# Hamiltonian simulation algorithms for near-term quantum hardware

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### Problem statement

#### Definition

Hamiltonian simulation:

"Given a description of a Hamiltonian H, and evolution time t, some initial state  $|\psi(0)\rangle$  produce the final state  $|\psi(t)\rangle$  (to some error  $\epsilon$ )"

 "The Hamiltonian of a system is the sum of the kinetic energies of all the particles, plus the potential energy of the particles associated with the system

# Why is this a difficult problem?

#### Definition

We assume that the quantum state is loaded into memory

- a classical computer can't store the state efficiently
- a classical computer cannot produce a complete description of the state

# Schrödinger equation

#### Definition

$$H|\psi(t)\rangle = i\hbar \frac{\delta}{\delta t} |\psi(t)\rangle$$

integrate both sides:

$$|\psi(t)\rangle = e^{-iHt/\hbar} |\psi(0)\rangle$$

break H down into potential and kinetic energy:

$$\hat{V}$$
 (E<sub>pot</sub>),  $E_{kin} = p^2 \frac{1}{2m}$ 

$$e^{rac{-i\hat{V}t}{\hbar}}$$
 and  $e^{rac{-ip^2t}{2m\hbar}}$  don't commute,  $e^{-iHt/\hbar} 
eq e^{rac{-i\hat{V}t}{\hbar}} e^{rac{-ip^2t}{2m\hbar}}$ 

# Lie-Trotter product formula

#### Definition

$$e^{A+B} = \lim_{n \to \inf} (e^{A/n} e^{B/n})^n$$

- product formula: simulate the sum-terms of a Hamiltonian by simulating each one separetly for a small time slice
- H = A + B + C
- $U = e^{-i(A+B+C)t} = (e^{-iCt/r}e^{-iBt/r}e^{-iCt/r})^r$
- switching between kinetic and potential energy terms

we arrive at: 
$$e^{\frac{-i\hat{H}t}{\hbar}}=\lim_{N\to\infty}(e^{\frac{-ip^2t}{2m\hbar N}}e^{\frac{-\hat{V}(\hat{x})t}{\hbar N}})^N$$

#### Definition

#### Split operator

#### Algorithm:

- **①** Apply a half step of the potential propagator to  $\psi(0)$
- Apply the Fourier transform: momentum basis
- Apply a full step of the kinetic propagator on the momentum basis
- Apply the Inverse Fourier transform: back to coordinate basis
- Solution
  Apply the second half step of the potential propagator

this algorithm results from splitting the propagator, substituting into the Schrödinger equation and projecting onto a coordinate basis  $|x\rangle$ 

End

# Simulating a wavefunction

• blue line:  $\langle \Psi | \Psi \rangle$ 

ullet orange line: potential energy,  $\hat{V}$ 

End

# Simulating a wavefunction

• blue line:  $\langle \Psi | \Psi \rangle$ 

ullet orange line: potential energy,  $\hat{V}$ 

### Problem statement

#### Definition

Find explicit gate count for near-term simulations

- gate count of near-term quantum hardware is very limited
- algorithms need to be optimized to use as few gates as possible
- even low gate counts of 250 gates are still too costly

# **Topics**

- Optimal fermion encoding
- Subscircuits: programming below the circuit model
- Non-asymptotic Trotter bounds

# Fermi-Hubbard Model

#### Hamiltonian:

$$H_{FH} = \sum_{i=1}^{N} h_{on-site}^{(i)} + \sum_{i < j, \sigma} h_{hopping}^{(i,j,\sigma)}$$

- Model existed before quantum computing
- Most practical applications involve fermionic system
- Classical approaches can be used as benchmark

# Sub-Circuit Model

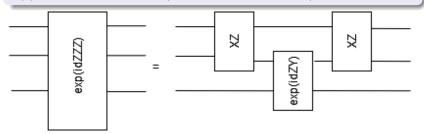
- Assumption: hardware can perform CNOT, Z, H .... gates
- Layer of abstraction: overhead
- Quantum circuit:  $C = \prod_{l=1}^{L} U_{l} V_{j}$
- $U_l = \prod_{i \in Q} u_i^l$  a layer of arbitrary single qubit unitary gates  $V_l = \prod_{ij \in \Gamma} v_{ij}(t_{ij}^l)$  a layer of non-overlapping, two-qubit gates
- two-qubit unitary gates:  $v_{ij}(t) = e^{ith_{ij}}, h_{ij} = Z_i Z_j$
- $h_{ij} = \sigma_i \sigma_j$  since  $\sigma_i \sigma_j = Z_i Z_j(*)$
- the cost  $T_{cost}(C) = L$ , L = circuit depth

#### Definition

how to decompse local Trotter steps  $e^{i\delta h}$ 

Let 
$$U(t) = e^{itH}, H = \frac{1}{2i}[h_1, h_2]$$

Let  $U(t) = e^{itH}$ ,  $H = \frac{1}{2i}[h_1, h_2]$   $U(t) = e^{it_1h_1}e^{it_2h_2}e^{it_2h_1}e^{it_1h_2}$  (Lemma 7, supplementary)



# Error models

- Per-gate error model
- Per-time error model
  - "Error budget required to execute a circuit"
- Single-qubit gates are considered free
- Obtain tight error expression  $\epsilon_p(\delta)$
- Guarantee  $\epsilon_p(\delta) < \epsilon_{target}$  by inverting the expression and derive a maximum possible Trotter step  $\delta_0 = \delta_0(\epsilon_{target})$

Table 1 Per-gate run-times.					
Fermion encoding	Trotter bounds	Standard decomposition	Sub-circuit decomposition		
VC	Ref. 23, Prop. F.4.	1,243,586	977,103		
	Analytic	121,478	95,447		
	Numeric	5391	4236		
Compact	Analytic	98,339	72,308		
	Numeric	4364	3209		

A comparison of the run-time  $T_{cost}$  for lattice size  $L \times L$  with L = 5, overall simulation time T = 7and target Trotter error  $\epsilon_{\text{target}} = 0.1$ , with  $\Lambda = 5$  fermions and coupling strengths  $|u|, |v| \le r = 1$ . Obtained by minimising over product formulas up to 4th order. T<sub>cost</sub> = circuit depth for pergate error model. In either gate decomposition case—standard and sub-circuit—we account single-qubit rotations as a free resource as explained in the Introduction; the value of  $T_{\rm cost}$ depends only on the two-gubit gates/interactions. Two-gubit unitaries are counted by unit time per gate in the per-gate error model. Here compact and VC denote the choice of fermionic encoding.

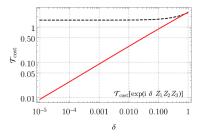
source: Hamiltonian simulation algorithms for near-term quantum hardware, Clinton et al.

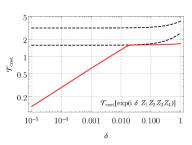
Table 2 Per-time run-times.				
Fermion encoding	Trotter bounds	Standard decomposition	Sub-circuit decomposition	
VC	Ref. <sup>23</sup> , Prop. F.4.	976,710	59,830	
	Analytic	95,409	17,100	
	Numeric	4234	1669	
compact	Analytic	77,236	1686	
	Numeric	3428	259	

A comparison of the run-time  $T_{cost}$  for lattice size  $L \times L$  with L = 5, overall simulation time T = 7and target Trotter error  $\epsilon_{\text{target}} = 0.1$ , with  $\Lambda = 5$  fermions and coupling strengths  $|u|, |v| \le r = 1$ . Obtained by minimising over product formulas up to 4th order.  $T_{cost} = T_{cost}(P_n(\delta_0)^{T/\delta_0})$  for per-time error model. In either gate decomposition case—standard and sub-circuit—we account single-qubit rotations as a free resource; the value of  $T_{cost}$  depends only on the two-qubit gates/interactions. Two-qubit unitaries are counted by their respective pulse lengths. Here compact and VC denote the choice of fermionic encoding.

source: Hamiltonian simulation algorithms for near-term quantum hardware, Clinton et al.

# Gate decomposition cost





source: Hamiltonian simulation algorithms for near-term quantum hardware. Clinton et al.

# Outlook

- Further optimization needed for quantum hardware
- Tighter error bounds might be reached
- Standard circuit decomposition will stay unfeasible on real hardware for some time
- The sub-circuit model might enable some algorithm to run on NISQ hardware

Hamiltonian simulation Implementation amiltonian simulation algorithms for near-term quantum hardware Results Outlook End

Thank you for your attention!