

The wave equation:- Maxwell's eqs. are

$$\begin{aligned} \text{(i)} \quad \nabla \cdot \mathbf{D} &= \rho & \text{(ii)} \quad \nabla \cdot \mathbf{B} &= 0 \\ \text{(iii)} \quad \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} & \text{(iv)} \quad \nabla \times \mathbf{H} &= \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \end{aligned}$$

In a uniform linear media, having permittivity  $\epsilon$  and permeability  $\mu$  then  $\mathbf{D} = \epsilon \mathbf{E}$ ,  $\mathbf{B} = \mu \mathbf{H}$ ,  $\mathbf{J} = \sigma \mathbf{E}$ .  $\rho = 0$   
Maxwell's eqs. for this medium can be written as

$$\begin{aligned} \text{(i)} \quad \nabla \cdot \mathbf{E} &= 0 & \text{(ii)} \quad \nabla \cdot \mathbf{H} &= 0 \\ \text{(iii)} \quad \nabla \times \mathbf{E} &= -\mu \frac{\partial \mathbf{H}}{\partial t} & \text{(iv)} \quad \nabla \times \mathbf{H} &= \sigma \mathbf{E} + \epsilon \frac{\partial \mathbf{E}}{\partial t} \end{aligned}$$

$$\text{Curl}(\nabla \times \mathbf{E}) = -\mu \frac{\partial}{\partial t}(\nabla \times \mathbf{H})$$

$\sigma$  = conductivity of medium

Substituting the value of  $\nabla \times \mathbf{H}$

$$\begin{aligned} \text{Curl curl } \mathbf{E} &= -\mu \frac{\partial}{\partial t} \left( \sigma \mathbf{E} + \epsilon \frac{\partial \mathbf{E}}{\partial t} \right) \\ &= -\mu \sigma \frac{\partial \mathbf{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \mathbf{E}}{\partial t^2} \end{aligned} \quad \text{--- (v)}$$

Similarly, taking curl of Maxwell's eq. (iv)

$$\text{Curl}(\nabla \times \mathbf{H}) = -\mu \sigma \frac{\partial \mathbf{H}}{\partial t} - \mu \epsilon \frac{\partial^2 \mathbf{H}}{\partial t^2} \quad \text{--- (vi)}$$

Using vector identity  $\text{curl curl } \mathbf{A} = \text{grad div } \mathbf{A} - \nabla^2 \mathbf{A}$  eqs. (v) and (vi) becomes

$$\nabla^2 \mathbf{E} - \mu \sigma \frac{\partial \mathbf{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0 \quad \text{--- (vii)}$$

$$\text{and } \nabla^2 \mathbf{H} - \mu \sigma \frac{\partial \mathbf{H}}{\partial t} - \mu \epsilon \frac{\partial^2 \mathbf{H}}{\partial t^2} = 0 \quad \text{--- (viii)}$$

Eqs. (vii) and (viii) are the wave eqs. of e.m. fields in a linear homogeneous medium where charge density is zero.

Its plane e.m. wave solution is of the form

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \cdot e^{i\mathbf{k} \cdot \mathbf{r} - i\omega t}, \quad \mathbf{r} \text{ is position vector}$$

$\mathbf{E}_0$  is the complex amplitude which is constant in space & time

$\mathbf{k}$  is the wave propagation vector defined as

$$\mathbf{k} = \frac{2\pi}{\lambda} \cdot \mathbf{n} = \frac{\omega}{v} \mathbf{n}, \quad \mathbf{n} \text{ being unit vector along } \mathbf{k}.$$

$v \rightarrow$  phase velocity of wave.

One of the simplest and familiar wave form is sine wave

$$f(z, t) = A \cos[k(z - vt) + S]$$

$A$  is the amplitude,  $S$  is the phase constant.

The time period of wave  $T = \frac{2\pi}{kv}$

$$\therefore \text{The frequency } \nu = \frac{1}{T} = \frac{kv}{2\pi} = \frac{v}{\lambda}$$

$$\text{angular frequency } \omega = 2\pi\nu = kv$$

In terms of  $\omega$ , the wave equation can be written as

$$f(z, t) = A \cos(kz - \omega t + S)$$

If the wave is travelling towards left, then

$$f(z, t) = A \cos(kz + \omega t - S)$$

The wave eq. can also be represented in complex notation using Euler's formula  $e^{i\theta} = \cos\theta + i\sin\theta$

$$\therefore f(z, t) = \text{Re} \left[ A e^{i(kz - \omega t + S)} \right]$$

where  $\text{Re}$  represents the real part of complex no.