VE 320 Summer 2019

Introduction to Semiconductor Devices

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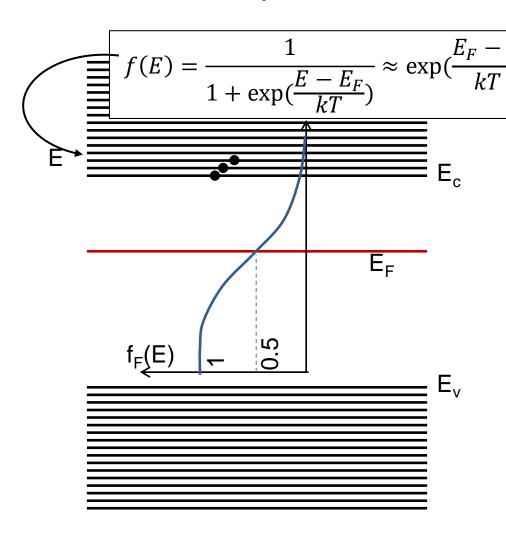
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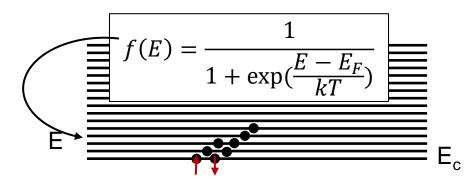
Lecture 4

The semiconductor in equilibrium (Chapter 4)



Non-degenerate Semiconductors:

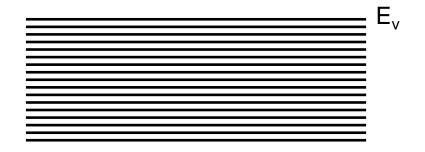
- Small number of impurity atoms
- Impurities introduce discrete, noninteracting donor energy states in n-type semiconductor
- Boltzmann approximation





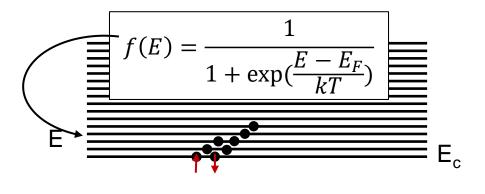
- High impurity concentration
- Donor electrons interact
- Split into a band of energies
- When n₀>N_c, then E_F>E_c:
 Degenerate n-type semiconductor

$$n_0 = N_c \exp\left[\frac{-(E_c - E_F)}{kT}\right]$$

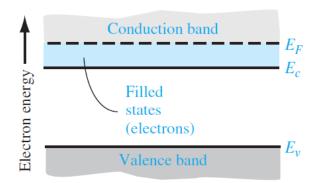


- When $p_0 > N_v$, the $E_F < E_v$:

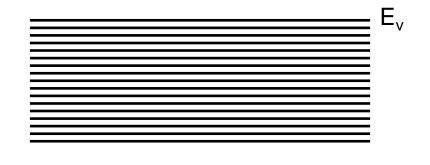
 Degenerate p-type semiconductor
- Fermi distribution, cannot use
 Boltzmann approximation

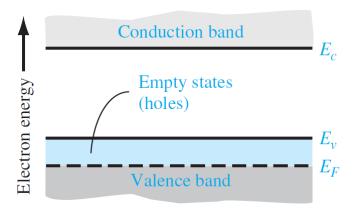


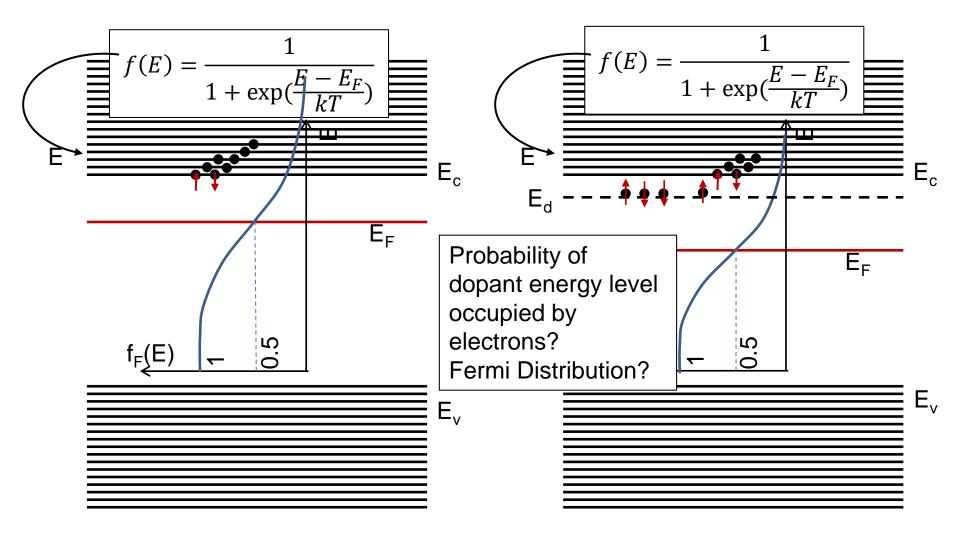
Degenerate n-type semiconductor: the states between $E_{\rm F}$ and $E_{\rm c}$ are mostly filled with electrons

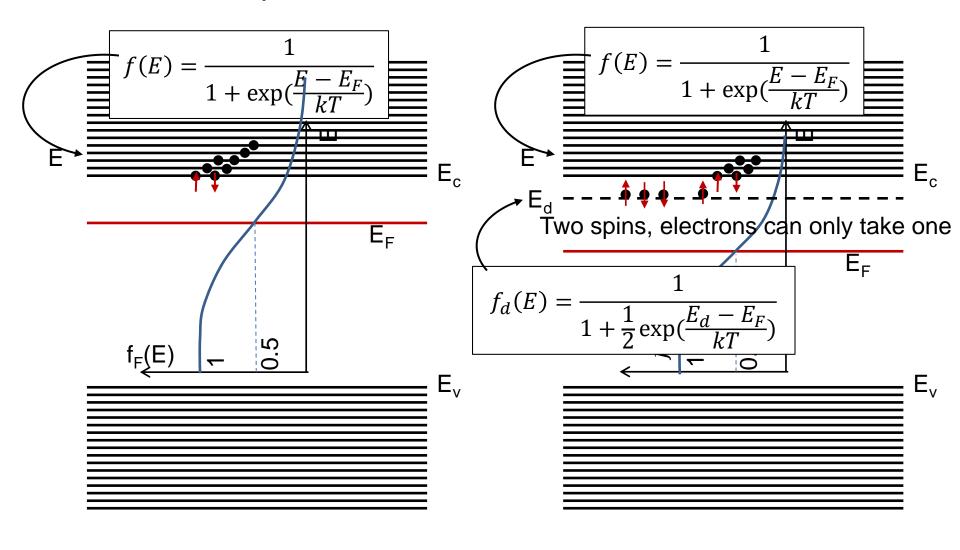


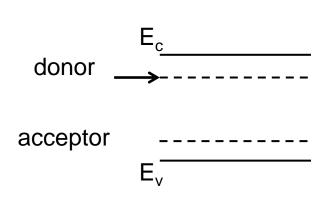
Degenerate p-type semiconductor: states between E_F and E_v are mostly empty











1 Probability of electrons occupying donor energy level

$$f_d(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_d - E_F}{kT})}$$

The density of electrons occupying the donor level:

$$1/2 = 1/g$$
, g is degeneracy factor

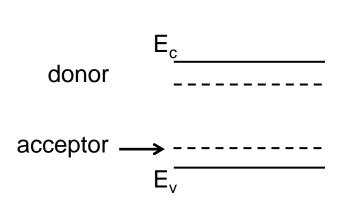
$$n_d = \frac{N_d}{1 + \frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)}$$

 $E_{\rm d}$: energy of the donor level

$$n_d = N_d - N_d^+$$

 N_{d}^{+} is the concentration of ionized donors--Important





1 Probability of electrons occupying acceptor energy level

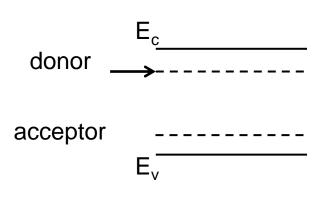
$$f_a(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_F - E_a}{kT})}$$

The density of holes occupying the acceptor level:

$$p_a = \frac{N_a}{1 + \frac{1}{g} \exp\left(\frac{E_F - E_a}{kT}\right)} = N_a - N_a^-$$

 $N_{\rm a}^{-}$ is the concentration of ionized acceptors--Important

The ground state degeneracy factor *g* is normally taken as 4 for the acceptor level in silicon and gallium arsenide because of the detailed band structure



1 Probability of electrons occupying dopant energy level

$$f_d(E) = \frac{1}{1 + \frac{1}{2} \exp(\frac{E_d - E_F}{kT})}$$

(2) Ionization rate

$$1 - f_d(E) = \frac{1}{1 + 2\exp(\frac{E_F - E_d}{kT})}$$

$$= \begin{cases} 1/3 & E_d = E_F \\ 1 & E_d \gg E_F + kT \\ 0 & E_d \ll E_F \end{cases}$$

Completer ionization

$$\begin{array}{c} & E_{c} \\ \hline \text{donor} & \longrightarrow \\ \\ \text{acceptor} & \longrightarrow \\ \hline E_{v} \\ \end{array}$$

Ratio of electrons that remains in dopant level to all electrons

$$(E_d - E_F) \gg kT$$

$$n_d \approx \frac{N_d}{\frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)} = 2N_d \exp\left[\frac{-(E_d - E_F)}{kT}\right]$$

Electrons in the conduction band:

$$n_0 = N_c \exp\left[\frac{-(E_c - E_F)}{kT}\right]$$

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$$\frac{n_d}{n_d+n_0} = \frac{2N_d \exp\left[\frac{-(E_d-E_F)}{kT}\right]}{2N_d \exp\left[\frac{-(E_d-E_F)}{kT}\right] + N_c \exp\left[\frac{-(E_c-E_F)}{kT}\right]} \quad \text{The donor states are essentially completely ionized for a typical do of 10$^{16} cm$^{-3}, almost all donor impurity atoms had donor donated an electron to$$

 E_c - E_d : ionization energy of the donor electrons

ionized for a typical doping of 10¹⁶ cm⁻³, almost all donor impurity atoms have donated an electron to the conduction band.



donor acceptor E_v E_c - E_d : ionization energy of the donor electrons

Ionization energy: the approximate energy required to elevate the donor electron into the conduction band

Bohr model:

Coulomb force = centripetal force

$$\frac{e^2}{4\pi \epsilon r_n^2} = \frac{m^* v^2}{r_n}$$

Angular momentum is also quantized $m^* r_n v = n \hbar$

$$m^* r_n v = n \hbar$$

Radius of the orbit is quantized

$$r_n = \frac{n^2 \hbar^2 4\pi \epsilon}{m^* e^2}$$

Bohr radius

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_0 e^2} = 0.53 \text{ Å} \qquad \frac{r_n}{a_0} = n^2 \epsilon_r \left(\frac{m_0}{m^*}\right)$$

*m** is the conductivity effective mass



acceptor

donor _____

 $E_{\rm c}$ - $E_{\rm d}$: ionization energy of the donor electrons

Total energy
$$E = T + V$$

T: kinetic energy, V: potential energy

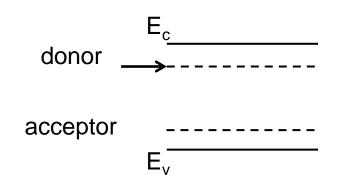
$$T = \frac{1}{2} m^* v^2 \qquad T = \frac{m^* e^4}{2(n\hbar)^2 (4\pi\epsilon)^2}$$

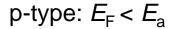
$$V = \frac{-e^2}{4\pi\epsilon r_n} = \frac{-m^* e^4}{(n\hbar)^2 (4\pi\epsilon)^2}$$

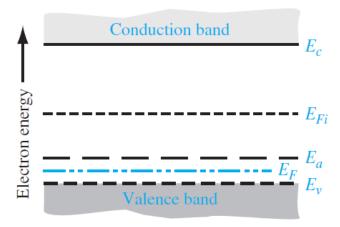
$$E = T + V = \frac{-m^* e^4}{2(n\hbar)^2 (4\pi\epsilon)^2}$$

Hydrogen atom in the lowest energy state: $E_{\rm H} = -13.6 \text{ eV}$

Silicon ionization energy: is $E_{\rm Si}=E_{\rm H}\times m^*/m_0/\epsilon_{\rm r}^2=$ -25.8 meV << $E_{\rm g}=1.12{\rm eV}$







Freeze-out:

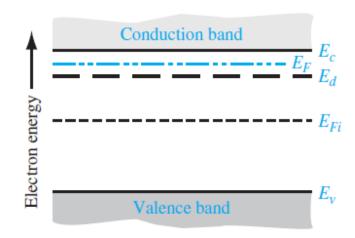
$$n_d = \frac{N_d}{1 + \frac{1}{2} \exp\left(\frac{E_d - E_F}{kT}\right)}$$

T=0K

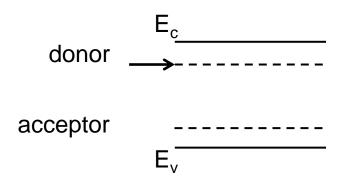
n-type semiconductor, each donor state must contain an electron, $n_d = N_d$ or $N_d^+ = 0$

$$\exp\left[(E_d - E_F)/kT\right] = 0$$

$$E_F > E_d$$



No electrons from the donor state are thermally elevated into the conduction band



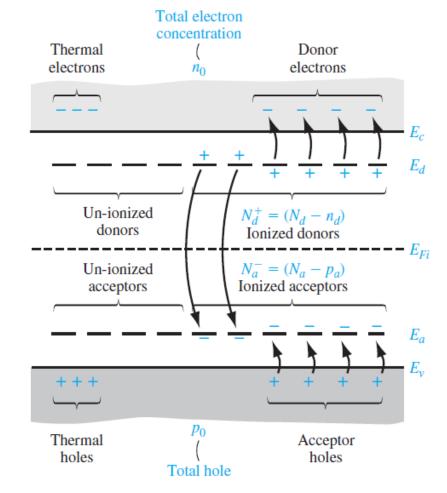
Compensated semiconductor: contains both donor and acceptor impurity atoms in the same region

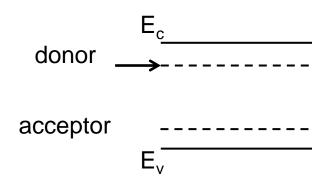
n-type compensated semiconductor: $N_d > N_a$ p-type compensated semiconductor: $N_a > N_d$ $N_a = N_d$: completely compensated semiconductor, like intrinsic.

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$
 See the

Compensated semiconductors at thermal equilibrium





Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Electron concentration in thermal equilibrium

Complete ionization: $n_d=0$, $p_a=0$

$$n_0 + N_a = N_d + p_0$$

$$n_0 + N_a = \frac{n_i^2}{n_0} + N_d$$

$$n_0^2 - (N_d - N_a)n_0 - n_i^2 = 0$$

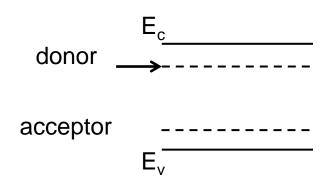
$$n_0 = \frac{(N_d - N_a)}{2} + \sqrt{\left(\frac{N_d - N_a}{2}\right)^2 + n_i^2}$$

Can calculate the electron concentration in an n-type semiconductor, $N_d > N_a$

$$p_0 = \frac{n_i^2}{n_0}$$

 $N_{\rm a}$ can be 0: only n doping





Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

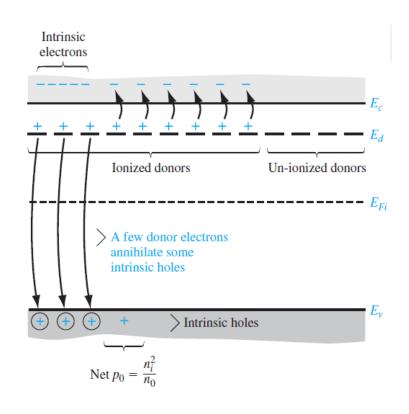
 $n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Electron concentration in thermal equilibrium

Complete ionization: n_d =0, p_a =0

$$n_0 > n_i$$
, $p_0 < n_i$



Electron concentration in thermal equilibrium

donor E_c acceptor E_v

Compensated semiconductors

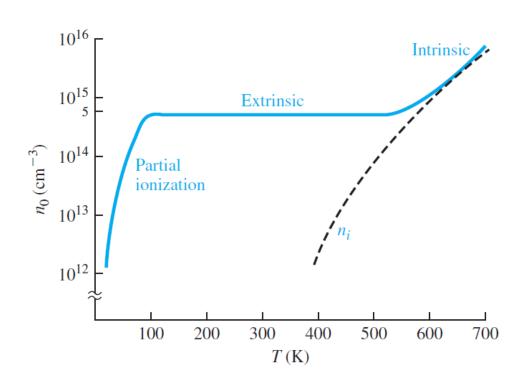
(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

High temperature: like intrinsic! – Usually bad for semiconductor devices



 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Fermi level position:

Special case: $N_a=0$, Only n-type doping

Boltzmann approximation

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



 $n_i >> N_d^+ \Rightarrow T \text{ very high}$

$$n_0 = p_0 = n_i = \sqrt{N_c N_v} \exp(-\frac{E_g}{2kT})$$

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Fermi level position:

Special case: N_a =0, Only n-type doping

Boltzmann approximation

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

$$E_c - E_F = kT \ln \left(\frac{N_c}{n_0}\right) = kT \ln \left(\frac{N_c}{N_d}\right)$$

 $N_{d}\uparrow$, E_{c} - $E_{F}\downarrow$, Fermi level closer to the conduction band

Fermi level position:

Special case: N_a =0, Only n-type doping

Boltzmann approximation

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

$$E_c - E_F = kT \ln \left(\frac{N_c}{n_0}\right) = kT \ln \left(\frac{N_c}{N_d}\right)$$

Compensated semiconductor:

$$N_{\rm d} \rightarrow N_{\rm d}$$
- $N_{\rm a}$

Fermi level position:

Special case: $N_a=0$, Only n-type doping

Boltzmann approximation

$$n_0 = N_d^+ + p_0$$

$$n_0 p_0 = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



 $n_i \leq N_d^+ \Rightarrow T \text{ not very high}$

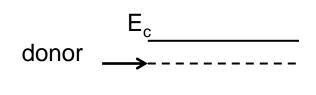
$$n_0 \approx N_d^+$$

(1) Assuming complete ionization:

$$n_0 \approx N_d$$

Also,
$$n_0 = n_i \exp [(E_F - E_{Fi})/kT]$$

$$E_F - E_{Fi} = kT \ln \left(\frac{n_0}{n_i}\right)$$



Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Fermi level position:

Special case: $N_d=0$, Only p-type doping

$$n_i << N_a^+ \Rightarrow T \text{ not very high}$$

$$p_0 \approx N_a^+$$

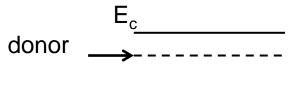
(1) Assuming complete ionization:

$$p_0 \approx N_a$$

Boltzmann approximation

$$p_0 = N_v \exp \left[-(E_F - E_v)/kT \right]$$
$$E_F - E_v = kT \ln \left(\frac{N_v}{N_a} \right)$$

$$E_{Fi} - E_F = kT \ln \left(\frac{p_0}{n_i}\right)$$



Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

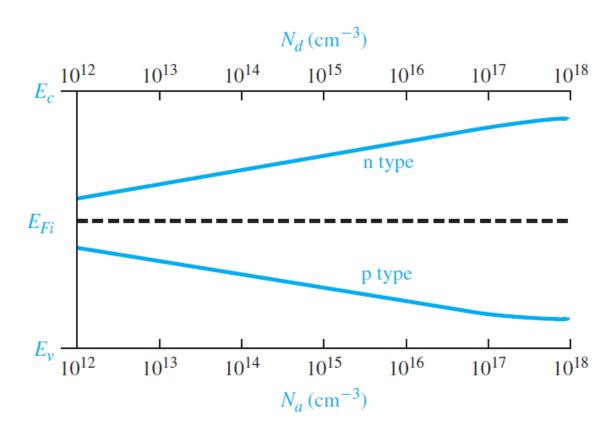
 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

Fermi level position:

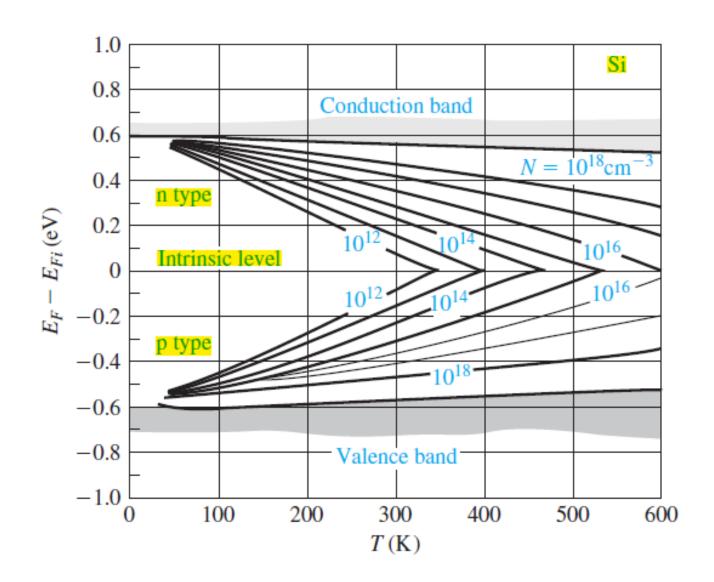
T not very high

Assuming complete ionization:

Boltzmann approximation



Fermi level position:



Fermi level position:

Special case: N_a =0, Only n-type doping

Boltzmann approximation

$$\begin{array}{c} & E_c \\ \hline \rightarrow & \end{array}$$

Compensated semiconductors

(3) Charge neutrality

$$n_0 + N_a^- = N_d^+ + p_0$$

$$n_0 = N_d^+ - N_a^- + p_0$$

$$n_0 + (N_a - p_a) = p_0 + (N_d - n_d)$$

 n_d : concentration of electrons in the donor states p_a : concentration of holes in the acceptor states

$$n_0 = N_d^+ + p_0 \qquad \qquad n_0 p_0^- = n_i^2$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)$$



$$n_i \leq N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 \approx N_d^+$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_F - E_d}{kT})}$$

$$n_i << N_d^+ \Rightarrow T \ not \ very \ high$$

$$n_0 = N_d^+$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_C}{kT}\right) = \frac{n_0}{N_C}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right)) = \underbrace{\begin{pmatrix} N_d \\ 1 + 2\exp\left(\frac{E_F - E_d}{kT}\right) \end{pmatrix}}_{N_d}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\frac{N_d}{1 + 2\exp(\frac{E_F - E_d}{kT})} = \frac{N_d}{1 + 2\exp(\frac{E_c - E_d}{kT})\exp(\frac{E_F - E_c}{kT})}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_C}{kT}\right) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c} \qquad \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)} = \frac{N_d}{1 + 2\exp\left(\frac{E_c - E_d}{kT}\right)\exp\left(\frac{E_F - E_c}{kT}\right)}$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_A}{kT})\frac{n_0}{N_c}}$$

$$n_i << N_d^+ \Rightarrow T \text{ not very high}$$

$$n_0 = N_d^+$$

$$(n_0 = N_c \exp\left(\frac{E_F - E_C}{kT}\right) = \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c}$$

$$\exp\left(\frac{E_F - E_c}{kT}\right) = \frac{n_0}{N_c} \qquad \frac{N_d}{1 + 2\exp\left(\frac{E_F - E_d}{kT}\right)} = \frac{N_d}{1 + 2\exp\left(\frac{E_c - E_d}{kT}\right)} \exp\left(\frac{E_c - E_d}{kT}\right) \exp\left(\frac{E_F - E_c}{kT}\right)$$

$$2\exp\left(\frac{E_A}{kT}\right)n_0^2 + N_c n_0 - N_d N_c = 0$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) = N_d^+ = \frac{N_d}{1 + 2\exp(\frac{E_A}{kT})\frac{n_0}{N_c}}$$

$$2\exp\left(\frac{E_A}{kT}\right)n_0^2 + N_c n_0 - N_d N_c = 0$$

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})}$$

$$n_0 = N_c \exp\left(\frac{E_F - E_c}{kT}\right) \Rightarrow \exp\left(\frac{E_F - E_c}{kT}\right) = \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})}$$

$$E_F = E_c + kT ln(\frac{\sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})} - 1}{4\exp(\frac{E_A}{kT})})$$

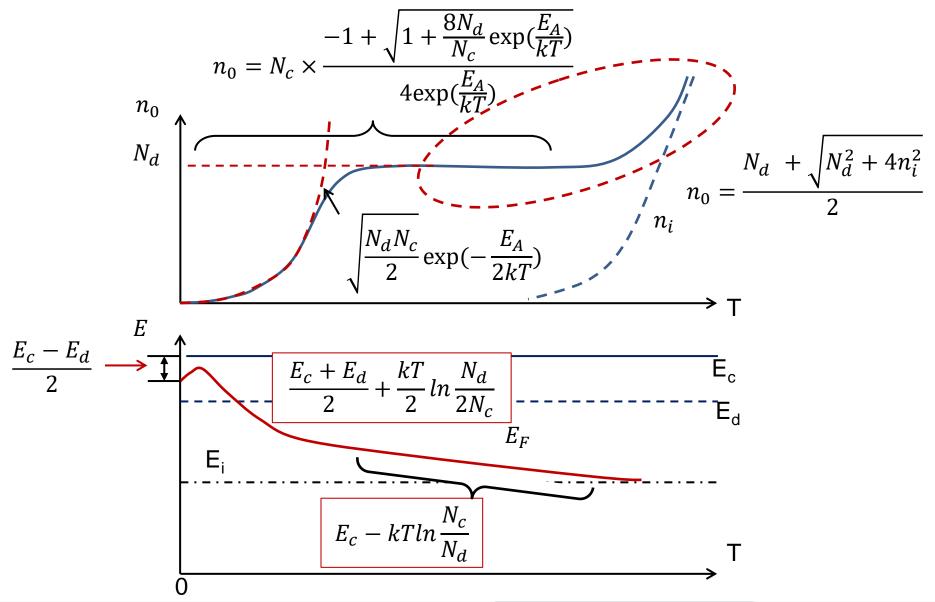
(2) Incomplete ionization, Boltzmann approximation, only n doping:

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})} = \begin{cases} \sqrt{\frac{N_d N_c}{2}} \exp(-\frac{E_A}{2kT}) & T \text{ small} \\ N_d & T \text{ big} \end{cases}$$

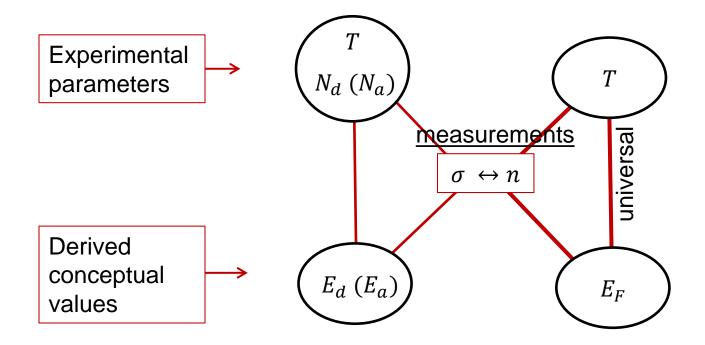
$$E_F = E_c + kT ln(\frac{\sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})} - 1}{4\exp(\frac{E_A}{kT})}) = \begin{cases} \frac{E_c + E_d}{2} + \frac{kT}{2} ln \frac{N_d}{2N_c} & T \text{ small} \\ E_c - kT ln \frac{N_c}{N_d} & T \text{ big} \end{cases}$$

$$n_0 = N_c \times \frac{-1 + \sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})}}{4\exp(\frac{E_A}{kT})} = \begin{cases} \sqrt{\frac{N_d N_c}{2}} \exp(-\frac{E_A}{2kT}) & T \text{ small} \\ N_d & T \text{ big} \end{cases}$$

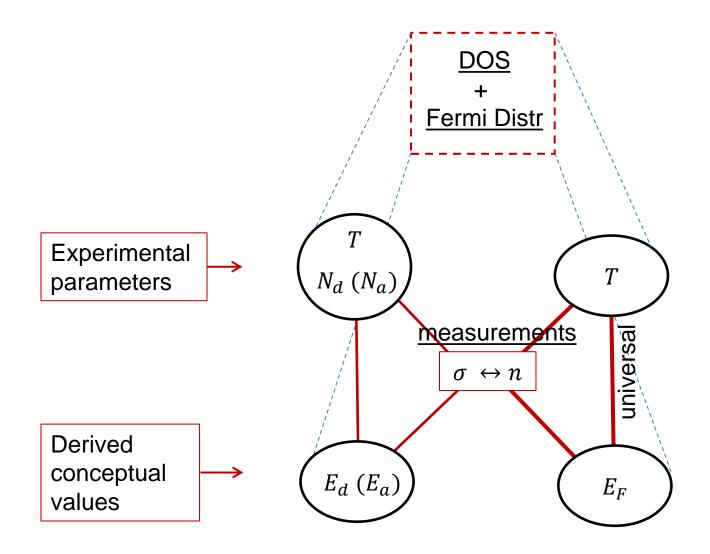
$$E_F = E_c + kT ln(\frac{\sqrt{1 + \frac{8N_d}{N_c} \exp(\frac{E_A}{kT})} - 1}{4 \exp(\frac{E_A}{kT})}) = \begin{cases} \frac{E_c + E_d}{2} + \frac{kT}{2} ln \frac{N_d}{2N_c}, & T \text{ small} \\ E_c - kT ln \frac{N_c}{N_d}, & T \text{ big} \end{cases}$$



Summary of Lecture 4



Summary of Lecture 4



Summary of Lecture 4

Why is E_F useful?

In thermal equilibrium, the Fermi energy level is a constant throughout a system

