

# Ayudantía 2. Física Estadística. 2.2 Pathria

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- ① dw del espacio fase de una partícula es invariante bajo una transformación de coordenadas cartesianas  $(x, y, z, p_x, p_y, p_z)$  a coordenadas esféricas.

Solución

$$x = r \cos \theta \sin \phi, \quad y = r \sin \theta \sin \phi, \quad z = r \cos \theta$$

$$|M| = \left| \frac{\partial(x, y, z, p_x, p_y, p_z)}{\partial(r, \theta, \phi, p_r, p_\theta, p_\phi)} \right| = \left| \frac{\partial(x, y, z)}{\partial(r, \theta, \phi)} \right| \left| \frac{\partial(p_x, p_y, p_z)}{\partial(p_r, p_\theta, p_\phi)} \right|$$

$$= r^2 \sin \theta \frac{\partial(p_x, p_y, p_z)}{\partial(p_r, p_\theta, p_\phi)}$$

$$p_x = m\dot{x} = m\dot{r} \sin \theta \cos \phi + m r \cos \theta \cos \phi \dot{\theta} - m r \sin \theta \sin \phi \dot{\phi}$$

$$p_y = m\dot{y} = m\dot{r} \sin \theta \sin \phi + m r \cos \theta \sin \phi \dot{\theta} + m r \sin \theta \cos \phi \dot{\phi}$$

$$p_z = m\dot{z} = m\dot{r} \cos \theta - m r \sin \theta \dot{\theta}$$

Recordamos que en esféricas:

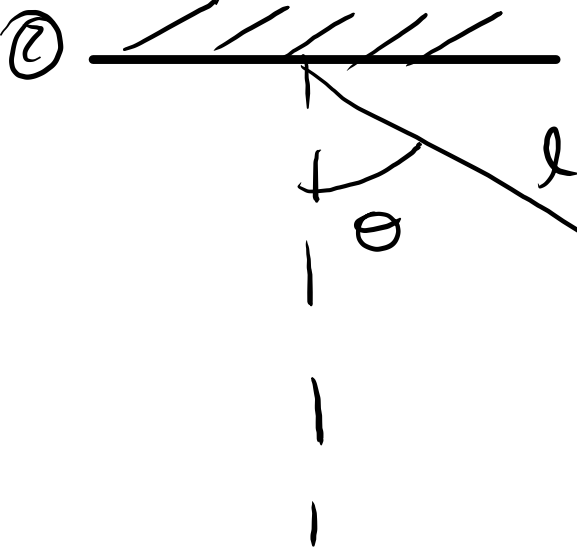
$$\underline{p_r = m\dot{r}, \quad p_\theta = m r \dot{\theta}, \quad p_\phi = m r^2 \sin^2 \theta \dot{\phi}}$$

$$p_x = \sin \theta \cos \phi p_r + \frac{\cos \theta \cos \phi}{r} p_\theta - \frac{\sin \phi}{r \sin \theta} p_\phi$$

$$p_y = \sin \theta \sin \phi p_r + \frac{\cos \theta \sin \phi}{r} p_\theta + \frac{\cos \phi}{r \sin \theta} p_\phi$$

$$p_z = \cos \theta p_r - \frac{\sin \theta}{r} p_\theta$$

$$\frac{\partial(p_x, p_y, p_z)}{\partial(p_r, p_\theta, p_\phi)} = \begin{vmatrix} \frac{\partial p_x}{\partial p_r} & \frac{\partial p_x}{\partial p_\theta} & \frac{\partial p_x}{\partial p_\phi} \\ \frac{\partial p_y}{\partial p_r} & \frac{\partial p_y}{\partial p_\theta} & \frac{\partial p_y}{\partial p_\phi} \\ \frac{\partial p_z}{\partial p_r} & \frac{\partial p_z}{\partial p_\theta} & \frac{\partial p_z}{\partial p_\phi} \end{vmatrix} = \frac{1}{r^2 \sin \theta} \Rightarrow |M| = 1$$



$$K = \frac{1}{2} m l^2 \dot{\theta}^2 \quad \leftarrow \theta \ll 1.$$

$$U = mgl(1 - \cos\theta) \approx mgl \frac{\theta^2}{2}$$

$$H = E = \frac{p_\theta^2}{2ml^2} + mgl \frac{\theta^2}{2}$$

$$1 = \frac{p_\theta^2}{a^2} + \frac{\theta^2}{b^2}, \quad \left. \begin{aligned} a^2 &\equiv 2ml^2 E \\ b^2 &\equiv 2E / mgl \end{aligned} \right\}$$

$$A(\text{area}) = \pi ab = 2\pi E \sqrt{\frac{l}{g}} = E\tau, \quad \tau = 2\pi \sqrt{\frac{l}{g}} \quad \leftarrow \frac{1}{\omega}$$

$$4 \quad x(t) = A \cos(\omega_0 t + \phi)$$

$$P(x) dx = P(x, x+dx)$$

Si  $T = \frac{2\pi}{\omega_0}$  es el tiempo requerido para que la partícula complete un ciclo,

$$P(x) dx = \frac{2 \frac{dt}{T}}{1} : \text{fracción del periodo que la partícula pasa en el intervalo } (x, x+dx)$$

Supongamos que  $\phi = \pi/2$ , entonces  $x(t) = A \sin(\omega_0 t)$

$$v = \frac{dx}{dt} \Rightarrow dt = \frac{dx}{v}, v = \omega_0 A \cos(\omega_0 t) = \omega_0 \sqrt{A^2 - x^2(t)}$$

$$P(x) dx = \frac{2 \frac{dt}{T}}{1} = \frac{2}{T} \frac{dx}{\omega_0 \sqrt{A^2 - x^2(t)}} = \frac{2\omega_0}{2\pi \omega_0 \sqrt{A^2 - x^2(t)}} = \frac{1}{\pi} \frac{dx}{\sqrt{A^2 - x^2(t)}}$$

$$\int_{-A}^A P(x) dx \stackrel{!}{=} 1.$$

$$\frac{1}{\pi} \int_{-A}^A \frac{dx}{\sqrt{A^2 - x^2(t)}} \stackrel{!}{=} 1$$

$$= \frac{1}{\pi} \arctan \left( \frac{x}{\sqrt{A^2 - x^2(t)}} \right) \Big|_{-A}^A = \frac{1}{\pi} \left( \frac{\pi}{2} - \left( -\frac{\pi}{2} \right) \right) = 1.$$

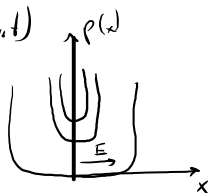
De la conservación de la energía,

$$E = \frac{mv^2}{2} + \frac{1}{2} m \omega_0^2 x^2$$

$$= \frac{1}{2} m \omega_0^2 A^2 \cos^2(\omega_0 t) + \frac{1}{2} m \omega_0^2 A^2 \sin^2(\omega_0 t)$$

$$E = \frac{1}{2} m \omega_0^2 A^2$$

$$\boxed{A^2 = \frac{2E}{m\omega_0^2}}$$



$$P(x) = \frac{1}{\pi} \left( \frac{k}{2E - kx^2} \right)^{1/2} \quad k = m\omega_0^2$$

$$P_n(x) = \frac{1}{\pi} \left( \frac{k}{2E_n - kx^2} \right)^{1/2}$$

$$E_n = \hbar \omega \left( n + \frac{1}{2} \right)$$

③  $\bar{p} = \langle \psi_k | \bar{\sigma} | \psi_k \rangle$ ,  $\bar{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$  Pathria  
Libgen

$$|\psi_k\rangle = a_+ \phi_{k, +1/2} + a_- \phi_{k, -1/2}$$

$$p_x = (a_+^*, a_-^*) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} = (a_+^* a_- + a_-^* a_+) \quad \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$p_y = (a_+^*, a_-^*) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} = i(a_-^* a_+ - a_+^* a_-) \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$p_z = (a_+^*, a_-^*) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} = (a_+^* a_+ - a_-^* a_-) \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$|\bar{p}| = \sqrt{p_x^2 + p_y^2 + p_z^2} = 1 \quad (a)$$

(b)  $\bar{p} = \underline{p_+ \langle \phi_{k, +1/2} | \bar{\sigma} | \phi_{k, +1/2} \rangle} + \underline{p_- \langle \phi_{k, -1/2} | \bar{\sigma} | \phi_{k, -1/2} \rangle}$

$$= p_+ \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + p_- \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ p_+ - p_- \end{pmatrix}, \quad p_+ = p_- = \frac{1}{2}.$$

$$\bar{p} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}, \quad |\bar{p}| = 0.$$

$$\rho_{ss'} = \begin{pmatrix} p_+ & 0 \\ 0 & p_- \end{pmatrix}, \quad \underline{p_+ + p_- = 1}$$

$(\bar{\sigma}, \mathbb{1})$  : form an new base.

$$|\psi_{k,s} \times \psi_{k,s'}| = \begin{pmatrix} a_+^* + a_+ & a_-^* - a_+ \\ a_+^* + a_- & a_-^* - a_- \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 + p_z & p_x - i p_y \\ p_x + i p_y & 1 - p_z \end{pmatrix}.$$

$$\underline{S} \quad (a) \quad \langle \hat{O} \rangle = \text{tr} [\hat{\rho}(t) \hat{O}] \quad \hat{\rho}(t) = \sum_i p_i |\psi_i(t)\rangle \langle \psi_i(t)|$$

$$\sum_i p_i = 1$$

$$p(t) \quad \{|\vec{r}\rangle\}$$

$$\hat{\rho}_{\vec{r}}(t) \rightarrow \langle \vec{r}' | \hat{\rho}(t) | \vec{r} \rangle$$

$$\hat{\phi}(p_1, p_2) = \hat{\rho}_{\vec{p}}(t) \rightarrow \langle \vec{p}' | \hat{\rho}(t) | \vec{p} \rangle$$

$$(b) \rightarrow \hat{\phi}(p_1, p_2) = f(p) \delta_{p_1 p_2}$$

$$\begin{aligned} \langle \vec{r}' | \hat{\rho}(t) | \vec{r} \rangle &= \sum_{\vec{p}', \vec{p}} \langle \vec{r}' | \vec{p} \times \vec{p}' | \hat{\rho}(t) | \vec{p} \times \vec{p} | \vec{r} \rangle \\ &= \frac{1}{V} \sum_{\vec{p}', \vec{p}} \phi(\vec{p}', \vec{p}) \exp \{ i(\vec{r}' \cdot \vec{p}' - \vec{r} \cdot \vec{p}) \} \\ &= \frac{1}{V} \sum_{\vec{p}', \vec{p}} f(p) \delta_{\vec{p} \vec{p}'} \exp \{ i(\vec{r}' \cdot \vec{p}' - \vec{r} \cdot \vec{p}) \} \\ &= \frac{1}{V} \sum_{\vec{p}} f(p) \exp \{ i(\vec{r}' - \vec{r}) \cdot \vec{p} \} \end{aligned}$$

$$\langle \vec{r} | \hat{\rho}(t) | \vec{r} \rangle = \frac{1}{V} \sum_p f(p) \quad \text{es una constante.}$$

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