

# Chemistry 3P51 – Fall 2013

## Quantum Chemistry

Lecture No. 16

Oct 21<sup>th</sup>, 2013

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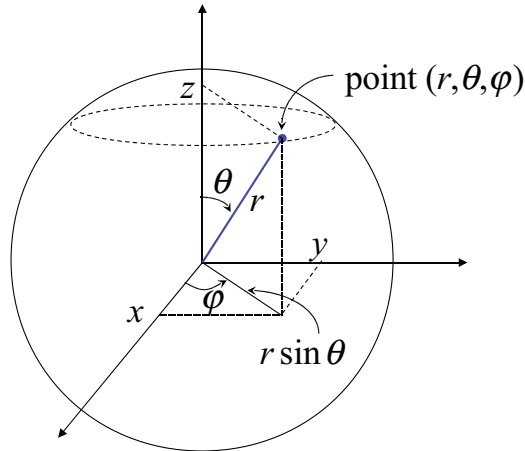
### *Objectives*

- To remind the student the spherical polar coordinates.
- To show the angular momentum operator components and square in spherical coordinates.
- To show the relationship between the kinetic energy operator and the angular momentum operator in rotational motion.
- To derive the Schrödinger equation for a particle in a ring.

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## Spherical polar coordinates

Cartesian coordinates are not ideal for describing orbital motion. Spherical coordinates are better suited for this purpose.



$$x = r \sin \theta \cos \varphi$$

$$y = r \sin \theta \sin \varphi$$

$$z = r \cos \theta$$

$$r^2 = x^2 + y^2 + z^2 \quad (1)$$

$$\cos \theta = \frac{z}{r} \quad (2)$$

$$\tan \varphi = \frac{y}{x} \quad (3)$$

$$0 \leq r < \infty, \quad 0 \leq \theta \leq \pi, \quad 0 \leq \varphi \leq 2\pi$$

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## Angular momentum in spherical coordinates

When expressed in spherical coordinates, the **components of the angular momentum operator** and **its square** take the following form

$$\begin{aligned} \hat{L}_x &= -i\hbar \left( -\sin \varphi \frac{\partial}{\partial \theta} - \frac{\cos \theta}{\sin \theta} \cos \varphi \frac{\partial}{\partial \varphi} \right) \\ \hat{L}_y &= -i\hbar \left( \cos \varphi \frac{\partial}{\partial \theta} - \frac{\cos \theta}{\sin \theta} \sin \varphi \frac{\partial}{\partial \varphi} \right) \\ \hat{L}_z &= -i\hbar \frac{\partial}{\partial \varphi} \end{aligned}$$

$$\hat{L}^2 = -\hbar^2 \left( \frac{\partial^2}{\partial \theta^2} + \frac{\cos \theta}{\sin \theta} \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right)$$

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## Angular momentum and kinetic energy operators

The angular momentum operator is very useful because is closely related to the **rotational motion** kinetic energy operator; that is, motion about a **fixed point** where the distance is  $r$

$$\hat{T} = \frac{\hat{L}^2}{2mr^2} \quad \text{provided } r = \text{const}$$

The quantity

$$I = mr^2$$

is called the **moment of inertia**.

Thus, the Schrödinger equation for a free ( $V=0$ ) particle in a ring or on the surface of a sphere of radius  $r$  can be written as

$$\frac{\hat{L}^2}{2I} \psi = E \psi$$

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## Particle in a ring

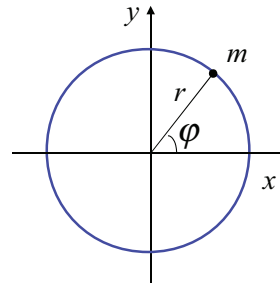
Consider a particle of mass  $m$  constrained to move in a ring of radius  $r$  and assume that the potential energy  $V(\varphi)$  of the particle is zero.

The Schrödinger equation for this system is

$$\hat{H} \psi = E \psi$$

where

$$\hat{H} = \frac{\hat{L}^2}{2I} = \frac{\hat{L}^2}{2mr^2}$$



Since the particle is constrained to remain in the  $z=0$  plane,

$$\hat{L}_x = \hat{L}_y = 0, \quad \text{so} \quad \hat{L}^2 = \hat{L}_z^2$$

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### ***Particle in a ring***

Recall that

$$\hat{L}_z = -i\hbar \frac{\partial}{\partial \varphi} \quad \hat{L}_z^2 = -\hbar^2 \frac{\partial^2}{\partial \varphi^2}$$

The Schrödinger equation for this particular problem is one-dimensional:

$$-\frac{\hbar^2}{2mr^2} \frac{\partial^2}{\partial \varphi^2} \psi(\varphi) = E\psi(\varphi)$$

Let us define

$$n^2 = \frac{2mr^2 E}{\hbar^2}$$

Therefore, we can rewrite the equation as  $\frac{d^2\psi}{d\varphi^2} = -n^2\psi(\varphi)$

The solution is

$$\psi(\varphi) = e^{in\varphi},$$

where  $n$  is any real number, positive or negative.

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