

# Physics Based Animation

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**Impulse based multi-body animation**

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# What's an impulse-based contact?

- We simulate contacts using collision impulses.
- The impulse is the integral of the force  $F$  over time. Since force is a vector quantity, impulse is also a vector quantity. Impulse applied to an object produces an equivalent vector change in its linear momentum (mass\*velocity), also in the resultant direction.
- Why do we like impulse-based contacts? Because it's fast!

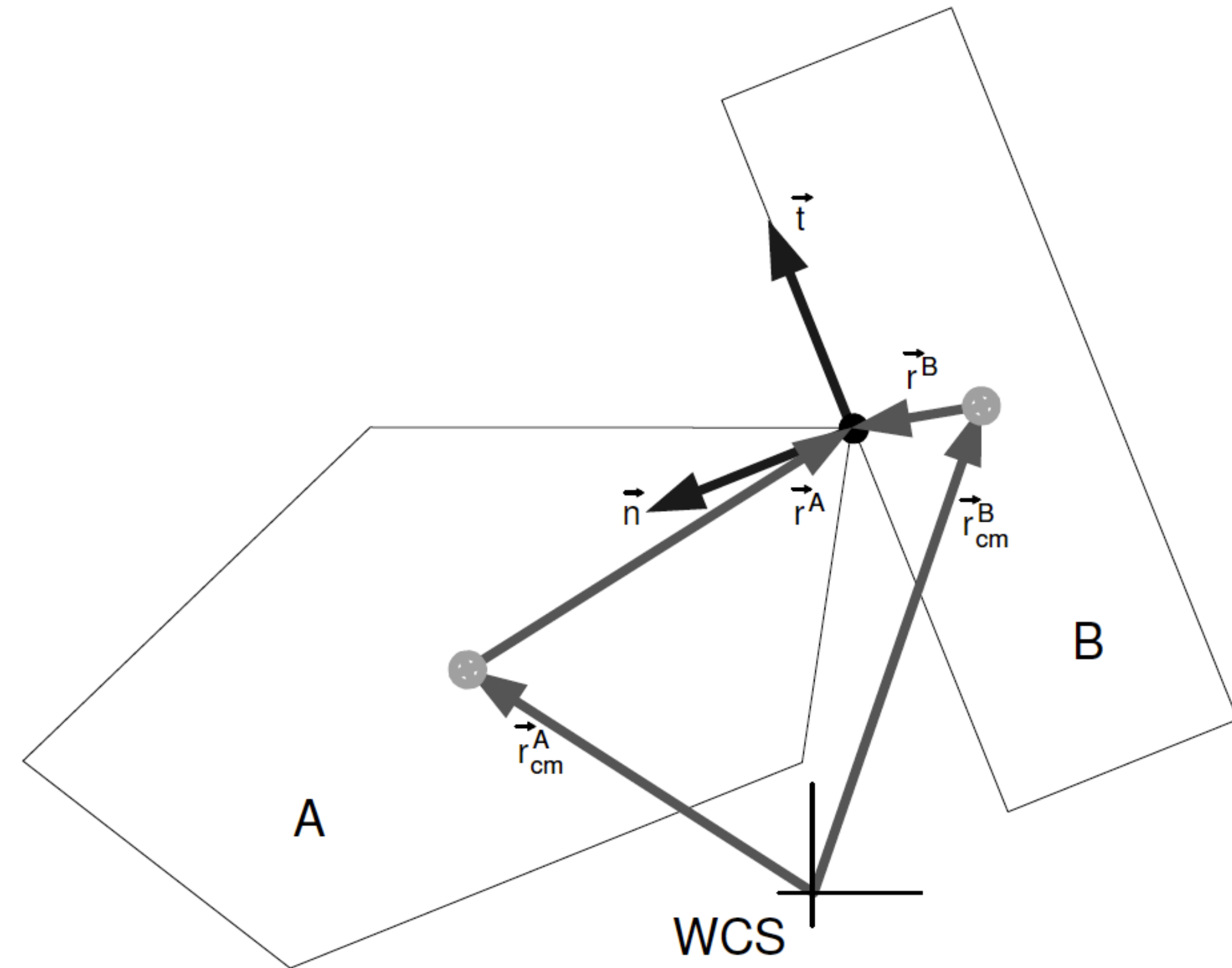
# Single point contact

- Single point contacts are not very realistic since collisions usually happens over a surface.
- We will assume we can compute the normal at collision point  $p$ .
- The normal by convention is always pointing from object B to A.
- We can always reach the collision point from either object's center of mass.

$$\mathbf{p}^A = \mathbf{r}^A + \mathbf{r}_{\text{cm}}^A,$$

$$\mathbf{p}^B = \mathbf{r}^B + \mathbf{r}_{\text{cm}}^B,$$

$$\mathbf{p}^A = \mathbf{p}^B.$$

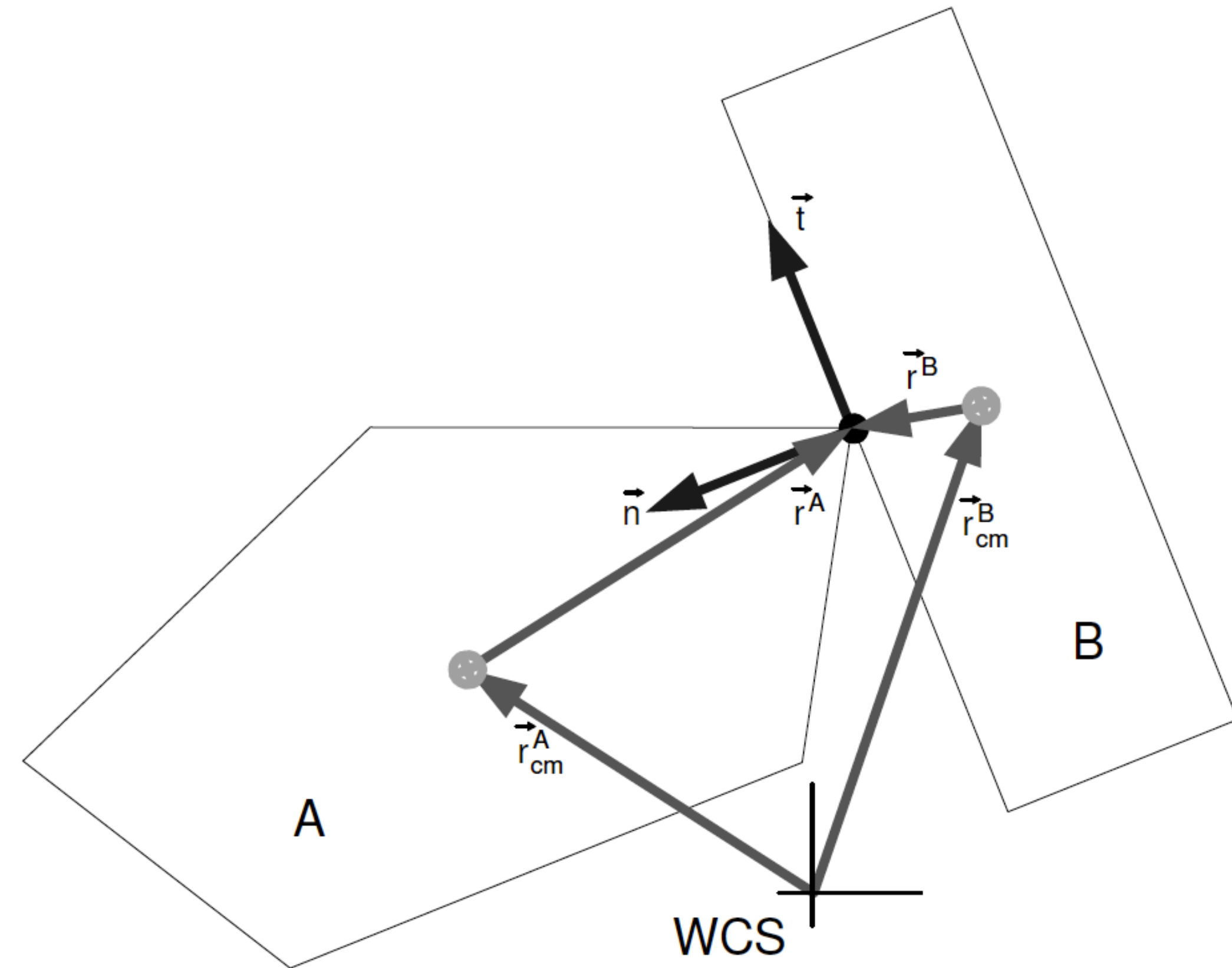


# Single point contact

- In 3D the contact point and contact normal define a contact plane. We can define the plane in 2 ways: explicitly for every point or using two orthogonal vectors w.r.t normal lying on the plane.

$$\forall x \in \mathbb{R}^3 \quad \text{where} \quad n \cdot x - n \cdot p = 0.$$

$$n = t_1 \times t_2.$$



# Single point contact - relative contact velocity

Recall, for a generic object  $X$   $\mathbf{p}^X = \mathbf{r}^X + \mathbf{r}_{\text{cm}}^X$ , We can use Body Frame notation  $\mathbf{r}^X = \mathbf{R}^X [\mathbf{r}^X]_{\text{BF}_X}$ ,

Obtaining  $\mathbf{p}^X = \mathbf{R}^X [\mathbf{r}^X]_{\text{BF}_X} + \mathbf{r}_{\text{cm}}^X$ . If we take the derivative we get:

$$\dot{\mathbf{p}}^X = \dot{\mathbf{R}}^X [\mathbf{r}^X]_{\text{BF}_X} + \mathbf{v}_{\text{cm}}^X,$$

$$\dot{\mathbf{p}}^X = \left[ \boldsymbol{\omega}^X \times \mathbf{R}_x^X \mid \boldsymbol{\omega}^X \times \mathbf{R}_y^X \mid \boldsymbol{\omega}^X \times \mathbf{R}_z^X \right] [\mathbf{r}^X]_{\text{BF}_X} + \mathbf{v}_{\text{cm}}^X,$$

$$\dot{\mathbf{p}}^X = \boldsymbol{\omega}^X \times \left( \mathbf{R}^X [\mathbf{r}^X]_{\text{BF}_X} \right) + \mathbf{v}_{\text{cm}}^X,$$

There's a trick! If we add zero  $\dot{\mathbf{p}}^X = \boldsymbol{\omega}^X \times \left( \underbrace{\mathbf{R}^X [\mathbf{r}^X]_{\text{BF}_X} + \mathbf{r}_{\text{cm}}^X}_{\mathbf{p}^X} - \mathbf{r}_{\text{cm}}^X \right) + \mathbf{v}_{\text{cm}}^X$ . Which is equal to:

$$\dot{\mathbf{p}}^X = \boldsymbol{\omega}^X \times (\mathbf{p}^X - \mathbf{r}_{\text{cm}}^X) + \mathbf{v}_{\text{cm}}^X,$$

Now we can compute the relative contact velocity  $\mathbf{u} = \dot{\mathbf{p}}^A - \dot{\mathbf{p}}^B$ .

And relative normal contact velocity  $u_n = \mathbf{n}^T (\dot{\mathbf{p}}^A - \dot{\mathbf{p}}^B),$

# Why the relative velocity is important?

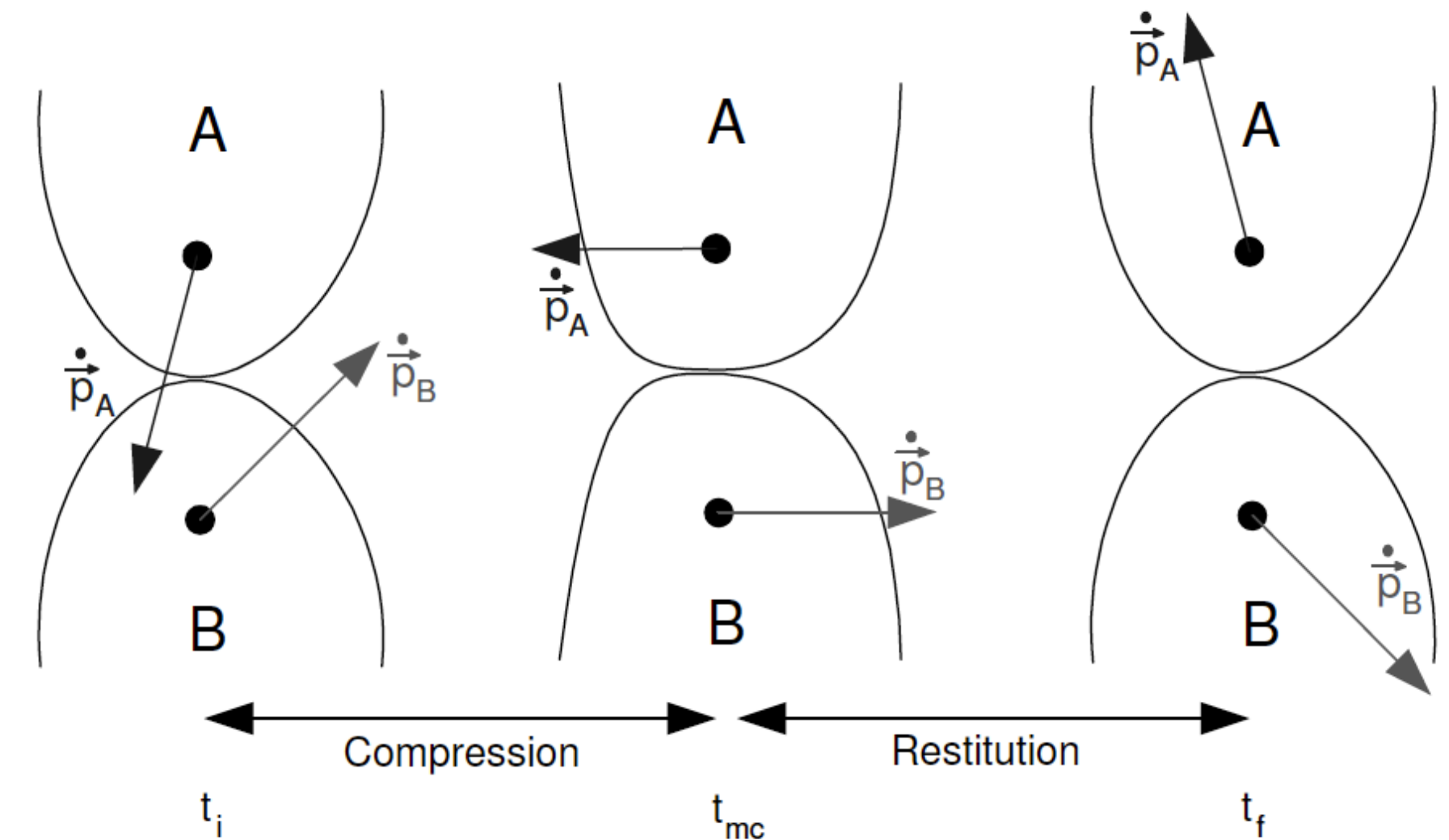
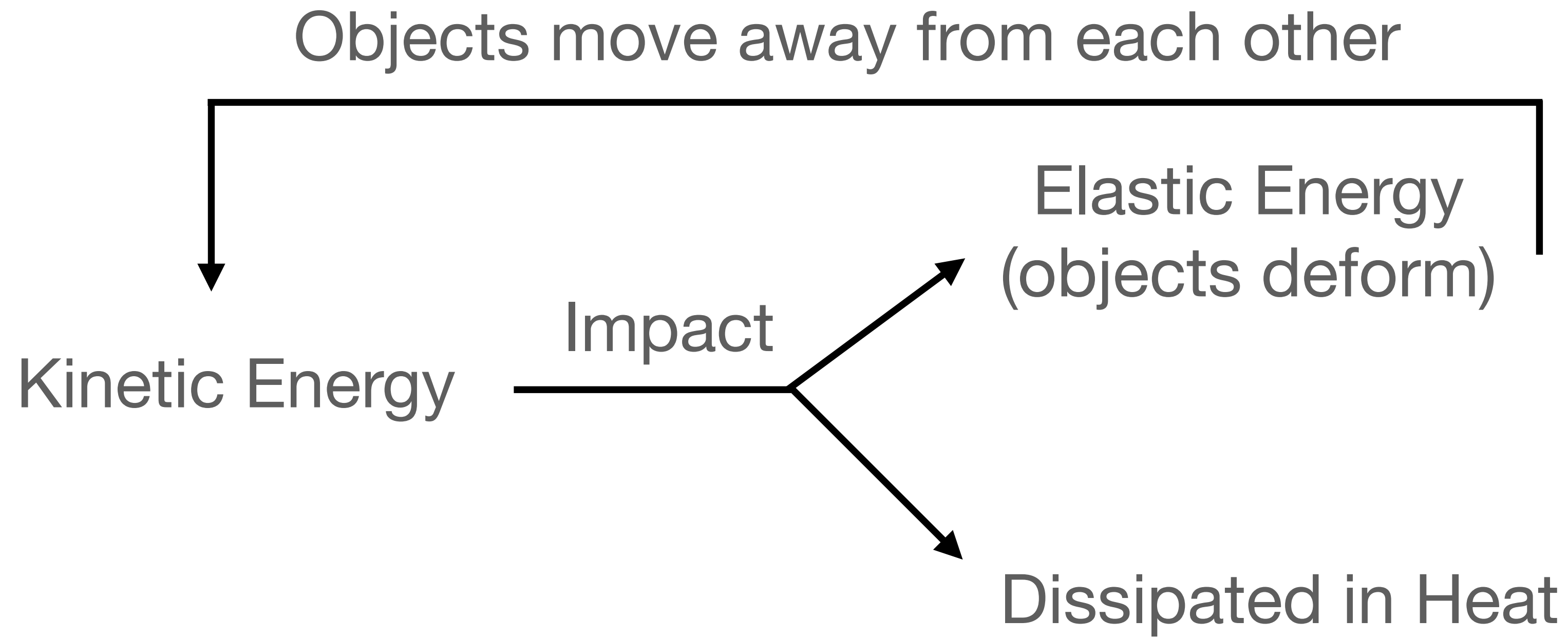
- Because now we can tell if the 2 objects are separating, resting or colliding.

$$\text{relative contact} = \begin{cases} \text{Separating} & \text{if } u_n > 0 \\ \text{Resting} & \text{if } u_n = 0 . \\ \text{Colliding} & \text{if } u_n < 0 \end{cases}$$

- And we can also derive our collision test:

$$u_n = \mathbf{n}^T \left( (\boldsymbol{\omega}^A \times \mathbf{r}^A + \mathbf{v}_{cm}^A) - (\boldsymbol{\omega}^B \times \mathbf{r}^B + \mathbf{v}_{cm}^B) \right) < 0.$$

# How contacts work in real life (sort of)





# Rigid bodies

## **Definition 6.1 (The Rigid Body Assumptions)**

*During a collision, the following four assumptions apply to rigid bodies:*

- *The duration of the collision is zero, that is  $t_f - t_i \rightarrow 0$ .*
- *Only impulsive forces must be used to avoid penetration.*
- *Positions and orientations are the same before and after the collision.*
- *Non impulsive forces have no effect during the collision.*



# Impulses

$$\mathbf{J} = \int_{t_i}^{t_f} \mathbf{F} dt.$$

Which can also be seen as the difference in momentum  $\mathbf{P} = m\mathbf{v}_{\text{cm}}$ :

$$\begin{aligned}\mathbf{F} &= \frac{d\mathbf{P}}{dt}, \\ \int_{t_i}^{t_f} \mathbf{F} dt &= \int_{t_i}^{t_f} d\mathbf{P}, \\ \mathbf{J} &= \mathbf{P}(t_f) - \mathbf{P}(t_i), \\ \mathbf{J} &= \Delta\mathbf{P}.\end{aligned}$$

Since  $t_f - t_i \rightarrow 0$ . We need to reparametrize w.r.t time

Time always increases  $\frac{d\gamma}{dt} > 0$ . So does the impulse in the normal direction  $\frac{dJ_n}{dt} > 0$ .

This has to hold:  $\frac{d\gamma}{dJ_n} > 0$ .

# Collision laws

- *Algebraic Collision Laws*
- *Incremental Collision Laws*
- *Full Deformation Collision Laws*
- *Compliant Contact Model Collision Laws*

*Collision Law  $\Leftarrow$  Physical Law + Contact Model*

# Physical Law — Impact

Newton's impact law

$$u_n(\gamma_f) = -e u_n(\gamma_i), \quad 0 \leq e \leq 1.$$

Poisson's hypothesis

$$J_n(\gamma_f) - J_n(\gamma_{mc}) = e J_n(\gamma_{mc}), \quad 0 \leq e \leq 1.$$

Stronge's hypothesis

$$W_n(\gamma_f) - W_n(\gamma_{mc}) = -e^2 W_n(\gamma_{mc}), \quad 0 \leq e \leq 1.$$

$$W = \int_{t_i}^{t_f} \mathbf{F}(t) \cdot \dot{\mathbf{x}}(t) dt. \quad \text{Reparametrize} \quad W = \int_{\gamma_i}^{\gamma_f} \mathbf{F}(\gamma) \cdot \dot{\mathbf{x}}(\gamma) \left( \frac{dt}{d\gamma} \right) d\gamma.$$

$$\mathbf{F}(\gamma) = \frac{d}{dt} \mathbf{J}(\gamma), \quad \dot{\mathbf{x}}(\gamma) = \ddot{\mathbf{u}}_A(\gamma), \quad W_A = \int_0^{\gamma_f} \mathbf{u}_A(\gamma) \cdot \frac{d}{d\gamma} \mathbf{J}(\gamma) d\gamma.$$

# Physical Law — Friction

Coulomb's friction law

$$\mathbf{F}_t = -\mu \frac{\mathbf{u}_t}{\|\mathbf{u}_t\|} \|\mathbf{F}_n\| ,$$

$$\dot{\mathbf{u}}_t = 0, \quad \|\mathbf{F}_t\| \leq \mu \|\mathbf{F}_n\| .$$

$$\dot{\mathbf{u}}_t \neq 0, \quad \|\mathbf{F}_t\| = \mu \|\mathbf{F}_n\| \quad \text{and} \quad \mathbf{F}_t \cdot \dot{\mathbf{u}}_t \leq 0 .$$

# Contact model — how the velocity changes

Newton's 2nd law  $\mathbf{F}^A(\gamma) = m_A \mathbf{a}_{\text{cm}}^A(\gamma),$

Euler's equation of motion  $\boldsymbol{\tau}^A(\gamma) = \mathbf{r}^A \times \mathbf{F}^A(\gamma) = \mathbf{I}_A \boldsymbol{\alpha}^A(\gamma) + \cancel{\boldsymbol{\omega}^A(\gamma) \times \mathbf{I}_A \boldsymbol{\omega}^A(\gamma)}.$  Alpha is angular acceleration

Integrating:

$$\int_{\gamma_i}^{\gamma} \mathbf{F}^A(\gamma) d\gamma = \int_{\gamma_i}^{\gamma} m_A \mathbf{a}_{\text{cm}}^A(\gamma) d\gamma,$$

$$\int_{\gamma_i}^{\gamma} \mathbf{r}^A \times \mathbf{F}^A(\gamma) d\gamma = \int_{\gamma_i}^{\gamma} \mathbf{I}_A \boldsymbol{\alpha}^A(\gamma) d\gamma.$$

$$\mathbf{J}^A(\gamma) = m_A \underbrace{(\mathbf{v}_{\text{cm}}^A(\gamma) - \mathbf{v}_{\text{cm}}^A(\gamma_i))}_{\Delta \mathbf{v}_{\text{cm}}^A(\gamma)},$$

$$\mathbf{r}^A \times \mathbf{J}^A(\gamma) = \mathbf{I}_A \underbrace{(\boldsymbol{\omega}^A(\gamma) - \boldsymbol{\omega}^A(\gamma_i))}_{\Delta \boldsymbol{\omega}^A(\gamma)}.$$

Now we have a relationship between impulses and change in velocities

$$\mathbf{J}^A(\gamma) = m_A \Delta \mathbf{v}_{\text{cm}}^A(\gamma),$$

$$\mathbf{r}^A \times \mathbf{J}^A(\gamma) = \mathbf{I}_A \Delta \boldsymbol{\omega}^A(\gamma).$$



# Contact model — how the momentum changes

$$\begin{aligned} \mathbf{J}^A(\gamma) &= m_A \Delta \mathbf{v}_{\text{cm}}^A(\gamma), \\ \mathbf{r}^A \times \mathbf{J}^A(\gamma) &= \mathbf{I}_A \Delta \boldsymbol{\omega}^A(\gamma). \end{aligned} \quad \longrightarrow \quad \begin{aligned} \Delta \mathbf{v} &= \frac{\mathbf{J}}{m}, \\ \Delta \boldsymbol{\omega} &= \mathbf{I}^{-1} (\mathbf{r} \times \mathbf{J}). \end{aligned}$$

$$\dot{\mathbf{p}}^A(\gamma) = \boldsymbol{\omega}^A(\gamma) \times \mathbf{r}^A + \mathbf{v}_{\text{cm}}^A(\gamma).$$

$$\Delta \dot{\mathbf{p}}^A(\gamma) = \Delta \boldsymbol{\omega}^A(\gamma) \times \mathbf{r}^A + \Delta \mathbf{v}_{\text{cm}}^A(\gamma),$$

By substituting and doing some algebra we get:

$$\Delta \mathbf{u}(\gamma) = \underbrace{\left( \left( \frac{1}{m_A} + \frac{1}{m_B} \right) \mathbf{1} - \left( (\mathbf{r}^A)^\times \mathbf{I}_A^{-1} (\mathbf{r}^A)^\times + (\mathbf{r}^B)^\times \mathbf{I}_B^{-1} (\mathbf{r}^B)^\times \right) \right)}_K \mathbf{J}^A(\gamma).$$



# Properties of the collision matrix

*The collision matrix  $\mathbf{K}$  is*

$$\Delta \mathbf{u}(\gamma) = \underbrace{\left( \left( \frac{1}{m_A} + \frac{1}{m_B} \right) \mathbf{1} - \left( (\mathbf{r}^A)^\times \mathbf{I}_A^{-1} (\mathbf{r}^A)^\times + (\mathbf{r}^B)^\times \mathbf{I}_B^{-1} (\mathbf{r}^B)^\times \right) \right)}_{\mathbf{K}} \mathbf{J}^A(\gamma).$$

- *constant*
- *symmetric*
- *positive definite*
- *invertible*

Differential Form of Impulse-Momentum Relation

$$\begin{aligned} \frac{d}{d\gamma} \mathbf{u}(\gamma) &= \mathbf{K} \frac{d}{d\gamma} \mathbf{J}(\gamma), \\ \frac{d}{d\gamma} \mathbf{J}(\gamma) &= \mathbf{K}^{-1} \frac{d}{d\gamma} \mathbf{u}(\gamma). \end{aligned}$$

This notation comes handy in the Stronge's hypothesis in the contact model.

$$W = \int_0^{\gamma_f} \mathbf{u}(\gamma) \cdot \mathbf{K}^{-1} \frac{d}{d\gamma} \mathbf{u}(\gamma) d\gamma,$$