

Relativity – Lecture 9

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Key concepts of lecture 8 - 1

- A physical vector quantity is represented by a four-vector in Special Relativity.
- A four-vector transforms between inertial frames under the Lorentz transformations.
- The norm of a four-vector is invariant.
- The four-velocity is $\mathbf{U} = \gamma_u(c, \mathbf{u})$.

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Key concepts of lecture 8 - 2

- The energy-momentum four-vector is $\mathbf{P} = (E/c, \mathbf{p})$.
- Here $\mathbf{p} = \gamma_u m \mathbf{u}$, and $E = \gamma_u mc^2$ is the total energy of the particle.
- The norm of \mathbf{P} is $m^2 c^2$. So $(mc^2)^2 = E^2 - (pc)^2$.
- The (rest-) mass is therefore invariant.
- The kinetic energy is $T = E - mc^2$.

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Example: rest energy of an electron

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Example: radioactive decay

Energy difference is 175 MeV.

Nuclear binding energy is large compared to particle masses!

All forms of energy contribute to the rest mass: electrostatic, nuclear, thermal, etc.

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Energy & momentum conservation

It has been shown experimentally that energy-momentum is conserved.

Therefore, in a particular frame E and p are separately conserved (e.g. in a collision), just as in classical mechanics.

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Example: energy & momentum conservation

Two protons ($m = 1\text{GeV}/c^2$) collide to form a pion ($m = 140\text{ MeV}/c^2$).

If all particles are at rest after collision, what was the initial velocity?

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Example: pion decay

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Useful relations

We know: $p = \gamma_u mu$, $E = \gamma_u mc^2$

Solve for u , γ_u :

Example: particle collision



Particle 1 moves with speed $u_1 = 15/17 c$ along the x-axis, and collides with stationary particle 2 to produce particle 3.
 Particles 1 and 2 have masses $m_1 = m_2 = 8/c^2$ units.

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Example: particle collision

Particle	$P_i = (E_i/c, p_i)$	β_i	m_i
1			
2			
3			

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Points to note from particle collision example

- The two incoming particles have the same energy-momentum vector length.
- The mass of particle 3 is *not* $m_3 = m_1 + m_2$. (m is frame-invariant, but not conserved!)

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Lorentz transform for E, p

$$\frac{E'}{c} = \gamma \left(\frac{E}{c} - \beta p_x \right)$$

$$p_x' = \gamma \left(p_x - \beta \frac{E}{c} \right)$$

Remember that E, p have γ_u factor in their definitions.
The γ in the transformation is γ_v .

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The centre-of-momentum frame

It is often easier to solve problems in the “centre-of-momentum” frame, where total momentum is zero. In other words, $p_{\text{before}} = p_{\text{after}} = 0$.



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The centre-of-momentum frame

Particle	$P'_i = (E'_i/c, \mathbf{p}'_i)$	β'_i	m'_i
1			
2			
3			

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Transforming frames: conclusion

Note that energy and momentum are conserved separately in any one frame.

However, when transforming frames, the energy and momentum change.

In other words, a Lorentz transformation changes energy into momentum, and vice versa.

However, the norm of the four-vector is invariant, so $E^2 = p^2 c^2 + (mc^2)^2$ is always true.

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Reminder: get the terminology right.

- **Conserved:** a quantity which is not changed by a physical process. This refers to one frame at a time, and a conserved quantity will typically have different numerical values in different frames.
- **Invariant:** a quantity which is not changed by a coordinate transformation. The term refers to more than one reference frame; an invariant quantity will not necessarily be conserved in a particular process.
- **Constant:** refers to a quantity which does not change in time, such as the mass of the Universe.
- The speed of light is conserved, invariant, and constant!

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Tip: solving problems

Try to solve problems first by using energy conservation alone. Some problems require you to use both energy and momentum conservation.

You can eliminate one variable using

$$E^2 = p^2 c^2 + (mc^2)^2$$

for example $p = \sqrt{(E/c)^2 - (mc)^2}$

For a massless particle, $E = pc$.

You can also leave out all of the c's and put them in at the end using dimensional analysis.