

Chemistry 3P51 – Fall 2013

Quantum Chemistry

Lecture No. 6
Sep 16th, 2013

1

Objectives

- To learn the requirements of a wave-function to be valid and the concept of “well-behaved ” function.
- To introduce the concept of expectation values show examples of how they are computed based on given wave-functions.
- To motivate the particle-in-a-box system by means of conjugated polyenes.
- To show the general strategy involved in the solution of the Schrödinger equation for a particle in a box and discuss such solution.

2

Requirements for a valid wave-function

1. The wave-function must be **continuous**.

2. The wave-function must be **single-valued**.

That is, there should not be two different values of probability density for finding the particle near the same point.

3. The wave-function should have **continuous derivatives** everywhere

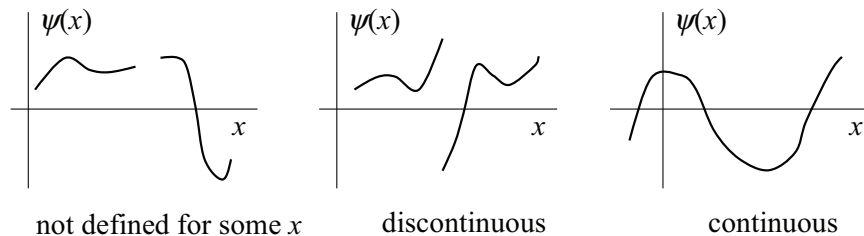
$$\frac{\partial \psi}{\partial x}, \frac{\partial^2 \psi}{\partial x^2}, \frac{\partial^3 \psi}{\partial x^3}, \dots, \frac{\partial^n \psi}{\partial x^n}$$

except, perhaps, at certain special points (at the nuclei or physical boundaries of the system)

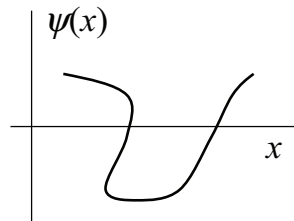
4. The wave-function should be **square-integrable**. In other words, **normalizable**

Functions satisfying 1 to 3 are usually termed as **well-behaved**. ₃

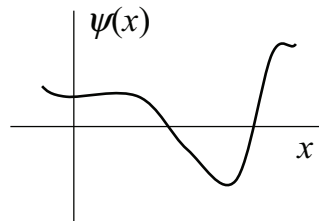
The wave-function must be continuous



The wave-function must be single-valued



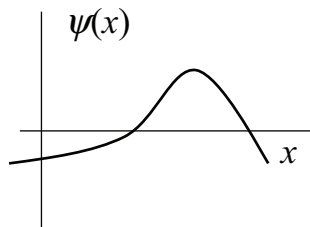
not single-valued



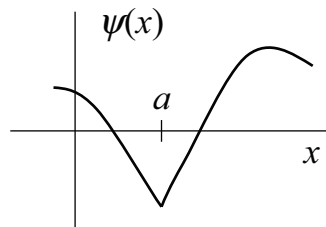
single-valued

5

The wave-function should have continuous derivatives everywhere



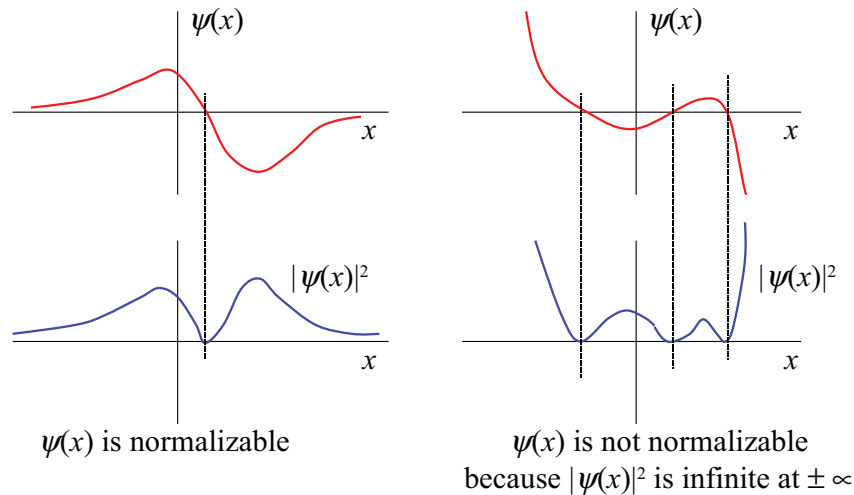
$\partial\psi/\partial x$ is continuous



$\partial\psi/\partial x$ is discontinuous at $x = a$

6

The wave-function should be square-integrable



7

Expectation values

- Based on the **probabilistic interpretation** of the wave-function, we introduce the **expectation value**, **expected value** or **average value**
- The **average value** of an observable A , when the system is in a state ϕ is given by

$$\langle A \rangle_{\phi}(t) = \frac{\int_{-\infty}^{\infty} \phi^*(x,t) \hat{A}(x,t) \phi(x,t) dx}{\int_{-\infty}^{\infty} \phi^*(x,t) \phi(x,t) dx}$$

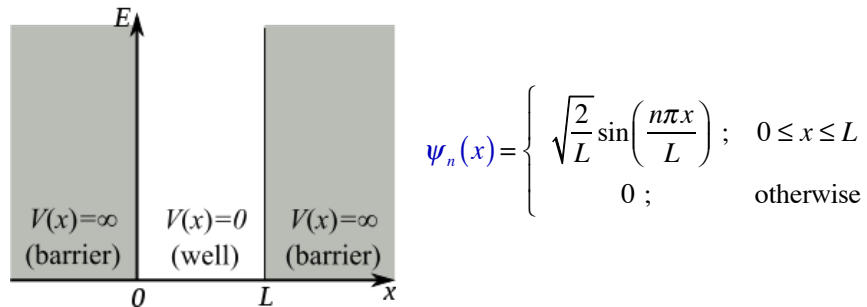
where A is the operator representing the observable A .

- The **average value** depends on the state, and in general depends on time

8

Example of expectation values

- Later we will see that the n^{th} state of the **particle-in-a-box** system represented by the following wave-function



- The **average position** of the particle, when at the n^{th} state is given by

$$\langle x \rangle_{\psi_n} = \frac{\int_0^L \psi_n^*(x) \hat{x} \psi_n(x) dx}{\int_0^L \psi_n^*(x) \psi_n(x) dx} = \frac{L}{2}$$

9

Example of expectation values

- The **average linear momentum** of the particle, when at the n^{th} state, is given by

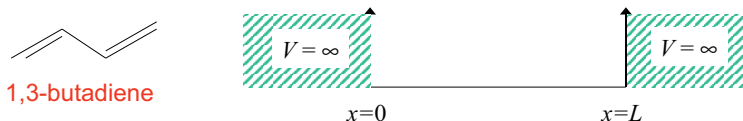
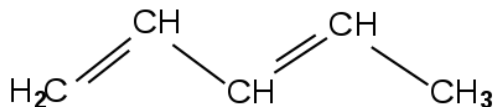
$$\langle p \rangle_{\psi_n} = \frac{\int_0^L \psi_n^*(x) \hat{p} \psi_n(x) dx}{\int_0^L \psi_n^*(x) \psi_n(x) dx} = 0$$

- Therefore the particle in a box **is on average** at the middle point of the box and **has on average** zero linear momentum.

10

Application problem No.1: Conjugated poleyenes

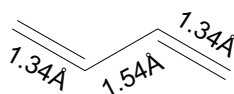
- A system of π -electrons in a linear conjugated poleyne can be thought of as a system of non-interacting electrons in a one-dimensional box



- Assumption 1.** The π -electrons do not interact with one another, but they obey the Pauli exclusion principle.
- Assumption 2.** The length L of the “box” is equal to the total length of the poleyne chain plus two radii of an sp^2 C atom

11

Example of the 1,3-butadine molecule



From standard reference tables:

$$\begin{aligned} R(\text{C}-\text{C}) &= 1.54 \text{ \AA} \\ R(\text{C}=\text{C}) &= 1.34 \text{ \AA} \\ r(sp^2 \text{ C}) &= 0.73 \text{ \AA} \end{aligned}$$

Using these values we find

$$L = 0.73 + 1.34 + 1.54 + 1.34 + 0.73 = 5.68 \text{ \AA}$$

$$V(x) = \begin{cases} 0 & ; \quad 0 \leq x \leq L \\ +\infty & ; \quad \text{otherwise} \end{cases}$$

12

Solving the Schrödinger equation for a particle in a box

- The potential for the system under discussion is given by

$$V(x) = \begin{cases} 0 ; & 0 \leq x \leq L \\ +\infty ; & \text{otherwise} \end{cases}$$

- Due to the piece-wise nature of the potential, the Schrödinger equation associated to this potential can be studied in three regions.
- Region I : $-\infty < x < 0 \rightarrow \psi_1(x)$
Region II : $0 \leq x \leq L \rightarrow \psi_2(x)$
Region III : $L < x < +\infty \rightarrow \psi_3(x)$
- In regions I and III the potential is “infinitely high”. Therefore the probability of finding the particle in those regions is zero

$$\psi_1(x) = \psi_3(x) = 0$$

13

Solving the Schrödinger equation for a particle in a box

- For region II the Schrödinger equation looks like

$$-\frac{\hbar^2}{2m} \frac{d^2\psi_2}{dx^2}(x) + 0 \cdot \psi_2(x) = E\psi_2(x)$$

which can be re-written as

$$\frac{d^2\psi_2}{dx^2}(x) + k^2\psi_2(x) = 0 \quad ; \quad k^2 = \frac{2mE}{\hbar^2}$$

- The general solution to this second order differential equation is

$$\psi_2(x) = A \sin(kx) + B \cos(kx)$$

- In order to determine the constants A and B we need to enforce conditions on the wave-function to make it well-behaved

14

Solving the Schrödinger equation for a particle in a box

- In order to make the wave-function continuous, the following conditions should be satisfied

$$\begin{aligned}\psi_2(0) = \psi_1(0) = 0 &\Rightarrow B = 0 \\ \psi_2(L) = \psi_3(L) = 0 &\Rightarrow A \sin(kL) = 0\end{aligned}$$

- The second condition implies that either $A = 0$ or $\sin(kL) = 0$. The former cannot be, otherwise would imply there is no wave-function in region II; **which would imply having probability zero of finding the particle in region III!**
- Therefore $\sin(kL) = 0$, which implies

$$k = \frac{n\pi}{L}$$

- The only part left to be determined is A ; which is computed by normalizing the wave-function.

15

Normalizing the wave-function for a particle in a box

- So far we have that the wave-function has the form

$$\psi(x) = \begin{cases} 0 & ; & x < 0 \\ A \sin\left(\frac{n\pi x}{L}\right) & ; & 0 \leq x \leq L \\ 0 & ; & x > L \end{cases}$$

- In order to determine the unknown A we normalize the wave-function. That is,

$$\int_{-\infty}^{+\infty} |\psi(x)|^2 dx = 1$$

16

Normalizing the wave-function for a particle in a box

- The previous integral leads to the following development

$$\begin{aligned}
 \int_{-\infty}^{+\infty} |\psi(x)|^2 dx &= \underbrace{\int_{-\infty}^0 |\psi_1(x)|^2 dx}_{\text{Region I}} + \underbrace{\int_0^L |\psi_2(x)|^2 dx}_{\text{Region II}} + \underbrace{\int_L^{+\infty} |\psi_3(x)|^2 dx}_{\text{Region III}} \\
 &= \underbrace{0}_{\text{Region I}} + \underbrace{A^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx}_{\text{Region II}} + \underbrace{0}_{\text{Region III}} \\
 &= A^2 \left(\frac{L}{2}\right)
 \end{aligned}$$

- Therefore, the normalization constant is

$$A = \sqrt{\frac{2}{L}}$$

17

Normalized wave-function and eigenvalues for a particle in a box

- The normalized **wave-function** for a particle in a box is then

$$\psi_n(x) = \begin{cases} 0 & ; \quad x < 0 \\ \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) & ; \quad 0 \leq x \leq L \\ 0 & ; \quad x > L \end{cases}$$

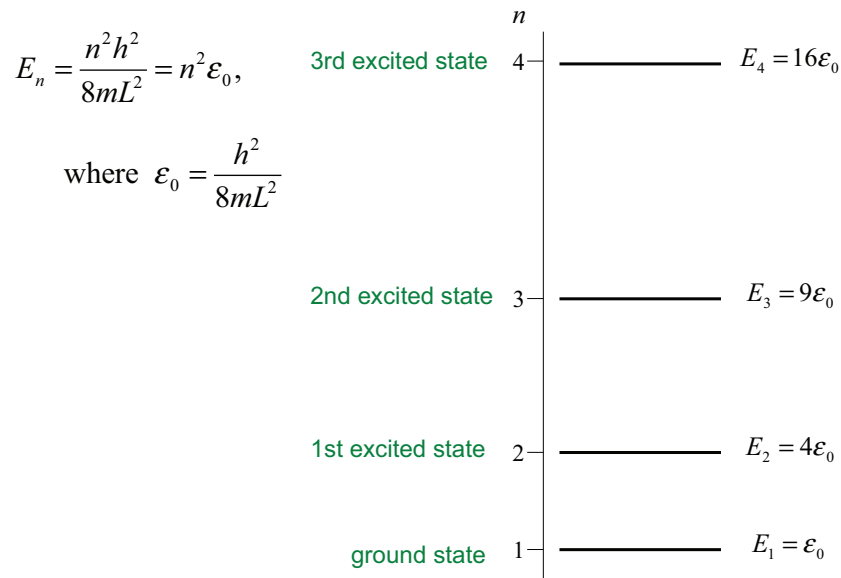
- From slides 11 and 12 of this lecture we conclude

$$k^2 = \frac{2mE}{\hbar^2} \quad ; \quad k = \frac{n\pi}{L} \quad \Rightarrow \quad E_n = \frac{\pi^2 \hbar^2}{2mL^2} n^2 = \frac{h^2}{8mL^2} n^2$$

- Notice that the energy levels are an integer multiple of a “basic energy” for the system. That is, **energy is quantized!**

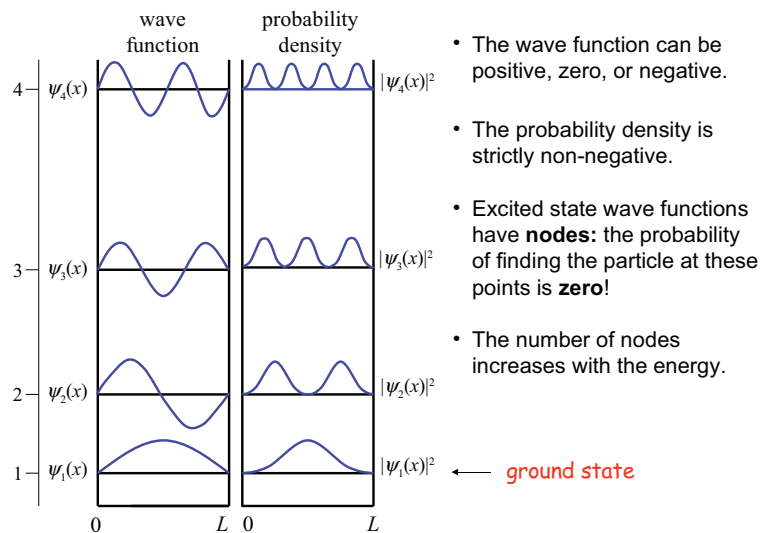
18

Energy level diagram for a particle in a box



19

Wave-functions and probability densities for a particle in a box



20