

## VP390 Problem Set 6

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### 1 Problem 1

For this problem, I choose the parameters of Gauss Wave Packet as:  
 $t_0 = 1, \alpha = 10^{-34}, m = 9 \times 10^{-31}, p_0 = 10^{30}$   
Then we get graphs when  $t = 0, 1, 2, 3$  as:

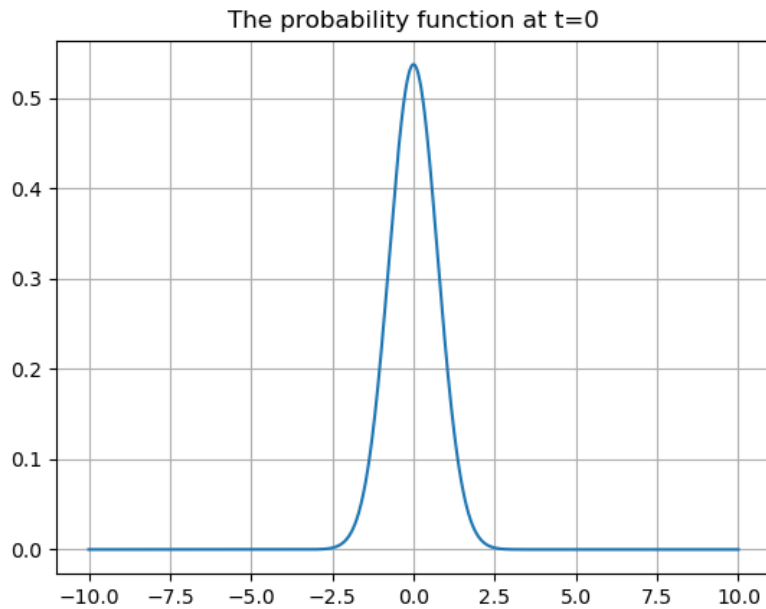


Figure 1:  $t = 0$

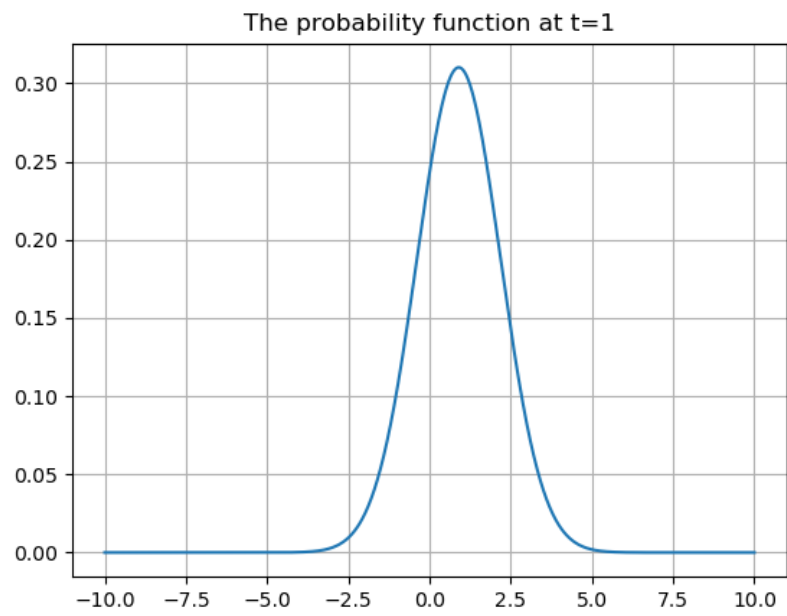


Figure 2:  $t = 1$

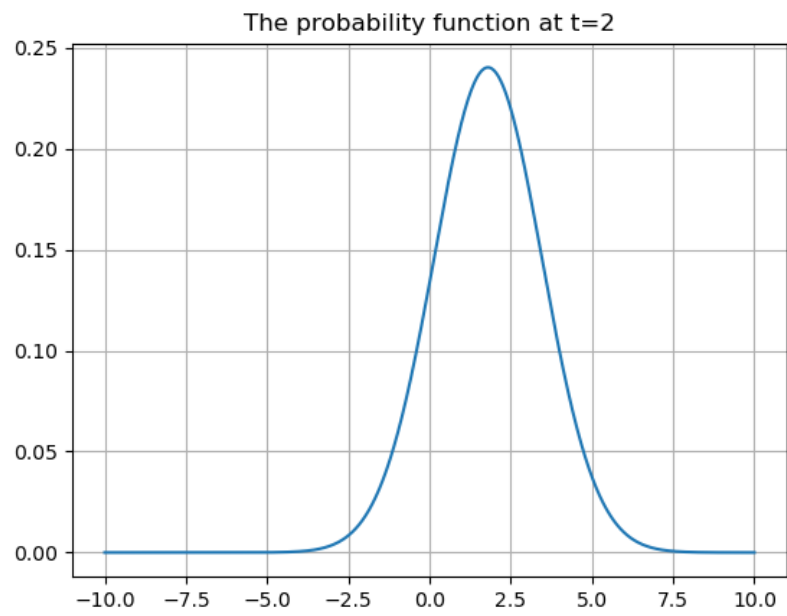


Figure 3:  $t = 2$

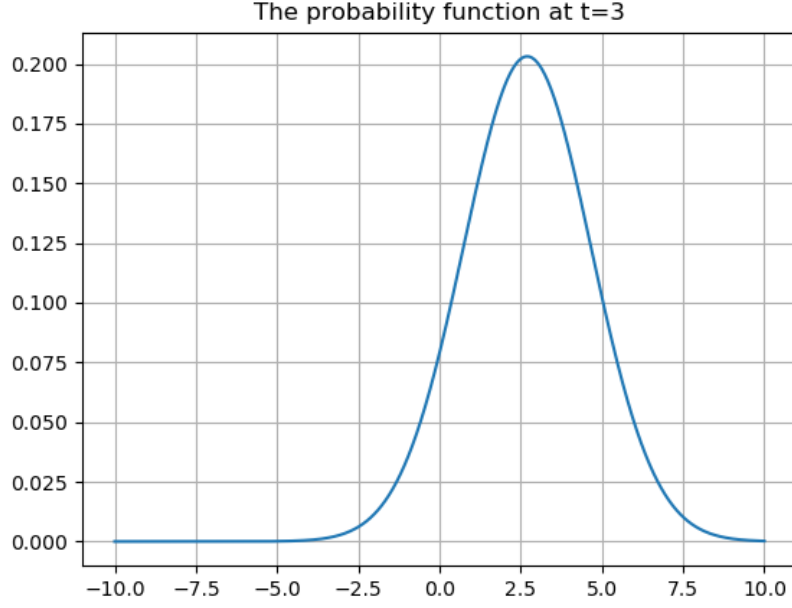


Figure 4:  $t = 3$

## 2 Problem 2

- (a) For the third maximum point, we get  $E = V_0 + \frac{\hbar^2 \pi^2}{8ma^2} 3^2$   
 $100 \times 1.6 \times 10^{-19} = 9 \frac{(1.054 \times 10^{-34})^2 \pi^2}{8 \times 9.1 \times 10^{-31} \times (10^{-10})^2} + V_0$   
 $V_0 = 15.28 \text{ eV}$

- (b) For  $E \rightarrow V_0$ ,  $T \rightarrow \frac{1}{1 + \frac{2ma^2 V_0}{\hbar^2}}$   
Then  $a = 1.0626 \times 10^{-9} \text{ m}$ , and the width is  $2.13 \times 10^{-9} \text{ m}$

## 3 Problem 3

For the barrier problem we get three equations:

$$\begin{cases} \psi_1(x) = Ae^{ik_1 x} + Be^{-ik_1 x} & x \leq -a \\ \psi_2(x) = Ce^{ik_2 x} + De^{-ik_2 x} & |x| < a \\ \psi_3(x) = Fe^{ik_1 x} & x \geq a \end{cases}$$

where  $k_1 = \frac{\sqrt{2mE}}{\hbar}$  and  $k_2 = \frac{\sqrt{2m(E-V_0)}}{\hbar}$

Since  $\psi(x)$  and  $\psi'(x)$  are continuous at  $x = \pm a$ , we get:

$$\begin{cases} Ae^{-ik_1 a} + Be^{ik_1 a} = Ce^{-ik_2 a} + De^{ik_2 a} & (1) \\ Ce^{ik_2 a} + De^{-ik_2 a} = Fe^{ik_1 a} & (2) \\ k_1(Ae^{-ik_1 a} - Be^{ik_1 a}) = k_2(Ce^{-ik_2 a} - De^{ik_2 a}) & (3) \\ k_2(Ce^{ik_2 a} - De^{-ik_2 a}) = k_1 Fe^{ik_1 a} & (4) \end{cases}$$

$$\begin{aligned}
k_2 * (2) - (4) \text{ we get } D &= \frac{k_2 - k_1}{2k_2} F e^{i(k_1 + k_2)a} \\
k_2 * (2) + (4) \text{ we get } C &= \frac{k_1 + k_2}{2k_2} F e^{i(k_1 - k_2)a} \\
k_1 * (1) + (3) \text{ we get } 2k_1 A &= (k_1 + k_2) C e^{i(k_1 - k_2)a} + (k_1 - k_2) D e^{i(k_1 + k_2)a} \\
\text{Then } \frac{A}{F} &= \frac{1}{4k_1 k_2} [(k_1 + k_2)^2 e^{i(k_1 - 2k_2)a} + (k_1 - k_2)^2 e^{i(k_1 + 2k_2)a}] \\
\frac{F}{A} &= \frac{4k_1 k_2}{(k_1 + k_2)^2 e^{i(k_1 - 2k_2)a} + (k_1 - k_2)^2 e^{i(k_1 + 2k_2)a}} = \frac{4k_1 k_2 e^{i(-k_1 + 2k_2)a}}{(k_1 + k_2)^2 + (k_1 - k_2)^2 e^{4ik_2 a}} \\
k_1 * (1) - (3) \text{ we get } 2k_1 B &= (k_1 - k_2) C e^{-i(k_1 + k_2)a} + (k_1 + k_2) D e^{i(k_2 - k_1)a} \\
\frac{B}{F} &= \frac{1}{4k_1 k_2} [(k_1^2 - k_2^2) e^{-i2k_2 a} + (k_2^2 - k_1^2) e^{i2k_2 a}] \\
\frac{B}{A} &= \frac{(k_1^2 - k_2^2) e^{-i2k_2 a} + (k_2^2 - k_1^2) e^{i2k_2 a}}{(k_1 + k_2)^2 e^{i(k_1 - 2k_2)a} + (k_1 - k_2)^2 e^{i(k_1 + 2k_2)a}} = \frac{(k_1^2 - k_2^2) e^{-ik_1 a} + (k_2^2 - k_1^2) e^{i(-k_1 + 4k_2 a)}}{(k_1 + k_2)^2 + (k_1 - k_2)^2 e^{4ik_2 a}} \\
\text{Hence } T + R &= \left(\frac{F}{A}\right)^2 + \left(\frac{B}{A}\right)^2 = 1
\end{aligned}$$

## 4 Problem 4

For the barrier problem we get three equations:

$$\begin{cases} \psi_1(x) = A e^{ik_1 x} + B e^{-ik_1 x} & x \leq -a \\ \psi_2(x) = C e^{ik_2 x} + D e^{-ik_2 x} & |x| < a \\ \psi_3(x) = F e^{ik_1 x} & x \geq a \end{cases}$$

$$\text{where } k_1 = \frac{\sqrt{2mE}}{\hbar} \text{ and } k_2 = \frac{\sqrt{2m(E+V_0)}}{\hbar}$$

We can plug them into the solution of rectangular barrier since the equations are in same forms by replacing  $V_0$  with  $-V_0$

$$\text{Then we get } \frac{1}{T} = 1 + \frac{1}{4} \frac{V_0^2}{E(E+V_0)} \sin^2(2\sqrt{\frac{2ma^2}{\hbar^2}}(E+V_0))$$

$$\text{Hence we get the result } \frac{1}{T} = 1 + \frac{1}{4} \frac{V_0^2}{E(E+V_0)} \sin^2(2k_2 a)$$

$$\text{For perfect transmission, we need } 2k_2 a = 2\sqrt{\frac{2ma^2}{\hbar^2}}(E+V_0) = n\pi$$

$$\text{Then } E = \frac{\pi^2 \hbar^2}{8ma^2} n^2 - V_0$$

## 5 Problem 5

$$\begin{aligned}
E &= V_0 + \frac{\hbar^2 \pi^2}{8ma^2} n^2 \\
a &= 0.5 \times 10^{-10}, n = 1, E = 0.9 \times 1.6 \times 10^{-19} \\
V_0 &= -5.885 \times 10^{-18} \text{ J}
\end{aligned}$$

## 6 Problem 6

- (a)  $-\hat{B}, \hat{A}] = -\hat{B}\hat{A} + \hat{A}\hat{B} = \hat{A}\hat{B} - \hat{B}\hat{A} = [\hat{A}, \hat{B}]$
- (b)  $\hat{A}[\hat{B}, \hat{C}] + [\hat{A}, \hat{C}]\hat{B} = \hat{A}\hat{B}\hat{C} - \hat{A}\hat{C}\hat{B} + \hat{A}\hat{C}\hat{B} - \hat{C}\hat{A}\hat{B} = (\hat{A}\hat{B})\hat{C} - \hat{C}(\hat{A}\hat{B}) = [\hat{A}\hat{B}, \hat{C}]$
- (c)  $[\hat{A}, \hat{C}] + [\hat{B}, \hat{C}] = \hat{A}\hat{C} - \hat{C}\hat{A} + \hat{B}\hat{C} - \hat{C}\hat{B} = (\hat{A} + \hat{B})\hat{C} - \hat{C}(\hat{A} + \hat{B}) = [\hat{A} + \hat{B}, \hat{C}]$
- (d)  $[\hat{A}, \gamma] = \hat{A}\gamma - \gamma\hat{A} = \gamma\hat{A} - \gamma\hat{A} = 0$

- (e)  $[\hat{A} - \alpha, \hat{B} - \beta] = (\hat{A} - \alpha)(\hat{B} - \beta) - (\hat{B} - \beta)(\hat{A} - \alpha) = \hat{A}\hat{B} - \hat{B}\hat{A} - \alpha(\hat{B} - \hat{B}) - \beta(\hat{A} - \hat{A}) = [\hat{A}, \hat{B}]$
- (f)  $[\hat{A}, f(\hat{A})] = \hat{A}f(\hat{A}) - f(\hat{A})\hat{A}$   
 Since  $f(\hat{A})$  is an analytic function, it can be expressed in the form  $f(\hat{A}) = \sum_{i=0}^{\infty} k_i \hat{A}^i$   
 Then  $\hat{A}f(\hat{A}) - f(\hat{A})\hat{A} = \sum_{i=0}^{\infty} k_i \hat{A}^{i+1} - \sum_{i=0}^{\infty} k_i \hat{A}^{i+1} = 0$   
 Hence  $[\hat{A}, f(\hat{A})] = 0$

## 7 Problem 7

- (a)  $[\hat{H}, \hat{p}]\psi(x) = (\hat{H}\hat{p} - \hat{p}\hat{H})\psi(x)$   
 Since  $\hat{H} = \hat{K} + \hat{V}$ ,  $\hat{V} = 0$ ,  $\hat{K} = \frac{\hat{p}^2}{2m}$   
 Then  $[\hat{H}, \hat{p}]\psi(x) = (\frac{\hat{p}^3}{2m} - \frac{\hat{p}^3}{2m})\psi(x) = 0$

Since  $\hat{V} = 0$ , we can consider  $\hat{H} = f(\hat{p})$  and  $f$  is analytic function, then according to  $[\hat{A}, f(\hat{A})] = 0$ , we know the result is 0.

- (b) For the momentum operator, the eigenfunction is  $\phi_{p_x}(x) = Ce^{\frac{i}{\hbar}p_x x}$   
 For Hamiltonian, since it's a free particle,  $\hat{H} = \hat{K} = \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$   
 $\hat{H}\phi_{p_x}(x) = -\frac{\hbar^2}{2m} \frac{\partial^2 \phi_{p_x}(x)}{\partial x^2} = \frac{C}{2m} p_x^2 \phi_{p_x}(x)$

Hence any eigenfunction of the 1D momentum operator is also an eigenfunction of the Hamiltonian of a free particle.

Conversely, for the eigenfunction of  $\hat{H}$  is  $\psi(x) = A \cos kx + B \sin kx$  where  $k = \sqrt{\frac{2mE}{\hbar}}$   
 $\hat{p}\psi(x) = -i\hbar \frac{\partial(A \cos kx + B \sin kx)}{\partial x} = ik\hbar(A \sin kx - B \cos kx)$

Hence the eigenfunction of the Hamiltonian of a free particle is not always the eigenfunction of the 1D momentum operator

- (c) In this case,  $\hat{H} = \hat{K} + \hat{V} = f(\hat{p}) + \hat{V}$   
 $[\hat{H}, \hat{p}]\psi(x) = [f(\hat{p}) + \hat{V}, \hat{p}]\psi(x) = ([f(\hat{p}), \hat{p}] + [\hat{V}, \hat{p}])\psi(x) = [\hat{V}, \hat{p}]\psi(x)$   
 $[\hat{H}, \hat{p}]\psi(x) = V(x) - i\hbar \frac{\partial \psi(x)}{\partial x} + i\hbar \frac{\partial V(x)\psi(x)}{\partial x} = -i\hbar V(x)\psi(x)' + [i\hbar V(x)\psi(x)' + i\hbar V(x)'\psi(x)] = i\hbar V(x)'\psi(x)$

For two operators  $\hat{A}, \hat{B}$  where one is a function while another takes one order derivative, their commutator will take the derivative of the operator's function instead of the input function, because it's cancelled out.

## 8 Problem 8

According to the general uncertainty principle  $\Delta_E \Delta_x \geq \frac{1}{2} |\langle \Psi, i[\hat{H}, \hat{x}] \Psi \rangle|$   
 $i[\hat{H}, \hat{x}] = i(\hat{H}\hat{x} - \hat{x}\hat{H}) = i\{-i\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)\}x - xV(x) + xi\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} = \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} x - x\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$   
 $i[\hat{H}, \hat{x}] = \frac{\hbar^2}{2m} (\frac{2\partial}{\partial x} + x\frac{\partial^2}{\partial x^2} - x\frac{\partial^2}{\partial x^2}) = \frac{\hbar^2}{m} \frac{\partial}{\partial x} = i\frac{\hbar}{m} (-i\hbar \frac{\partial}{\partial x}) = i\frac{\hbar}{m} \hat{p}$   
 Hence  $\Delta_E \Delta_x \geq \frac{1}{2} |\langle \Psi, i\frac{\hbar}{m} \hat{p} \Psi \rangle| = \frac{\hbar}{2m} \langle \hat{p} \rangle$

## 9 Problem 9

1. No, since the polarization direction of  $A$  and  $B$  are perpendicular to each other, the vertically polarized light can't penetrate through  $B$
2. (i) According to Malus's Law,  $I = I_0 \cos^2 \frac{\pi}{2} \cos^2 \frac{\pi}{2}$ , behind the  $B$  we can get light with intensity  $I = \frac{1}{4} I_0$   
(ii) With the increase of  $\theta$ , the light intensity behind  $C$  is  $I = I_0 \cos^2 \theta \sin^2 \theta$ , which is max when  $\theta = \frac{\pi}{2}$  and minimum when  $\theta = 0$

According to our intuition, we should never get light behind  $B$  because the polarization direction of  $A$  and  $B$  are perpendicular to each other, however, it fails when we put another polarizer  $C$  between then when  $\theta \neq 0$

In this guess, I guess the polarizer  $C$  plays a role like measurement or observation for the polarization of the light when  $\theta \neq 0$ , and it disturb the light's original polarization. Hence, Malus's law is an example showing how observation will cause the collapse of wave function

## Python Code for Problem 1

```
import numpy as np
import matplotlib.pyplot as plt
import math
t0=1
a=1e34
h=1.05e-34
m=9e-31
p0=1e30
t=30
x=np.linspace(-100,100,10000)
b=a*h*(1+t**2/(t0**2))**0.5
psi=math.e**(-(x-p0*m*t)**2/(b**2))/(b*math.pi**0.5)
title='The probability function at t=' + str(t)
plt.title(title)
plt.grid(1)
plt.plot(x,psi)
plt.show()
```