

6.2.17 GOING THROUGH THE AVAILABLE PLOTS IN TAUOLA
LOOKING UP VARIABLE NAMES IN THE DOCUMENTATION
FOR INTERESTING/ UNIQUE PLOTS

2) THE uxp_z VARIABLE DEPICTS THE Z-POSITION
OF THE PRIMARY VERTEX

THE NUMBER OF ENTRIES 7 917 540 IS THE NUMBER
OF COLLISIONS OBSERVED.

THE MEAN GIVES THE AVERAGE POSITION OF THESE
COLLISIONS, WHICH IS NOT QUITE CENTERED IN THE
DETECTOR, BUT RATHER A BIT BIASED TO THE LEFT
BY -9.163 PRESUMABLY mm.

3) OF THE TWO BEAMS IN THE LHC. EACH HAS
3564 BUNCHES OF WHICH 28 OF CAN BE
FILLED WITH PROTONS.
EVERY BUNCH SLOT CONTAINS $N = \mathcal{O}(10^{11})$ PROTONS
SLOTS ARE SPACED BY 25 ns

VARIATIONS IN SPACING BETWEEN & WITHIN BUNCHES WILL
DIRECTLY AFFECT WHERE IN THE DETECTOR THE
COLLISIONS OCCUR.

A MISMATCH IN BUNCH SIZE COULD LEAD TO INTERACTIONS
AT DIFFERENT LOCATIONS WHICH WOULD SHOW UP
AS DISTINCT SEPARATE PEAKS

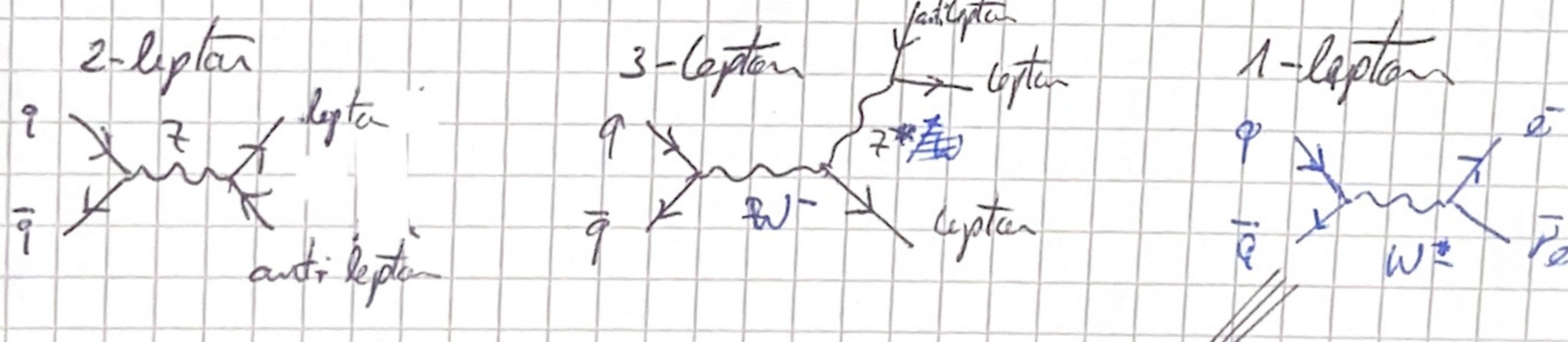
CONSIDERING THE CURRENT BINNING THERE SEEMS
TO BE ONLY ONE PEAK.

SMALLER DEVIATIONS COULD BE OBSERVED WHEN
COMPARING TO GENERATED DATA USING THE
MEAN & STD DEVIATION

whole transition from
calorimeter to —
and we lose leptons there

4) lepn DEPICTS THE NUMBER OF PRESELECTED LEPTONS,
I.E. NUMBER OF LEPTONS PRODUCED SINCE DATA IS
PREFILTERED.
THE PEAK IN THE 1. BIN SUGGESTS THAT MOST INTERACTIONS
YIELD A SINGLE LEPTON WITH A DRAMATIC DECREASE
IN MULTILEPTON PRODUCTION.
MORE THAN 2 LEPTONS & 0 LEPTONS ARE NOT OBVIOUSLY
DEPICTED IN THE PLOT.

5) Initially we expected 2 leptons in the final state, since it correspond to the tree-level Feynman diagram. In this case, due to the high energy, it's more complex diagrams play a role and we get more, or less, than 2 leptons.



6.2 [B] LEP_{PT} VARIANCE DESCRIBING THE TRANSVERSE MOMENTUM.

THE RISE AROUND 25 GeV IS A RESULT OF THE DATA PRESELECTION, SO THAT FEW LEPTONS UNDER THAT ENERGY ARE OBSERVED.

WE OBSERVE A SOMEWHAT EXPONENTIAL DROP IN COUNT AT HIGHER TRANSVERSE MOMENTA. THE MEAN $\langle \text{LEP}_{\text{PT}} \rangle$ FALLS AROUND ~~30 GeV~~ 30 GeV .

AT LEAST ONE HAS TO BE ABOVE 25 GeV
THE FEW BELOW 25 GeV ARE NEEDED IN MULTILEPTON PRODUCTION

THE LEP_ETA VARIABLE CONTAINS THE PSEUDO RAPIDITY DUE TO THE DETECTOR SETUP THERE ARE NO MEASUREMENTS PAST $\eta \approx 2.5$ WHICH IS THE CUTOFF TO THE LEFT & RIGHT. THE CAUSE OF THE GAPS BETWEEN $\pm 1.4 - \pm 1.5$ ARE THE

THE LEP_ETA PHI VARIABLE DEPICTS THE AZIMUTHAL ANGLE OF THE LEPTONS.

DUE TO THE UNIFORMAL ACCEPTANCE OF THE DETECTOR WE OBSERVE A FAIRLY FLAT DISTRIBUTION. THUS WE OBSERVE VALUES EVENLY BETWEEN $\pm \pi \text{ rad}$.

THE NUMBER OF ENTRIES DESCRIBE ALL THE MEASURED LEPTONS. DUE TO THE PRESENCE OF SOME COLLISIONS WHERE MULTIPLE LEPTONS ARE PRODUCED WE FIND THAT THE TOTAL NUMBER OF ENTRIES IS HIGHER THAN IN ?

6.3 [B] WHAT THE SCRIPT DOES

- IMPORT NECESSARY PACKAGES
- EXAMINE COMMAND LINE ARGUMENTS PROVIDED
IE. INPUT FILE & NUMBER OF EVENTS
- CHECK VALIDITY OF INPUT FILE & OPEN IT
- ACCES THE TREE CONTAINING THE VARIABLES
- CREATE A NEW OUTPUT FILE
- INITIALIZE A HISTOGRAM
- LOOPING OVER DATA FILLING THE HISTOGRAM & UPDATING PROGRESS OUTPUT
 - HANDLING DATA THAT IS STORED IN VECTORS
- CREATING A CANVAS TO DRAW THE HISTOGRAM
- WRITING TO OUTPUT FILE

WE LIMIT THE NUMBER OF EVENTS TO 50 000.

WHEN RUNNING THE SCRIPT WE GET A HISTOGRAM OF THE VXP_Z VARIABLE. DUE TO THE LIMITED NUMBER OF ENTRIES OBSERVED & THE VARIA~~T~~ DIFFERENCE IN BINNING WE OBSERVE DIFFERENT VALUES FOR THE MEAN & STD DEVIATION.

WE TRY PRINTING OUT SOME OF THE VARIABLES IN THE TREE. BOTH VXP_Z & LEP_N ENTRIES ARE NUMERIC BEING FLOATS FOR THE INTERACTION VERTEX, AND INTEGERS FOR THE NUMBER OF LEPTONS PRODUCED.

THE OUTPUT OF LEP_{PT} IS GIVEN BY A LIST, WHERE THE NUMBER OF ENTRIES CORRESPOND TO THE NUMBER OF LEPTON PRODUCED IN THAT COLLISION. THE FLOAT VALUES CONTAINED BEING TRANSVERSE MOMENTUM

THE CONVERSION FROM MEV TO GEV CAN BE DONE IN THE PRINTOUT "print(lpx_pt)". [pt/10000 for pt in myCHM.lep_pt]]

- TRYING TO PRINT OUT VPX-ETA WE GET POINTERS TO THE SAME LOCATION FOR ALL VALUES SEEMINGLY

Q - we use the definition of θ from spherical coordinates $\cot\theta = \frac{p_z}{p_t}$

we now want def. of y in $\sinh(y)$: $\sinh(y) = \left(e^{-\text{Im}(\tan(\frac{\theta}{2}))} e^{\text{Re}(\tan(\frac{\theta}{2}))} \right) \frac{1}{2}$

$$\Leftrightarrow \frac{1}{2} (\cot \frac{\theta}{2} - \tan \frac{\theta}{2}) = \frac{1}{2} \frac{\cos^2(\theta/2) - \sin^2(\theta/2)}{\cos(\theta/2) \cdot \sin(\theta/2)} \Rightarrow \text{now we use trig. identities}$$

$$= \frac{\cos \theta}{\sin \theta} = \cot \theta \Rightarrow \sinh(y) = \cot \theta = \frac{p_z}{p_t} \Rightarrow p_z = p_t \sinh(y)$$

WHEN IMPLEMENTING THIS, WE ITERATE THROUGH A NUMBER OF ENTRIES.
WE LIMIT THE m_{ll} VALUES TO THOSE > 0 BY ONLY ADDING
THESE TO THE PLOT.

EXAMINING THE PLOT WE OBSERVE A MAJOR PEAK AROUND 90 GeV WITH A ROUGHLY GAUSSIAN SHAPE TOWARDS THE
RIGHT, WITH FEW TO NONE MAJOR MASSES MEASURED BEYOND
~~140 GeV~~ BUT THE MEAN VALUE OF THE HIST IS AT 73.74 GeV
WITH A STD DEV OF 29.65. THIS IS DUE TO A NUMBER
OF ENTRIES & SMALLER PEAKS ~~located~~ BELOW THE 90 GeV PEAK.
MOST OF THESE ENTRIES AROUND 30 GeV & 6 GeV WITH
A STEADY BASELINE OF ~ 220 ENTRIES

Since for now, we haven't add strong constraints to the events used, we suspect that the peaks with $m_{ll} < 50$ GeV are results from leptons produced by other decay processes and not directly by Z-Boson decays.

COMPARING - NOW EXAMINING THE MONTE CARLO PLOT
WE ARE NOT OBSERVING ~~for~~ THE PREVIOUS PEAKS
IS BELOW THE EXPECTED VALUE & THE MAIN PEAK IS
HIGHER IN AMPLITUDE

From the log & log-log scale, we don't get much more information,

we rather get a better comparison of the peaks with $m_{ll} < 50$ GeV.
here we see that the few low m_{ll} entries reach ~ 200 in measurement,
while ~~for~~ in MC they barely go over 10.

64 PT THE MOST INFLUENTIAL CUTS IN THE HISTOGRAM
SEEM TO BE THE PRIMARY VERTEX &
THE OPPOSITE CHARGE CUT.

② Event-Level cut are more general constraints which are not Z-Boson specific, such that particles not coming from the desire decay pass-through. Whereas more specific conditions, which are specifically tailored for the decay tend to play a more important role, e.g. charge neutrality.

→ 30 GeV max is effe of
pt-cut.

10 GeV Y-Golem resonance
5 GeV T' - charm resonance
q-q likely decay into leptons



we only see the leptons
which come from the
tau decay. They are unstable
and decay before they
get detected. We can
study or identify them
from the p_T & R and V_T

6.3 [4] WE CALCULATE THE INVARIANT MASS USING TLorentzVector BY INITIALIZING LEPT1 & LEPT2 VECTORS, APPLYING
• SET PTEta Phi M WITH THE VALUES
AND PASSING VALUES FROM THE TREE.
LASTLY APPLYING THE .M METHOD TO THE SUM
OF THE LEPTONS.

WE ADJUSTED THE LINE WIDTH OF THE MANUAL
VALUES TO BE THICKER, SINCE THE HISTOGRAMS
COVERED EACH OTHER IDENTICALLY, AS FAR AS
WE CAN TELL.

CONSIDERING THE LINES NEEDED TLorentzVector
IS EASIER TO USE & LESS PRONE TO USER ERROR

⑤ SINCE THE INVARIANT MASS OF THE BOSON IS CONSERVED

THE THE INVARIANT MASS OF THE OBSERVED LEPTONS SHOULD
BE THE SAME. PREVIOUS OBSERVATIONS OF
THE PLOTS HOLD TRUE.

WITH SOME SURPRISE WITH RESPECT TO THE HIGH NUMBER
OF ENTRIES AT LOWER INV. MASSES

THE SMALLER PEAKS AROUND $\sim 5\text{ GeV}$ & $\sim 30\text{ GeV}$.

6.4 [5] THE RESULTING HISTOGRAMS OF THE EGAMMA & MUON DATA NOW APPROXIMATE A BECK CURVE MORE CLEANLY.
THE HARD CUT OFF IMPLEMENTED AROUND $80 - 100\text{ GeV}$ IS SLIGHTLY NOTICABLE.
FOR THE EGAMMA & MUON WE USE ETIET, FOR MUON & TAUON THE ETIET IS USED. WE OMIT THIS TRIGGER FOR ZTAUTAU SINCE IT IS NOT THAT INFLUENTIAL & THERE IS NO ETIET.

THE TAU MC SET SHOWS A PEAK AROUND 46 GeV AND IS THUS ALMOST ENTIRELY CUT BY THE INVARIANT MASS TRIGGER.

6.5 [7] WE COMBINE THE EGAMMA & MUON DATA INTO ANALYSIS MERGED AND USED THE TBROWSER TO QUICKLY VALIDATE THE INTEGRITY OF THE PLOT.
~~the other~~

[2] We proceed by finishing the code plot.py which allows us to plot the invariant mass distribution generated with the measurements and with the MC. Here we take into account the proper scaling for the simulations by multiplying them with $\alpha = \text{inv} \cdot x_{\text{sec}} \cdot i / \text{sum} \cdot i$

6.6 [1] First we remind ourselves, that according to the central limit theorem the a statistical distribution converges to a gaussian distribution if the number of samples $n \rightarrow \infty$ while expectation value μ and variance σ^2 remain finite. We see that ~~the~~ our data set fulfills said requirements, thus the Gauss distribution takes into account the statistical nature (random & Gaussian) of the experiment.

On the other hand, the Breit-Wigner distribution is related to the physical properties of the system. The former distribution describes the probability of finding a particle with a certain energy at a given ~~tree~~ cross section. In short, it describes unstable particles, e.g. Z -boson, in vacuum.

In conclusion by taking both the physical and statistical nature of the process is that we can get a satisfactory fit

[2]

From PDG we read $M_{Z, \text{PDG}} = 91,188 \pm 0,0020 \text{ GeV}$ and $\Gamma_{Z, \text{PDG}} = 2,4955 \pm 0,0023 \text{ GeV}$.

Whereas from our analysis fitting the convolution of a Breit-Wigner and a Gauss distribution we find determine $M_Z = 90,629 \pm 0,005$, $\Gamma_Z = 3,258 \pm 0,021$. We calculate the deviation from both values

$$\sigma_{\frac{\Delta M}{M}} = \frac{|M_{Z, \text{PDG}} - M_Z|}{(\Delta M_{Z, \text{PDG}}^2 + \Delta M_Z^2)^{1/2}} = 27,11. \text{ Analogously we find } \sigma_{\frac{\Delta \Gamma}{\Gamma}}.$$

$\sigma_{\frac{\Delta M}{M}} = 36,09$. From this comparison we can clearly see that our values deviate significantly from the ones measured at the LHC.

→ e emit Brems-emission which makes the detection unclear

μ have a clearer peak due to the clearer path that the leave in our detector

ARBITRARITY OF THE
SOLUTION $\chi^2 \rightarrow \text{TRK}$ FOR DIFFERENT

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6.8 |

Systematic errors,

→ fit model, it does not match the peak.

→ e missing energy due to Bremsstrahlung. Could correct for missing energy.

② modify the Breit-Wigner add extra parameters to move the peak and model the asymmetry

6.7 E1 WE IMPLEMENT THE TAG & PROBE METRIC
FILTERING FOR TIGHT OPTION & RUNNING
WITH ALL DATA
SAVING INTERMITTENT VALUES IN DENOMINATOR
(ENUMERATOR), WHICH WILL ALLOW US
LATER TO CALCULATE THE ERROR
USING A BINOMIAL DISTRIBUTION SINCE A
PARTICLE EITHER PASSES OR FAILS A
TEST

$$E = \frac{\text{NUMERATOR}}{\text{DENOMINATOR}}$$

WITH THE EFFICIENCY

$$\sigma_E = \sqrt{\frac{E(1-E)}{\text{DENOMINATOR}}}$$

THIS CAN
EASILY BE IMPLEMENTED USING THE
OPTION B IN ROOTS TH1::DIVIDE

LIMIT OF 6 PAGES