



## ***Robotics 2***

# **Dynamic model of robots: Newton-Euler approach**

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# Approaches to dynamic modeling

(reprise)



## energy-based approach (Euler-Lagrange)



## Newton-Euler method (balance of forces/torques)

- multi-body robot seen as a whole
- constraint (internal) reaction forces between the links are automatically eliminated: in fact, they do not perform work
- closed-form (symbolic) equations are directly obtained
- best suited for study of dynamic properties and **analysis** of control schemes
- dynamic equations written separately for each link/body
- **inverse dynamics in real time**
  - equations are evaluated in a **numeric** and **recursive** way
  - best for **synthesis** (=implementation) of model-based control schemes
- by elimination of reaction forces and back-substitution of expressions, we still get closed-form dynamic equations (identical to those of Euler-Lagrange!)



# Derivative of a vector in a moving frame

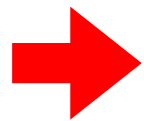
... from velocity to acceleration

$${}^0v_i = {}^0R_i {}^iv_i$$

$${}^0\dot{R}_i = S({}^0\omega_i) {}^0R_i$$

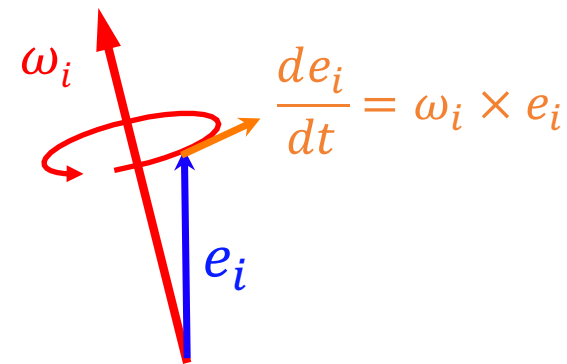
$${}^0\dot{v}_i = {}^0a_i = {}^0R_i {}^ia_i = {}^0R_i {}^i\dot{v}_i + {}^0\dot{R}_i {}^iv_i$$

$$= {}^0R_i {}^i\dot{v}_i + {}^0\omega_i \times {}^0R_i {}^iv_i = {}^0R_i ({}^i\dot{v}_i + {}^i\omega_i \times {}^iv_i)$$



$${}^ia_i = {}^i\dot{v}_i + {}^i\omega_i \times {}^iv_i$$

derivative of "unit" vector





# Dynamics of a rigid body

## ■ Newton dynamic equation

- **balance**: sum of forces = variation of **linear** momentum

$$\sum f_i = \frac{d}{dt}(mv_c) = m\dot{v}_c$$

## ■ Euler dynamic equation

- **balance**: sum of torques = variation of **angular** momentum

$$\begin{aligned}\sum \mu_i &= \frac{d}{dt}(I\omega) = I\dot{\omega} + \frac{d}{dt}(R\bar{I}R^T)\omega = I\dot{\omega} + (\dot{R}\bar{I}R^T + R\bar{I}\dot{R}^T)\omega \\ &= I\dot{\omega} + S(\omega)R\bar{I}R^T\omega + R\bar{I}R^T\cancel{S^T(\omega)}\omega = I\dot{\omega} + \omega \times I\omega\end{aligned}$$

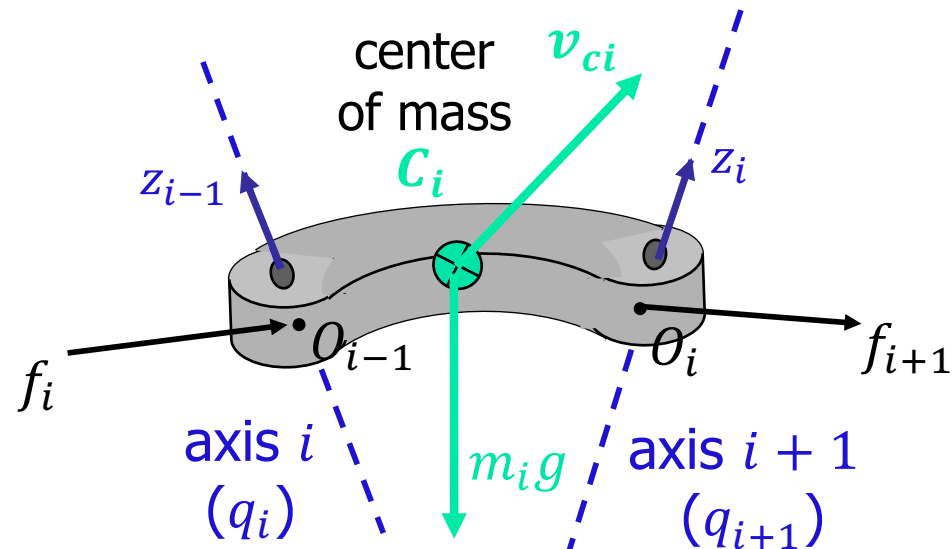
## ■ principle of **action and reaction**

- forces/torques: applied **by** body ***i*** **to** body ***i* + 1**  
= **—** applied **by** body ***i* + 1** **to** body ***i***



# Newton-Euler equations - 1

link  $i$



FORCES

$f_i$  force applied  
from link  $i - 1$  on link  $i$

$f_{i+1}$  force applied  
from link  $i$  on link  $i + 1$

$m_i g$  gravity force

all vectors expressed in the  
same RF (better RF <sub>$i$</sub> )

Newton equation

$$f_i - f_{i+1} + m_i g = m_i a_{ci}$$

N

linear acceleration of  $C_i$



# Newton-Euler equations - 2

link  $i$

## TORQUES

$\tau_i$  torque applied  
from link  $(i - 1)$  on link  $i$

$\tau_{i+1}$  torque applied  
from link  $i$  on link  $(i + 1)$

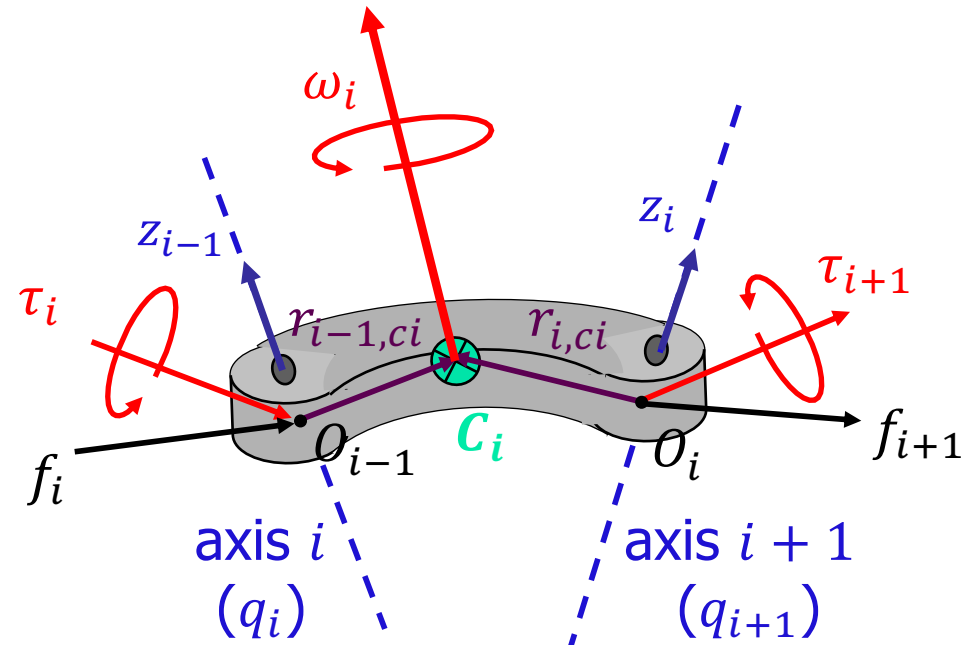
$f_i \times r_{i-1,ci}$  torque due to  $f_i$  w.r.t.  $C_i$

$-f_{i+1} \times r_{i,ci}$  torque due to  $-f_{i+1}$  w.r.t.  $C_i$

Euler equation

$$\tau_i - \tau_{i+1} + f_i \times r_{i-1,ci} - f_{i+1} \times r_{i,ci} = I_i \dot{\omega}_i + \omega_i \times (I_i \omega_i)$$

angular acceleration of body  $i$



all vectors expressed in  
the same RF (RF<sub>i</sub> !!)



# Forward recursion

## Computing velocities and accelerations

- “moving frames” algorithm (as for velocities in Lagrange)
- for simplicity, only revolute joints here  
(see [textbook](#) for the more general treatment)

### initializations

$${}^i\omega_i = {}^{i-1}R_i^T [{}^{i-1}\omega_{i-1} + \dot{q}_i {}^{i-1}z_{i-1}]$$

←  ${}^0\omega_0$

$${}^i\dot{\omega}_i = {}^{i-1}R_i^T [{}^{i-1}\dot{\omega}_{i-1} + \ddot{q}_i {}^{i-1}z_{i-1}] + {}^{i-1}\dot{R}_i^T [{}^{i-1}\omega_{i-1} + \dot{q}_i {}^{i-1}z_{i-1}]$$

AR 
$$= {}^{i-1}R_i^T [{}^{i-1}\dot{\omega}_{i-1} + \ddot{q}_i {}^{i-1}z_{i-1} + \dot{q}_i {}^{i-1}\omega_{i-1} \times {}^{i-1}z_{i-1}]$$

←  ${}^0\dot{\omega}_0$

$${}^i a_i = {}^{i-1}R_i^T {}^{i-1}a_{i-1} + {}^i\dot{\omega}_i \times {}^i r_{i-1,i} + {}^i\omega_i \times ({}^i\omega_i \times {}^i r_{i-1,i})$$

←  ${}^0a_0 - {}^0g$

$${}^i a_{ci} = {}^i a_i + {}^i\dot{\omega}_i \times {}^i r_{i,ci} + {}^i\omega_i \times ({}^i\omega_i \times {}^i r_{i,ci})$$

the gravity force term can be skipped in Newton equation, if added here



# Backward recursion

## Computing forces and torques

from  $N_i$   $\longrightarrow$  to  $N_{i-1}$  eliminated, if inserted in forward recursion ( $i=0$ ) initializations

$${}^i f_i = {}^i R_{i+1}^{i+1} f_{i+1} + m_i ({}^i a_{ci} - \cancel{{}^i g}) \quad \leftarrow f_{N+1} \quad \tau_{N+1}$$

$\downarrow$

**F/TR**

$${}^i \tau_i = {}^i R_{i+1}^{i+1} \tau_{i+1} + ({}^i R_{i+1}^{i+1} f_{i+1}) \times {}^i r_{i,ci} - {}^i f_i \times ({}^i r_{i-1,i} + {}^i r_{i,ci}) + {}^i I_i {}^i \dot{\omega}_i + {}^i \omega_i \times {}^i I_i {}^i \omega_i$$

from  $E_i$   $\longrightarrow$  to  $E_{i-1}$

at each step of this recursion, we have two vector equations ( $N_i + E_i$ ) at the joint providing  $f_i$  and  $\tau_i$ : these contain ALSO the **reaction forces/torques** at the joint axis  $\Rightarrow$  they should be "projected" next along/around this axis

**FP**  $u_i = \begin{cases} {}^i f_i^T {}^i z_{i-1} + F_{vi} \dot{q}_i & \text{for prismatic joint} \\ {}^i \tau_i^T {}^i z_{i-1} + F_{vi} \dot{q}_i & \text{for revolute joint} \end{cases}$

$\uparrow$  **generalized forces** (in rhs of Euler-Lagrange eqs)

$\nwarrow$  **add any dissipative term** (here, viscous friction only)

$\Rightarrow$  **N scalar equations at the end**





# Comments on Newton-Euler method

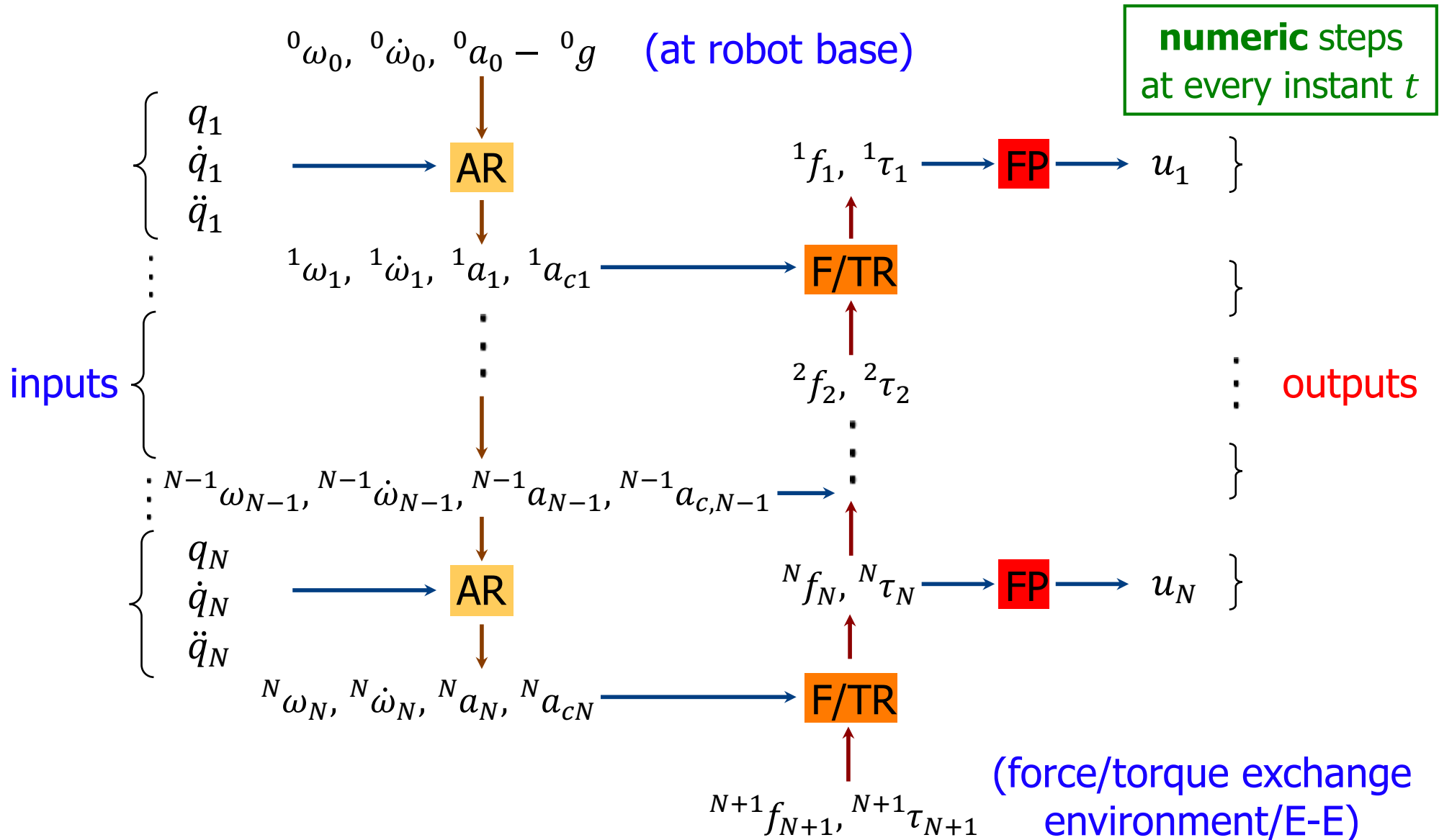
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- the previous forward/backward recursive formulas can be evaluated in symbolic or numeric form
  - **symbolic**
    - substituting expressions in a recursive way
    - at the end, a closed-form dynamic model is obtained, which is identical to the one obtained using Euler-Lagrange (or any other) method
    - there is **no** special convenience in using N-E in this way
  - **numeric**
    - substituting numeric values (numbers!) at each step
    - **computational complexity** of each step remains constant  $\Rightarrow$  grows **in a linear fashion** with the number  **$N$**  of joints ( **$O(N)$** )
    - strongly recommended for real-time use, especially when the number  **$N$**  of joints **is large**



# Newton-Euler algorithm

efficient computational scheme for inverse dynamics





# Matlab (or C) script

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general routine  $NE_{\alpha}(\arg_1, \arg_2, \arg_3)$

- data file (of a specific robot)
  - number  $N$  and types  $\sigma = \{0,1\}^N$  of joints (revolute/prismatic)
  - table of DH kinematic parameters
  - list of **ALL** dynamic parameters of the links (and of the motors)
- input
  - vector parameter  $\alpha = \{^0g, 0\}$  (presence or absence of gravity)
  - three ordered **vector arguments**
    - typically, samples of joint **position, velocity, acceleration** taken from a desired trajectory
- output
  - generalized force  $u$  for the **complete** inverse dynamics
  - ... or **single terms** of the dynamic model



# Examples of output

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- complete inverse dynamics

$$u = NE_0(q_d, \dot{q}_d, \ddot{q}_d) = M(q_d)\ddot{q}_d + c(q_d, \dot{q}_d) + g(q_d) = u_d$$

- gravity term

$$u = NE_0(q, 0, 0) = g(q)$$

- centrifugal and Coriolis term

$$u = NE_0(q, \dot{q}, 0) = c(q, \dot{q})$$

- $i$ -th column of the inertia matrix

$$u = NE_0(q, 0, e_i) = M_i(q)$$

$e_i$  =  $i$ -th column  
of identity matrix

- generalized momentum

$$u = NE_0(q, 0, \dot{q}) = M(q)\dot{q} = p$$



## A further example of output

- **factorization** of centrifugal and Coriolis term

$$u = NE_0(q, \dot{q}, 0) = c(q, \dot{q}) = S(q, \dot{q})\dot{q}$$

- for later use, what about a “mixed” velocity term?

$$S(q, \dot{q})\dot{q}_r \not\Rightarrow \begin{cases} u = NE_0(q, \dot{q}_r, 0) = S(q, \dot{q}_r)\dot{q}_r \\ u = NE_0(q, e_i \dot{q}_{ri}, 0) = S_i(q, e_i \dot{q}_{ri})\dot{q}_{ri} \end{cases} \text{no good!}$$

a)  $S(q, \dot{q})\dot{q}_r = S(q, \dot{q}_r)\dot{q}$  , when using Christoffel symbols

$$\text{b) } S(q, \dot{q} + \dot{q}_r)(\dot{q} + \dot{q}_r) = S(q, \dot{q})\dot{q} + S(q, \dot{q}_r)\dot{q}_r + 2S(q, \dot{q})\dot{q}_r$$

$$\Rightarrow u = \frac{1}{2} (NE_0(q, \dot{q} + \dot{q}_r, 0) - NE_0(q, \dot{q}, 0) - NE_0(q, \dot{q}_r, 0)) \\ = S(q, \dot{q})\dot{q}_r \quad (\text{i.e., with 3 calls of standard NE algorithm})$$

[Kawasaki et al., IEEE T-RA 1996]



# Modified NE algorithm

modified routine  $\widehat{NE}_\alpha(\text{arg}_1, \text{arg}_2, \text{arg}_3, \text{arg}_4)$  with 4 arguments

[De Luca, Ferrajoli, ICRA 2009]

$$\widehat{NE}_\alpha(x, y, y, z) = NE_\alpha(x, y, z) \quad \text{consistency property}$$

e.g.,  $u = \widehat{NE}_g(q, 0, 0, 0) = NE_g(q, 0, 0) = g(q)$

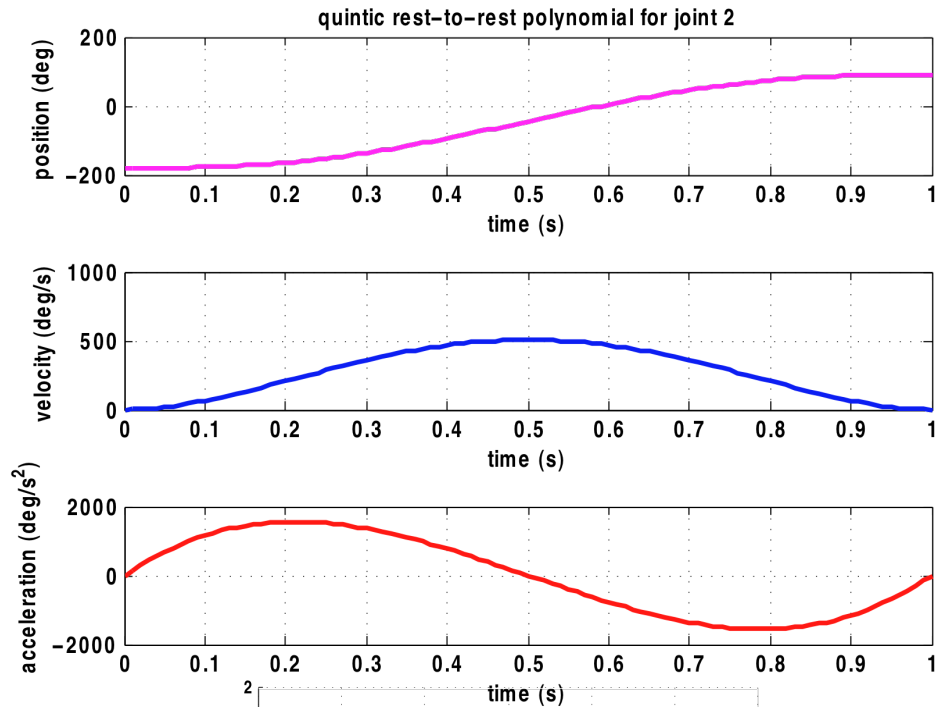
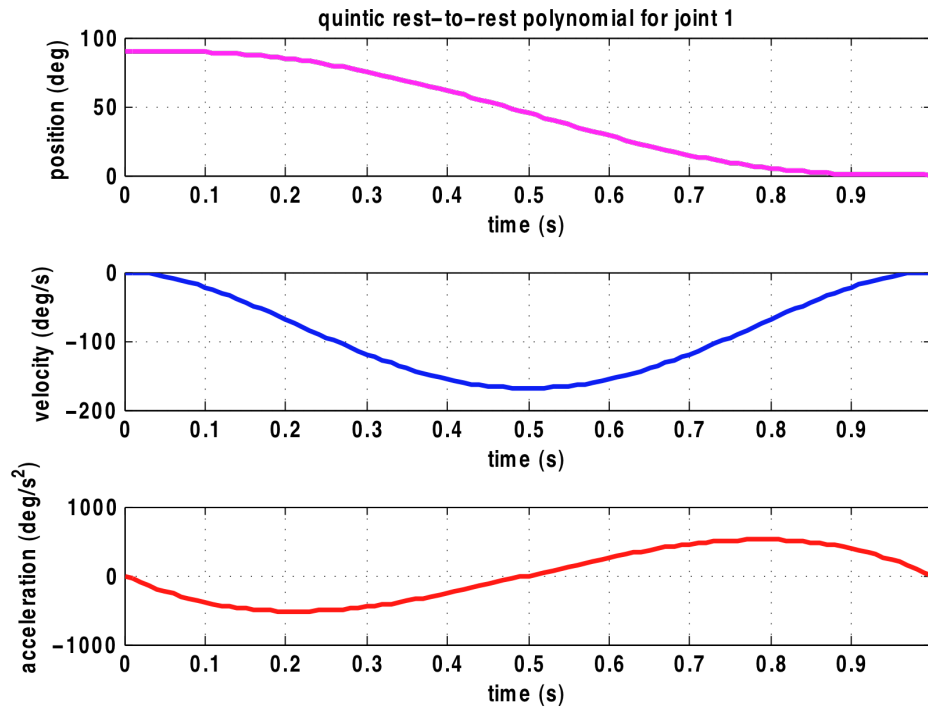
$$u = \widehat{NE}_0(q, \dot{q}, \dot{q}, 0) = NE_0(q, \dot{q}, 0) = c(q, \dot{q}) = S(q, \dot{q})\dot{q}$$

$$\Rightarrow u = \widehat{NE}_0(q, \dot{q}, \dot{q}_r, 0) = S(q, \dot{q})\dot{q}_r \quad \text{with } \dot{M} - 2S \text{ skew-symmetric} \\ \text{(i.e., with 1 call of modified NE algorithm)}$$

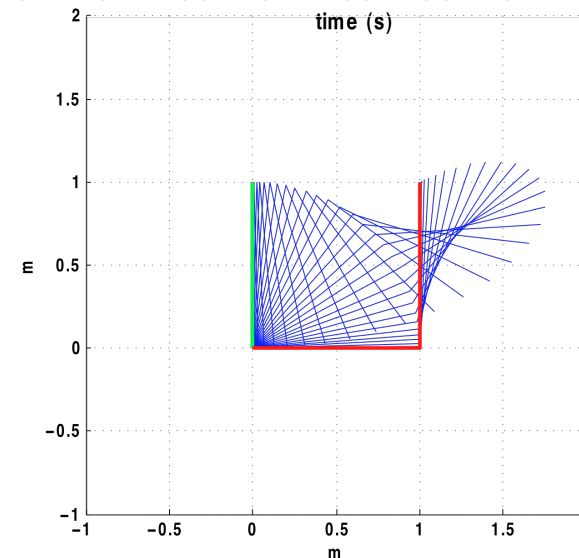
$$\Rightarrow u = \widehat{NE}_0(q, \dot{q}, e_i, 0) = S_i(q, \dot{q}) \\ \text{(i.e., the full matrix } S \text{ with } N \text{ calls of modified NE algorithm)}$$



# Inverse dynamics of a 2R planar robot

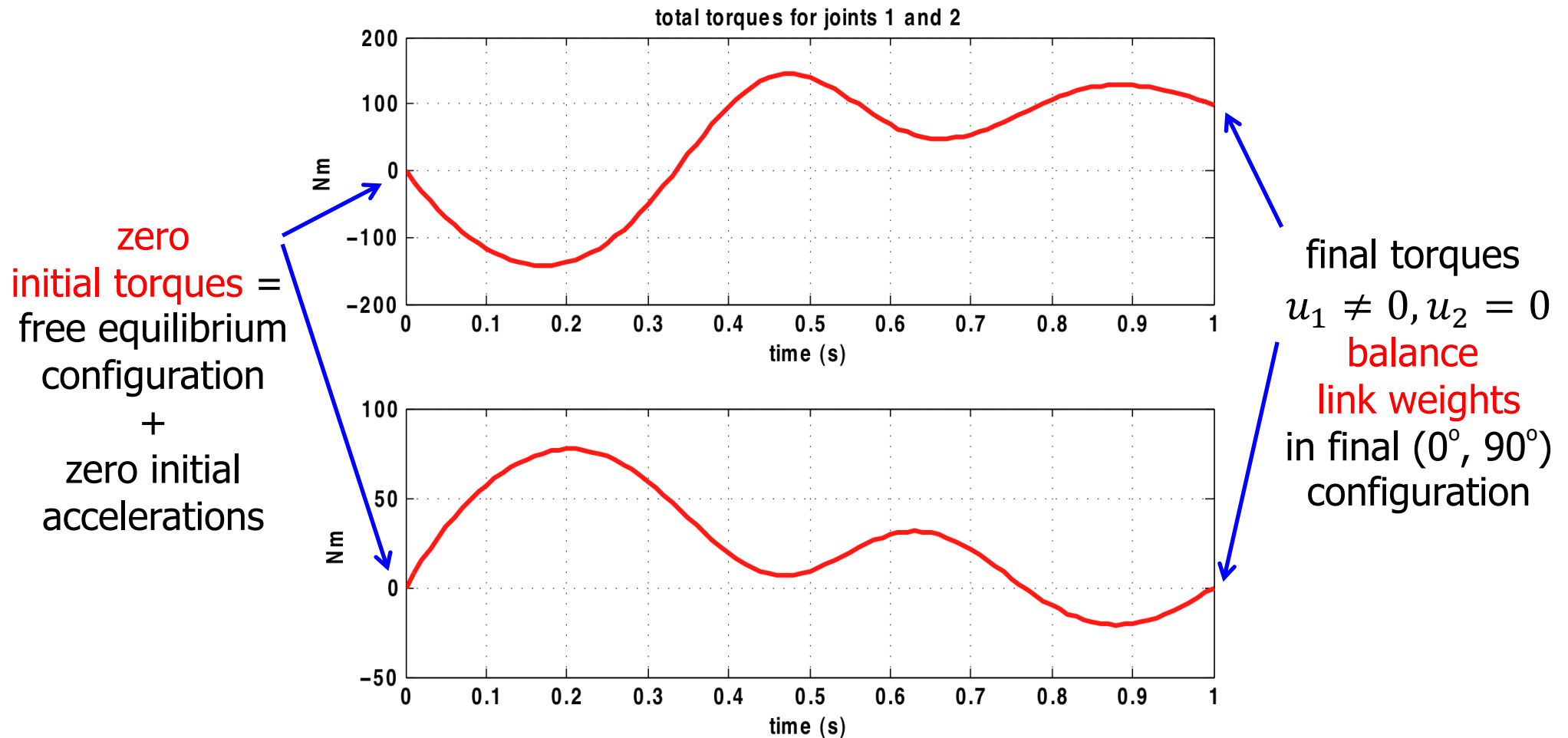


desired (smooth) joint motion:  
quintic polynomials for  $q_1, q_2$  with  
zero vel/acc boundary conditions  
from  $(90^\circ, -180^\circ)$  to  $(0^\circ, 90^\circ)$  in  $T = 1$  s





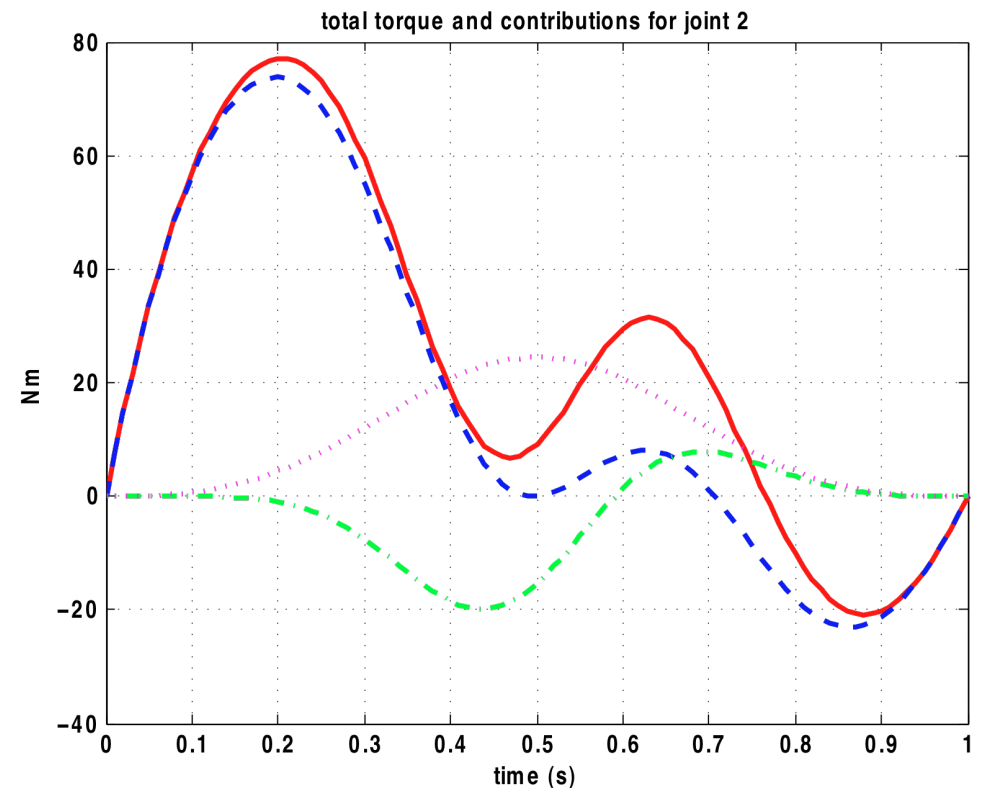
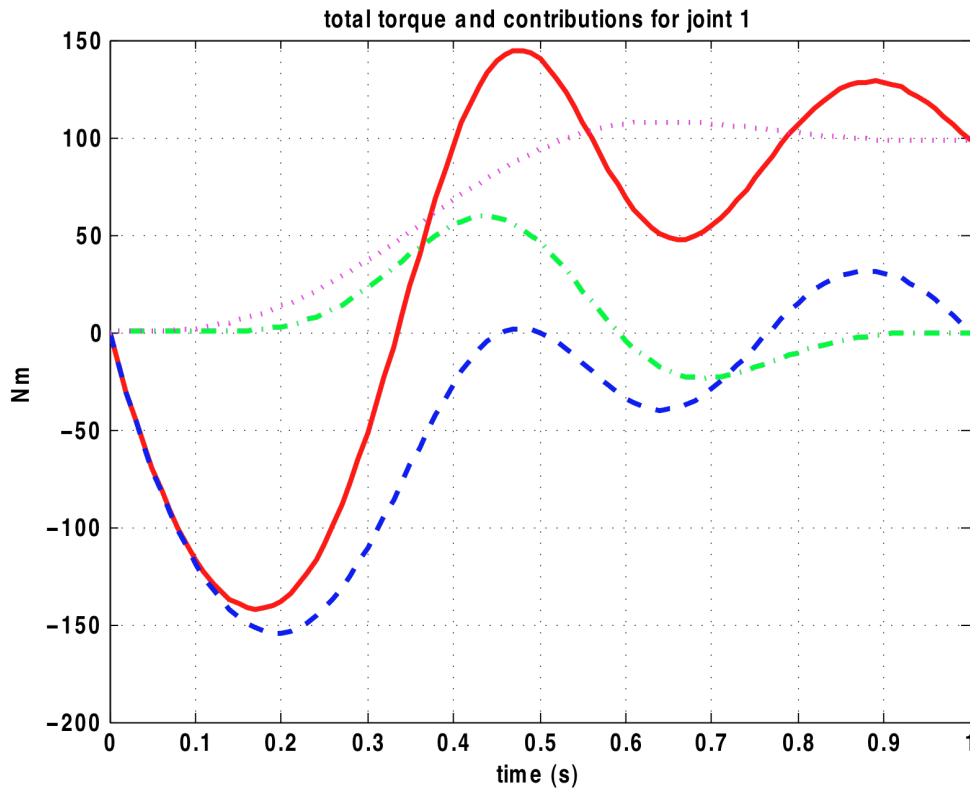
# Inverse dynamics of a 2R planar robot



motion in vertical plane (under gravity)  
both links are thin rods of uniform mass  $m_1 = 10$  kg,  $m_2 = 5$  kg



# Inverse dynamics of a 2R planar robot



torque contributions at the two joints for the desired motion

— = total, --- = inertial  
-.-.- = Coriolis/centrifugal, ..... = gravitational

# Use of NE routine for simulation

## direct dynamics



- numerical integration, at **current** state  $(q, \dot{q})$ , of
$$\ddot{q} = M^{-1}(q)[u - (c(q, \dot{q}) + g(q))] = M^{-1}(q)[u - n(q, \dot{q})]$$

- Coriolis, centrifugal, and gravity terms

$$n = NE_g(q, \dot{q}, 0) \quad \text{complexity } O(N)$$

- $i$ -th column of the inertia matrix, for  $i = 1, \dots, N$

$$M_i = NE_0(q, 0, e_i) \quad O(N^2)$$

- numerical inversion of inertia matrix

$$InvM = \text{inv}(M) \quad O(N^3) \quad \text{but with small coefficient}$$

- given  $u$ , integrate acceleration computed as

$$\ddot{q} = InvM * [u - n] \quad \longrightarrow \quad \text{new state } (q, \dot{q}) \quad \text{and repeat over time ...}$$