

Problem Set 4

To be handed in by 11.30am, April 24th

1. Show that the anti-commutation relations for creation and annihilation operators in eq. (10.9) lead to the correct anti-commutation relations for the fields $\psi_\alpha(\mathbf{x}, t)$ and $\pi_\beta(\mathbf{x}, t)$. It will be useful to use the relationships in eq. (10.10). [8]

2. Verify the form of the Hamiltonian

$$H = \int \frac{d^3\mathbf{k}}{(2\pi)^3} \frac{1}{2E} E \sum_{s=1,2} [b^\dagger(s, \mathbf{k})b(s, \mathbf{k}) + d^\dagger(s, \mathbf{k})d(s, \mathbf{k})].$$

Explain why commutation relations for ψ and π and therefore for b and d would have led to a disaster. [7]

3. By considering a photon moving along the z -axis show that

$$\langle \chi | a^\dagger(3, \mathbf{k})a(3, \mathbf{k}) - a^\dagger(0, \mathbf{k})a(0, \mathbf{k}) | \chi \rangle = 0$$

results from

$$\langle \chi | \partial_\mu A^\mu | \chi \rangle = 0. \quad [6]$$

4. Derive the trace results

$$\begin{aligned} \text{Tr}(\gamma^{\mu_1} \dots \gamma^{\mu_n}) &= 0 \quad \text{for } n \text{ odd} \\ \text{Tr}(\gamma^{\mu_1} \dots \gamma^{\mu_n}) &= g^{\mu_1 \mu_2} \text{Tr}(\gamma^{\mu_3} \dots \gamma^{\mu_n}) - g^{\mu_1 \mu_3} \text{Tr}(\gamma^{\mu_2} \gamma^{\mu_4} \dots \gamma^{\mu_n}) + \dots \\ &\quad + g^{\mu_1 \mu_n} \text{Tr}(\gamma^{\mu_2} \dots \gamma^{\mu_{n-1}}) \\ \text{Tr}(\not{a} \not{b}) &= 4 a \cdot b, \\ \text{Tr}(\not{a} \not{b} \not{c} \not{d}) &= 4(a \cdot b c \cdot d - a \cdot c b \cdot d + a \cdot d b \cdot c). \end{aligned}$$

(Hint: for the first one use $(\gamma^5)^2 = 1$.) [7]

5. Prove the result in eq. (11.13). It will be helpful first to prove

$$\begin{aligned} \gamma^\alpha \gamma^\mu \gamma_\alpha &= -2\gamma^\mu \\ \gamma^\alpha \gamma^\mu \gamma^\nu \gamma_\alpha &= 4g^{\mu\nu} \\ \gamma^\alpha \gamma^\mu \gamma^\nu \gamma^\rho \gamma_\alpha &= -2\gamma^\rho \gamma^\nu \gamma^\mu, \end{aligned}$$

and you may make use of the result in Section 11.1. [10]

6. Show that the spin summed/averaged squared matrix element for Compton scattering in the massless limit is given by

$$|\mathcal{M}_{fi}|^2 = 2e^4 \left(-\frac{u}{s} - \frac{s}{u} \right).$$

You may use the result that the matrix element squared for the s -channel diagram is $2e^4 \left(-\frac{u}{s} \right)$, but must derive explicitly the interference term contribution. Evaluate the total cross section using the expressions in the centre of mass frame in Section 8.4. Why does this create a problem? [9]

7. Why is it not possible in QCD to use the on-shell scheme? [3]