

# Vv256 Lecture 16

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- A crucial step of the Laplace method is finding the inverse Laplace transform

$$Y(s) \longrightarrow y(t)$$

- Definition:

$$\mathcal{L}[y(t)] = \int_0^{\infty} e^{-st} y(t) dt = Y(s)$$

- Linear property:

$$Y(s) = c_1 Y_1(s) + c_2 Y_2(s)$$

- Transform of a derivative:

$$\mathcal{L}[f^{(n)}(t)] = s^n F(s) - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - s^0 f^{(n-1)}(0),$$

- Translation Theorems:

$$\mathcal{L}[e^{at} f(t)] = F(s - a) \quad \text{and} \quad \mathcal{L}[f(t - a)u(t - a)] = e^{-as} F(s)$$

Q: What is the inverse Laplace transform of a product?

$$H(s) = F(s)G(s) \quad \text{where} \quad \begin{aligned} \mathcal{L}[f(t)] &= F(s) \\ \mathcal{L}[g(t)] &= G(s) \end{aligned}$$

- If one of those function is

$$e^{-as}$$

then we can use the second translation theorem

$$\mathcal{L}[f(t-a)u(t-a)] = e^{-as}F(s)$$

Q: What happens otherwise? For example,

$$H(s) = \frac{1}{s^2(s-1)} = \frac{1}{s-1} - \frac{1}{s} - \frac{1}{s^2} \quad \text{or} \quad R(s) = \frac{s}{(s^2+k^2)^2}$$

- We could use a PFD for  $H$ . However, it is not always possible. e.g.  $R(s)$ .

- Firstly notice

$$\mathcal{L}^{-1}[H(s)] \neq \mathcal{L}^{-1}[F(s)] \cdot \mathcal{L}^{-1}[G(s)]$$

- For that, consider

$$\mathcal{L}^{-1}\left[\frac{1}{s^2(s-1)}\right] = \mathcal{L}^{-1}\left[\frac{1}{s-1} - \frac{1}{s} - \frac{1}{s^2}\right] = e^t - 1 - t$$

whereas

$$\mathcal{L}^{-1}\left[\frac{1}{s^2}\right] \cdot \mathcal{L}^{-1}\left[\frac{1}{s-1}\right] = te^t$$

- Hence

$$\mathcal{L}^{-1}\left[\frac{1}{s^2(s-1)}\right] \neq \mathcal{L}^{-1}\left[\frac{1}{s^2}\right] \cdot \mathcal{L}^{-1}\left[\frac{1}{s-1}\right]$$

## Convolution Theorem

If  $f(t)$  and  $g(t)$  are piecewise continuous on  $[0, \infty)$  and of exponential order, then

$$\int_0^{\infty} e^{-st} \left( \int_0^t f(\tau) g(t - \tau) d\tau \right) dt = \mathcal{L}[f(t)] \mathcal{L}[g(t)] = F(s)G(s)$$

### Proof

- Let

$$F(s) = \int_0^{\infty} e^{-s\tau} f(\tau) d\tau, \quad \text{and} \quad G(s) = \int_0^{\infty} e^{-s\mu} g(\mu) d\mu$$

- Take the right-hand side, we have

$$\begin{aligned} F(s)G(s) &= \left( \int_0^{\infty} e^{-s\tau} f(\tau) d\tau \right) \left( \int_0^{\infty} e^{-s\mu} g(\mu) d\mu \right) \\ &= \int_0^{\infty} \int_0^{\infty} e^{-s(\tau+\mu)} f(\tau) g(\mu) d\mu d\tau \end{aligned}$$

## Proof

$$\begin{aligned} F(s)G(s) &= \int_0^\infty \left( \int_0^\infty e^{-s(\tau+\mu)} f(\tau) g(\mu) d\mu \right) d\tau \\ &= \int_0^\infty f(\tau) \left( \int_0^\infty e^{-s(\tau+\mu)} g(\mu) d\mu \right) d\tau \end{aligned}$$

- To compute the inner integral, we hold  $\tau$  as fixed, and integrate w.r.t.  $\mu$ ,

$$\int_0^\infty e^{-s(\tau+\mu)} g(\mu) d\mu$$

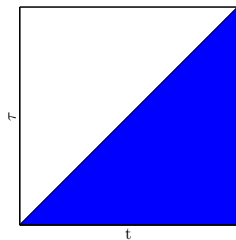
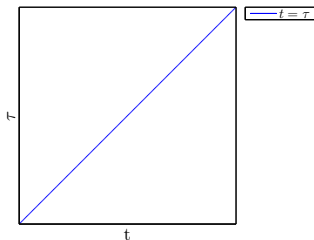
- However, we don't evaluate it directly, let  $t = \tau + \mu$ , and thus  $\frac{dt}{d\mu} = 1$

$$\begin{aligned} F(s)G(s) &= \int_0^\infty f(\tau) \left( \int_0^\infty e^{-s(\tau+\mu)} g(\mu) d\mu \right) d\tau \\ &= \int_0^\infty f(\tau) \left( \int_{t=\tau}^\infty e^{-st} g(t-\tau) dt \right) d\tau \end{aligned}$$

## Proof

- Consider what region we are integrating over,

$$\int_0^\infty \int_{t=\tau}^\infty dt d\tau = \iint dA = \int_0^\infty \int_0^t d\tau dt$$



- $$F(s)G(s) = \iint e^{-st} f(\tau) g(t - \tau) dA = \int_0^\infty e^{-st} \left( \int_0^t f(\tau) g(t - \tau) d\tau \right) dt$$

## Definition

If  $f(t)$  and  $g(t)$  are piecewise continuous on the interval  $[0, \infty)$ , then the function

$$\int_0^t f(\tau)g(t - \tau) d\tau$$

is known as the **convolution** of  $f$  and  $g$ . This special function is denoted as

$$(f * g)(t) = \int_0^t f(\tau)g(t - \tau) d\tau$$

- The convolution of  $f$  and  $g$ , which is a function of  $t$ , is often considered as a product between  $f$  and  $g$  since it has some properties of a typical product.

$$f * g = g * f \quad \text{commutative}$$

$$f * (g * h) = (f * g) * h \quad \text{associative}$$

$$f * (g + h) = f * g + f * h \quad \text{distributive}$$



- Recall we had

$$\mathcal{L}^{-1}\left[\frac{1}{s^2(s-1)}\right] = \mathcal{L}^{-1}\left[\frac{1}{s-1} - \frac{1}{s} - \frac{1}{s^2}\right] = e^t - 1 - t$$

- Let us apply the convolution theorem

$$\begin{aligned}\mathcal{L}^{-1}\left[\frac{1}{s^2(s-1)}\right] &= \mathcal{L}^{-1}\left[\frac{1}{s^2}\right] * \mathcal{L}^{-1}\left[\frac{1}{(s-1)}\right] \\ &= (t) * (e^t) = (e^t) * (t) = \int_0^t f(\tau)g(t-\tau) d\tau \\ &= \int_0^t e^\tau(t-\tau) d\tau \\ &= e^t - t - 1\end{aligned}$$

## Exercise

Find the inverse Laplace transform of  $\frac{1}{(s^2 + k^2)^2}$ .

## Solution

- Let  $F(s) = G(s) = \frac{k}{s^2 + k^2}$ , so

$$f(t) = g(t) = \mathcal{L}^{-1}\left[\frac{k}{s^2 + k^2}\right] = \sin kt$$

- The convolution theorem states that  $\mathcal{L}[f * g] = F(s)G(s)$

$$\begin{aligned}\mathcal{L}^{-1}[F(s)G(s)] &= f * g = \int_0^t \sin k\tau \sin k(t - \tau) d\tau \\ &= \frac{1}{2k} \sin kt - \frac{t}{2} \cos kt\end{aligned}$$

- $$\begin{aligned}\mathcal{L}^{-1}\left\{\frac{1}{(s^2 + k^2)^2}\right\} &= \frac{1}{k^2} \mathcal{L}^{-1}\left\{\frac{k^2}{(s^2 + k^2)^2}\right\} \\ &= \frac{1}{k^2} \left(\frac{1}{2k} \sin kt - \frac{t}{2} \cos kt\right) = \frac{\sin kt - kt \cos kt}{2k^3}\end{aligned}$$

# Series Circuits

- In a series circuit, Kirchhoff's second law states that the sum of the voltage drops across an inductor, resistor, and capacitor is equal to the voltage

$$E(t)$$

- It is known that voltage drops across an inductor, resistor, and capacitor are,

$$\Delta V_L = L \frac{di}{dt}, \quad \Delta V_R = Ri(t), \quad \text{and} \quad \Delta V_C = \frac{1}{C} \int_0^t i(\tau) d\tau,$$

respectively, where  $i(t)$  is the current and  $L$ ,  $R$ , and  $C$  are constants.

- So the current in a  $LRC$  series circuit is governed by the following equation,

$$L \frac{di}{dt} + Ri(t) + \frac{1}{C} \int_0^t i(\tau) d\tau = E(t)$$

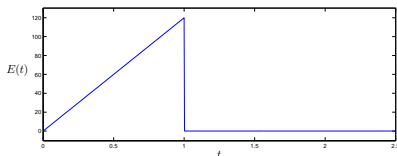
## Exercise

Determine the current  $i(t)$  in a series  $LRC$  circuit when

$$L = 0.1, \quad R = 2, \quad C = 0.1, \quad \text{and} \quad i(0) = 0,$$

and voltage is given by

$$E(t) = 120t - 120tu(t - 1).$$



## Solution

- To find the current  $i(t)$ , we need to solve the following

$$\frac{1}{10} \frac{di}{dt} + 2i + 10 \int_0^t i(\tau) d\tau = 120t - 120tu(t - 1)$$

- This is known as **integrodifferential equation**.

## Solution

- First notice that, in general,

$$f * g = \int_0^t f(\tau)g(t - \tau) d\tau = \int_0^t f(\tau) d\tau \quad \text{for } g(t - \tau) = 1$$

- According to the convolution theorem,

$$\mathcal{L}[f * g] = \mathcal{L}[f] \cdot \mathcal{L}[g] = \frac{F(s)}{s}, \quad \text{where } \mathcal{L}[f] = F(s)$$

for the case when  $g(t - \tau) = 1$ .

- Back to the current  $i(t)$ ,

$$\mathcal{L}\left[\int_0^t f(\tau) d\tau\right] = \frac{F(s)}{s} \implies \mathcal{L}\left[\int_0^t i(\tau) d\tau\right] = \frac{I(s)}{s}$$

where  $I(s) = \mathcal{L}[i(t)]$ .

## Solution

- Thus the Laplace transform of the integrodifferential equation is

$$\mathcal{L}\left[\frac{1}{10}\frac{di}{dt} + 2i + 10\int_0^t i(\tau) d\tau\right] = \mathcal{L}\left[120t - 120tu(t-1)\right]$$

$$\Rightarrow \frac{1}{10}sI(s) + 2I(s) + 10\frac{I(s)}{s} = 120\left(\frac{1}{s^2} - e^{-s}\mathcal{L}[t+1]\right)$$

$$\Rightarrow \frac{1}{10}sI(s) + 2I(s) + 10\frac{I(s)}{s} = 120\left(\frac{1}{s^2} - e^{-s}\frac{1}{s^2} - e^{-s}\frac{1}{s}\right)$$

- Rearranging and collecting  $I(s)$ ,

$$I(s) = 1200\left[\frac{1}{s(s+10)^2} - \frac{1}{s(s+10)^2}e^{-s} - \frac{1}{(s+10)^2}e^{-s}\right]$$

## Solution

- By partial fractions,  $I(s) = 1200 \left[ \underbrace{\frac{1/100}{s}}_{(1)} - \underbrace{\frac{1/100}{s+10}}_{(2)} - \underbrace{\frac{1/10}{(s+10)^2}}_{(3)} - \underbrace{\frac{1/100}{s}}_{(4)} e^{-s} + \underbrace{\frac{1/100}{s+10}}_{(5)} e^{-s} - \underbrace{\frac{9/10}{(s+10)^2}}_{(6)} e^{-s} \right]$

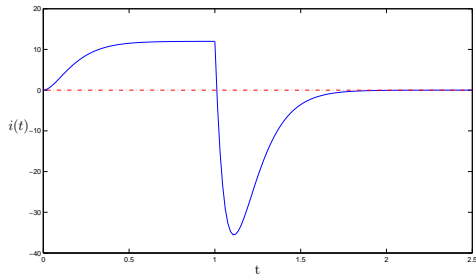
- Transforming back to the time domain,

$$i(t) = 1200 \left[ \underbrace{\frac{1}{100}}_{(1)} - \underbrace{\frac{1}{100} e^{-10t}}_{(2)} - \underbrace{\frac{1}{10} t e^{-10t}}_{(3)} - \underbrace{\frac{1}{100} u(t-1)}_{(4)} + \underbrace{\frac{1}{100} e^{-10(t-1)} u(t-1)}_{(5)} - \underbrace{\frac{9}{10} (t-1) e^{-10(t-1)} u(t-1)}_{(6)} \right]$$

## Solution

- Written as a piecewise-defined function,

$$i(t) = \begin{cases} 12 - 12e^{-10t} - 120te^{-10t}, & 0 \leq t < 1 \\ -12e^{-10t} + 12e^{-10(t-1)} - 120te^{-10t} - 1080(t-1)e^{-10(t-1)} & 1 \leq t. \end{cases}$$



- The input  $E(t)$  is discontinuous, however, the output  $i(t)$  is a continuous



- Consider a periodic function  $\mathcal{L}[f(t)]$  where  $f(t+T) = f(t)$ .

### Theorem

Let  $f(t)$  be piecewise continuous on  $[0, \infty)$  and of exponential order, and periodic with **period  $T$** , then

$$\mathcal{L}[f(t)] = \frac{1}{1 - e^{-sT}} \int_0^T e^{-st} f(t) dt$$

### Proof

- Break the interval into two parts,

$$\mathcal{L}[f(t)] = \int_0^T e^{-st} f(t) dt + \int_T^\infty e^{-st} f(t) dt$$

- When we let  $t = u + T$ , then the second integral becomes,

$$\int_0^\infty e^{-s(u+T)} f(u+T) du = e^{-sT} \int_0^\infty e^{-su} f(u) du = e^{-sT} \mathcal{L}[f(t)]$$

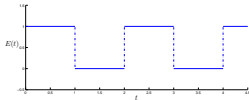
## Exercise

The current  $i(t)$  in a single-loop  $LR$  series circuit is governed by

$$L \frac{di}{dt} + Ri = E(t)$$

Find the current  $i(t)$  when  $i(0) = 0$  and  $E(t)$  is a square wave function

$$E(t) = \begin{cases} 1, & 0 \leq t \leq 1 \\ 0, & 1 \leq t \leq 2 \end{cases}, \quad E(t) = E(t+2)$$



## Solution

- Find the Laplace transform of input function  $E(t)$ ,

$$\begin{aligned} \mathcal{L}[E(t)] &= \frac{1}{1 - e^{-2s}} \int_0^2 e^{-st} E(t) dt \\ &= \frac{1}{1 - e^{-2s}} \int_0^1 e^{-st} \cdot 1 dt = \frac{1}{1 - e^{-2s}} \frac{1 - e^{-s}}{s} = \frac{1}{s(1 + e^{-s})} \end{aligned}$$

## Solution

- Applying the Laplace transform to the equation, we have

$$LsI(s) + RI(s) = \frac{1}{s(1 + e^{-s})}$$

- Rearranging and collecting  $I(s)$ , we have

$$I(s) = \frac{1/L}{s(s + R/L)} \cdot \frac{1}{1 + e^{-s}}$$

- To find the current in the  $t$ -domain, we first make use of power series,

$$1 - r + r^2 - r^3 \dots = \frac{1}{1 - (-r)} \implies \frac{1}{1 + e^{-s}} = 1 - e^{-s} + e^{-2s} - e^{-3s} \dots$$

- By partial fraction,

$$\frac{1/L}{s(s + R/L)} = \frac{1}{L} \left( \frac{L/R}{s} - \frac{L/R}{s + R/L} \right)$$

## Solution

- Together, we have

$$\begin{aligned} I(s) &= \frac{1}{R} \left( \frac{1}{s} - \frac{1}{(s + R/L)} \right) (1 - e^{-s} + e^{-2s} - e^{-3s} + \dots) \\ &= \frac{1}{R} \left( \frac{1}{s} - \frac{e^{-s}}{s} + \frac{e^{-2s}}{s} - \frac{e^{-3s}}{s} + \dots \right) \\ &\quad - \frac{1}{R} \left( \frac{1}{s + R/L} - \frac{e^{-s}}{s + R/L} + \frac{e^{-2s}}{s + R/L} - \frac{e^{-3s}}{s + R/L} + \dots \right) \end{aligned}$$

- Now we are ready to back transform!!!

$$\begin{aligned} i(s) &= \frac{1}{R} \left( 1 - u(t-1) + u(t-2) - u(t-3) + \dots \right) \\ &\quad - \frac{1}{R} \left( e^{-Rt/L} - e^{-R(t-1)/L} u(t-1) + e^{-R(t-2)/L} u(t-2) - \dots \right) \end{aligned}$$

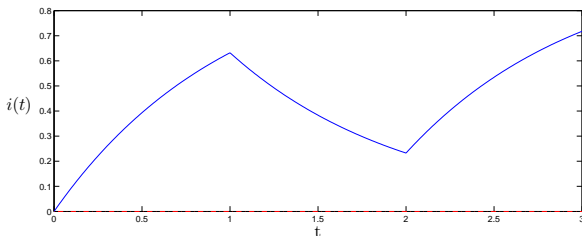
## Solution

- Tidy it up, we have

$$i(t) = \frac{1}{R} \left(1 - e^{-Rt/L}\right) + \frac{1}{R} \sum_{k=1}^{\infty} (-1)^k \left(1 - e^{-R(t-k)/L}\right) u(t-k)$$

- For  $R = 1$ ,  $L = 1$ , and let's consider  $0 \leq t \leq 3$ , then

$$i(t) = 1 - e^{-t} - (1 - e^{-(t-1)})u(t-1) + (1 - e^{-(t-2)})u(t-2)$$



- Recall we had our theorem on the Laplace transform of a derivative function

$$\mathcal{L}\left[\frac{df}{dt}\right] = s\mathcal{L}[f(t)] - f(0) = sF(s) - f(0)$$

Q: What will be the inverse Laplace transform of

$$\frac{d}{ds}F(s) \quad \text{where} \quad F(s) = \mathcal{L}[f(t)]$$

- Suppose  $f(t)$  is well-behaved, thus we can interchange the order of

integration and differentiation

$$\begin{aligned} \frac{d}{ds}F(s) &= \frac{d}{ds} \int_0^\infty e^{-st} f(t) dt = \int_0^\infty \frac{d}{ds} \left( e^{-st} f(t) \right) dt \\ &= \int_0^\infty -te^{-st} f(t) dt = -\mathcal{L}[tf(t)] \end{aligned}$$

## Exercise

Find the Laplace transform of  $te^{3t}$  using the last theorem.

## Solution

- We can treat it as a **translation**

$$\mathcal{L}\left[e^{at}f(t)\right] = F(s - a)$$

- Then

$$\begin{aligned}\mathcal{L}[te^{3t}] &= \frac{n!}{(s - a)^{n+1}} \\ &= \frac{1}{(s - 3)^2}\end{aligned}$$

for  $s > 3$ .

- We can treat it as a **derivative**

$$\mathcal{L}[tg(t)] = -\frac{d}{ds}F(s)$$

- Then

$$\begin{aligned}\mathcal{L}[te^{3t}] &= -\frac{d}{ds} \left( \frac{1}{s - a} \right) \\ &= -\frac{d}{ds} \left( \frac{1}{s - 3} \right) \\ &= \frac{1}{(s - 3)^2}\end{aligned}$$

for  $s > 3$ .

- We can use the last result to find the Laplace transform of

$$\begin{aligned}\mathcal{L}[t^2 f(t)] &= \mathcal{L}[t \cdot t f(t)] = -\frac{d}{ds} \mathcal{L}[t f(t)] = -\frac{d}{ds} \left( -\frac{d}{ds} \mathcal{L}[f(t)] \right) \\ &= \frac{d^2}{ds^2} \left( \mathcal{L}[f(t)] \right)\end{aligned}$$

- The preceding two cases suggest the following general result:

### Theorem

If  $F(s) = \mathcal{L}[f(t)]$  and  $n = 1, 2, 3, \dots$ , then

$$\mathcal{L}[t^n f(t)] = (-1)^n \frac{d^n}{ds^n} F(s)$$

Q: When is this theorem useful?



## Exercise

Solve  $t\ddot{y} - t\dot{y} + y = 2$ ,  $y(0) = 2$ ,  $\dot{y}(0) = -4$  using the method of Laplace.

## Solution

- Take the Laplace transform and apply the last theorem, we have

$$\begin{aligned} & -\frac{d}{ds}\mathcal{L}[\ddot{y}] + \frac{d}{ds}\mathcal{L}[\dot{y}] + \mathcal{L}[y] = \mathcal{L}[2] \\ & -\frac{d}{ds}\left(s^2Y(s) - sy(0) - y'(0)\right) + \frac{d}{ds}\left(sY(s) - y(0)\right) + Y(s) = \frac{2}{s} \\ & -2sY(s) - s^2Y'(s) + y(0) + Y(s) + sY'(s) + Y(s) = \frac{2}{s} \\ & Y'(s) + \frac{2}{s}Y(s) = \frac{2}{s^2} \end{aligned}$$

- From this first order linear equation, and the initial condition  $\dot{y}(0) = -4$

$$Y(s) = 2s^{-1} + Cs^{-2} \implies y(t) = 2 + ct \implies y(t) = 2 - 4t$$

- A linear equation with variable coefficients can be solved by using the method of Laplace transform, but it is not a universal method.

Using series is still the standard method for variable coefficients

- However, a second approach offers a way to compute various things, e.g. the Laplace transform of special functions

Q: How can we find the closed-form for the Laplace transform of  $J_0$ .

$$\begin{aligned}\mathcal{L}[J_0] &= \mathcal{L}\left[\sum_{n=0}^{\infty} \frac{(-1)^n}{n!\Gamma(1+\lambda+n)} \left(\frac{t}{2}\right)^{2n+\lambda}\right] \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n}{n!\Gamma(1+n)} \left(\frac{1}{2}\right)^{2n} \mathcal{L}[t^{2n}] \\ &= \sum_{n=0}^{\infty} \frac{(-1)^n}{n!\Gamma(1+n)} \left(\frac{1}{2}\right)^{2n} \frac{(2n)!}{s^{2n+1}}\end{aligned}$$

Q: The closed form is not obvious, what shall we do next?

- Let us consider the values of  $J_0$  and  $J'_0$  at  $t = 0$ .

$$J_0(0) = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \Gamma(1+n)} \left(\frac{t}{2}\right)^{2n} \Big|_{t=0} = 1 + 0 + 0 + \cdots = 1$$

and

$$J'_0(0) = \sum_{n=1}^{\infty} \frac{(-1)^n}{n! \Gamma(1+n)} \frac{2n}{2} \left(\frac{t}{2}\right)^{2n-1} \Big|_{t=0} = 0$$

- Thus the following initial-value problem defines  $J_0$ .

$$t^2 \ddot{y} + t \dot{y} + (t^2 - \lambda^2)y = 0 \quad y(0) = 1 \quad \dot{y}(0) = 0$$

where  $\lambda = 0$ .

Q: Why is this useful? What shall we do next?

- Take the Laplace transform, we have

$$\mathcal{L}[t\ddot{y} + \dot{y} + ty] = \mathcal{L}[0]$$

- Apply the last theorem, we have

$$-\frac{d}{ds}\mathcal{L}[\ddot{y}] + \mathcal{L}[\dot{y}] - \frac{d}{ds}\mathcal{L}[y] = \mathcal{L}[0]$$

- Apply the theorem of the Laplace transform of a derivative,

$$-\frac{d}{ds}\left(s^2Y(s) - sy(t_0) - y'(t_0)\right) + sY(s) - y(t_0) - \frac{dY}{ds} = 0$$

$$-2sY(s) - s^2Y'(s) + y(t_0) + sY(s) - y(t_0) - Y'(s) = 0$$

$$-(s^2 + 1)Y'(s) - sY(s) = 0$$

- Solve this first order linear equation, we have

$$Y(s) = \frac{C}{\sqrt{s^2 + 1}} \quad \text{where } C \text{ is an arbitrary constant.}$$

Q: How can we determine the arbitrary constant  $C$ ?

$$Y(s) = \frac{C}{\sqrt{s^2 + 1}}$$

- We need some connection between initial values in  $t$ -domain and  $s$ -domain.

### Theorem

If  $f(t)$  and  $f'(t)$  are continuous, and of exponential order, then

$$\lim_{t \rightarrow 0} f(t) = \lim_{s \rightarrow \infty} sF(s) \quad \text{where} \quad F(s) = \int_0^{\infty} e^{-st} f(t) dt$$

Furthermore,

$$\lim_{t \rightarrow \infty} f(t) = \lim_{s \rightarrow 0} sF(s) \quad \text{where} \quad F(s) = \int_0^{\infty} e^{-st} f(t) dt$$

provided the limit of  $f(t)$  exists when  $t \rightarrow \infty$ .

- This is often known as the initial-value/final-value theorem.

- Since  $J_0$  converges for  $t \in \mathbb{R}$ , and it is a solution to a second-order equation

$$\lim_{t \rightarrow 0} J_0 = J_0(0) = 1$$

- So the initial-value theorem implies

$$1 = \lim_{s \rightarrow \infty} s \mathcal{L}[J_0] = \lim_{s \rightarrow \infty} s \frac{C}{\sqrt{s^2 + 1}} = \lim_{s \rightarrow \infty} \frac{C}{\sqrt{1 + \frac{1}{s^2}}} = C$$

- Hence the Laplace transform of Bessel function of the first kind of order 0 is

$$\mathcal{L}[J_0] = \frac{1}{\sqrt{s^2 + 1}}$$

### Exercise

Solve  $t\ddot{y} + \dot{y} + 4ty = 0$ ,  $y(0) = 3$ ,  $\dot{y}(0) = 0$  using the method of Laplace.

## Solution

- Taking the Laplace transform, we obtain

$$-\frac{d}{ds} \left( s^2 Y(s) - sy(0) - y'(0) \right) + sY(s) - y(0) - 4 \frac{dY}{ds} = 0$$

- Use the initial conditions and simplify, we have

$$(s^2 + 4)Y'(s) + sY(s) = 0$$

- Solve this first-order linear equation, and use the following theorem, we have

$$Y = \frac{C}{\sqrt{s^2 + 4}} = \frac{C}{2\sqrt{\left(\frac{s}{2}\right)^2 + 1}} \implies y(t) = C J_0(2t) \implies y(t) = 3 J_0(2t)$$

## Theorem

Let  $F(s) = \int_0^\infty e^{-st} f(t) dt$ , then  $\frac{1}{a} F\left(\frac{s}{a}\right) = \int_0^\infty e^{-st} f(at) dt$ , where  $a > 0$ .