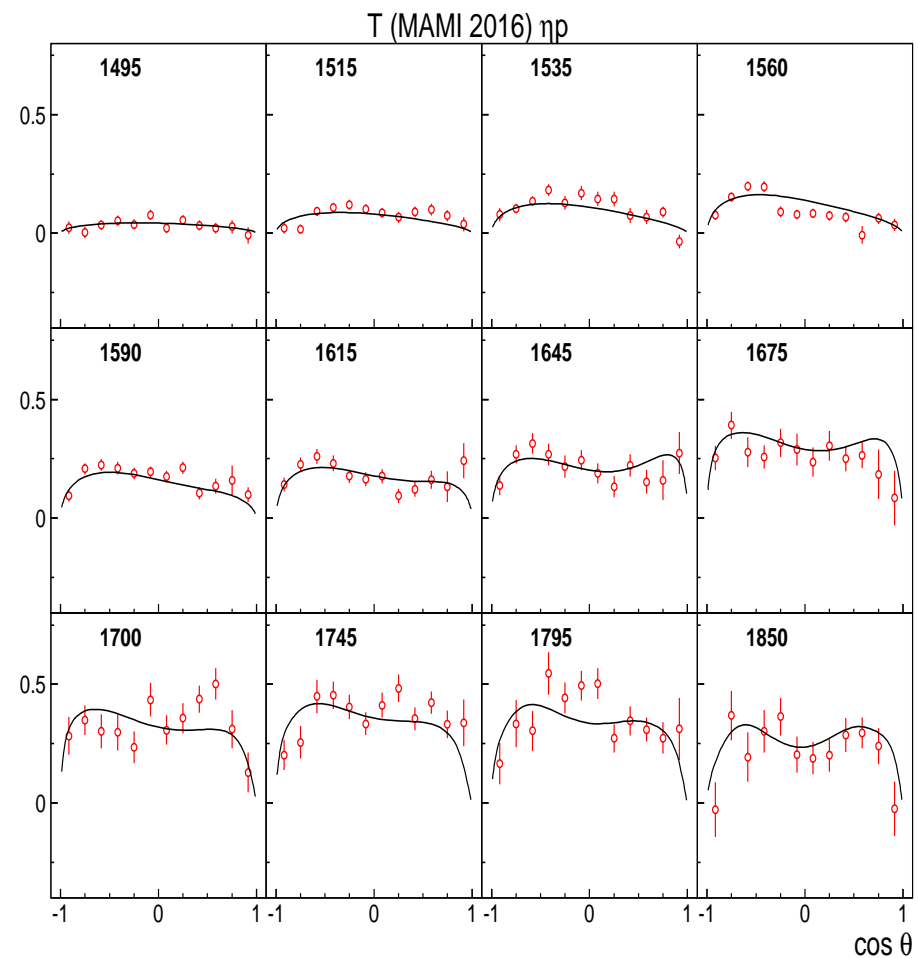
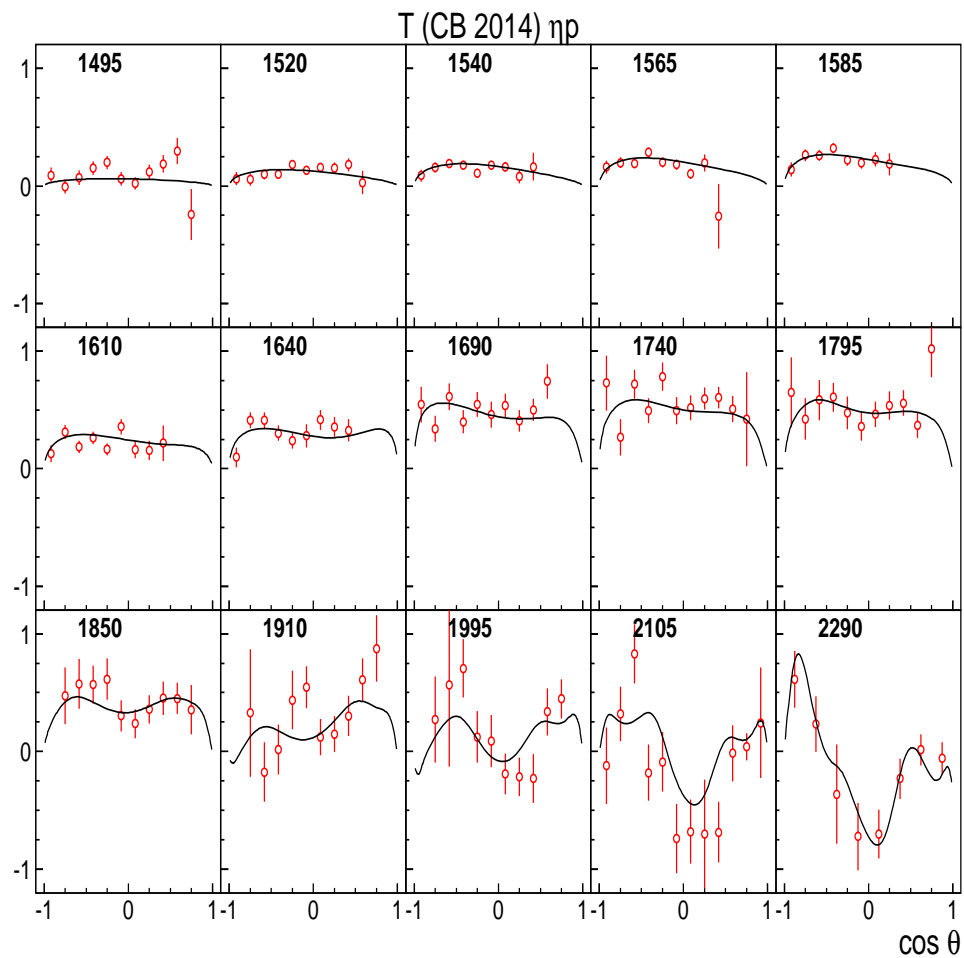
The analysis of the new $\gamma p \rightarrow \eta p$ data. $d\sigma/d\Omega$ (MAMI)



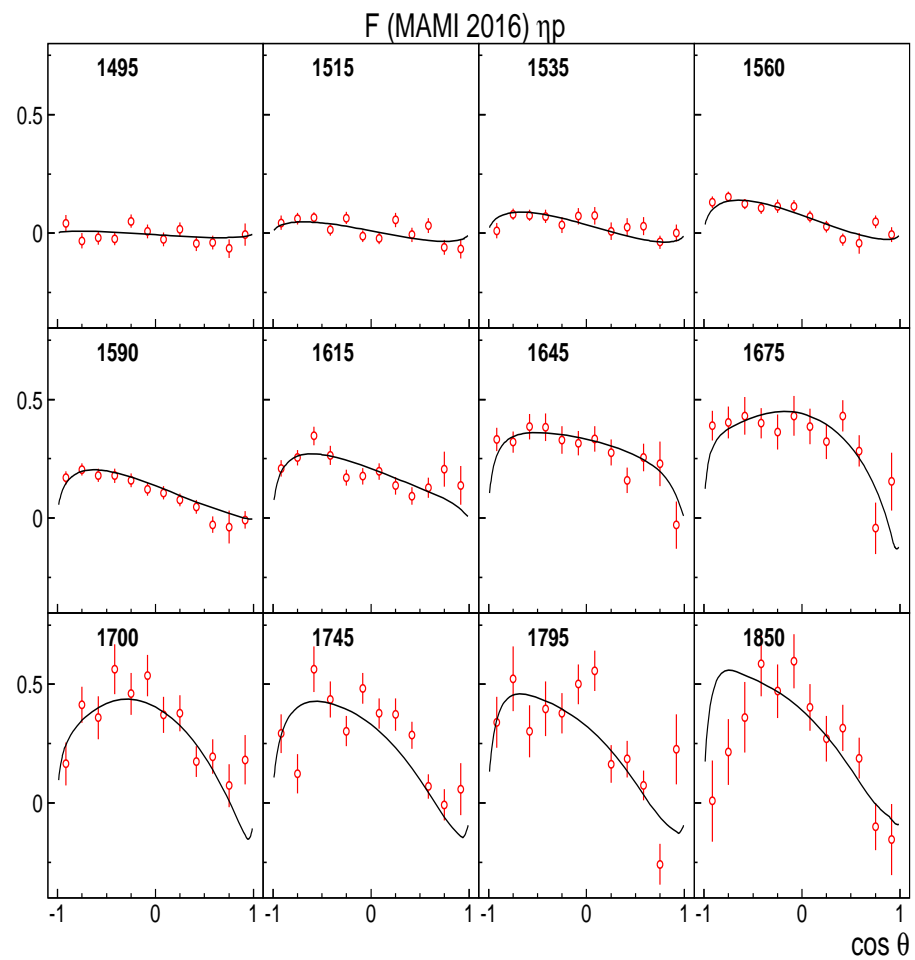
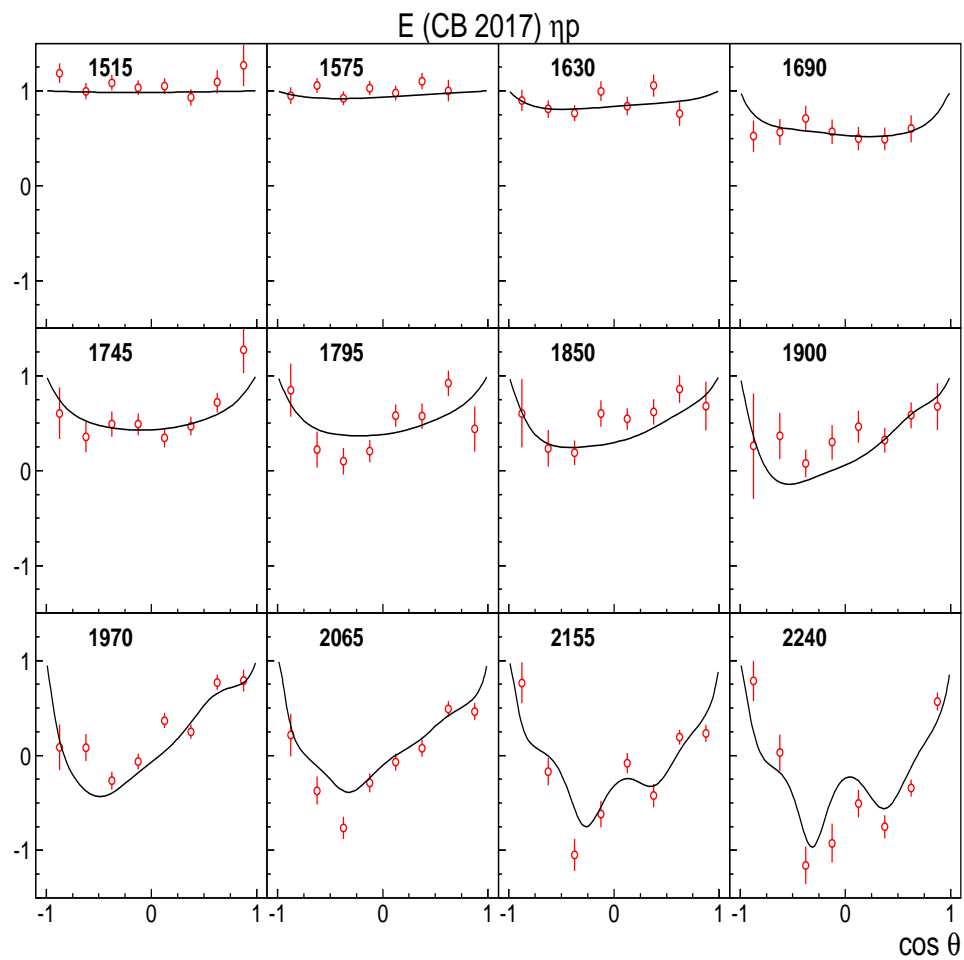
The analysis of the new $\gamma p \rightarrow \eta p$ data. H, P, T (CB-ELSA)



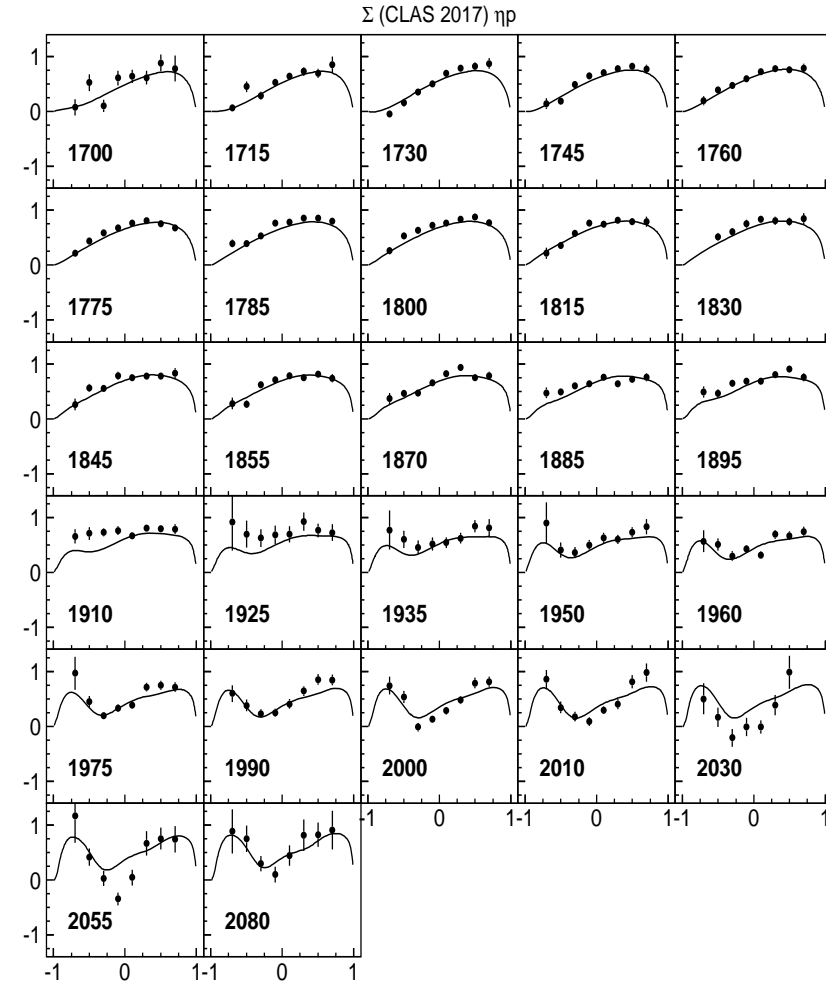
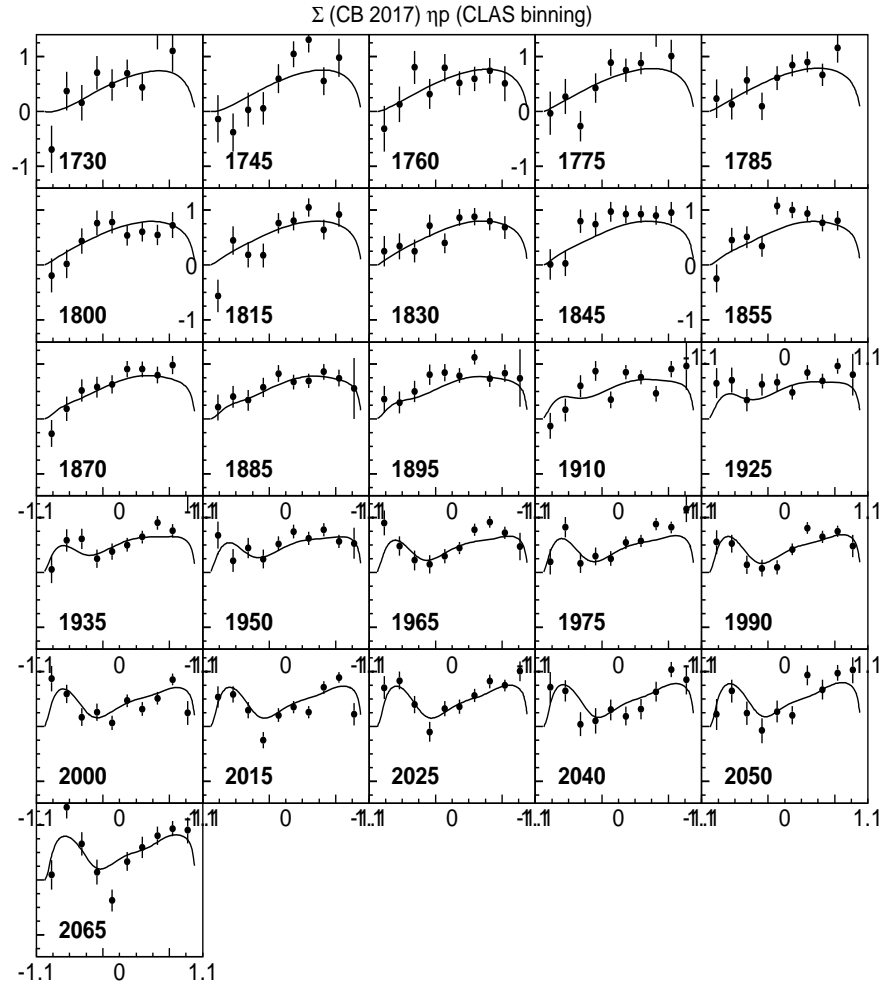
The analysis of the new $\gamma p \rightarrow \eta p$ data. T (CB-ELSA), (MAMI scale 1.4)



The analysis of the new $\gamma p \rightarrow \eta p$ data. E (CB-ELSA), F (MAMI) (scale 1.4)



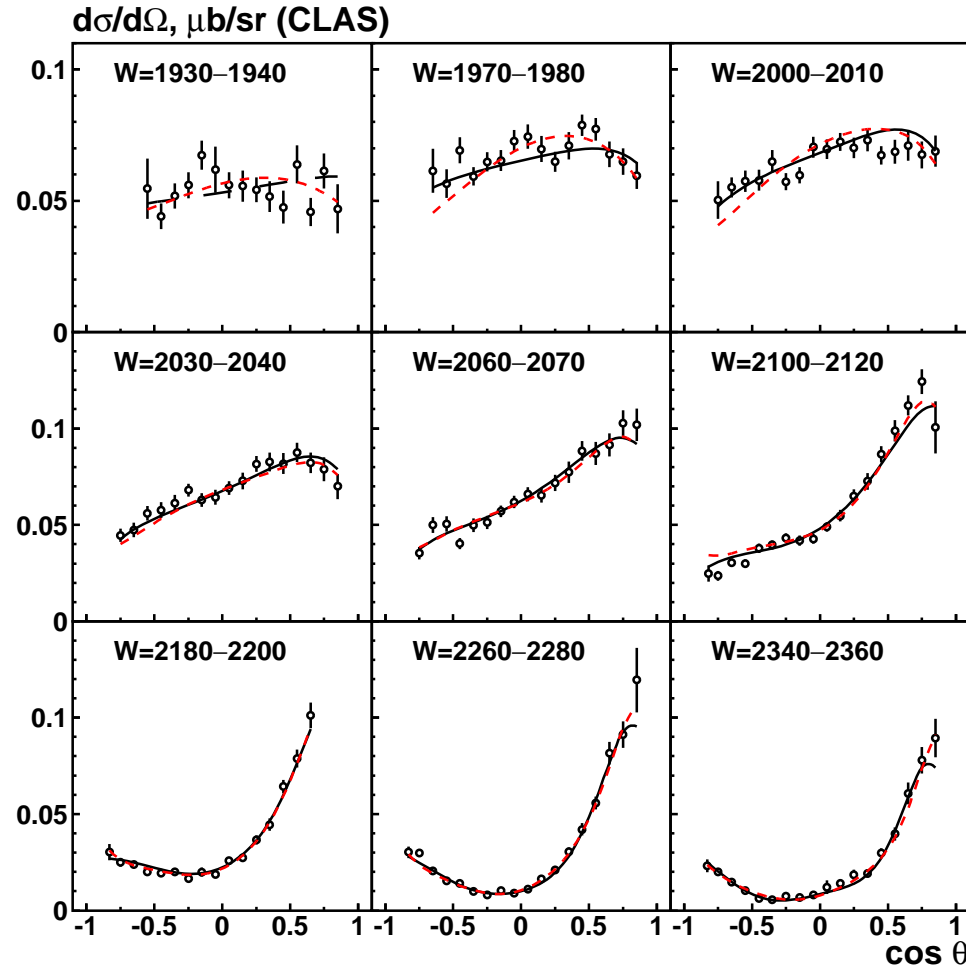
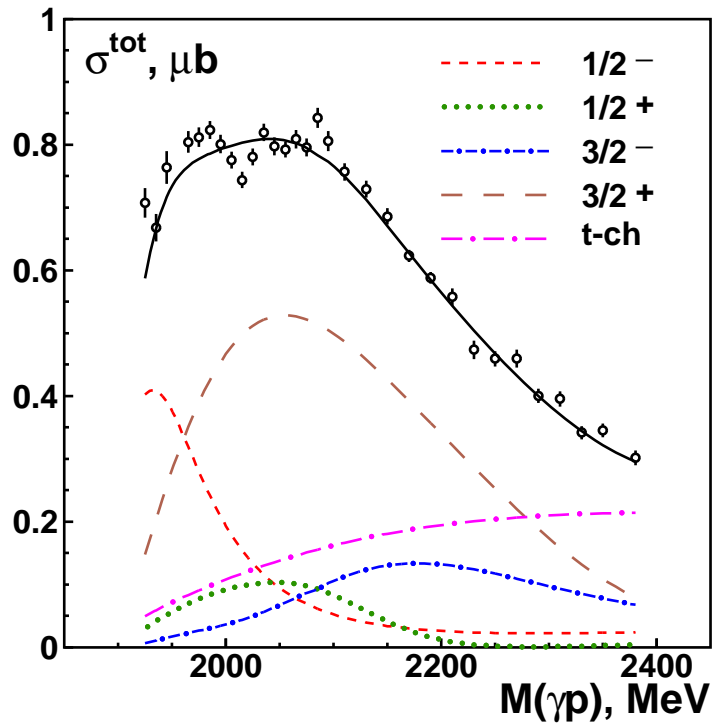
The analysis of the new $\gamma p \rightarrow \eta p$ data. Σ (CB-ELSA and CLAS)



Resonance branchings to the ηN channel

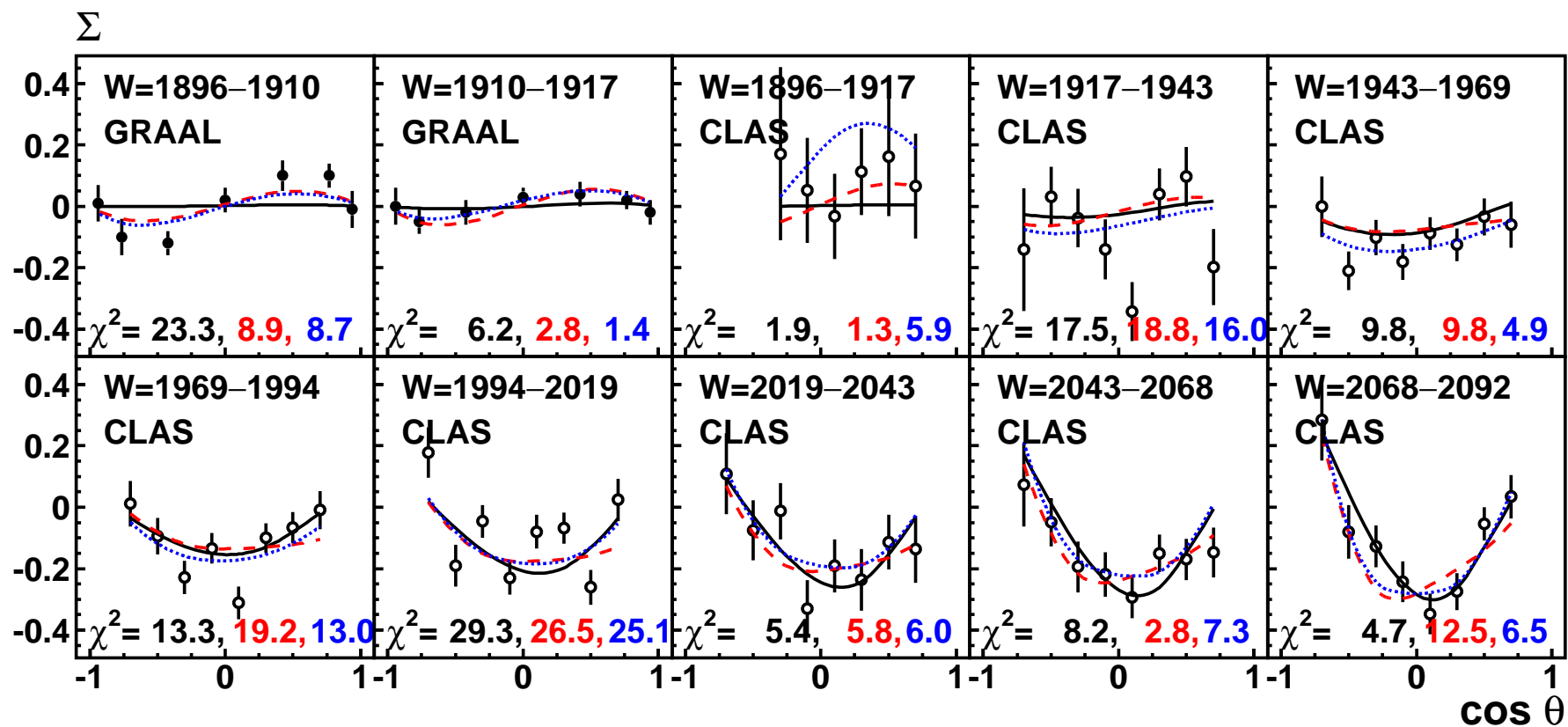
Res.	BR	Res.	BR	Res.	BR
$N(1535)$ $1/2^-$	0.42 ± 0.04 0.42 ± 0.10	$N(1650)$ $1/2^-$	0.32 ± 0.04 $0.05 - 0.15$	$N(1895)$ $1/2^-$	0.10 ± 0.05 (0.21 ± 0.06)
$N(1710)$ $1/2^+$	0.25 ± 0.09 $0.10 - 0.30$	$N(1880)$ $1/2^+$	0.19 ± 0.07 $(0.25^{+0.30}_{-0.20})$	$N(2100)$ $1/2^+$	0.25 ± 0.10 0.61 ± 0.61
$N(1520)$ $3/2^-$	< 0.001 0.0023 ± 0.0004	$N(1700)$ $3/2^-$	0.01 ± 0.01 0 ± 0.01	$N(1875)$ $3/2^-$	0.02 ± 0.01 0.012 ± 0.018
$N(1720)$ $3/2^+$	0.03 ± 0.02 0.021 ± 0.014	$N(1900)$ $3/2^+$	0.03 ± 0.01 ~ 0.12	$N(2120)$ $3/2^-$	≤ 0.01 -
$N(1675)$ $5/2^-$	0.005 ± 0.005 0 ± 0.007	$N(2060)$ $5/2^-$	0.04 ± 0.01 0.04 ± 0.02	$N(2190)$ $7/2^-$	0.025 ± 0.005 0 ± 0.01
$N(1680)$ $5/2^+$	0.002 ± 0.001 0 ± 0.007	$N(2000)$ $5/2^+$	0.002 ± 0.001 0.002 ± 0.002	$N(1990)$ $7/2^+$	≤ 0.01 -

The analysis of the $\gamma p \rightarrow \eta' p$ data.



Strong contribution from the $S_{11}(1895)$, $P_{13}(1900)$, $P_{11}(2100)$ and $D_{13}(2120)$ states.

The beam asymmetry on $\gamma p \rightarrow \eta' p$



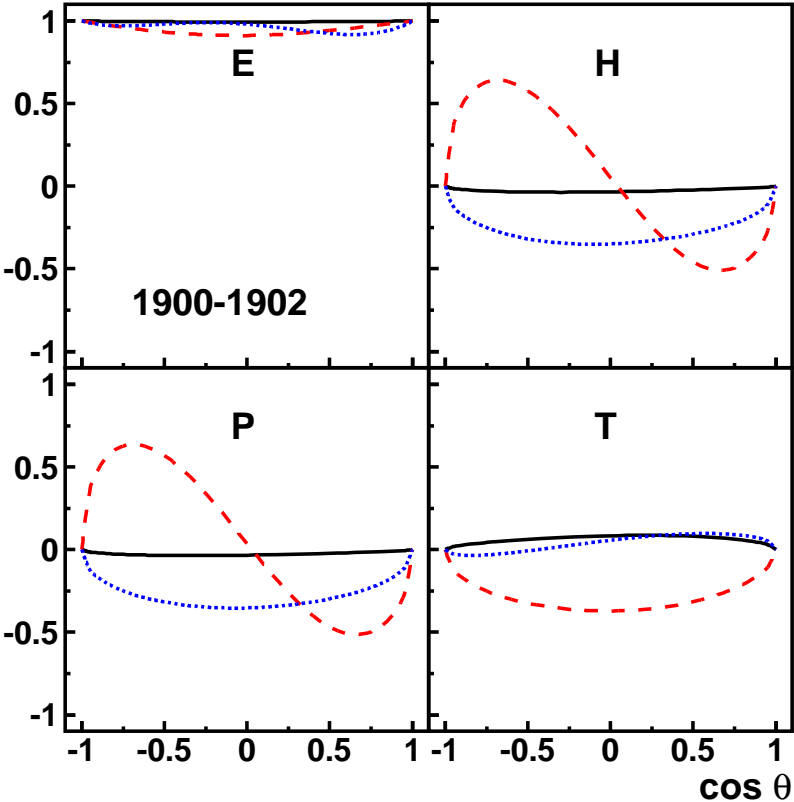
No narrow states

$D_{15}(1903)$

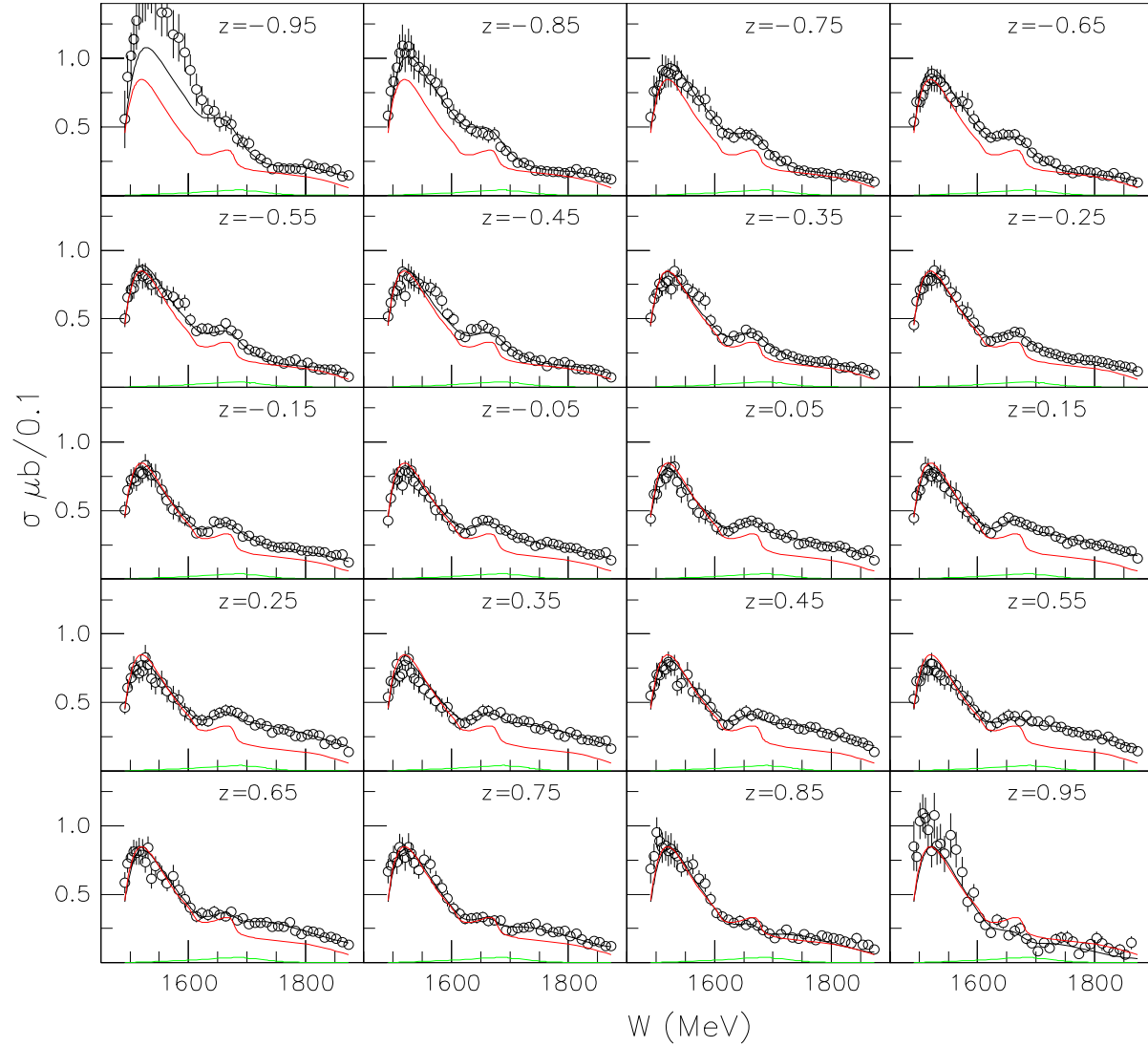
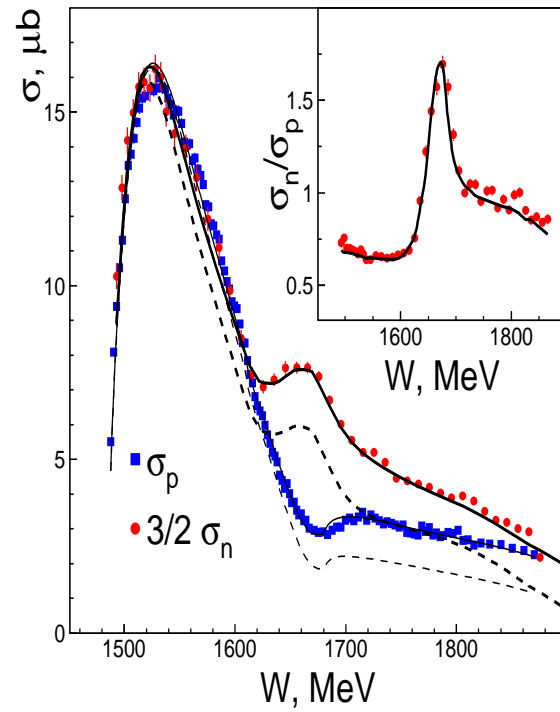
$D_{13}(1900)$

The description of the data below W=1917 MeV and the prediction of other observables

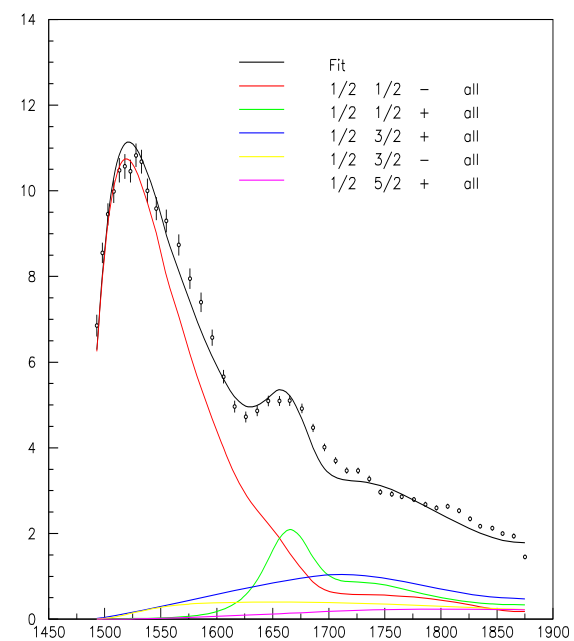
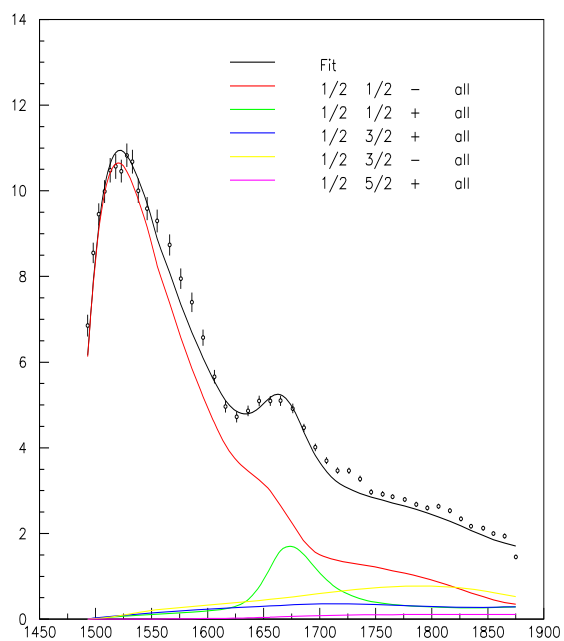
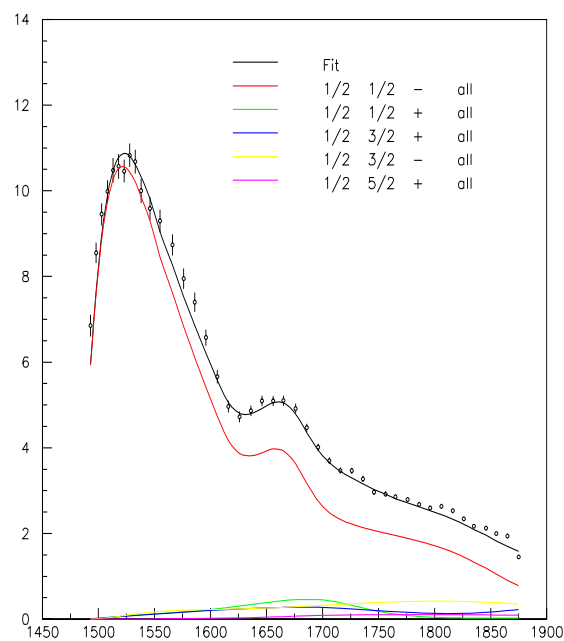
Resonance	N	Basic	D_{13}	D_{15}
M (MeV)			1900	1903
Γ (MeV)			1	1
χ^2 (Σ)	13	29.5	11.7	10.1
χ^2 ($\frac{d\sigma}{d\Omega}$)	50	120.3	59.9	129.0



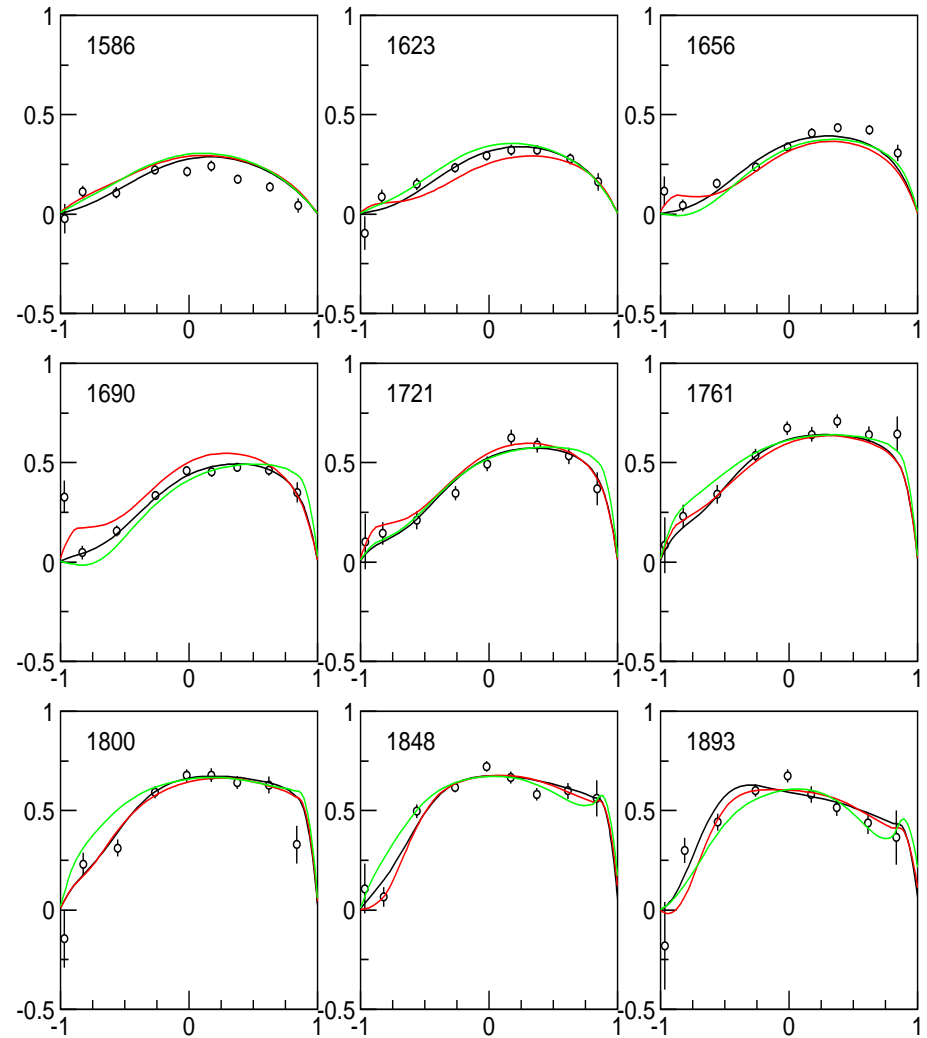
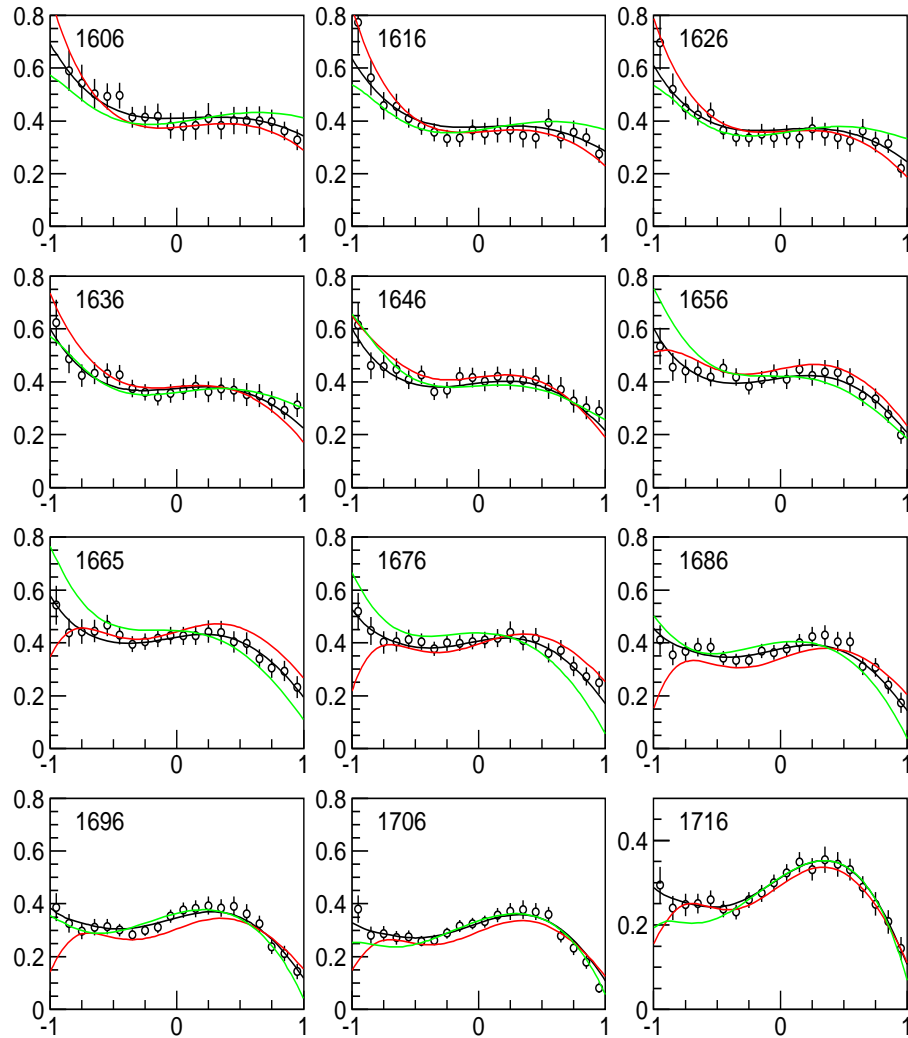
Solution with interference between S_{11} states



Solutions with the $P_{11}(1680)$ states

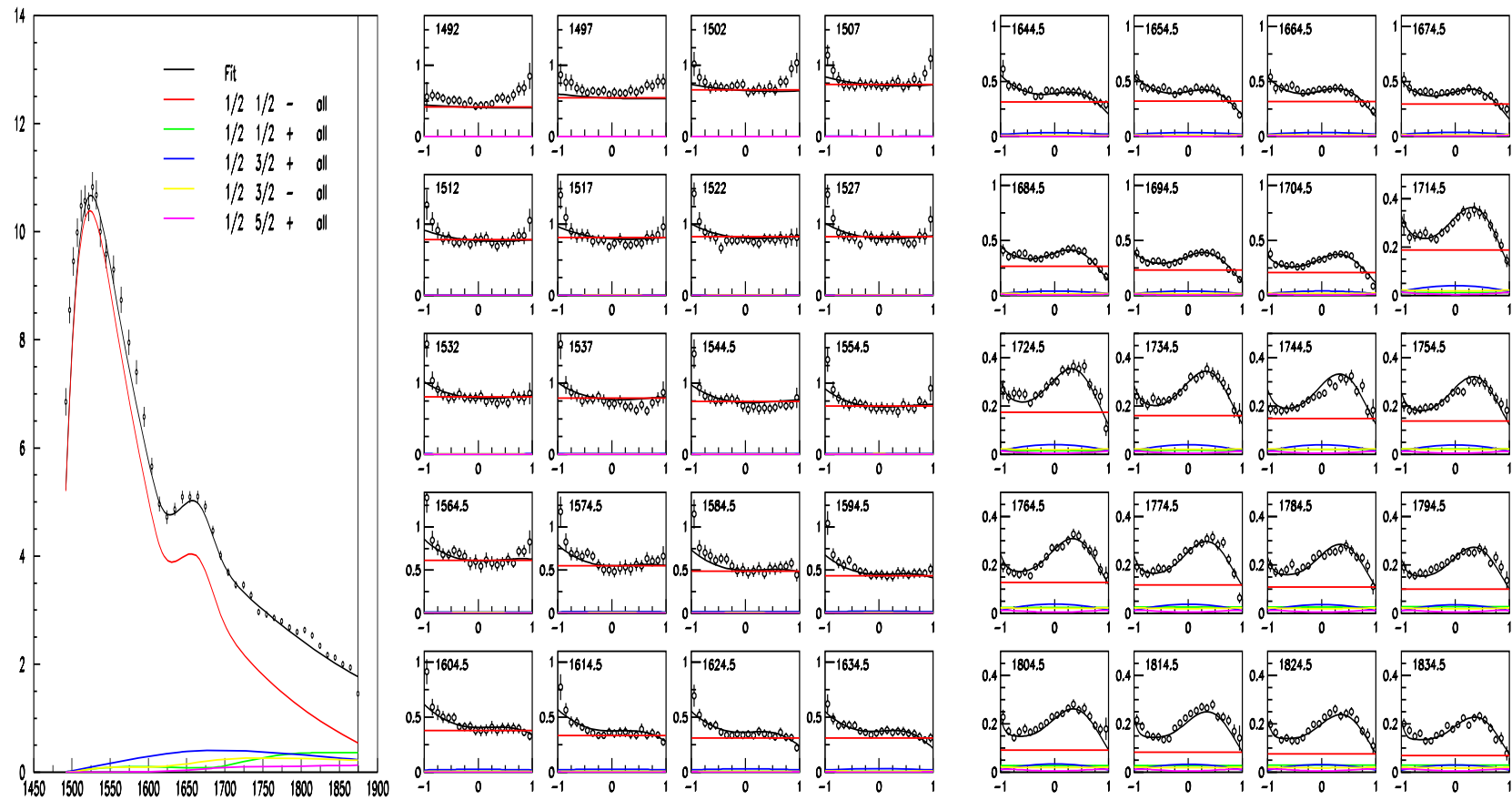


The description of the new data as well as GRAAL data is notably worse

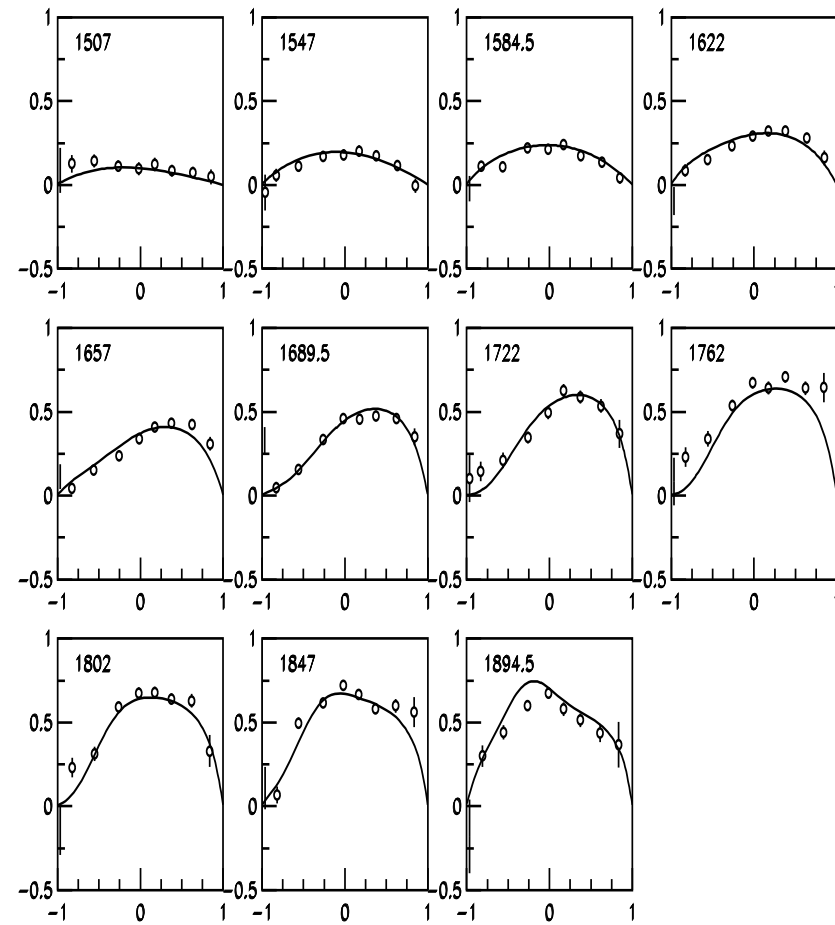
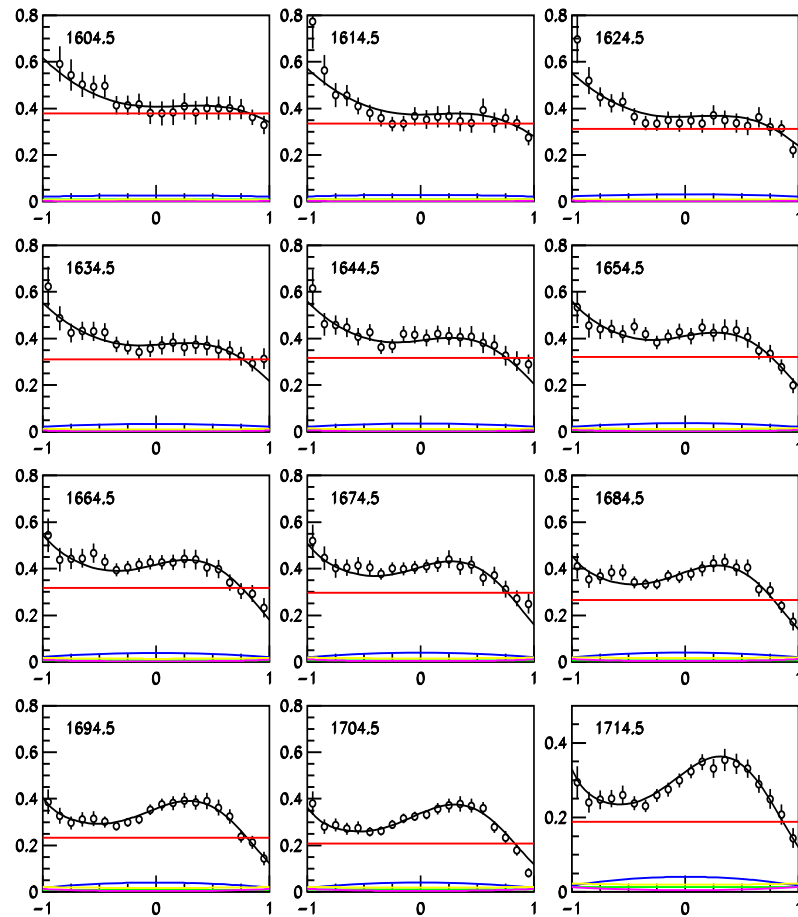


Limit for the production of $P_{11}(1680)$: $|A^{\frac{1}{2}}| Br(\eta n) < 5 \text{ GeV}^{-\frac{1}{2}} 10^{-3}$

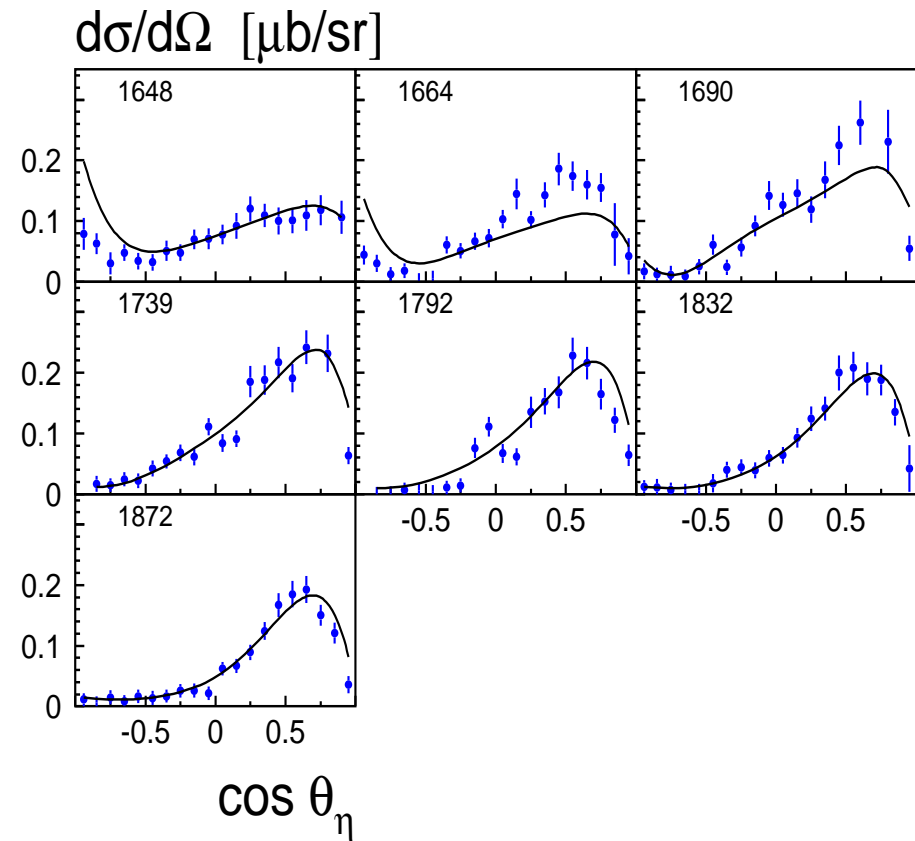
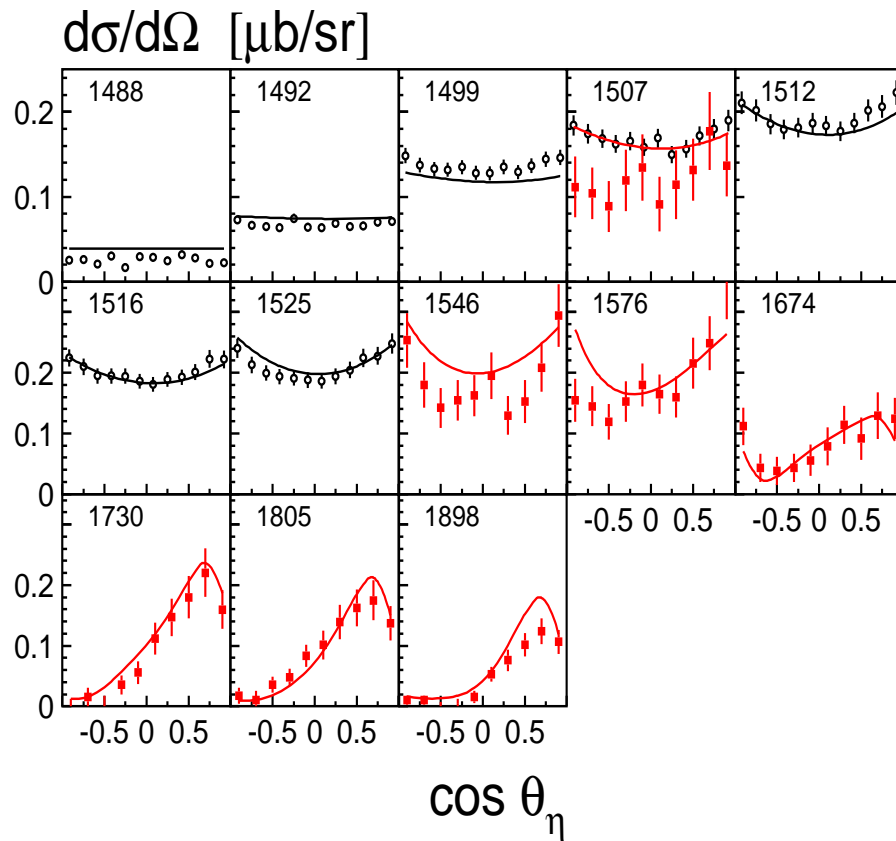
The description of the data with N/D-based approach



The description of the differential cross section and beam asymmetry in the selected energy region



The description of the $\pi p \rightarrow \eta p$ data with N/D-based approach

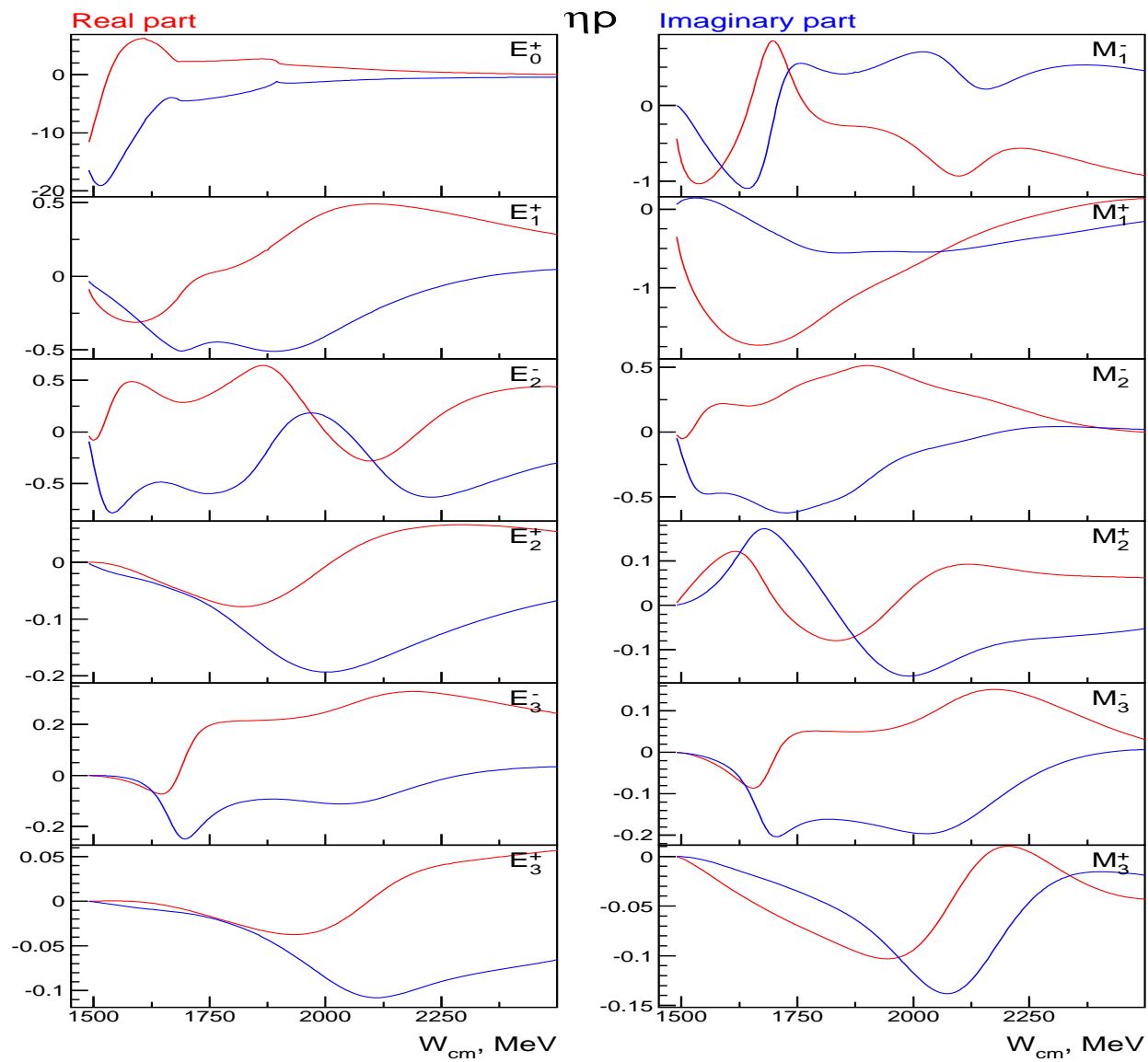


S. Prakhov et al., Phys. Rev. C72(2005) 015203

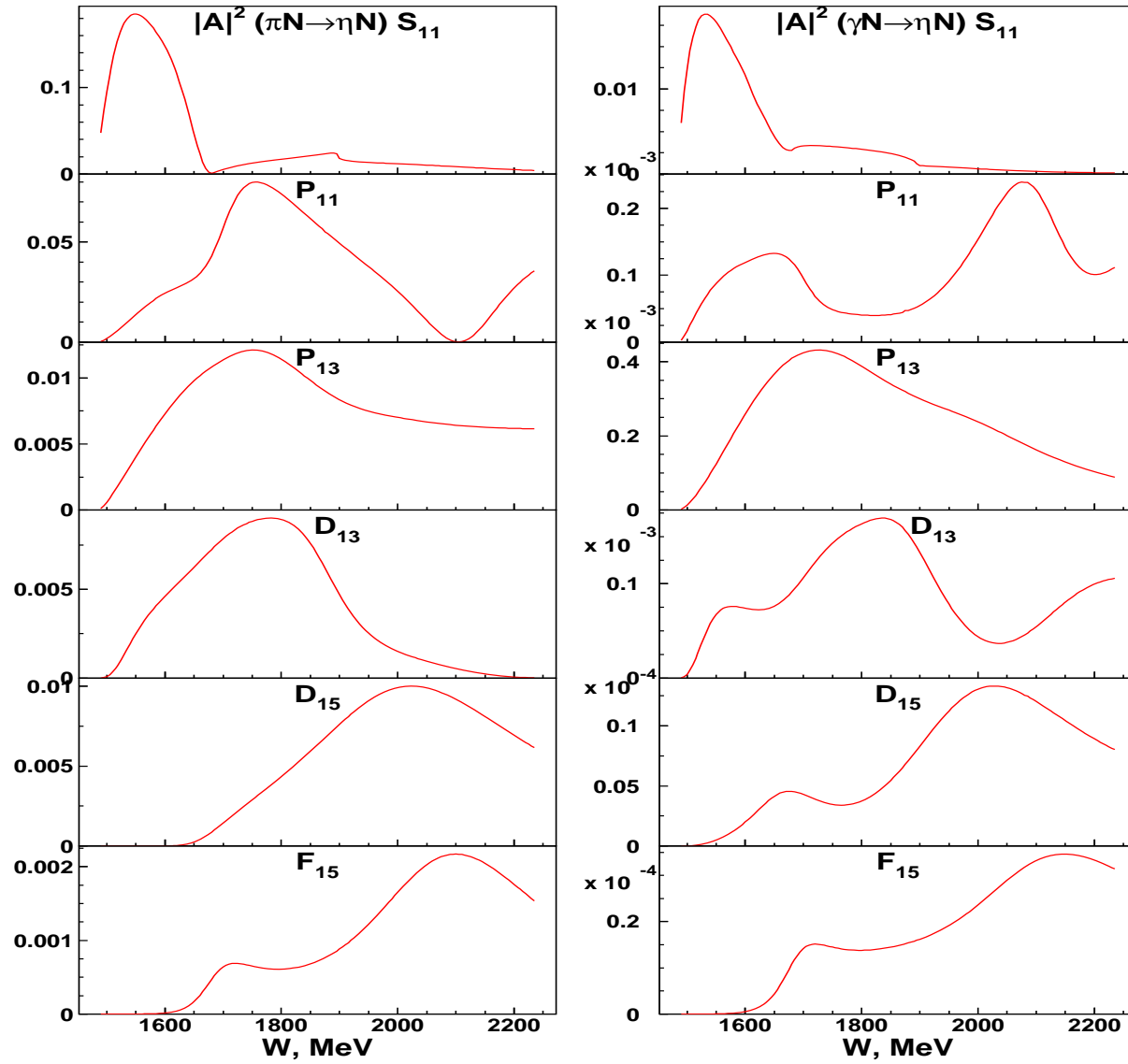
R.M. Brown et al. Nucl. Phys. B153, 89 (1979)

W. B. Richards et al., Phys. Rev. D1, 10 (1970)

$\gamma p \rightarrow \eta p$ multipoles



Amplitudes squared for the $\pi N \rightarrow \eta N$ and $\gamma p \rightarrow \eta p$ reactions



Residues in the pole for $\pi N \rightarrow \eta N$ (MeV) and $\gamma N \rightarrow \eta N$ ($10^{-3} \text{ GeV}^{-\frac{1}{2}}$) reactions amplitudes.

Res	$N(1535)1/2^-$	$N(1650)1/2^-$	$N(1895)1/2^-$
$\pi N \rightarrow \eta N$	(29±2) -(84±3)^o	-(22±3) (2±12)^o	(5±3) (25±20)^o
$\gamma p \rightarrow \eta p (\mathbf{E}, -A^{\frac{1}{2}})$	-(19±2) -(60±3)⁰	(4±0.3) (31±10)^o	(1.7±0.8) (20±20)^o
Res	$N(1440)1/2^+$	$N(1710)1/2^+$	$N(1880)1/2^+$
$\pi N \rightarrow \eta N$	-(20±6) (0±30)^o	(7±2) (54±15)^o	(12±5) (60±18)^o
$\gamma p \rightarrow \eta p (\mathbf{M}, A^{\frac{1}{2}})$	(5±2) (15±30)^o	(2.2±0.7) -(73±16)^o	(0.9±0.35) (60±17)^o
Res	$N(1520)3/2^-$	$N(1720)3/2^+$	$N(1900)3/2^+$
$\pi N \rightarrow \eta N$	(1.7±0.4) -(88±15)^o	(8.1±2.6) (45±10)^o	(2.9±0.8) (55±30)^o
$\gamma p \rightarrow \eta p (A^{\frac{1}{2}})$	(0.22±0.06) (86±18)^o	(8.7±3.0) (49±10)^o	-(1.3±0.4) -(40±25)^o
$\gamma p \rightarrow \eta p (A^{\frac{3}{2}})$	(1.18±0.20) -(74±17)^o	-(4.0±2.1) -(53±12)^o	(12±5) (14±10)^o
$\gamma p \rightarrow \eta p (\mathbf{E})$	-(0.91±0.15) -(71±16)⁰	-(4.3±1.7) (35±15)^o	(4.0±1.3) (6±12)^o
$\gamma p \rightarrow \eta p (\mathbf{M})$	-(0.44±0.10) -(79±14)⁰	-(6.1±2.7) -(98±14)^o	-(10.0±4) (17±12)^o