3.4.1 Hadamard Test

In a previous section we discussed the *Swap Test* to measure the similarity between two unknown states $|\psi\rangle$ and $|\phi\rangle$ without having to measure these state directly. Note that we use the words *similarity* and *overlap* interchangeably. We derived that for the swap test circuit, the probability $Pr(|0\rangle)$ of measuring state $|0\rangle$ was:

$$Pr(|0\rangle) = \frac{1}{2} + \frac{1}{2} \langle \psi | \phi \rangle^2.$$

We can invert this and express the overlap as:

$$\langle \psi | \phi \rangle^2 = 1 - Pr(|0\rangle).$$

In this section we present another test of this nature, the *Hadamard Test*.

Both the Swap test and the Hadamard test can be visualized with an analogy using real-valued vectors. The numbers come out differently, but the principle is the same. Think about how we compute the inner product of the sum $(\vec{a} + \vec{b})$ of two *normalized*, real-valued vectors \vec{a} and \vec{b} (they have to be normalized, else the math doesn't work out):

$$(\vec{a} + \vec{b})^{T}(\vec{a} + \vec{b}) = \sum_{i} (a_{i} + b_{i})^{2}$$

$$= \sum_{i} a_{i}^{2} + \sum_{i} b_{i}^{2} + 2 \sum_{i} a_{i}b_{i}$$

$$= 2 + 2\vec{a}^{T}\vec{b}.$$
(3.1)

Note the three extreme cases where \vec{a} and \vec{b} point in the same direction, are orthogonal, or are anti-parallel. In these three cases, Equation 3.1 yields:

parallel:
$$\vec{a} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 and $\vec{b} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, then $\vec{a}^T \vec{b} = 1$.
orthogonal: $\vec{a} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ and $\vec{b} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, then $\vec{a}^T \vec{b} = 0$.
anit-parallel: $\vec{a} = \begin{bmatrix} -1 \\ 0 \end{bmatrix}$ and $\vec{b} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, then $\vec{a}^T \vec{b} = -1$.

Now let's apply this principle for the Hadamard test. Remember how the Swap test used two quantum registers to hold the states ψ and ϕ ? The Hadamard test is different, it uses only one quantum register which will hold the superposition of the two states $|a\rangle$ and $|b\rangle$ for which we want to determine overlap. Hence, as a precondition, we need to prepare this initial state:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle|a\rangle + |1\rangle|b\rangle).$$
 (3.2)

How can we generate such a state? First, let's see how the partial expressions

look as state vectors:

$$|0\rangle|a
angle = egin{bmatrix} 1 \ 0 \end{bmatrix} \otimes egin{bmatrix} a_0 \ a_1 \end{bmatrix} = egin{bmatrix} a_0 \ a_1 \ 0 \ 0 \end{bmatrix},$$

and

$$|1
angle|b
angle = egin{bmatrix} 0 \ 1 \end{bmatrix} \otimes egin{bmatrix} b_0 \ b_1 \end{bmatrix} = egin{bmatrix} 0 \ 0 \ b_0 \ b_1 \end{bmatrix}.$$

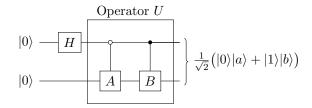
As a vector, state $|\psi\rangle$ in Equation 3.2 would be:

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle|a\rangle + |1\rangle|b\rangle) = \frac{1}{\sqrt{2}} \begin{bmatrix} a_0 \\ a_1 \\ b_0 \\ b_1 \end{bmatrix}.$$

We need operators A and B that produce the states $|a\rangle$ and $|b\rangle$ when applied to state $|0\rangle$. Note that A and B have to be a unitary and that we name the matrix element in a way to produce the desired output vectors:

$$\begin{split} A\left|0\right\rangle &=A\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}a_0&a_2\\a_1&a_3\end{bmatrix}\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}a_0\\a_1\end{bmatrix}=\left|a\right\rangle,\\ B\left|0\right\rangle &=B\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}b_0&b_2\\b_1&b_3\end{aligned}\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}b_0\\b_1\end{bmatrix}=\left|b\right\rangle. \end{split}$$

To construct the circuit, we use qubit 0 as an ancilla and qubit 1 should hold the superposition of states a and b. We can use an initial Hadamard gate to produce an equal superposition of $|0\rangle$ and $|1\rangle$ with which we control the gates A and B on qubit 1, as shown in this circuit:



Note that in the literature the two controlled operators A and B are often referred to as a combined single operator U. Let's verify this in code! First we create random unitary operators A and B and apply them to state $|0\rangle$ to extract the relevant state components a_0, a_1, b_0 , and b_1 :

```
def make_rand_operator():
    """Make a unitary operator U, derive u0, u1."""

U = ops.Operator(unitary_group.rvs(2))
    if not U.is_unitary():
        raise AssertionError('Error: Generated non-unitary operator')
    psi = U(state.bitstring(0))
    u0 = psi[0]
    u1 = psi[1]
    return (U, u0, u1)

def hadamard_test():
    """Perform Hadamard Test."""

A, a0, a1 = make_rand_operator()
B, b0, b1 = make_rand_operator()
```

With these parameters, we can construct the state in two different ways. First we compute it explicitly, following Equation 3.2:

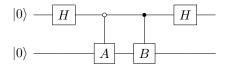
```
# Construct the desired end state psi as an explicit expression. # psi = 1/sqrt(2)(|0>|a> + |1>|b>) psi = (1 / cmath.sqrt(2) * (state.bitstring(0) * state.State([a0, a1]) + state.bitstring(1) * state.State([b0, b1])))
```

To compare, we construct the state with a circuit and confirm that the result matches the closed form above:

```
# Let's see how to make this state with a circuit.
qc = circuit.qc('Hadamard test - initial state construction.')
qc.reg(2, 0)
qc.h(0)
qc.applyc(A, [0], 1) # Controlled-by-0
qc.applyc(B, 0, 1) # Controlled-by-1

# The two states should be identical!
if not np.allclose(qc.psi, psi):
    raise AssertionError('Incorrect result')
```

Now let's add another Hadamard gate to the ancilla qubit 0:



This changes the state to:

$$\begin{split} \frac{1}{\sqrt{2}}H\big(|0\rangle|a\rangle + |1\rangle|b\rangle\big) &= \frac{1}{\sqrt{2}}\big(H|0\rangle|a\rangle + H|1\rangle|b\rangle\big) \\ &= \frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\big(|0\rangle|a\rangle + |1\rangle|a\rangle + |0\rangle|b\rangle - |1\rangle|b\rangle\big) \\ &= \frac{1}{2}|0\rangle\left(|a\rangle + |b\rangle\right) + \frac{1}{2}|1\rangle\left(|a\rangle - |b\rangle\right). \end{split}$$

The probability of measuring state $|0\rangle$ in the first qubit is the norm squared of the probability amplitude:

$$\left|\frac{1}{2}(|a\rangle + |b\rangle)\right|^{2} = \frac{1}{2}(\langle a| + \langle b|)\frac{1}{2}(|a\rangle + |b\rangle)$$

$$= \frac{1}{4}\underbrace{(\langle a|a\rangle + \langle a|b\rangle + \langle b|a\rangle + \underbrace{\langle b|b\rangle}_{=1})}_{=1}$$

$$= \frac{1}{4}(2 + \langle a|b\rangle + \langle a|b\rangle^{*}). \tag{3.3}$$

The two inner products are complex numbers. For a given complex number z, adding $z + z^* = a + ib + a - ib = 2a$. Hence we can write Equation 3.3, the probability of measuring $|0\rangle$ on qubit 0, as:

$$Pr(|0\rangle) = \frac{1}{2} + \frac{1}{2} \text{Re}(\langle a|b\rangle)$$

or, alternatively:

$$2Pr(|0\rangle) - 1 = \text{Re}(\langle a|b\rangle)$$

We can quickly verify this in code as well:

```
# Now let's apply a final Hadamard to ancilla.
qc.h(0)

# At this point, this inner product estimation should hold:
# P(|0>) = 1/2 + 1/2 Re(<a|b>)

# Or
# 2 * P(|0>) - 1 = Re(<a|b>)
dot = np.dot(np.array([a0, a1]).conj(), np.array([b0, b1]))
p0 = qc.psi.prob(0, 0) + qc.psi.prob(0, 1)
if not np.allclose(2 * p0 - 1, dot.real, atol = 1e-6):
    raise AssertionError('Incorrect inner product estimation')
```

Can we obtain an estimate for the imaginary part of the inner product as well? Yes we can. For this, we start with a slightly modified initial state:

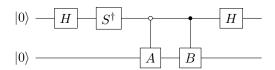
$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle|a\rangle - i|1\rangle|b\rangle).$$

The construction is similar to the one above, but we have to apply a factor

of -i to the $|1\rangle$ part of the state by adding an S^{\dagger} gate right after the initial Hadamard gate:

$$\begin{array}{c|c} |0\rangle & \hline & H \\ \hline & S^{\dagger} \\ \hline & |0\rangle & \hline & A \\ \hline & B \\ \end{array} \right\} \frac{1}{\sqrt{2}} \Big(|0\rangle |a\rangle - i |1\rangle |b\rangle \Big)$$

Similar to above, we add a final Hadamard gate to the ancilla:



which changes the state to:

$$\begin{split} \frac{1}{\sqrt{2}}H\big(|0\rangle|a\rangle - i|1\rangle|b\rangle\big) &= \frac{1}{\sqrt{2}}\big(H|0\rangle|a\rangle - Hi|1\rangle|b\rangle\big) \\ &= \frac{1}{\sqrt{2}}\frac{1}{\sqrt{2}}\big(|0\rangle|a\rangle + |1\rangle|a\rangle - i|0\rangle|b\rangle + i|1\rangle|b\rangle\big) \\ &= \frac{1}{2}|0\rangle\left(|a\rangle - i|b\rangle\right) + \frac{1}{2}|1\rangle\left(|a\rangle + i|b\rangle\right). \end{split}$$

The probability of measuring state $|0\rangle$ on the ancilla is, again, the norm squared of the probability amplitude:

$$\left|\frac{1}{2}(|a\rangle - i|b\rangle)\right|^{2} = \frac{1}{2}(\langle a| + i\langle b|)\frac{1}{2}(|a\rangle - i|b\rangle)$$

$$= \frac{1}{4}\underbrace{(\langle a|a\rangle - i\langle a|b\rangle + i\langle b|a\rangle + \underbrace{\langle b|b\rangle}_{=1})}_{=1}$$

$$= \frac{1}{4}(2 - i\langle a|b\rangle + i\langle a|b\rangle^{*}). \tag{3.4}$$

The inner products are complex numbers. For a complex number z = a + ib, $z^* = a - ib$ and these relations hold:

$$-iz = -i(a+ib) = -ia+b,$$

$$iz^* = i(a-ib) = ia+b,$$

$$\Rightarrow -iz+iz^* = -ia+b+ia+b$$

$$= 2b$$

$$= 2\operatorname{Im}(z).$$

Using this in Equation 3.4, we obtain the final result:

$$Pr(|0\rangle) = \frac{1}{2} + \frac{1}{2}\operatorname{Im}(\langle a|b\rangle).$$

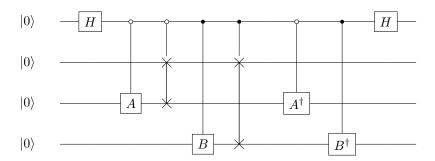
or, alternatively:

$$2Pr(|0\rangle) - 1 = \operatorname{Im}(\langle a|b\rangle).$$

We can confirm this in code as well:

```
# Now let's try the same to get to the imaginary parts.
    psi = 1/sqrt(2)(|0>|a> - i|1>|b>)
psi = (1 / cmath.sqrt(2) *
       (state.bitstring(0) * state.State([a0, a1]) -
        1.0j * state.bitstring(1) * state.State([b0, b1])))
# Let's see how to make this state with a circuit.
qc = circuit.qc('Hadamard test - initial state construction.')
qc.reg(2, 0)
qc.h(0)
qc.sdag(0)
                       # <- this gate is new.
qc.applyc(A, [0], 1) # Controlled-by-0
qc.applyc(B, 0, 1)
                       # Controlled-by-1
# The two states should be identical!
if not np.allclose(qc.psi, psi):
    raise AssertionError('Incorrect result')
# Now let's apply a final Hadamard to ancilla.
qc.h(0)
# At this point, this inner product estimation should hold:
\# P(|0>) = 1/2 + 1/2 Im(\langle a|b \rangle)
# Or
\# 2 * P(|0>) - 1 = Im(\langle a/b \rangle)
dot = np.dot(np.array([a0, a1]).conj(), np.array([b0, b1]))
p0 = qc.psi.prob(0, 0) + qc.psi.prob(0, 1)
if not np.allclose(2 * p0 - 1, dot.imag, atol = 1e-6):
  raise AssertionError('Incorrect inner product estimation')
```

In cases where the construction of A and B is more complex there is the potential that the end state is entangled with the construction of A and B. In such cases, we should introduce an explicit result quantum register. After computation of A and B we superimpose the results onto this register, before uncomputing the construction of A and B, as described in the section on uncomputation. For example:



3.4.2 Inversion Test

In this chapter we will discuss the *inversion test*, a third way to approximate the similarity between states via estimating their scalar product. So far we have learned about the swap test, which utilized a register for each input $|a\rangle$ and $|b\rangle$ together with an ancilla. We also learned about the Hadamard test, which uses the ancilla but just one register, assuming there are operators A and B to construct the states $|a\rangle$ and $|b\rangle$.

The inversion test takes this one step further. It no longer needs an ancilla, just one quantum register, but it needs the ability to construct B^{\dagger} . We again assume operators A and B produce states $|a\rangle$ and $|b\rangle$, with:

$$\begin{split} A\left|0\right\rangle &=A\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}a_0&a_2\\a_1&a_3\end{bmatrix}\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}a_0\\a_1\end{bmatrix}=\left|a\right\rangle,\\ B\left|0\right\rangle &=B\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}b_0&b_2\\b_1&b_3\end{bmatrix}\begin{bmatrix}1\\0\end{bmatrix}=\begin{bmatrix}b_0\\b_1\end{bmatrix}=\left|b\right\rangle, \end{split}$$

and construct this simple circuit:

$$|0\rangle$$
 — A — B^{\dagger} — \nearrow —

The expectation value of the projective measurement $M = |0\rangle\langle 0|$ is given by

$$\begin{split} \left(\langle 0|A^{\dagger}B|0\rangle \right) \left(\langle 0|B^{\dagger}A|0\rangle \right) &= \langle 0|A^{\dagger}B|0\rangle \langle 0|B^{\dagger}A|0\rangle \\ &= |\langle 0|B^{\dagger}A|0\rangle|^2 \\ &= |\underbrace{\langle 0|B^{\dagger}A|0\rangle}_{=\langle b|}|^2 \\ &= |\langle b|a\rangle|^2 \\ &= |\langle a|b\rangle \langle b|a\rangle = \langle b|a\rangle \langle a|b\rangle \\ &= |\langle a|b\rangle|^2. \end{split}$$

Note that the norm of the inner product is symmetric. In code, we reuse the mechanism introduced in the section on the Hadamard test to construct random unitaries A and B with function <code>make_rand_operator()</code>. The inversion test itself is then straightforward:

```
def inversion_test():
  """Perform Inversion Test."""
  # The inversion test allows keeping the number of qubits to a minimum
  # when trying to determine the overlap between two states. However, it
  # requires a precise way to generate states a and b, as well as the
  # adjoint for one of them.
  # If we have operators A and B (similar to the Hadamard Test),
  # to determine the overlap between a and be (\langle a|b \rangle), we run:
        B_adjoint A |0>
  # and determine the probability p0 of measuring |0>. p0 is an
  # a precise estimate for <a/b>.
 A, a0, a1 = make_rand_operator()
 B, b0, b1 = make_rand_operator()
  # For the inversion test, we will need B^\dagger:
 Bdag = B.adjoint()
  # Compute the dot product <a|b>:
 dot = np.dot(np.array([a0, a1]).conj(), np.array([b0, b1]))
  # Here is the inversion test. We run B^{\circ} \to A / 0 and find
  # the probability of measuring |0>:
  qc = circuit.qc('Hadamard test - initial state construction.')
 qc.reg(1, 0)
  qc.apply1(A, 0)
 qc.apply1(Bdag, 0)
  # The probability amplitude of measuring |0> should be the
  # same value as the dot product.
 p0, _ = qc.measure\_bit(0, 0)
 if not np.allclose(dot.conj() * dot, p0):
    raise AssertionError('Incorrect inner product estimation')
```

Bibliography