

Vv256 Lecture 17

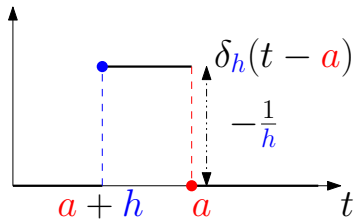
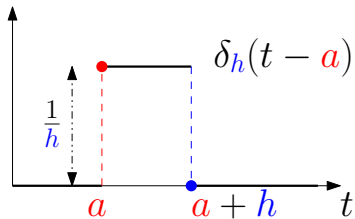
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- Mechanical systems are often acted on by external forces of **large magnitude** that act only for a **very short** period of time. e.g. golf ball hit by a club.
- Such force can be modelled by step functions,

$$\delta_h(t - a) = \frac{u(t - a) - u(t - (a + h))}{h}$$



Definition

This step function $\delta_h(t - a)$ is sometimes called a **unit impulse function**, since

$$\int_0^{\infty} \delta_h(t - a) dt = 1$$

Exercise

Find the Laplace transform of

$$\delta_h(t - a)$$

Solution

- By the definition of the unit impulse function, we have

$$\begin{aligned}\mathcal{L}[\delta_h(t - a)] &= \mathcal{L}\left[\frac{u(t - a) - u(t - (a + h))}{h}\right] \\ &= \frac{1}{h} \left(\mathcal{L}[1 \cdot u(t - a)] - \mathcal{L}[1 \cdot u(t - (a + h))] \right)\end{aligned}$$

- By the alternative second translation theorem, we have

$$\mathcal{L}[\delta_h(t - a)] = \frac{e^{-as} - e^{-(a+h)s}}{hs} = e^{-as} \left(\frac{1 - e^{-hs}}{hs} \right), \quad s > 0$$

Exercise

Solve the following initial-value problem

$$\ddot{y} + y = \delta_h(t - 2\pi), \quad y(0) = 1, \quad y'(0) = 0$$

Solution

- Taking the Laplace transform,

$$\begin{aligned}\mathcal{L}[\ddot{y}] + \mathcal{L}[y] &= \mathcal{L}[\delta_h(t - 2\pi)] \\ s^2 Y(s) - s + Y(s) &= e^{-2\pi s} \left(\frac{1 - e^{-hs}}{hs} \right) \\ Y(s) &= \frac{s}{s^2 + 1} + \left[\frac{1}{s^2 + 1} \right] \cdot \left[e^{-2\pi s} \left(\frac{1 - e^{-hs}}{hs} \right) \right]\end{aligned}$$

- Thus the solution according to the convolution theorem is given by

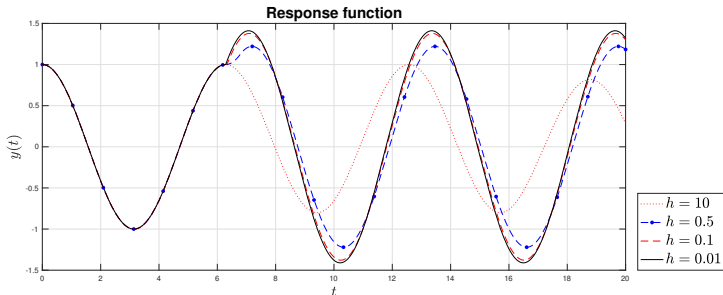
$$y = \cos t + \int_0^t \sin(\tau) \delta_h(t - \tau - 2\pi) d\tau$$

Solution

- Computing the integral, we have

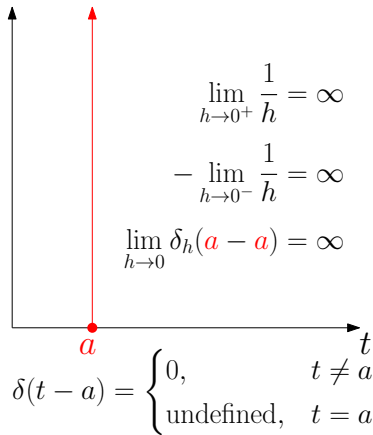
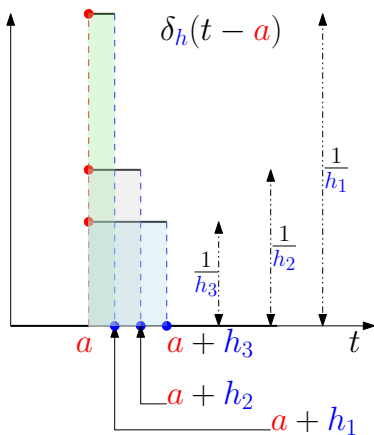
$$\begin{aligned} y &= \cos t + \int_0^t \sin(\tau) \delta_h(t - \tau - 2\pi) d\tau \\ &= \cos t + \frac{1}{h} \left[\left(\cos(t - h) - 1 \right) u\left(t - (2\pi + h)\right) - \left(\cos(t) - 1 \right) u\left(t - 2\pi\right) \right] \end{aligned}$$

- The following shows how the solution changes as h approaches zero.



- In practice it is convenient to work with the limit of the unit impulse function

$$\delta(t - a) = \lim_{h \rightarrow 0} \delta_h(t - a)$$



Definition

This generalization $\delta(t - a) = \lim_{h \rightarrow 0} \delta_h(t - a)$ is called the **Dirac delta function**.

- The Dirac delta function $\delta(t - a)$ is **not** a function in the usual sense, and it is an idealization of what happens in practice. e.g. elastic collision.
- From the definition of $\delta_h(t - a)$, we can see that $\delta(t - a)$ is essentially,

$$\delta(t - a) = \lim_{h \rightarrow 0} \delta_h(t - a) = \begin{cases} \infty, & t = a \\ 0, & t \neq a \end{cases}$$

with the following property by construction,

$$\begin{aligned} \int_{-\infty}^{\infty} \delta(t - a) dt &= \int_{-\infty}^{\infty} \lim_{h \rightarrow 0} \delta_h(t - a) dt \\ &= \lim_{h \rightarrow 0} \int_{-\infty}^{\infty} \delta_h(t - a) dt = \lim_{h \rightarrow 0} 1 = 1 \end{aligned}$$

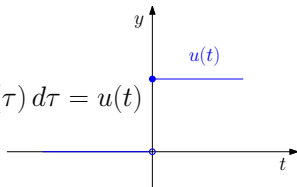
- Of course for an interval $[c, d]$ such that $c < a < d$, we have the same result

$$\int_c^d \delta(t - a) dt = \int_{-\infty}^{\infty} \delta(t - a) dt = 1$$

- In fact, it is also true for smaller and smaller interval around a , that is,

$$\begin{aligned} \int_{a^-}^{a^+} \delta(t - a) dt &= \lim_{c \rightarrow a^-} \int_c^{a^+} \delta(t - a) dt = 1 \\ \int_c^{a^+} \delta(t - a) dt &= 1 \\ \Rightarrow \int_{a^-}^{a^+} \delta(t - a) dt &= 1 \end{aligned}$$

- In some sense $\dot{u}(t - a) = \delta(t - a)$ since $\int_{-\infty}^t \delta(\tau) d\tau = u(t)$



- With the definition we are using,

$$\delta(t - a) = \lim_{h \rightarrow 0} \delta_h(t - a)$$

- We have the following properties:

1. Because $\delta(t - a)$ is the limit of graphs of area 1, the area under its graph is 1

$$\int_c^d \delta(t - a) dt = \begin{cases} 1 & c < a < d, \\ 0 & \text{otherwise.} \end{cases}$$

2. For any continuous function $f(t)$, we have $f(t)\delta(t - a) = f(a)\delta(t - a)$ and

$$\int_c^d f(t)\delta(t - a) dt = \begin{cases} f(a) & c < a < d, \\ 0 & \text{otherwise.} \end{cases}$$

3. For a unit step function $u(t - a)$, the Dirac delta function

$$\delta(t - a) = u'(t - a)$$

is defined to be the generalized derivative.

4. Suppose it is valid to interchange the order of limiting operations,

$$\mathcal{L}[\delta(t - a)] = \mathcal{L}\left[\lim_{h \rightarrow 0} \delta_h(t - a)\right] = \lim_{h \rightarrow 0} \mathcal{L}[\delta_h(t - a)]$$

then we can obtain the Laplace transform of the Dirac delta function

Theorem

$$\mathcal{L}[\delta(t - a)] = e^{-sa} \quad \text{where } a > 0$$

Proof

• Assuming it is valid to interchange the order of limiting operations, we have

$$\begin{aligned} \mathcal{L}[\delta(t - a)] &= \lim_{h \rightarrow 0} \mathcal{L}[\delta_h(t - a)] = \lim_{h \rightarrow 0} e^{-as} \left(\frac{1 - e^{-hs}}{hs} \right), \quad s > 0 \\ &= e^{-as} \cdot 1 \\ &= e^{-as}, \quad s > 0 \end{aligned}$$

according to L'Hospital's rule.

Exercise

Consider the following IVP again, this time with the Dirac delta function instead.

$$\ddot{y} + y = \delta(t - 2\pi), \quad y(0) = 1, \quad y'(0) = 0$$

Solution

- Taking the Laplace transform, we have

$$s^2 Y(s) - s + Y(s) = e^{-2\pi s} \implies Y(s) = \frac{s}{s^2 + 1} + e^{-2\pi s} \frac{1}{s^2 + 1}$$

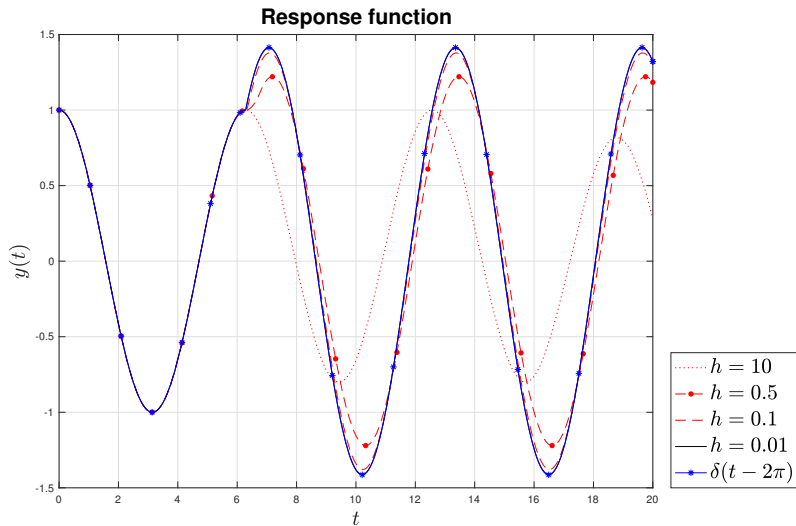
- Thus according to the convolution theorem, we have

$$y = \cos(t) + \int_0^t \delta(\tau - 2\pi) \sin(t - \tau) d\tau$$

- Applying property 4 of the Dirac delta function, we have

$$y(t) = \cos t + \sin(t - 2\pi)u(t - 2\pi) = \cos t + \sin(t)u(t - 2\pi)$$

- Notice the solution to this IVP is fairly close to the cases with small h ,



- So far, we have avoided having discontinuity at

$$t = 0$$

however, there is a need to distinguish the followings

$$0^-, \quad 0, \quad 0^+$$

when comes to the unilateral Laplace transform of a function in general.

- In general, we actually need the following

$$\mathcal{L}_- [f(t)] = \int_{0^-}^{\infty} f(t) e^{-st} dt$$

when the singularity happens at $t = 0$ to avoid inconsistency, for example

$$\delta(t)$$