RADIAL VELOCITIES FOR 889 LATE-TYPE STARS¹

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ABSTRACT

We report radial velocities for 844 FGKM–type main-sequence and subgiant stars and 45 K giants, most of which had either low-precision velocity measurements or none at all. These velocities differ from the standard stars of Udry et al. by 0.035 km s⁻¹ (rms) for the 26 FGK standard stars in common. The zero point of our velocities differs from that of Udry et al.: $\langle V_{\text{Present}} - V_{\text{Udry}} \rangle = +0.053$ km s⁻¹. Thus, these new velocities agree with the best known standard stars both in precision and zero point, to well within 0.1 km s⁻¹.

Nonetheless, both these velocities and the standards suffer from three sources of systematic error, namely, convective blueshift, gravitational redshift, and spectral type mismatch of the reference spectrum. These systematic errors are here forced to be zero for $G2\ V$ stars by using the Sun as reference, with Vesta and day sky as proxies. But for spectral types departing from solar, the systematic errors reach $0.3\ km\ s^{-1}$ in the F and K stars and $0.4\ km\ s^{-1}$ in M dwarfs.

Multiple spectra were obtained for all 889 stars during 4 years, and 782 of them exhibit velocity scatter less than 0.1 km s⁻¹. These stars may serve as radial velocity standards if they remain constant in velocity. We found 11 new spectroscopic binaries and report orbital parameters for them.

Subject headings: binaries: spectroscopic — catalogs — stars: fundamental parameters — stars: kinematics — stars: late-type — techniques: radial velocities — techniques: spectroscopic

On-line material: machine-readable tables

1. INTRODUCTION

The radial velocity of a star is ideally the component of the velocity vector of its center of mass that lies along the line of sight. Radial velocities are valuable for a variety of astrophysical investigations, including studies of the structure of the Milky Way Galaxy, the orbits of long-period binary stars, and the distances to star clusters (see, e.g., Binney & Merrifield 1998). "Barycentric" radial velocities (sometimes referred to as "absolute" radial velocities), such as reported here, are measured relative to the barycenter, or center of mass, of the solar system. Such velocities are often (incorrectly) termed "heliocentric," though the Sun moves with a speed of $\sim\!13~{\rm m~s^{-1}}$ relative to the barycenter.

Radial velocities of stars in the Galaxy are often measured with an accuracy of only ~ 0.5 km s⁻¹. With advances in the accuracy of proper motion measurements to ~ 1 mas yr⁻¹ (e.g., Perryman et al. 1996) for many stars, a corresponding increase in the accuracy of radial velocities is required. Meanwhile, the best *relative* radial velocities have precisions of 3 m s⁻¹ (Butler et al. 2000) and new instruments (e.g., HARPS) are designed to achieve a precision of 1 m s⁻¹ (Bouchy et al. 2001; Pepe et al. 2000). These relative velocities have proved useful in the detection of extrasolar

The precision-velocity technology has been applied to the establishment of barycentric radial velocities, most notably by the Geneva team (Udry, Mayor, & Queloz 1999a; Udry et al. 1999b). They have measured velocities for 38 stable dwarf stars with precision better than 0.05 km s⁻¹.

Here we provide barycentric radial velocities with typical accuracies of $0.3~\rm km~s^{-1}$ (and precise to $0.03~\rm km~s^{-1}$ for a given spectral type) for 889 stars. Our intent is to provide a velocity measurement at the current epoch for a variety of purposes. Velocity variations with a timescale of hundreds of years may be detected by comparison of present and future velocities. We also hope to establish radial velocity standard stars, by identifying a subset that exhibit no significant velocity variation above $10~\rm m~s^{-1}$.

2. BARYCENTRIC RADIAL VELOCITIES

The Doppler searches for planets have been successful because of the relative ease with which the change in radial velocity may be measured with high internal precision. Such relative velocities circumvent the technical challenges associated with the determination of an accurate velocity zero point, and they avoid the physical interpretation of a barycentric Doppler shift which becomes bewildering at levels below 1 km s⁻¹.

Barycentric Doppler shifts carry an unclear interpretation for several reasons. Stellar lines suffer a gravitational redshift (Misner, Thorne, & Wheeler 1973) upon leaving the

planets (e.g., Marcy, Cochran, & Mayor 2000) but they are not necessarily tied to a velocity zero point. Nonetheless, the precision-velocity instruments have overcome many observational and technical hurdles related to spectroscopic Doppler-shift measurements, either by using a gas absorption cell or a fiber-fed comparison lamp spectrum (Valenti, Butler, & Marcy 1995; Butler et al. 1996; Baranne 1999).

¹ Based on observations obtained at the W. M. Keck Observatory, which is operated jointly by the University of California and the California Institute of Technology, and on observations obtained at the Lick Observatory, which is operated by the University of California.

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stellar photosphere, yielding effective redshifts of $V_{\rm grav} = GM/Rc$ (Dravins et al. 1999). This redshift varies from 680 m s⁻¹ for F5 V to 500 m s⁻¹ for M5 V stars, the range of spectral types considered here. Uncertainties in stellar masses and radii prevent an accurate removal of this effect. Furthermore, stellar lines suffer a transverse Doppler effect (essentially time dilation), the removal of which requires knowledge of the full velocity vector of the star's space motion. This effect is $\simeq 50$ m s⁻¹ for the fastest moving stars (Lindegren, Dravins, & Madsen 1999).

More importantly, stellar Doppler shifts are affected by subphotospheric convection (granulation), macroturbulence, stellar rotation, pressure shifts, oscillations, and activity cycles. The most important of these effects is granulation. The textbook explanation is that a larger contribution to the stellar flux emerges from hot, rising gas than from falling gas in the convective cells. These motions yield an overall blueshift of spectral lines (Dravins 1999). However, the exact convective blueshift depends on the full three-dimensional hydrodynamics and radiative transfer in each spectral line as a function of depth in the photosphere (Dravins 1999). The blueshift clearly depends on spectral type and is expected to be $\sim -1000 \text{ m s}^{-1}$ for F5 V, -400 m s^{-1} for G2 V, and -200 m s^{-1} for K0 V (Dravins et al. 1999). Effects due to pressure shifts are less than 100 m s^{-1} for ordinary stars (Dravins et al. 1999; Allende Prieto et al. 1997). Stellar rotation also imposes minor radial velocity effects (Gray 1999).

Moreover, the measurement of spectroscopic barycentric radial velocities usually requires a reference stellar spectrum. Due to constraints of telescope time and available standard stars, only a small number of reference spectra can be used. Typically, the Doppler measurements require use of reference spectra having different spectral types than the program star, which leads to spectral mismatch errors. Since the strengths of spectral lines vary with temperature and metallicity, the spectra of the reference and program stars will be significantly different. In effect, the relative displacement in wavelength between two nonidentical spectra is not uniquely defined and therefore is dependent on the algorithm used. Such spurious Doppler effects are minimized, but not eliminated, by using high-resolution spectra with many lines resolved, as adopted in this present survey.

The effects of convection, stellar gravitational redshift, transverse Doppler shift, and the other effects mentioned above cannot be determined for a star with an accuracy that is comparable to the internal measurement errors of ~ 10 m s⁻¹. Removal of such effects by a model of each star might introduce more model-dependent errors. Nonetheless, with considerable modeling, the measurement of a Doppler shift may be used to determine the "true" velocity component of the center of mass of the star. Alternatively (and more traditionally), Doppler measurements may be left merely as an observable, namely the spectroscopic shift in wavelength, often quoted as $z = \Delta \lambda/\lambda_0$ (Lindegren et al. 1999).

Here we adopt the philosophy that spectroscopic bary-centric radial velocities should first be corrected for local effects, such as that caused by the observer's motion relative to the solar system barycenter (\sim 30 km s⁻¹) and by the solar gravitational blueshift (\sim 3 m s⁻¹). In addition, we will correct all of our velocities of FGK stars for gravitational redshift and convective blueshift to first order by using the known radial velocity of the Sun to set the zero point for the stellar velocity measurements.

Our quoted Doppler shifts represent velocities as if measured at the solar system barycenter but with the Sun and its potential well removed. Clearly, the velocity measurements presented here are amenable to future corrections for the spectral-type dependence relative to G2 V for gravitational redshift and convective blueshift, in order to yield the most accurate velocities possible for them (e.g., Gullberg 1999; Saar & Fischer 2000).

3. SPECTROSCOPIC OBSERVATIONS

The spectra were obtained with the HIRES echelle spectrometer (Vogt et al. 1994) on the 10 m Keck I telescope and with the "Hamilton" echelle spectrometer fed by either the 3 m Shane or the 0.6 m Coude Auxilliary (CAT) Telescopes (Vogt 1987). During an observation the starlight is sent through a glass cell that is filled with iodine vapor (Marcy & Butler 1992) before entering the spectrometer, which superimposes iodine lines on the stellar lines. These iodine absorption lines are used to calibrate the wavelength scale of the spectrum from 5000 to 6000 Å.

Our Doppler planet search project contains 889 stars at the Keck and Lick Observatories (Butler et al. 2000). Until now only relative radial velocities have been computed from these spectra, and they have a precision of $\sim 3 \text{ m s}^{-1}$ (Butler et al. 1996; Vogt et al. 2000), which has allowed the discovery of Jovian and sub-Jovian sized extrasolar planets (Marcy et al. 2000).

To achieve barycentric velocities, we adopt two standard spectra. We use the National Solar Observatory (NSO) FTS solar spectrum (Kurucz et al. 1984) and an M dwarf composite spectrum (see § 4) as reference spectra. The 889 stars each typically have ~12 spectra obtained during 4 years from 1997 to 2001. The distribution of spectral types in our sample is: 14% F, 46% G, 27% K, and 13% M stars. Except for 45 K giants all stars are main-sequence dwarfs or subgiants. All stars are void of a visible companion within 2", though some were subsequently revealed to be spectroscopic binaries (see § 9).

4. DOPPLER METHOD

The barycentric radial velocities reported here are found in a manner similar to that used to find the relative radial velocities for the planet search. An observed spectrum is fitted with a synthetic spectrum that is composed of the individual stellar and iodine spectra. In detail, the synthetic spectrum is the product of the deconvolved stellar "template" spectrum (with the spectrometer instrumental profile removed) with a high-resolution spectrum of molecular iodine. The product of these two is convolved with the instrumental profile of the spectrometer (at the time of the observation), to produce the final synthetic spectrum, as described in Butler et al. (2000).

The observed spectrum to be synthesized is broken into "chunks" of length 40 pixels corresponding to roughly \sim 2 Å. In total 14 free parameters are fitted in the Doppler analysis; 11 devoted to the instrumental profile, along with the wavelength zero point of each chunk, the wavelength dispersion across each chunk, and the Doppler shift of the stellar spectrum relative to the stellar template of that star, $z = (\Delta \lambda/\lambda)$. All of the parameters, except for z, are extracted primarily from the iodine portion of the observed spectrum. After the best parameters are found for all the chunks of a

spectrum they are saved for further analysis, notably the weighted average of z. We apply a correction to all velocities for our topocentric motion relative to the barycenter (McCarthy 1995). For a more in-depth discussion of this standard Doppler analysis to obtain relative velocities, see Marcy & Butler (1992), Butler et al. (1996), and Valenti et al. (1995).

To obtain barycentric radial velocities, the approach was similar to the standard analysis described above and indeed we used some parameters derived from that analysis. Here we used the National Solar Observatory (NSO) solar spectrum (Kurucz et al. 1984) as the deconvolved template, for all F, G, and K stars. For M dwarfs, we constructed an inhouse M star composite spectrum as the template (see below). Since all parameters except for z are extracted from the iodine portion of the spectrum previously, the use of a different stellar template spectrum does not affect these parameters significantly. Therefore, in our new fit, we simply adopt the values of the 13 non-z parameters that were previously obtained in the Doppler analysis for relative velocities (i.e., for the planet search) and we fit here only for the barycentric radial velocity, z.

Ordinarily, when the deconvolved stellar spectrum of the individual star is used as the template, the χ^2_{ν} statistic is near unity (usually less than 1.7), because the stellar contribution to the model is a previously obtained deconvolved spectrum of the same star. But in our fit for barycentric velocities, the value of reduced χ^2_{ν} is much larger (poorer) because of the mismatch between the template spectrum (NSO or M dwarf composite) and that of the program star. The program star differs from the solar or M-dwarf template in spectral type, metallicity, and $v \sin i$. Thus, the spectral fits are not as good, yielding χ^2_{ν} of 3–7.

The internal error per observation, as defined by the weighted standard deviation of the mean of the velocities from all (\sim 400) chunks, is higher for our barycentric radial velocities than for the relative radial velocities used in the planet search. We find that while the relative velocities carry errors of \sim 3 m s⁻¹ (for the planet search) the average internal error per observation for the barycentric velocities reported here is \sim 20 m s⁻¹. A more conservative estimate of our velocity precision is given in § 6 as 0.03 km s⁻¹ from comparison with the standard stars.

Due to the fact that only one stellar template is used for a large range of stellar types, we expect that our errors (internal error per observation) will depend on B-V. The greater the difference between template and program star, the greater the expected error due to spectral mismatch. We plot this internal error in our barycentric velocities versus B-V in Figure 1. Since the NSO solar spectrum, B-V=0.64 (Carroll & Ostlie 1996), is used for F, G, and K stars we expect the errors to be minimized for solar type stars. This is confirmed in Figure 1. The M dwarf composite template spectrum was created from five different M dwarfs having average spectral type of M3 (see below). The internal velocity errors for stars analyzed with this stellar template are minimized at $B-V\approx 1.5$, as can also be seen in Figure 1.

Each program star has an average of 12 observations. As many observations as possible are analyzed per star, up to 30, in order to minimize the uncertainty in the mean of the barycentric radial velocities. If the radial velocity of the star were stable we should obtain an uncertainty in the mean of $\sim 20/\sqrt{12} \approx 6 \text{ m s}^{-1}$. A majority of stars indeed exhibit velocity scatter of $\sim 10 \text{ m s}^{-1}$ (rms), while others have an rms

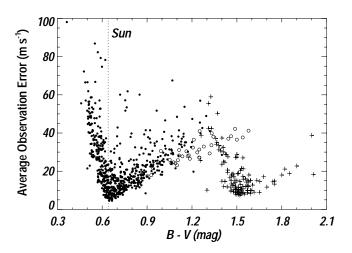


Fig. 1.—Internal velocity error per observation (averaged per star) vs. B-V. The dots represent the stars for which the NSO solar spectrum was used as the reference (filled dots represent dwarfs and open dots represent giants), and the pluses represent the stars for which the M star composite spectrum was used as the reference. The B-V for the Sun (B-V=0.64) is shown for clarity. The spectral type dependence of errors is apparent.

scatter larger than $100~{\rm m~s^{-1}}$. Scatter in the latter stars is almost always caused by companions or large chromospheric activity. Overall, 782 of the 889 stars have an rms velocity scatter less than $0.1~{\rm km~s^{-1}}$, and $107~{\rm stars}$ have an rms velocity scatter larger than $0.1~{\rm km~s^{-1}}$. The barycentric radial velocities for these stars are reported in Tables 1 and 2, respectively.

As mentioned above, the NSO solar spectrum could not be used as the template for the M dwarf program stars because the spectra are too different. Therefore, a separate reference spectrum was required for the M stars. For this purpose a composite spectrum of five M stars was produced in the following manner.

Five M stars were selected having barycentric radial velocities reported by Marcy, Lindsay, & Wilson (1987) with low uncertainties: GJ 251 (M4), GJ 411 (M2), GJ 526 (M4), GJ 752A (M3.5), and GJ 908 (M2). One spectrum of high S/N ratio was used from each of these M dwarfs. The spectra were shifted back, to remove the Doppler shifts caused by the barycentric motion of the observatory and by the barycentric radial velocity of the star itself relative to the barycenter, taken from Marcy et al. (1987). To check for any residual Doppler shifts, these corrected spectra were then cross-correlated with respect to one of them, GJ 251. Assuming the resulting displacements were due to random errors, the mean of the residual velocities was taken to be the barycentric velocity zero point. Using this reference point the spectra were again corrected for their Doppler shifts. Finally, the spectra were put on the same wavelength scale and co-added to create a M star composite spectrum.

5. VELOCITY ZERO POINT

Using observations of the day sky and the minor planet Vesta we found the zero point of our velocities for FGK stars. These references were used because they have essentially solar spectra, and the radial velocities of Vesta and the Sun relative to a topocentric observer are easily determined. We used the on-line "JPL Ephemeris Generator" to find

 $\label{eq:table 1} TABLE \ 1$ Radial Velocities of Stable Stars a

						Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$
Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$	Name	Name	(NSO/M)	(-2440000)	(days)	$(m s^{-1})$
Name	Name	(NSO/M)	(-2440000)	(days)	$(m s^{-1})$	(1)	(2)	(3)	(4)	(5)	(6)
(1)	(2)	(3)	(4)	(5)	(6)		HID 10220	NIGO	10050	245	
HD 166	HID 544	NICO	0702	2662	(527	HD 13531 HD 13579	HIP 10339 HIP 10531	NSO NSO	10979 11339	245 293	7203 -12723
HD 166 HD 283	HIP 544 HIP 616	NSO NSO	8783 11002	2662 1185	-6537 -43102	HD 13612B	HIP 10303	NSO	11076	1184	-5337
HD 377	HIP 682	NSO	11417	208	1184	HD 13825	HIP 10505	NSO	11270	432	-2237
HD 400	HIP 699	NSO	11137	507	-15141	HD 13931	HIP 10626	NSO	11265	748	30538
HD 531	BD +07 9	NSO	11419	208	13460	HD 14412	HIP 10798	NSO	11006	1391	7383
HD 1326A	GJ 15A	M	9928	4656	11814	HD 15176	HIP 11432	NSO	11494	145	-41783
HD 1326B	GJ 15B	M	11294	509	10976	HD 15335	HIP 11548	NSO	11031	1219	41202
HD 1388	HIP 1444	NSO	11016	1338	28498	HD 16141	HIP 12048	NSO	11067	1008	-50971
HD 1461	HIP 1499	NSO	11083	1184	-10166	HD 16160	HIP 12114	NSO	9351	3608	25766
HD 1832	HIP 1813	NSO	11250	950	-30550	HD 16397 HD 16623	HIP 12306 HIP 12364	NSO NSO	11060 11088	1082 829	-99660 17502
HD 1835	HIP 1803	NSO	9123	3026	-2405	HD 16895	HIP 12777	NSO	8930	2937	24453
HD 2025 HD 2774	HIP 1936 HIP 2497	NSO NSO	11218 11478	1389 117	3241 -51566	HD 17190	HIP 12926	NSO	11146	1186	14138
HD 3074	HIP 2663	NSO	10955	1390	29539	HD 17230	HIP 12929	NSO	11443	386	11061
HD 3079	HIP 2712	NSO	11193	390	-12347	HD 17332	HIP 13027	NSO	11481	140	4427
HD 3651	HIP 3093	NSO	9708	4708	-32961	HD 17660	HIP 13258	NSO	11396	324	-28904
HD 3674	HIP 3119	NSO	11161	514	166	HD 17925	HIP 13402	NSO	9385	4084	18068
HD 3765	HIP 3206	NSO	11151	1294	-63202	HD 18143	HIP 13642	NSO	11079	1186	31950
HD 3861	HIP 3236	NSO	11055	105	-14796	HD 18144	HIP 13601	NSO	11418	386	-1329
HD 4203	HIP 3502	NSO	12020	431	-14140	HD 18449 HD 18632	HIP 13905 HIP 13976	NSO NSO	11487 11263	118 540	-37021 28826
HD 4208	HIP 3479	NSO	11345	1427	56726	HD 18803	HIP 13976 HIP 14150	NSO	11203	1426	28826 9878
HD 4256 HD 4307	HIP 3535 HIP 3559	NSO NSO	11069 11303	1177 1390	9460 -10349	HD 18907	HIP 14086	NSO	11172	2	42718
HD 4614	HIP 3821	NSO	9380	4080	8314	HD 19034	HIP 14241	NSO	11082	1186	-20344
HD 4614B		M	10036	1860	11293	HD 19308	HIP 14532	NSO	11234	714	32723
HD 4628	HIP 3765	NSO	9262	3747	-10230	HD 19373	HIP 14632	NSO	7328	532	49449
HD 4915	HIP 3979	NSO	11348	291	-3742	HD 19467	HIP 14501	NSO	10996	1217	6936
HD 5065	HIP 4127	NSO	11336	510	-73844	HD 19994	HIP 14954	NSO	9069	4758	19331
HD 5133	HIP 4148	NSO	11470	96	-13067	HD 20165	HIP 15099	NSO	11070	1186	-16676
HD 5372	HIP 4393	NSO	11325	950	596	HD 20619	HIP 15442	NSO	10941	1178	22689
HD 6101	HIP 4849	NSO	11041	58	22170	HD 20630 HD 21019	HIP 15457 HIP 15776	NSO NSO	7834 10991	1847 1178	19021 41630
HD 6611 HD 6734	HIP 5276	NSO	11027	1426	-6068 -94511	HD 21197	HIP 15919	NSO	11145	1186	-13107
HD 7047	HIP 5315 HIP 5534	NSO NSO	11130 11128	1426 418	9264	HD 21313	HIP 16107	NSO	11901	182	-20017
HD 7228	HIP 5682	NSO	11450	509	-20538	HD 21847	HIP 16517	NSO	11263	714	30033
HD 7590	HIP 5944	NSO	11425	27	-13073	HD 22049	HIP 16537	NSO	7147	177	16332
HD 7727	HIP 5985	NSO	11262	514	5898	HD 22072	HIP 16641	NSO	11070	1463	11050
HD 8262	HIP 6405	NSO	11290	507	5636	HD 22484	HIP 16852	NSO	8391	2301	28080
HD 8389		NSO	10869	1177	34647	HD 22879	HIP 17147	NSO	11076	1427	120356
HD 8574	HIP 6643	NSO	11328	1161	18886	HD 23249 HD 23356	HIP 17378 HIP 17420	NSO NSO	11162 11220	1330 501	-6295 25287
HD 8648		NSO	11441	987	923	HD 23439	HIP 17666	NSO	11102	1119	50704
HD 8763 HD 8941	HIP 6732 HIP 6869	NSO NSO	11487 11171	117 419	-42532 9132	HD 24040	HIP 17960	NSO	11278	954	-9423
HD 9224	HIP 7090	NSO	11174	507	14937	HD 24213	HIP 18106	NSO	11181	712	-39608
HD 9280	HIP 7080	NSO	11172	3	41861	HD 24238	HIP 18324	NSO	11038	1088	38809
HD 9331	HIP 7221	NSO	11933	246	-20086	HD 24341	HIP 18309	NSO	11038	1088	142798
HD 9407	HIP 7339	NSO	11248	538	-33291	HD 24365	HIP 18208	NSO	11092	1088	19278
HD 9562	HIP 7276	NSO	11071	1390	-14989	HD 24451	HIP 18774	NSO	11404	354	17670
HD 9826	HIP 7513	NSO	8547	2576	-28674	HD 24496	HIP 18267	NSO	11236	954	18936
HD 9986	HIP 7585	NSO	11276	987	-21047	HD 24727 HD 24892	HIP 18388 HIP 18432	NSO NSO	11304 10952	954 1217	-18140 45410
HD 10002 HD 10086	HIP 7539 HIP 7734	NSO NSO	10973 11297	1119 514	11562 2143	HD 24916	HIP 18512	NSO	11245	1427	3540
HD 10086 HD 10126	HIP 7733	NSO	11301	293	56156	HD 25069	HIP 18606	NSO	11133	1427	38484
HD 10120	HIP 7902	NSO	11098	1121	17838	HD 25665	HIP 19422	NSO	11436	29	-13546
HD 10436	HIP 8070	NSO	11469	184	-50957	HD 25680	HIP 19076	NSO	9509	4084	24034
HD 10476	HIP 7981	NSO	10921	1387	-33647	HD 25723	HIP 19011	NSO	11490	144	26748
HD 10697	HIP 8159	NSO	11044	1044	-46022	HD 25918	HIP 19301	NSO	11538	5	-36198
HD 10700	HIP 8102	NSO	9838	5053	-16619	HD 26151	HIP 19232	NSO	11416	412	-6804
HD 10780	HIP 8362	NSO	11404	384	2764	HD 26161 HD 26162	HIP 19428 HIP 19388	NSO NSO	11236 11497	714 116	12924 24776
HD 11020	HIP 8346	NSO	10999	1081	22786	HD 26767	HIP 19386	NSO	11728	564	38288
HD 11226 HD 11505	HIP 8548 HIP 8798	NSO NSO	11310 11279	513 420	10314 -16470	HD 26794	HIP 19788	NSO	10980	1131	56573
HD 11964	HIP 9094	NSO	11109	1391	-10470 -9366	HD 26965	HIP 19849	NSO	7657	1845	-42331
HD 12051	HIP 9269	NSO	11097	1125	-35102	HD 28005	HIP 20800	NSO	11231	714	34768
HD 12235	HIP 9353	NSO	9727	3965	-18434	HD 28187	HIP 20638	NSO	10990	1215	18321
HD 12414	HIP 9473	NSO	10997	301	15898	HD 28343	HIP 20917	M	11476	356	-35073
HD 12661	HIP 9683	NSO	11463	944	-47310	HD 28344	HIP 20899	NSO	11524	171	39145
HD 12846	HIP 9829	NSO	11435	595	-4662	HD 28676	HIP 21158	NSO	11305	361	6671 46353
HD 13043	HIP 9911	NSO	11030	1043	-39333	HD 28946 HD 29150	HIP 21272 HIP 21436	NSO NSO	11395 11396	386 386	-46353 -6615
HD 13507	HIP 10321	NSO	11141	397	6172	111/2/100	-111 21 130	1,50	11570	200	0013

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Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$	Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$
Name	Name	(NSO/M)	(-2440000)	(days)	$(m s^{-1})$	Name	Name	(NSO/M)	(-2440000)	(days)	$(m \ s^{-1})$
(1)	(2)	(3)	(4)	(5)	(6)	(1)	(2)	(3)	(4)	(5)	(6)
HD 29528	HIP 21703	NSO	11901	182	-18922	HD 56274	HIP 35139	NSO	10929	1125	66663
HD 29883	HIP 21988	NSO	11359	385	17843		HIP 35209	NSO	11164	705	8150
HD 30562	HIP 22336	NSO	11370	293	77189		HIP 36249	NSO	11285	744	5013
HD 30708	HIP 22576	NSO	11209	714	-55738	HD 59747	HIP 36704	NSO	11441	472	-15744
HD 31253	HIP 22826	NSO	11345	954	12184	HD 60491	HIP 36827	NSO	11385	410	-9667
HD 31560	HIP 22907	NSO	11162	1217	6203		HIP 37349	NSO	11255	1163	-18210
HD 31966	HIP 23286	NSO	11239	743	-18058		HIP 37853	NSO	10415	97	106159
HD 32147	HIP 23311	NSO	9858	4085	21552		HIP 38228	NSO	11232	703	-15888
HD 32923 HD 32963	HIP 23835 HIP 23884	NSO NSO	11305 11235	388 715	20558 -62435		HIP 38216 HIP 38541	NSO NSO	10867 10968	1132 416	44973 -234268
HD 33021	HIP 23852	NSO	11070	1216	-02433 -22300		HIP 38931	NSO	11005	1120	-234208 -4417
HD 33632	HIP 24332	NSO	11182	708	-1707		HIP 39157	NSO	10958	1133	14832
HD 33636	HIP 24205	NSO	11411	954	5714	HD 66171	HIP 39822	NSO	10944	1119	36412
HD 33793	HIP 24186	M	11376	410	245194	HD 66428	HIP 39417	NSO	11926	92	44140
HD 34411	HIP 24813	NSO	7970	1859	66511		HIP 39780	NSO	11431	456	-36012
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HD 34721		NSO	11041	1217	40448		HIP 40118	NSO	11068	1218	29575
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HD 35681	HIP 25580	NSO	11195	702	12136		HIP 40761	NSO	11743	422	17534
HD 35974		NSO	11017	1162	76683		HIP 40843	NSO	8392	2250	32733
HD 36003	HIP 25623	NSO	11214	1215	-55527		HIP 41226	NSO	11267	795	-13087
HD 36395	HIP 25878	M	10194	4537	8665		HIP 41484	NSO	11177	798	-32342
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HD 37213	HIP 26273	NSO	11088	1163	12464		HIP 41479	NSO	11255	742	60217
HD 37394 HD 37588	HIP 26779 HIP 26689	NSO NSO	9701 10857	2486 0	1207 -58276		HIP 41844 HIP 42030	NSO NSO	11046 11255	745 742	13629 -18253
HD 37962	HIP 26737	NSO	10899	862	3133		HIP 41926	NSO	10792	808	14785
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HD 39715	HIP 27918	NSO	11461	386	-33797		HIP 42499	NSO	10981	1091	-12088
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HD 40397 HD 40650	HIP 28267 HIP 28634	NSO NSO	11113 11030	713 343	143621 -76840		HIP 43410 HIP 43587	NSO NSO	11173 9301	469 2513	4256 27351
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HD 43523	HIP 30023	NSO	11261	709	-16172		HIP 44897	NSO	9485	3286	26120
HD 43745	HIP 29843	NSO	10964	1187	-2716		HIP 45343	M	11045	974	11142
HD 43947	HIP 30067	NSO	11181	714	40545		HIP 120005	M	11077	974	12495
HD 44420 HD 44985	HIP 30243 HIP 30552	NSO NSO	11745 11227	464 714	-531 32253		HIP 45737 HIP 46580	NSO NSO	11416 10970	412 323	51009 29836
HD 45067	HIP 30545	NSO	11118	1166	47280		HIP 47202	NSO	11982	167	28994
HD 45184	HIP 30503	NSO	10940	1220	-3856		HIP 47690	NSO	10989	1121	-12225
HD 45350	HIP 30860	NSO	11819	431	-20727	HD 84737	HIP 48113	NSO	8233	1919	4900
HD 45391	HIP 30862	NSO	11459	472	-5388	HD 85725	HIP 48468	NSO	10726	752	61610
HD 45588	HIP 30711	NSO	10911	1187	35856		HIP 49081	NSO	8649	2716	55956
HD 46375	HIP 31246	NSO	11491	514	-1032		HIP 49350	NSO	11113	737	-260
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HD 48682	HIP 32480	NSO	8280	2301	-44233 -23933		HIP 49099	NSO	11062	737	-17886
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HD 50806	HIP 332// HIP 33094	NSO NSO	10965 10870	1119 1187	-15023 72388		HIP 50473 HIP 50478	NSO NSO	11074 11167	358 1119	23011 29607
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HD 52456	HIP 33848	NSO	11325	514	-11909		HIP 51258	NSO	11703	423	39984
HD 52711	HIP 34017	NSO	9614	4502	24604		HIP 51459	NSO	8584	3023	8529
HD 53665	HIP 34239	NSO	11232	747	-14703		HIP 51468	NSO	11381	383	4948
HD 55575	HIP 35136	NSO	11386	476 685	84809		HIP 51579	NSO	11853	432	-8875 -4751
HD 56124	HIP 35265	NSO	11324	685	22526	HD 91638	HIP 51784	NSO	11056	381	-4751

Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$	Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$
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HD 92945	HIP 52462	NSO	11323	1218	22856		HIP 66621	NSO	11950	307	16460
HD 93745	HIP 52888	NSO	11019	1218	37947		HIP 67155	M	10336	2856	15809
HD 94765	HIP 53486	NSO	11599	93	5525	HD 120066	HIP 67246	NSO	11205	1241	-30561
HD 95128	HIP 53721	NSO	9800	4241	11235	HD 120467	HIP 67487	NSO	11252	1157	-37806
HD 95650	HIP 53985	M	11816	423	-13899		HIP 67422	NSO	11492	392	-20380
HD 95735		M	10213	4748	-84689		HIP 68030	NSO	11016	23	-10532
HD 96418		NSO	11101	411	-8013		HIP 68184	NSO	11958	548	-26471
HD 96574		NSO	11271	57	36055		HIP 68337	NSO	11130	1157	-57444
HD 96700 HD 97004		NSO	10995	1261 293	12769		HIP 68469	M	11497	172	-25813
HD 97004 HD 97037		NSO NSO	11361 11294	709	5352 -15914		HIP 68593 HIP 69357	NSO NSO	11177 11117	675 1396	1409 3288
HD 97101		NSO	10384	2902	-16376		HIP 69414	NSO	11104	474	37773
HD 97101B		M	10307	1930	-15333		HIP 69526	NSO	11585	89	-16242
HD 97334	HIP 54745	NSO	9429	3273	-3662		HIP 69518	NSO	11016	22	-14644
HD 97343	HIP 54704	NSO	11051	1261	39794	HD 125184	HIP 69881	NSO	11184	1403	-12377
HD 97658	HIP 54906	NSO	10965	1088	-1654	HD 125455	HIP 70016	NSO	11154	1404	-9806
HD 98281		NSO	11044	1121	13330		HIP 70319	NSO	10208	4746	-19241
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HD 98697		NSO	11441	453	-14312		HIP 70782	NSO	11050	202	-5839
HD 99109 HD 99491		NSO	11478	534	32997		HIP 70873	NSO	11129	533	-405
HD 99491 HD 99492	HIP 55846 HIP 55848	NSO NSO	11047 11110	1286 1243	4190 3726		HIP 71284 HIP 71395	NSO NSO	8333 11392	2080 696	141 -9569
HD 100180		NSO	11017	1243	-4854		HIP 71462	NSO	11063	795	-42074
HD 100623		NSO	11003	1162	-21959		HIP 71774	NSO	11149	1158	-8137
HD 101177		NSO	11461	386	-16912		HIP 71803	NSO	11827	479	12685
HD 101259	HIP 56830	NSO	11077	1121	96905	HD 130087	HIP 72190	NSO	11835	483	-15673
HD 101501	HIP 56997	NSO	7992	2213	-5565	HD 130307	HIP 72312	NSO	11218	1429	12864
HD 102158	HIP 57349	NSO	11009	1121	28122		HIP 72339	NSO	12038	407	-12441
HD 102634		NSO	11099	381	1254		HIP 72577	NSO	11147	1429	-32159
HD 102870	HIP 57757	NSO	7466	1058	4448		HIP 72567	NSO	9090	3698	-2502
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HD 103432 HD 103932		NSO NSO	11033 11114	1241 1218	6079 48499		HIP 72772 GJ 566A	NSO NSO	11105 7325	1242 1157	-28834 1303
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HD 104800	HIP 58843	NSO	11820	547	9989	HD 132142	HIP 73005	NSO	11094	1158	-14771
HD 105113	HIP 59021	NSO	11070	1218	31871	HD 132375	HIP 73309	NSO	11139	624	-24370
HD 105405	HIP 59175	NSO	11412	325	2196		HIP 73593	NSO	11181	796	-32628
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HD 105631	HIP 59280	NSO	11033	1240	-2428		HIP 74234	NSO	11365	692	310614
HD 106116 HD 106156	HIP 59532 HIP 59572	NSO NSO	11063 11063	1241 1240	14543 -7415		HIP 74500 HIP 74432	NSO NSO	10962 11259	952 1480	5038 -38921
HD 107148	HIP 60081	NSO	11890	548	25216		HIP 74702	NSO	11108	468	-3149
HD 107705	HIP 60353	NSO	11278	797	4759		HIP 75101	NSO	11483	720	-46935
HD 108510	HIP 60816	NSO	11207	0	6888		HIP 75104	NSO	11020	21	-27029
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HD 109358		NSO	11735	480	6259		HIP 75266	NSO	10982	1403	-26374
HD 110315	HIP 61901	NSO	11108	1242	24604		HIP 75277	NSO	11196	798	-7042
HD 110537	HIP 62039	NSO	11386	748	35461		HIP 75281	NSO	11253	842	-48942
HD 111031	HIP 62345	NSO	11051	1243	-20458		HIP 75722	NSO	11234	1479	7752
HD 111066 HD 111398	HIP 62349	NSO	11351	330	6407		HIP 76200 HIP 76114	NSO NSO	11100	1158	12024
HD 111398 HD 111484	HIP 62536 HIP 62596	NSO NSO	11411 11408	331 721	3154 -20643		HIP 76114 HIP 76228	NSO NSO	11221 11509	625 368	-35691 10674
HD 111484B	SAO 119612	NSO	11387	721	-19377		HIP 76375	NSO	11092	1134	-67101
HD 111515	HIP 62607	NSO	11059	1243	2548		HIP 76543	NSO	11427	338	37708
HD 111631	HIP 62687	M	11901	517	5041	HD 139477	HIP 76315	NSO	11387	453	-8893
HD 112257	HIP 63048	NSO	11044	473	-39428	HD 141004	HIP 77257	NSO	8789	3560	-66416
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HD 114710	HIP 64394	NSO	7397	1058	5295		HIP 77760	NSO	8418	3561	-56107
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HD 116442	HIP 65352	NSO	11301	1242	28421		HIP 79214	NSO	11112	1208	-14067 -53211
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HD 117207	HIP 65808	NSO	11199	1240	-17476		HIP 79152	NSO	11086	359	-42685

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HD 145809	HIP 79524	NSO	11088	1428	21146		HIP 90004	NSO	12060	406	-25603
HD 145897	HIP 79540	NSO	11457	301	-23611	HD 169191	HIP 90067	NSO	11414	127	-19268
HD 145934	SAO 102017	NSO	11342	1101	-27954	HD 169830	HIP 90485	NSO	12036	406	-17304
HD 145958	HIP 79492	NSO	11183	1518	18422		SAO 123515	NSO	11177	1208	-28841
HD 146233		NSO	11245	1088	11748		HIP 90593	NSO	11954	457	-59423
HD 146362		NSO	11121	1134	-14855		HIP 90656	NSO	11147	1480	-54752
HD 146775		NSO	11116	1245	-30218		HIP 90790	NSO	10812	2173	-43131
HD 147044		NSO	11339	746	-14560		HIP 90586	NSO	11307	744	-22670
HD 147231		NSO	11182	1157	-16461		HIP 90864	NSO	11058	1479	-46252
HD 147379 HD 147379B		M M	12021 12040	482 482	-18789 -18530		HIP 91287 HIP 91332	NSO NSO	11040 11902	1421 394	-23259 -67256
HD 147776		NSO	11541	265	7341		HIP 91438	NSO	11149	1421	37103
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HD 149200		NSO	11017	22	-7176		HIP 91507	NSO	11016	23	-11419
HD 149652	HIP 81279	NSO	11080	359	-31167	HD 173701	HIP 91949	NSO	11071	1158	-45602
HD 149661	HIP 81300	NSO	10107	3304	-12857	HD 173739	HIP 91768	M	11078	1104	-834
HD 149806	HIP 81375	NSO	11358	747	10391	HD 173740	HIP 91772	M	11078	1104	1187
HD 150433		NSO	11122	1065	-40215		HIP 92200	NSO	11499	443	15544
HD 150698		NSO	11073	1131	48229		HIP 92532	NSO	11301	418	-13083
HD 150933		NSO	11287	747	-27880		HIP 92747	NSO	11424	106	-7100
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HD 151541		NSO	11067	1158	9475		HIP 93185	NSO	11232	825	-40611
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HD 152391		NSO	10746	2588	45066		HIP 93746	NSO	11183	1213	-72173
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HD 157347	HIP 85042	NSO	11256	748	-35901	HD 185720	HIP 96813	NSO	11027	41	15992
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HD 163153	HIP 87710	NSO	11234	750	-73024		HIP 98677	NSO	11009	1517	20385
HD 163489	GJ 4035	NSO	11135	1163	-49364		HIP 98767	NSO	11189	1340	-45308
HD 164507	HIP 88217	NSO	11265	824	5370		HIP 98792	NSO	10990	1479	-2527
HD 164595	HIP 88194	NSO	11249	748	2020		HIP 98978	NSO	11093	441	-12842
HD 164922	HIP 88348	NSO	11196	1520	20248	HD 191785	HIP 99452	NSO	10947	1427	-49286
HD 165222	HIP 88574	M	10299	4747	32671	HD 192020	GJ 4138	NSO	11118	1207	-11366
HD 165438	HIP 88684	NSO	11457	274	-26747		HIP 99711	NSO	11346	810	-10738
HD 165567	HIP 88533	NSO	11202	382	3853		HIP 99727	NSO	11781	1767	-526
HD 165634	HIP 88839	NSO	11394	83	-4867		HIP 99729	NSO	12075	407	-451
HD 165683	HIP 88636	NSO	11426	32	-546		HIP 100072	NSO	11084	443	15584
HD 166620	HIP 88972	NSO	9370	3649	-19418		HIP 199427	NSO	11355	26	-1623
HD 167389		NSO	11254	749	-5545		HIP 100363	NSO	11976	369	6921
HD 168009	HIP 89474	NSO	11329	747	-64595	HD 193901	HIP 100568	NSO	11593	296	-171455

Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$	Primary	Alternate	Ref.	$\langle \mathrm{JD} \rangle$	ΔT	$\langle RV \rangle$
Name	Name	(NSO/M)	(-2440000)	(days)	$(m s^{-1})$	Name	Name	(NSO/M)	(-2440000)	(days)	$(m \ s^{-1})$
(1)	(2)	(3)	(4)	(5)	(6)	(1)	(2)	(3)	(4)	(5)	(6)
HD 194035	HIP 100500	NSO	11252	444	16665	HD 217877	HIP 113896	NSO	11182	442	-12676
HD 194766	HIP 100895	NSO	11101	412	-15131	HD 217987	HIP 114046	M	11139	385	8809
HD 195104	HIP 101059	NSO	11022	22	-14936	HD 218029	HIP 113864	NSO	11427	106	-8708
HD 195564	HIP 101345	NSO	11141	1428	9535	HD 218133	HIP 114028	NSO	11147	443	-48799
HD 196201		NSO	11431	426	-19821	HD 218209	HIP 113989	NSO	11114	1040	-15895
HD 196761 HD 196850	HIP 101997 HIP 101875	NSO NSO	11023 11196	1478 1207	-41987 -21045	HD 218566 HD 218730	HIP 114322 HIP 114424	NSO NSO	11028 11160	1186 442	-37804 2009
HD 196885	HIP 101966	NSO	11150	441	-30189	HD 218739	HIP 114385	NSO	11006	0	-5681
HD 197076		NSO	11035	1389	-35409	HD 218792	HIP 114449	NSO	11423	105	2548
HD 197139	HIP 101986	NSO	11437	125	-24024	HD 218868	HIP 114456	NSO	11193	442	-30622
HD 198089		NSO	11158	441	-33434	HD 219134	HIP 114622	NSO	10206	672	-18558
HD 198802 HD 199305	HIP 103077 HIP 103096	NSO M	11104 11164	1101 1077	-3171 -17161	HD 219172 HD 219538	HIP 114670 HIP 114886	NSO NSO	11137 11096	419 1090	-2828 9990
HD 199476	HIP 102970	NSO	11040	765	-30230	HD 219834B	HIP 115125	NSO	11301	1427	10782
HD 199598	HIP 103455	NSO	11354	747	-25301	HD 219953	HIP 115194	NSO	10989	1517	-48118
HD 199960	HIP 103682	NSO	11054	1338	-17605	HD 220339	HIP 115445	NSO	11132	1389	34001
HD 200538	HIP 104071	NSO	10991	1389	15377	HD 220957	HIP 115839	NSO	10427	97	-81264
HD 200746 HD 201091	HIP 104075 HIP 104214	NSO NSO	11408 7369	67 1066	13807 -65726	HD 221146 HD 221354	HIP 115951 HIP 116085	NSO NSO	11100 11410	439 336	-15137 -25113
HD 201091	HIP 104214	M	8258	2297	-64023	HD 221356	HIP 116106	NSO	10925	1479	-23113 -12713
HD 202108		NSO	11258	747	2502	HD 221830	HIP 116421	NSO	11211	444	-112299
HD 202573	HIP 105000	NSO	10823	799	-26179	HD 222033	HIP 116542	NSO	11195	443	-13088
HD 202575		NSO	11175	1337	-18065	HD 222143	HIP 116613	NSO	11152	793	-169
HD 202751 HD 203644	HIP 105152 HIP 105497	NSO NSO	11054 11436	1338 75	-27427 -4356	HD 222368 HD 222582	HIP 116771 HIP 116906	NSO NSO	7945 11315	1787 950	5656 12067
HD 204587	HIP 106147	M	11004	1338	-4336 -84186	HD 223238	HIP 117367	NSO	11216	950	-15420
HD 206332		NSO	11254	746	-44542	HD 223498	HIP 117526	NSO	11031	1185	-23985
HD 206387	HIP 107107	NSO	11807	821	-7761	HD 223559	HIP 117567	NSO	11465	117	-62839
HD 206860		NSO	8929	2894	-16833	HD 223691	HIP 117668	NSO	10999	704	1656
HD 207804	HIP 107840	NSO	10636	63	-14848	HD 223807	HIP 117756	NSO	11483	116	-15825
HD 207874 HD 208313	HIP 107941 HIP 108156	NSO NSO	11243 11509	720 335	-26702 -13251	HD 224156 HD 224383	HIP 117953 HIP 118115	NSO NSO	11482 11146	117 633	14883 -31205
HD 208801		NSO	10595	2552	-50179	HD 225216	HIP 379	NSO	11476	117	-28860
HD 209128	HIP 108691	NSO	11432	144	8115	HD 225261	HIP 400	NSO	11105	1389	7511
HD 209290	HIP 108782	M	11425	30	18363	HD 230409	HIP 93341	NSO	11578	346	-2238
HD 209458 HD 209747	HIP 108859	NSO	11586	414 127	-14759 -18889	HD 230999 HD 232979	HIP 94615	NSO M	11930	335	-81527 34401
HD 209747 HD 209761	HIP 109068 HIP 109023	NSO NSO	11404 11443	105	-18889 -28712	HD 232641	HIP 21553 HIP 46639	NSO	11693 11935	356 92	36182
HD 209875	HIP 109144	NSO	11208	750	-40896	HD 239960	HIP 110893	M	11146	833	-33937
HD 210277	HIP 109378	NSO	10756	766	-20873	HD 245409	HIP 26335	M	11002	1161	22046
HD 210302	HIP 109422	NSO	11472	363	-16259	HD 260655	HIP 31635	M	11181	713	-58178
HD 210392 HD 210460	HIP 109428	NSO	11972	369	-1496	HD 265866 HD 281540	HIP 33226	M	10225	2907	22914
HD 210460 HD 210667	HIP 109439 HIP 109527	NSO NSO	10843 11175	1093 435	20429 -19443	HD 285968	HIP 19143 HIP 21932	NSO M	11515 11238	170 742	110217 26219
HD 210752	HIP 109646	NSO	11157	390	57225	HD 349726	HIP 93873	M	11304	723	32407
HD 210762	HIP 109602	NSO	11438	93	-8799	GJ 2	HIP 428	M	12055	405	-240
HD 211038	HIP 109822	NSO	10807	796	10407	GJ 4A	HIP 473	M	11979	405	1669
HD 212801	HIP 110853	NSO	11021	1389	-8475	GJ 4B	BD +45 4408B	M	12067	432	-1655
HD 213119 HD 213519	HIP 110986 HIP 111148	NSO NSO	11407 11247	81 949	-30288 -31630	GJ 14 GJ 26	HIP 1368 G132-11	M M	9537 12030	3745 608	2957 -347
HD 213575	HIP 111274	NSO	11282	449	-21544	GJ 47	G243-50	M	11852	551	7574
HD 213628	HIP 111349	NSO	10939	1003	-50445	GJ 48	HIP 4856	M	11159	866	1500
HD 214557	HIP 111748	NSO	11177	443	-38534	GJ 49	HIP 4872	M	11083	877	-5966
HD 214749	HIP 111960	NSO	10984	1165	1	GJ 54.1	HIP 5643	M	11263	532	28089
HD 214868 HD 214995	HIP 111944 HIP 112067	NSO NSO	11404 11439	128 93	-10881 -28855	GJ 70 GJ 83.1	HIP 8051 G73-12	M M	11867 11566	552 383	-25883 -28567
HD 215152	HIP 112190	NSO	11345	783	-13795	GJ 96	HIP 11048	M	11890	607	-37941
HD 215648	HIP 112447	NSO	8320	1809	-5856	GJ 107B	BD +48 746B	M	9713	3608	25765
HD 216259	HIP 112870	NSO	10981	1517	1291	GJ 109	HIP 12781	M	11522	173	30568
HD 216625	HIP 113086	NSO	11020	23	10993	GJ 156	HIP 18280	M	10978	1089	62597
HD 216646 HD 216899	HIP 113084 HIP 113296	NSO M	11435 11166	93 1127	-7670 -27317	GJ 173 GJ 226	HIP 21556 HIP 29277	M M	11788 11126	319 741	-6768 -1622
HD 217004	HIP 113296 HIP 113386	NSO	11015	1072	371	GJ 273	HIP 36208	M	10129	3167	18216
HD 217014	HIP 113357	NSO	10005	9	-33225	GJ 273.1	HIP 36357	NSO	11155	0	-3951
HD 217107	HIP 113421	NSO	11199	700	-13399	GJ 357	HIP 47103	M	11204	743	-34581
HD 217165	HIP 113438	NSO	11496	452	15391	GJ 361	HIP 47513	M	11753	391	11511
HD 217357	HIP 113576	M	11087	1074	16420	GJ 362	HIP 47650	M M	11777	393 422	6606 15493
HD 217459 HD 217563	HIP 113622 HIP 113686	NSO NSO	11446 11448	137 96	18630 4661	GJ 373 GJ 382	HIP 48714 HIP 49986	M M	11703 11053	422 974	15493 7932
HD 217618	HIP 113695	NSO	11515	334	-12730	GJ 388	SAO 81292	M	9343	3135	12420
HD 217813	HIP 113829	NSO	11423	31	2030	GJ 390	HIP 51007	M	11818	428	21590

TABLE 1—Continued

Primary Alternate Ref. $\langle JD \rangle$ ΛT $\langle RV \rangle$ (NSO/M) (-2440000) $(m s^{-1})$ Name (days) Name (2) (4) GJ 393..... HIP 51317 11020 936 8335 M GJ 397..... HIP 51525 Μ 11323 351 21052 GJ 402..... HIP 53020 M 11776 513 -1042GJ 406..... G45-20 M 11316 380 19482 GJ 408..... HIP 53767 M 11133 975 3151 GJ 412A..... HIP 54211 M 10386 2903 68886 GJ 413.1.... HIP 54532 M 11759 456 -3830GJ 424..... HIP 55360 Μ 11796 513 60401 GJ 433..... HIP 56528 M 11419 17973 GJ 436..... HIP 57087 M 11778 513 9607 GJ 445..... HIP 57544 M 11227 863 -111654GJ 447..... HIP 57548 Μ 11362 722 -31087GJ 450..... HIP 57802 Μ 11247 863 273 GJ 465..... HIP 60559 M 11764 455 51174 GJ 486..... HIP 62452 M 11146 867 19090 GJ 514..... HIP 65859 Μ 11085 1158 14556 -14557 GJ 519..... HIP 66459 M 11827 482 GJ 528B BD +27 2296B NSO 11519 369 -22896GJ 553.1..... HIP 70975 11882 520 -1756Μ HIP 71253 171 GJ 555..... M 11496 -1453GJ 569..... HIP 72944 M 9424 1336 -7209GJ 581..... HIP 74995 Μ 11498 176 -9398 GJ 615.1B BD +13 3091B 18427 NSO 11244 1520 GJ 625..... HIP 80459 M 11224 842 -13034GJ 628..... HIP 80824 M 11759 688 -21222GJ 649..... 12002 HIP 83043 M 780 4316 GJ 667C..... Μ 11982 458 6353 457 GJ 671..... HIP 84790 M 12056 -19528 GJ 678.1A..... HIP 85665 749 -12457M 11218 GJ 686..... HIP 86287 M 11149 1099 -9515GJ 687..... -28779 HIP 86162 M 11173 1101 GJ 694..... HIP 86776 M 11250 748 -14269GJ 699..... -110506HIP 87937 M 10418 4022 HIP 92403 Μ -10499GJ 729..... 11397 721 GJ 745B HIP 93899 Μ 11304 723 32171 GJ 793..... HIP 101180 M 11209 401 10599 HIP 102401 Μ 9736 4037 -24702 HIP 104432 12001 -58267GJ 821..... Μ 430

these instantaneous velocities.⁶ Four observations of the day sky and two observations of Vesta were analyzed with the before mentioned Doppler method using the NSO solar spectrum. These references showed that our raw velocities, from Keck and Lick, were consistently large by 522 ± 5 m s⁻¹.

There are various possible sources for the 522 m s⁻¹ offset in our raw velocities. The absolute wavelength scales of both the NSO solar spectrum and our FTS iodine spectrum are not well known. According to Kurucz et al. (1984), the solar lines are broadened by 200 m s⁻¹ due to the change in the radial velocity during the observation, and the wavelengths may have errors as large as 100 m s⁻¹. The wavelength scale of the iodine FTS spectrum comes from the calibration made at the McMath telescope at Kitt Peak. We have no independent way to verify the integrity of the zero point in wavelength of this FTS iodine spectrum. The third concern stems from the instrumental profile of the HIRES and the Hamilton spectrometers. It is known that the PSF is asymmetric to some degree (Valenti et al. 1995) and may give rise to systematic velocity shifts. However, it is unlikely that the same asymmetry in the same direction would be found at

TABLE 1—Continued

Primary Name	Alternate Name	Ref. (NSO/M)	(JD) (-2440000)	ΔT (days)	$\langle RV \rangle$ (m s ⁻¹)
(1)	(2)	(3)	(4)	(5)	(6)
GJ 849	HIP 109388	M	10972	834	-15256
GJ 905	G171-10	M	11248	542	-77949
GJ 908	HIP 117473	M	10076	4505	-71147
GJ 1005	HIP 1242	M	11044	60	-26425
GJ 1148	HIP 57050	M	11834	484	-9095
GJ 2066	HIP 40501	M	11042	744	62205
GJ 2130A	HIP 86961	M	12007	460	-28990
GJ 3126	G244-47	M	11872	551	-84115
GJ 3709B		M	11786	426	-9226
GJ 3804	HIP 67164	M	11910	517	4969
GJ 3897B	HIP 74434	NSO	11213	1306	-38829
GJ 3992	HIP 84099	M	11928	419	-44443
GJ 4048A	G204-58	M	11994	458	477
GJ 4062	G205-28	M	11940	423	-19003
GJ 4063		M	11519	296	12495
GJ 4070	HIP 91699	M	12021	456	-31764
GJ 4098	G207-19	M	12025	423	-1730
GJ 4333	HIP 115332	M	12042	229	-6507
GJ 9492	HIP 71898	M	11184	864	18704
HIP 795	BD + 07.9s	NSO	11419	208	14685
HIP 5004	G269-87	NSO	11865	547	45733
HIP 10449	SAO 129772	NSO	11481	140	28116
HIP 15904	BD +11 468	NSO	11668	249	86739
HIP 52942	SAO 99310	NSO	11425	536	24574
HIP 59406		M	11745	392	-9100
HIP 80295	BD -11 4126	NSO	11455	395	-17315
HIP 89215	BD +05 3640	NSO	11795	689	-1050
HIP 103039		M	11239	747	16300
HIP 103269	BD +41 3931	NSO	11531	336	-130673
HIP 106924	BD +59 2407	NSO	11640	296	-244634
BD +18 4505C		NSO	11999	722	-89693
BD -10 3166		NSO	11499	533	26679
G161-29		NSO	11448	413	22370
G195-59		M	11782	401	-3451
G60-06		NSO	11418	533	-18105

NOTE.— Table 1 is also available in machine-readable form in the electronic edition of the *Astrophysical Journal Supplement*.

two different spectrometers. Therefore, we consider this possibility less plausible.

All Doppler measurements which used the NSO solar spectrum as the reference template were corrected for this offset of 522 m s⁻¹. Since currently no M dwarfs exist having a definitive barycentric radial velocity, the velocity zero point for the M stars in our sample was set using the previously published velocities for the five standards from Marcy et al. (1987), which have errors of ~0.4 km s⁻¹.

6. COMPARISON OF PRESENT VELOCITIES WITH STANDARD STARS

To get an external measure of the accuracy of our bary-centric radial velocities we have compared our velocities to published velocities of supposed radial velocity standard stars. We carry out this comparison separately for the FGK stars and for the M stars.

6.1. F. G. and K Stars

We compared our velocities to those of the 26 standard stars that were measured by Udry et al. (1999a). The results of this comparison are shown in Figures 2 and 3 and yield $\langle V_{\rm Udry} - V_{\rm Present} \rangle = -53~{\rm m~s^{-1}}$, with an rms scatter of

⁶ See http://ssd.jpl.nasa.gov/cgi-bin/eph.

^a Stars with $\sigma_{\rm rms}$ < 100 m s⁻¹.

 $\label{eq:table 2} TABLE~2$ Radial Velocities: Stars with rms $>0.1~km~s^{-1a}$

		Ref.		RV		ΔT	$\sigma_{ m rms}$		
Primary Name	Alternate Name	(NSO/M)	JD (-2440000)	$(m s^{-1})$	$\langle JD \rangle (-2440000)$	(days)	$(m s^{-1})$	N	Comment
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
HD 1854	SAO 192490	NSO	10984.112	24327	10714	119	3651	4	
HD 3346	HIP 2900	NSO	10012.747	-33377	9317	2220	172	29	
HD 3770	HIP 3169	NSO	11774.903	-8124	11555	1132	178	13	L
HD 3795	HIP 3185	NSO	11756.021	-46595	11131	1535	156	24	L
HD 4271	HIP 3540	NSO	11047.927	32566	11037	21	7799	2	
HD 4747	HIP 3850	NSO	11756.051	10399	11092	1731	315	21	CO
HD 7483	HIP 5881	NSO	12129.043	-17717	12015	515	2602	11	CO
HD 8331 HD 9770	HIP 6442	NSO NSO	11154.627 11070.080	-12270 33562	11091 11057	128 27	2349 122	2 2	
HD 11909	HIP 7372 HIP 9110	NSO	11438.923	-6173	11430	13	137	4	L
HD 14802	HIP 11072	NSO	11170.808	18824	10863	709	418	7	L
HD 17037	HIP 12764	NSO	11581.755	16915	11452	1295	560	14	L
HD 18445	HIP 13769	NSO	11580.723	49877	11491	1147	222	10	C
HD 21017	HIP 15861	NSO	11779.008	8561	11455	74	371	5	L
HD 25329	HIP 18915	M	11412.131	-25177	11179	361	338	3	
HD 25535	HIP 18824	NSO	10786.825	10205	10576	421	203	2	
HD 28388	HIP 20802	NSO	11170.888	17748	10942	457	498	6	L
HD 29461	HIP 21654	NSO	11581.849	40518	11763	563	116	8	L
HD 29587	HIP 21832	NSO	11073.077	111729	10856	610	786	4	00
HD 30339	HIP 22429	NSO	12003.721	12208	11911	211	1538	10	CO
HD 30649	HIP 22596	NSO	11582.778	28904	11391	1545	108	14	L
HD 31412 HD 32387	HIP 22919 HIP 23550	NSO NSO	11581.861 11072.132	47494 54468	11317 10759	1062 706	144 1819	12	L
HD 32450	HIP 23452	M	11072.132	-11004	10759	233	380	2	
HD 34101	HIP 24419	NSO	11073.046	33664	10950	1942	2543	7	CO
HD 35956	HIP 25662	NSO	11551.891	9942	11346	1822	2457	14	C
HD 39587	HIP 27913	NSO	10068.858	-12171	8846	4580	1294	38	CO
HD 43587	HIP 29860	NSO	11582.848	5776	11160	1638	3361	14	C
HD 50639	HIP 33109	NSO	11552.956	-2947	11412	1169	207	8	L
HD 54563	HIP 34608	NSO	10784.061	-15291	10553	365	8953	4	
HD 64468	HIP 38657	NSO	11581.909	-10624	11267	1541	3991	13	C
HD 65430	HIP 39064	NSO	11581.938	-28709	11214	1846	470	26	CO
HD 68988	HIP 40687	NSO	12064.768	-69520	11919	513	120	13	C
HD 72780	HIP 42112	NSO	11626.698	28059	11693	412	169	10	L
HD 73512 HD 73668	HIP 42418 HIP 42488	NSO NSO	11171.030 11229.858	10737 -21456	11199 11270	56 775	29558 112	2 24	L
HD 81133	HIP 45995	NSO	10546.835	11635	10624	400	2456	3	L
HD 86680	HIP 49060	NSO	11974.896	7317	11935	91	4728	8	L
HD 89744	HIP 50786	NSO	11698.722	-5401	11675	737	166	45	C
HD 94340	HIP 53217	NSO	10545.893	-15203	10518	84	5126	3	
HD 100167	HIP 56257	NSO	11242.931	-27596	11037	411	4649	2	
HD 101206	HIP 56829	M	11551.097	40597	11107	945	22398	5	C
HD 101563	HIP 57001	NSO	10955.788	-9602	10745	494	656	5	L
HD 102540	HIP 57574	NSO	10955.792	-4260	10745	494	4055	5	L
HD 103829	HIP 58318	NSO	12101.779	-552	11987	202	128	8	L
HD 106252	HIP 59610	NSO	11628.817	15368	11767	1209	114	15	С
HD 111312	HIP 62505	NSO	11582.041	623	11186	1119	1094	5	C
HD 114762	HIP 64426 HIP 65708	NSO	9858.733 11026.672	49576	9723	4069	434	66	С
HD 117126 HD 117176	HIP 65708 HIP 65721	NSO NSO	9124.792	-7637 5036	10975 9853	195 4928	1685 182	5 106	С
HD 117635	HIP 65982	NSO	10954.918	-51998	10752	4928	336	3	C
HD 120136	HIP 67275	NSO	8779.738	-16542	9558	5097	340	93	С
HD 120690	HIP 67620	NSO	10606.790	7159	10535	144	984	2	~
HD 122676	HIP 68634	NSO	11004.712	-8017	10955	196	1419	3	
HD 122742	HIP 68682	NSO	10504.970	-8081	8986	3545	3712	24	CO
HD 129814	HIP 72043	NSO	11680.036	6727	11474	1266	101	13	L
HD 131511	HIP 72848	NSO	9914.784	-31806	8966	2866	8832	20	CO
HD 131976	HIP 73182	M	10181.904	35628	9298	2958	10056	9	C
HD 136118	HIP 74948	NSO	11627.902	-3062	11660	1326	136	40	C
HD 136580	HIP 75039	NSO	11627.962	-25626	11463	1291	204	19	L
HD 140913	HIP 77152	NSO	11304.900	-18679	10474	1511	742	12	CO
HD 142229	HIP 77810	NSO	11704.927	-21443	11660	758	486	9	L
HD 142267	HIP 77801	NSO	10603.902	35666	10440	328	3931	2	

TABLE 2—Continued

Primary Name (1)	Alternate Name (2)	Ref. (NSO/M) (3)	JD (-2440000) (4)	RV (m s ⁻¹) (5)	⟨JD⟩ (−2440000) (6)	Δ <i>T</i> (days) (7)	$\sigma_{\rm rms} \atop ({\rm m \ s^{-1}}) \atop (8)$	N (9)	Comment (10)
HD 145206	HIP 79195	NSO	11407.673	-36942	11379	53	2180	3	
HD 152311	HIP 82621	NSO	10713.733	-21385	10622	166	354	3	
HD 157681	HIP 84950	NSO	11811.653	-7723	11449	299	133	6	
HD 158222	HIP 85244	NSO	10982.968	-14301	11103	413	2698	3	
HD 160346	HIP 86400	NSO	10656.768	15947	9951	1456	3724	7	
HD 161198	HIP 86722	NSO	10666.816	23951	10516	391	351	3	
HD 161797	HIP 86974	NSO	11372.852	-17004	9872	4324	210	53	L
HD 166435	HIP 88945	NSO	11026.782	-14403	11017	23	105	4	
HD 167665	HIP 89620	NSO	11702.962	7693	11320	1879	322	16	L
HD 168443	HIP 89844	NSO	11071.770	-48636	10976	825	182	37	С
HD 169822	HIP 90355	NSO	11793.810	-18274	11824	817	312	22	C
HD 171115	HIP 91004	NSO	11780.777	-2291	11425	106	175	5	
HD 173667	HIP 92043	NSO	8437.936	23043	9353	5165	124	91	
HD 174457	HIP 92418	NSO	11755.929	-25820	11700	755	908	9	CO
HD 175518	HIP 92918	NSO	11014.815	-66763	11019	22	1946	6	
HD 183255	HIP 95575	M	10666.899	-68922	10606	119	10961	3	
HD 184860	HIP 96471	NSO	11793.808	63944	11373	1845	514	21	С
HD 186704	HIP 97255	NSO	11411.863	-15117	11383	70	3021	4	Ĺ
HD 188376	HIP 98066	NSO	10713.746	-39136	10492	430	11891	4	
HD 190406	HIP 98819	NSO	11411.872	4757	9477	4364	121	52	L
HD 190771	HIP 98921	NSO	11826.657	-25063	11478	1072	136	13	L
HD 195019	HIP 100970	NSO	11792.783	-91582	11214	788	188	50	С
HD 197214	HIP 102264	NSO	10713.839	-21884	10565	437	4096	4	
HD 199918	HIP 103735	NSO	10957.090	52560	10661	591	3443	5	
HD 200565	HIP 103983	NSO	11792.788	-2425	11702	793	137	24	L
HD 208527	HIP 108296	NSO	10304.870	4154	10102	2550	133	30	
HD 208776	HIP 108473	NSO	11440.695	31491	11150	435	601	10	CO
HD 209779	HIP 109110	NSO	11751.921	-17310	11543	335	1722	6	L
HD 215578	SAO 108160	NSO	11439.882	-16217	11208	1823	1323	21	L
HD 217303	HIP 113562	NSO	11782.895	-36049	11439	96	205	7	
HD 219420	HIP 114834	NSO	11049.844	-22295	11027	44	2813	5	L
HD 220077	HIP 115279	NSO	11707.121	1521	11727	821	132	20	
HD 223084	HIP 117258	NSO	11706.125	2164	11310	701	180	33	L
GJ 84	HIP 9724	M	11583.710	23335	11419	1294	867	11	CO
GJ 285	HIP 37766	M	11629.714	26531	10322	2280	299	7	
GJ 494	HIP 63510	M	12127.800	-11228	12004	546	110	15	A
GJ 595	HIP 76901	M	11584.166	84921	11886	898	4624	10	CO
GJ 623	HIP 80346	M	10181.925	-26284	9208	3222	1238	12	CO
GJ 873	HIP 112460	M	11447.748	413	9714	4401	124	17	
GJ 876	HIP 113020	M	11072.938	-1591	10514	4016	181	53	С
HIP 35519	SAO 115271	NSO	11552.969	11636	11625	803	1529	5	-
HIP 52940	BD +13 2311B	NSO	11706.829	25212	11574	1137	1310	15	CO
HIP 57450	BD +51 1696	NSO	11679.832	64092	11520	810	1211	10	20
CD -32 8503B		NSO	11679.866	20940	11264	724	319	5	

Note.— Table 2 is also available in machine-readable form in the electronic edition of the Astrophysical Journal Supplement.

35 m s⁻¹. The corresponding uncertainty in the mean is $35/\sqrt{26} = 7$ m s⁻¹. Thus, the formal difference in zero points is

$$\langle V_{\text{Present}} - V_{\text{Udry}} \rangle = +53 \pm 7 \text{ m s}^{-1}. \tag{1}$$

Thus, there appears to be a statistically significant difference in the zero point of the velocities reported here compared to those of Udry et al. (1999a) of 53 m s⁻¹. We do not know the origin of this difference, nor do we know which scale is more "accurate." Few studies will be affected by such small differences. But future highly precise proper motion measurements and precise orbit calculations may require such accurate velocities, including proper treatment of gravitational redshift and convective blueshift.

However, the difference between the present velocities and those of Udry et al. (1999a) do exhibit a significant B-V dependence as seen in Figure 4. The slope of the linear trend is -182 ± 24 m s⁻¹ per mag with an rms scatter around the fit of 23 m s⁻¹. This dependence is similar to the B-V dependence seen between the CfA and CORAVEL data (Stefanik, Latham, & Torres 1999; Udry et al. 1999a). This color dependence is likely caused by some spectral type mismatch in one or all of the radial-velocity scales. For solar type stars (near G2 V) where our zero point is well set, we are in good agreement with the CORAVEL velocities, with an offset of only 25 m s⁻¹ as seen in Figure 4.

According to Udry et al. (1999a) the ELODIE measurements ensure temporal stability of better than 15 m s⁻¹ dur-

^a Stars with $\sigma_{\rm rms} \ge 100 \, {\rm m \ s^{-1}}$.

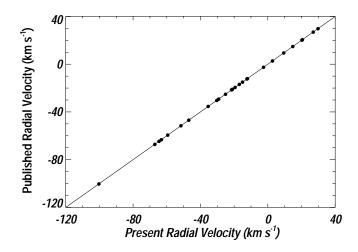


Fig. 2.—Velocities of standard stars vs. velocities measured here, for F, G, and K stars. The standard stars are on the CORAVEL scale (Udry et al. 1999a). The present velocities agree well with the standards, with no visible nonlinear departure of velocity scale.

ing timescales of years for their standard stars. However, their quoted velocities have been rounded off at the 50 m s⁻¹ level. From our measurements of these stars in common, the temporal variability is less than 10 m s⁻¹ during \sim 5 yr. Therefore, it is not certain whether the rms scatter of 35 m s⁻¹ between the two sets of velocities is due to our errors or the rounding of the ELODIE data plus their errors. We conclude that the barycentric velocities reside on the same scale with an offset of \sim 50 m s⁻¹, and a scatter of \sim 35 m s⁻¹.

Comparing our velocities for 29 common stars with those reported by Stefanik et al. (1999), with a correction of +136 m s⁻¹ added to their native velocities in Table 1 and 2 (R. P. Stefanik, 2002, private communication), we obtain $\langle V_{\rm CfA} - V_{\rm Present} \rangle = 15$ m s⁻¹, which is only marginally different from zero, with an rms scatter of 123 m s⁻¹ as seen in Figure 5. The differences exhibit a significant B-V dependence as seen in Figure 6. The slope of the linear trend is $+448 \pm 111$ m s⁻¹ per mag with an

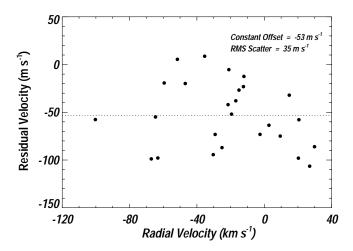


Fig. 3.—Standard-star velocities (Udry et al. 1999a) minus present velocities for all 26 FGK stars in common (as in Fig. 2). The differences reveal that the present velocities are higher than those of Udry et al. by 53 m s $^{-1}$ and exhibit an rms scatter of 35 m s $^{-1}$. Thus, the present velocities and those of Udry et al. (1999a) each have internal accuracy better than 35 m s $^{-1}$ and differ in zero point by \sim 53 m s $^{-1}$.

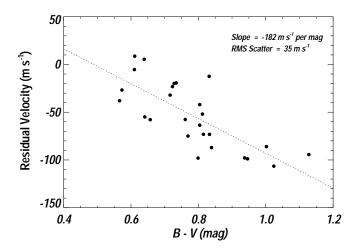


Fig. 4.—Standard-star velocities (Udry et al. 1999a) minus the present velocities as a function of B-V. A dependence is apparent, suggesting systematic errors in at least one set of velocities. The slope is -182 ± 24 m s⁻¹ per mag. The rms scatter is 35 m s⁻¹ before fitting and 23 m s⁻¹ after fitting a line to the data.

rms scatter around the fit of 109 m s⁻¹. This color dependence is likely caused by some spectral type mismatch in one or all of the radial-velocity scales. For solar type stars, we are in good agreement with the CfA velocities, with no offset as seen in Figure 6. This is likely due to the CfA velocity zero point being set by observations of minor planets as was also done for the present velocities.

Interestingly, the sign of the slope in Figure 6 is opposite of that in Figure 4 between the CORAVEL and present velocities. Since synthetic spectra were used to derive both the CfA and ELODIE velocities, which were used as the reference system in Udry et al. (1999a), (Stefanik et al. 1999; Baranne et al. 1996; Gullberg 1999) this seems to give credence to the notion that using synthetic spectra does not entirely solve the problem of spectral dependent systematic errors.

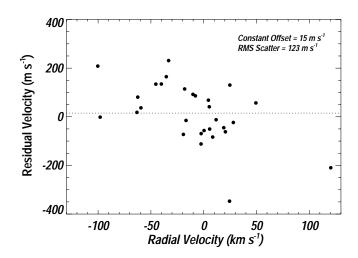


Fig. 5.—Standard-star velocities (Stefanik et al. 1999) minus present velocities for all 29 FGK stars in common. The differences reveal that the present velocities are lower than those of Stefanik et al. by 15 m s⁻¹ and exhibit an rms scatter of 123 m s⁻¹. Thus, the present velocities and those of Stefanik et al. (1999) differ in zero point by $\sim\!15$ m s⁻¹.

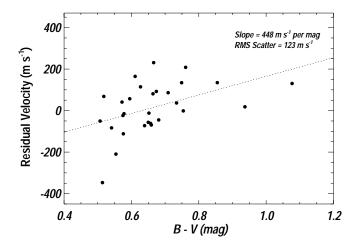


Fig. 6.—Standard-star velocities (Stefanik et al. 1999) minus the present velocities as a function of B-V. A dependence is apparent, suggesting systematic errors in at least one set of velocities. The slope is $+448 \pm 111$ m s⁻¹ per mag. The rms scatter is 123 m s⁻¹ before fitting and 109 m s⁻¹ after fitting a line to the data.

6.2. M Stars

The radial-velocity standard stars listed by Udry et al. (1999a) do not include any M dwarfs, and we do not have any M dwarfs in common with Stefanik et al. (1999) or the older CORAVEL standard stars (Udry et al. 1999b). Therefore, we compared our present velocities for M dwarfs to those given in Marcy et al. (1987). The results of this comparison, for 21 stars in common, are shown in Figures 7 and 8 and yield $\langle V_{\rm MLW} - V_{\rm Present} \rangle = -21 \, \rm m \, s^{-1}$ with an rms scatter of 164 m s⁻¹. The offset is obviously very low since the M star reference spectrum was created using the velocities quoted by Marcy et al. (1987). Since the average internal error for the velocities in Marcy et al. (1987) is \sim 200 m s⁻¹, most of the scatter in the differences is due to them. No B-Vdependence is seen in the residuals. It is therefore difficult to ascertain the uncertainty in our present velocities for M dwarfs and similarly difficult to ascertain a zero-point error (whatever that would mean). But a conservative estimate of the errors would be the ~ 0.4 km s⁻¹ uncertainty of the

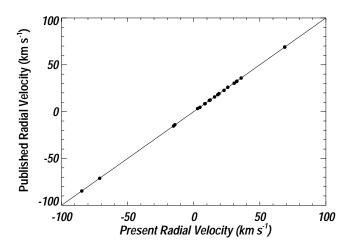


Fig. 7.—Velocities of standard stars (Marcy et al. 1987) vs. present velocities for M dwarfs. The velocities agree well with no apparent nonlinear dependence.

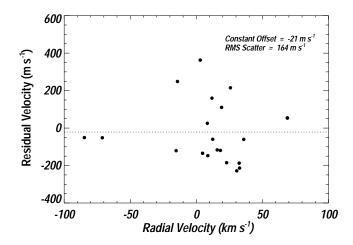


Fig. 8.—Difference between standard stars velocities (Marcy et al. 1987) and present velocities for all 21 M stars in common (as in Fig. 7). There is an rms scatter of 164 m s⁻¹, and a constant offset of -21 m s⁻¹, which is not statistically significant. Thus, our velocities for the stars having B-V>1.3 reside on the velocity scale set by Marcy et al. (1987). No B-V dependence is seen in the residuals.

Marcy et al. (1987) velocities which were used in setting the velocity zero point.

Since the observed offset (21 m s⁻¹) for the M dwarfs is within a factor of 2 of the internal scatter of \sim 10 m s⁻¹, we have not applied any correction to the radial velocities for these stars. There is some concern that our two sets of stars, the F, G, and K stars and the M stars, might not be on the same velocity zero point. A comparison between them is difficult due to the significant spectral type mismatch errors.

7. UNCERTAINTY ESTIMATES

Due to the several systematic errors affecting radial velocities on the order of $0.1~\rm km~s^{-1}$ it is difficult to ascertain the true uncertainties of the velocities reported here. Normally, two different methods are used to estimate uncertainties: (1) Standard deviation of the mean (i.e., the internal rms scatter of points), and (2) comparison with published values. We have done both here. On average the standard deviation of the mean, due to the internal scatter of points, for the velocities in Table 1 is $\sim 10~\rm m~s^{-1}$. We also compared our velocities with the best known published standard star velocities of Udry et al. (1999a). This comparison, as shown in the previous section, gave an rms scatter of the differences of 35 m s⁻¹. Therefore, using the traditional methods of estimating uncertainties our velocities are accurate to $\sim 35~\rm m~s^{-1}$.

In this case the traditional methods fail due to the astrophysical sources of errors that affect all spectroscopic measurements of radial velocity. The values of these errors are not known adequately, otherwise we would have corrected for them. It is likely that these systematic errors were also not taken into account by Udry et al. (1999a) or Stefanik et al. (1999). This means that a comparison between their data and ours will not yield the true uncertainty of our velocities or theirs

What the comparison does show is the precision of velocities within a given spectral type. The velocities of all stars of a given spectral type will have a nearly constant offset from their true kinematic velocities, because the systematic errors are dependent on spectral type. Relative to that constant offset the velocities in that spectral type are very precise. The

comparison with Udry et al. (1999a) shows our precision to be no worse than 0.035 km s^{-1} (see Fig. 3).

This high precision within a spectral type can be effectively used to look for moving groups. Moving groups have velocity dispersions of typically $\sim 0.5 \, \mathrm{km \ s^{-1}}$ such as the Pleiades group (Jones 1970). The present velocities are precise enough within a spectral type to judge whether a star belongs to the moving group or not.

Even though the exact values of the systematic errors are not known we can estimate our true uncertainties. There are three major systematic errors in our velocities, namely, convective blueshift, gravitational redshift, and spectral type mismatch. These errors change systematically with spectral type. Since our velocity zero point was set here using the day sky and the minor planet Vesta, which have well known radial velocities due to solar system dynamics, the systematic errors were forced to zero for solar-type stars. The errors will rise as the spectral type departs from solar type. Because the zero-point calibration was used for all FGKtype stars there will be differential errors due to convective blueshift and gravitational redshift that increase with departure from solar type. Similarly, the spectral type mismatch errors only occur for non-solar-type stars, since the NSO solar spectrum was used for the reference, and is therefore unaffected by the velocity zero-point calibration.

The approximate values of the systematic errors are as follows. According to Dravins (1999) the convective blueshift is approximately $-1000~\rm m~s^{-1}$ for F5 V (B-V=0.4), $-400~\rm m~s^{-1}$ for G2 V (solar-type, B-V=0.64), and $-200~\rm m~s^{-1}$ for K0 V (B-V=0.9). The effective velocity error caused by gravitational redshift can be computed from, $V_{\rm grav}=GM/Rc$ (Dravins et al. 1999). We find, with masses and radii given by Allen (2000), that the redshift is +680 m s^{-1} for F5 V, +636 m s^{-1} for G2 V, and +590 m s^{-1} for K0 V. The sum of the two effects shows that the overall systematic error is approximately $-320~\rm m~s^{-1}$ for F5 V, +236 m s^{-1} for G2 V, and +390 m s^{-1} for K0 V. Due to our zero-point calibration these errors are here forced to zero for solar-type stars. Therefore, we expect that our velocities are low by $\sim\!556~\rm m~s^{-1}$ for F5 V stars, true for solar-type stars, and high by $\sim\!154~\rm m~s^{-1}$ for K0 V stars.

With these estimates of the systematic errors in hand we can compute a correction for the velocities as a function of B-V. We fit a line to the first two points (F5 V and G2 V) and to the second two points (G2 V and K0 V) to obtain two linear interpolations. The first equation (eq. [2]) is for main-sequence stars earlier than solar type (B-V < 0.64), and the second equation (eq. [3]) is for main-sequence stars later than solar type (B-V > 0.64), not including M type stars. Equation (3) may be used for stars with 1.3 > B-V > 0.9, but represents an extrapolation of the points given above and should be used with caution. Since most of our main-sequence stars have B-V < 1.1 this should not cause too many problems. These equations do not account for spectral type mismatch errors because we do not know enough about these effects yet for our velocities:

$$RV_{kin} = RV_{spec} - 2.317 \times (B - V) + 1.483 \text{ km s}^{-1},$$

 $B - V < 0.64;$ (2)

$$RV_{kin} = RV_{spec} - 0.642 \times (B - V) + 0.411 \text{ km s}^{-1} ,$$

$$0.64 < B - V < 1.3 . \tag{3}$$

These corrections are valid only for main-sequence stars for which the NSO solar spectrum was used as the reference. The velocities for the M type stars, with B-V>1.3 and for which the M composite spectrum was used as the reference, have a zero-point uncertainty of ~ 0.4 km s⁻¹, and therefore a correction for gravitational redshift or convective blueshift is not warranted at this time. However, we expect the velocities of the M stars to be very precise due to their low rms velocity scatter and their small spectral type range which minimizes the systematic errors. We expect them to also be precise to 0.03 km s⁻¹, just as the FGK main-sequence stars. If the "true" radial velocity of only one of them were known then they all could be corrected for their zero-point error and be very accurate.

We expect that the corrections given in equations (2) and (3) account for convective blueshift and gravitational redshift to within \sim 0.1 km s⁻¹. The errors due to spectral type mismatch are likely to be \sim 0.1 km s⁻¹ (Griffin et al. 2000). From the errors expressed in equations (2) and (3), an additional \sim 0.1 km s⁻¹ for the spectral type mismatch error, along with the distribution of spectral types of our survey stars, the typical uncertainty of the uncorrected radial velocities in Tables 1 and 2 is \sim 0.3 km s⁻¹. For the few subgiants, the errors are somewhat larger, but not easily estimated without models of subphotospheric convection in such stars.

We cannot at this time give a correction for the velocities of our 45 giants. Giants have a gravitational redshift on the order of 0.1 km s⁻¹, which is much smaller than that for dwarfs due to the large radius of giants, \sim 15 R_{\odot} (Allen 2000). The convective blueshift is not known for giants and hinders us from giving a rough velocity correction. Hopefully, more in-depth future studies of the sources of error discussed here, such as recent work by Pourbaix et al. (2002), will allow for more accurate corrections of radial velocities and eventually yield "true" radial velocities. Until that time these corrections may be used for the present velocities.

8. FINAL RADIAL VELOCITIES AND DESCRIPTION OF TABLES

The barycentric radial velocities for all 889 stars are reported in Tables 1 and 2. The 782 stars that exhibit an rms velocity scatter less than 100 m s⁻¹ are reported in Table 1. Primary and alternate star names are given in the first two columns, and the stellar spectrum used as the template (either NSO or M dwarf composite) is listed in column (3). The mean time of the observations is given in column (4) under $\langle JD \rangle$ to establish the characteristic epoch for the velocity measurement. The span of observations in days is given in column (5), and the mean barycentric radial velocity of all observations for a star is in column (6). Stars with only one observation were put in Table 1 even though their rms scatter is not defined. They can be distinguished by $\Delta T = 0$.

The 107 stars with an rms velocity scatter greater than 100 m s⁻¹ are reported in Table 2. Primary and alternate star names are given in the first two columns. The stellar template spectrum (either NSO or M dwarf) is given in column (3). The Julian Date (JD) of one specific observation is given in column (4). The barycentric radial velocity of that one observation is given in column (5). The mean date of all observations and span of observations are given in columns

TABLE 3
ORBITAL PARAMETERS

Star	P (days)	$K (km s^{-1})$	в	ω (deg)	$T_0 = (-2450000)$	$M_1 \ (M_{\odot})$	$M_{2,\mathrm{min}} \ (M_{\mathrm{Jup}})$	a_{\min} (AU)	$f(M)$ (M_{\odot})	Other References
HD 4747	6832 (653)	0.65(0.1)	0.64 (0.06)	257 (5)	453 (473)	0.83	42.3	6.7	0.000087	
HD 7483	701.42 (0.01)		0.12 (0.003)	313 (0.5)	1791.1(0.7)	0.92	136	1.6	0.0020	
HD 30339	15.0778 (0.0003)	5.94 (0.01)	0.25(0.001)	43(1)	1881.199 (0.0003)	1.10	77.8	0.13	0.00030	
HD 34101	803.51 (0.03)	3.59 (0.01)	0.08(0.001)	275 (1)	(2) (2)	0.87	167	1.7	0.0038	
HD 39587	5136 (12)	1.85(0.02)	0.45(0.01)	111 (1)	1463 (17)	0.89	143	5.9	0.0024	
HD 65430	3138 (342)	1.11 (0.2)	0.32(0.02)	77 (1)	3267 (302)	0.78	8.79	4.0		
HD 122742	3617 (7)	6.41(0.01)	0.48(0.001)	183(0.1)	2030(3)	96.0	545	5.3		1, 2, 3, 4, 5
HD 131511	125.396 (0.001)	19.10 (0.01)	0.51(0.001)	219(0.1)	203.407 (0.004)	0.78	455	0.52		6,7
HD 140913	147.968 (0.001)	1.94(0.01)	0.54 (fixed)	18(1)	1321.42 (0.02)	0.98	43.2	0.55		8, 9, 10, 11
HD 174457	840.80 (0.05)	1.25(0.01)	0.23(0.01)	139(1)	2020(4)	1.19	65.8	1.9	0.00016	
HD 208776	2624 (371)	5.46 (1.3)	0.27(0.04)	245 (12)	(100)	1.24	511	4.2	0.0400	
GJ 84	6818 (2491)	2.18 (0.2)	0.44(0.08)	238 (10)	1777 (41)	0.39	115	5.6	0.0054	
GJ 595	62.6277 (0.0001)	6.57(0.01)	0.26(0.001)	253(0.1)	1999.50(0.02)	0.28	0.09	0.22	0.0017	
GJ 623	1366.1 (0.4)	2.08 (0.04)	0.67(0.01)	251 (1)	1298 (10)	0.31	42.0	1.7	0.00053	12, 13, 14, 15, 16, 17
HIP 52940	1393.5 (0.2)	2.15 (0.01)	0.37 (0.003)	199(1)	1541(2)	1.12	126	2.6	0.0011	

REFERENCES.—(1) Duquennoy & Mayor 1991; (2) Martin et al. 1998; (3) Blazit et al. 1987; (4) Kamper 1987; (5) Wagman 1949; (6) Beavers & Salzer 1983; (7) Kamper & Lyons 1981; (8) Halbwachs et al. 2000; (9) Mazeh et al. 1996; (10) Oetiker et al. 2001; (11) Mayor et al. 1997; (12) McCarthy 1986; (13) Barbieri et al. 1996; (14) Henry & McCarthy 1993; (15) Marcy & Moore 1989; (16) McCarthy & Henry 1987; (17) Lippincott & Worth 1978.

(6) and (7). The rms scatter of the velocities of all observations is given in column (8), and the number of observations in column (9). The last column is for comments where stars with companion orbits or linear trends are noted. "L" indicates that velocities vary linearly with time (see Table 4), "CO" indicates that a companion and its orbit were found, "C" indicates that a published companion exists but an orbit is not given here, and "A" indicates that the star is chromospherically active based on the emission reversal seen at the Ca II H and K lines in our spectra.

The orbital parameters for 15 stars with companions are reported in Table 3. The orbital period P, velocity semiamplitude K, eccentricity e, longitude of periastron ω , and time of periastron T_0 are given as well as the primary mass M_1 , minimum secondary mass $M_{2,\min}$, minimum semimajor axis a_{\min} , and the mass function, f(M). References are also given to other sources which have information on these stars, their companions, and orbital parameters. The data for stars with linear trends are given in Table 4. The slope of the radial velocity curve and the number of observations is given for 30 stars.

9. ORBITS OF BINARIES

We found 107 stars in this program that exhibit an rms velocity scatter greater than 100 m s⁻¹. We attempted to fit these with a Keplerian orbit. We found 29 stars that yield

TABLE 4
STARS WITH LINEAR TRENDS

Primary Name	Alternate Name	Slope $(m s^{-1} day^{-1})$	Number of Observations
HD 3770	HIP 3169	0.41 (0.02)	13
HD 3795	HIP 3185	0.363 (0.003)	24
HD 11909	HIP 9110	23.7 (0.8)	4
HD 14802	HIP 11072	-1.73(0.02)	7
HD 17037	HIP 12764	-1.5(0.1)	14
HD 21017	HIP 15861	-11.3(0.3)	5
HD 28388	HIP 20802	-2.7(0.3)	6
HD 29461	HIP 21654	0.58 (0.02)	9
HD 30649 ^a	HIP 22596	0.27 (0.02)	15
HD 31412	HIP 22919	0.414 (0.009)	12
HD 50639	HIP 33109	-0.51(0.01)	8
HD 72780	HIP 42112	-0.91(0.03)	10
HD 73668	HIP 42488	-0.69(0.02)	24
HD 86680 ^b	HIP 49060	112(7)	8
HD 101563	HIP 57001	-3.22(0.07)	5
HD 102540	HIP 57574	-19.9(0.5)	5
HD 103829	HIP 58318	-2.28(0.07)	8
HD 129814	HIP 72043	-0.224(0.004)	13
HD 136580	HIP 75039	-0.47(0.02)	18
HD 142229	HIP 77810	-1.47(0.06)	9
HD 161797	HIP 86974	-0.147(0.006)	42
HD 167665	HIP 89620	-0.59(0.02)	16
HD 186704 ^b	HIP 97255	-88(8)	4
HD 190406	HIP 98819	-0.066(0.001)	58
HD 190771	HIP 98921	0.30 (0.02)	13
HD 200565	HIP 103983	-0.54(0.01)	24
HD 209779	HIP 109110	-10.7(0.1)	6
HD 215578	SAO 108160	-2.51(0.03)	21
HD 219420	HIP 114834	-135.5(0.5)	5
HD 223084	HIP 117258	-0.85(0.04)	33

^a Exhibits some additional positive curvature.

good Keplerian orbital fits to their relative radial velocities. Many of these binaries have been previously published from our velocities and therefore will not be duplicated here. Several papers contain these previously reported single-line spectroscopic binaries, namely Marcy et al. (1999), Butler et al. (2000), Vogt et al. (2002), Fischer et al. (2002), Cumming et al. (1999), and Marcy et al. (2001b).

For 15 stars, our velocities provide unpublished orbital solutions or reveal unknown companions. These orbits are listed in Table 3 which gives the usual orbital parameters for single-line spectroscopic binaries. Plots of the velocities and the associated Keplerian fits are shown in Figures 9–23. The mass for the primary star for each system was estimated from the catalog by Prieto & Lambert (1999) or by using B-V and Allen (2000). From the orbital parameters and primary masses, we determined the minimum companion masses (M_{\min}) , mass functions, f(M), and minimum semimajor axes, a_{\min} , which are also listed in Table 3. The values of $M_{\rm min}$ range from 42 $M_{\rm Jup}$ to 545 $M_{\rm Jup}$, and therefore the companions are all candidate brown dwarfs or H-burning stellar companions. The uncertainties in the orbital parameters were found by a Monte Carlo technique in which Gaussian velocity noise was added to the best-fit theoretical velocity curve at the times of observation. We ran 50 trials for each case, with orbital parameters being rederived for each trial. The standard deviations of the resulting orbital parameters were taken as the uncertainties. For some stars the Monte Carlo method underestimated the uncertainties in the orbital parameters. For these stars a more conservative estimate of the uncertainties was made by fitting many different Keplerian orbits to the observed relative radial velocities, without any modeled Gaussian noise, and then looking at the scatter of the parameters produced by the best orbital fits.

The stars HD 4747, HD 65430, and GJ 84 have large orbital uncertainties due to incomplete phase coverage. There were insufficient velocity measurements to constrain the orbit of HD 18445 even though it is known to have a companion (Halbwachs et al. 2000). Thus, it is not listed in Table 3. The eccentricity for HD 140913 was fixed to e=0.54 (Latham et al. 1989) in performing the Keplerian fit, since not enough points were available to constrain the eccentricity. Only the other four parameters

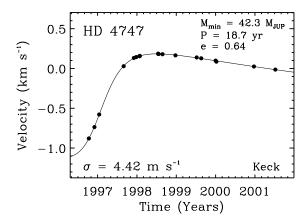


Fig. 9.—Doppler velocities for HD 4747 (G8/K0 V). The solid line is a Keplerian orbital fit with a period of 18.7 yr, a semiamplitude of 0.65 km s⁻¹, and an eccentricity of 0.64, yielding a minimum ($M_{\rm min}$) of 42.3 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 4.42 m s⁻¹.

^b Exhibits some additional negative curvature.

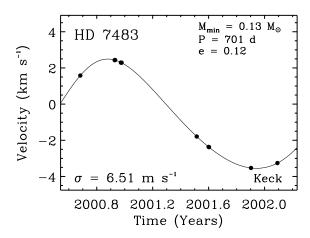


Fig. 10.—Doppler velocities for HD 7483 (G5 V). The solid line is a Keplerian orbital fit with a period of 701 days, a semiamplitude of 3.02 km s⁻¹, and an eccentricity of 0.12, yielding a minimum ($M_{\rm min}$) of 0.13 M_{\odot} for the companion. The rms of the Keplerian fit is 6.51 m s⁻¹.

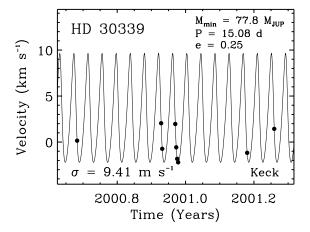


Fig. 11.—Doppler velocities for HD 30339 (F8 V). The solid line is a Keplerian orbital fit with a period of 15.08 days, a semiamplitude of 5.94 km s⁻¹, and an eccentricity of 0.25, yielding a minimum ($M_{\rm min}$) of 77.8 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 9.41 m s⁻¹.

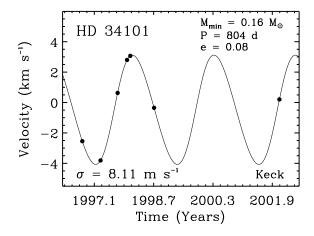


Fig. 12.—Doppler velocities for HD 34101 (G8 V). The solid line is a Keplerian orbital fit with a period of 804 days, a semiamplitude of 3.76 km s⁻¹, and an eccentricity of 0.08, yielding a minimum ($M_{\rm min}$) of 0.16 M_{\odot} for the companion. The rms of the Keplerian fit is 8.11 m s⁻¹.

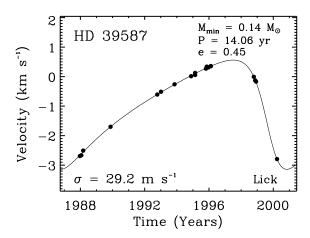


Fig. 13.—Doppler velocities for HD 39587 (G0 V). The solid line is a Keplerian orbital fit with a period of 14.06 yr, a semiamplitude of 1.85 km s⁻¹, and an eccentricity of 0.45, yielding a minimum ($M_{\rm min}$) of 0.14 M_{\odot} for the companion. The rms of the Keplerian fit is 29.2 m s⁻¹.

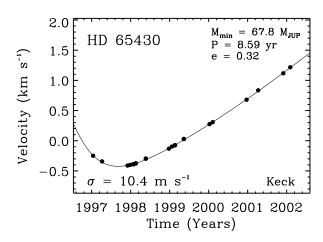


Fig. 14.—Doppler velocities for HD 65430 (K0 V). The solid line is a Keplerian orbital fit with a period of 8.59 yr, a semiamplitude of 1.11 km s⁻¹, and an eccentricity of 0.32, yielding a minimum ($M_{\rm min}$) of 67.8 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 10.4 m s⁻¹.

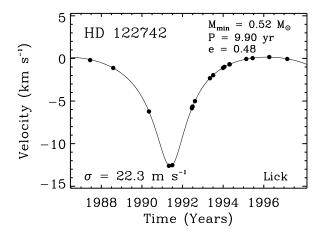


Fig. 15.—Doppler velocities for HD 122742 (G8 V). The solid line is a Keplerian orbital fit with a period of 9.90 yr, a semiamplitude of 6.41 km s⁻¹, and an eccentricity of 0.48, yielding a minimum ($M_{\rm min}$) of 0.52 M_{\odot} for the companion. The rms of the Keplerian fit is 22.3 m s⁻¹.

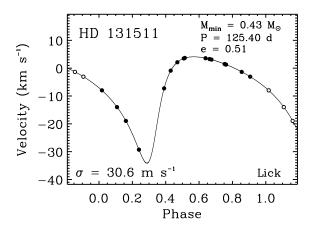


Fig. 16.—Doppler velocities for HD 131511 (K2 V). The solid line is a Keplerian orbital fit with a period of 125.40 days, a semiamplitude of 19.10 km s⁻¹, and an eccentricity of 0.51, yielding a minimum ($M_{\rm min}$) of 0.43 M_{\odot} for the companion. The rms of the Keplerian fit is 30.6 m s⁻¹.

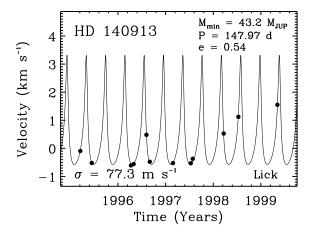


Fig. 17.—Doppler velocities for HD 140913 (G0 V). The solid line is a Keplerian orbital fit with a period of 147.97 days, a semiamplitude of 1.94 km s⁻¹, yielding a minimum ($M_{\rm min}$) of 43.2 $M_{\rm JUP}$ for the companion. The eccentricity was fixed to e=0.54 given by Latham et al. (1989) since not enough points were available to constrain the eccentricity. The rms of the Keplerian fit is 77.3 m s⁻¹.

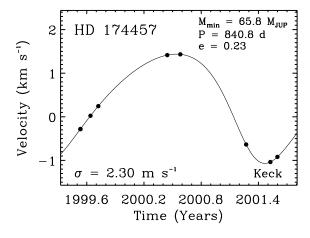


Fig. 18.—Doppler velocities for HD 174457 (F8 V). The solid line is a Keplerian orbital fit with a period of 840.8 days, a semiamplitude of 1.25 km s⁻¹, and an eccentricity of 0.23, yielding a minimum ($M_{\rm min}$) of 65.8 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 2.30 m s⁻¹.

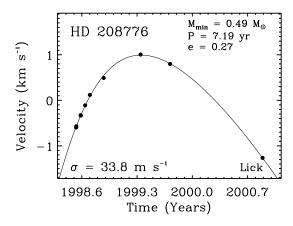


Fig. 19.—Doppler velocities for HD 208776 (G0 V). The solid line is a Keplerian orbital fit with a period of 7.19 yr, a semiamplitude of 5.46 km s⁻¹, and an eccentricity of 0.27, yielding a minimum ($M_{\rm min}$) of 0.49 M_{\odot} for the companion. The rms of the Keplerian fit is 33.8 m s⁻¹.

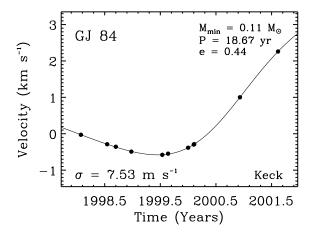


Fig. 20.—Doppler velocities for GJ 84 (M3 V). The solid line is a Keplerian orbital fit with a period of 18.67 yr, a semiamplitude of 2.18 km s⁻¹, and an eccentricity of 0.44, yielding a minimum ($M_{\rm min}$) of 0.11 M_{\odot} for the companion. The rms of the Keplerian fit is 7.53 m s⁻¹.

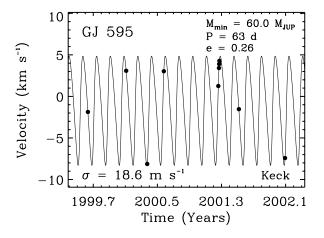


Fig. 21.—Doppler velocities for GJ 595 (M3 V). The solid line is a Keplerian orbital fit with a period of 63 days, a semiamplitude of 6.57 km s⁻¹, and an eccentricity of 0.26, yielding a minimum ($M_{\rm min}$) of 60.0 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 18.6 m s⁻¹.

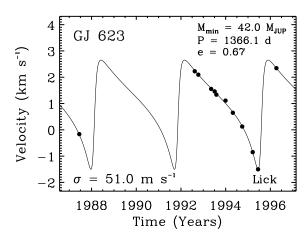


Fig. 22.—Doppler velocities for GJ 623 (M3 V). The solid line is a Keplerian orbital fit with a period of 1366.1 days, a semiamplitude of 2.08 km s⁻¹, and an eccentricity of 0.67, yielding a minimum ($M_{\rm min}$) of 42.0 $M_{\rm JUP}$ for the companion. The rms of the Keplerian fit is 51.0 m s⁻¹.

were fitted for this star. The Keplerian fit for HD 208776 is particularly poor as insufficient velocities are available to constrain the orbital period to better than a factor of 2. Eleven of the spectroscopic binary stars appear to be newly discovered here: HD 4747, HD 7483, HD 30339, HD 34101, HD 39587, HD 65430, HD 174457, HD 208776, GJ 84, GJ 595, HIP 52940. Velocities are available upon request of G. M. The companions all have $M_{\rm min}$ in the substellar range or low-mass stellar range and thus offer interesting targets for studies with adaptive optics or interferometry.

10. CONCLUSION

We have provided barycentric radial velocities with an internal precision of 0.03 km s⁻¹ for 889 stars. The error estimates stem both from the internal errors found from the spectral chunks within each spectrum and from the comparison with accurate velocities on the CORAVEL scale (Udry et al. 1999a). The radial velocities of the F, G, and K dwarfs reside on a velocity zero point defined by the observations of the Sun, using the day sky and Vesta as proxies. Our velocity scale differs by only 0.053 km s⁻¹ from that of Udry et al. (1999a) and $0.015~\rm km~s^{-1}$ from that of Stefanik et al. (1999), thus adding confidence to the zero points of all three sets of velocities. The radial velocities of the M dwarfs reside on the velocity system defined by Marcy et al. (1987) and have not been further corrected, nor is such a correction known to be necessary. These M dwarf velocities are probably accurate to within $200 \,\mathrm{m \, s^{-1}}$.

The Doppler shifts reported here have such high accuracy that gravitational redshift and convective blueshift impose comparable (or greater) wavelength shifts. These effects were somewhat removed from our velocity measurements by using the Sun for the velocity zero point. We expect therefore that for G2 V stars the present velocities represent their "true" kinematic velocities within 0.03 km s $^{-1}$. However, for stars departing from solar type the sum of the two astrophysical effects will produce systematic errors dependent on spectral type. From F to K type dwarfs this variation will be as large as $\sim\!0.3$ km s $^{-1}$ and will cause our velocities to be low for F type stars and high for K type stars. Using estimates for these effects we give rough velocity corrections

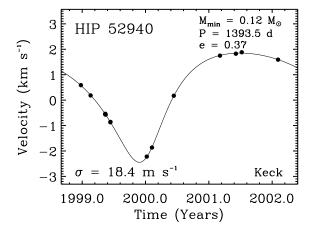


Fig. 23.—Doppler velocities for HIP 52940 (F8 V). The solid line is a Keplerian orbital fit with a period of 1393.5 days, a semiamplitude of 2.15 km s⁻¹, and an eccentricity of 0.37, yielding a minimum ($M_{\rm min}$) of 0.12 M_{\odot} for the companion. The rms of the Keplerian fit is 18.4 m s⁻¹.

in equations (2) and (3). We presume that these corrections bring the velocities within ~ 0.15 km s⁻¹ of their "true" kinematic values.

These precise radial velocities can be used to complement future highly precise proper motion measurement, such as those projected to be obtained by the GAIA mission of the ESA. The radial velocities and proper motions will give the three components of space motion of stars. These precise space motions may be useful for discerning the membership of moving groups, since young moving groups have velocity dispersions of $\sim 0.5 \, \mathrm{km \, s^{-1}}$ (Jones 1970).

The 782 stars listed in Table 1 exhibited a velocity scatter of less than $100 \, \mathrm{m \, s^{-1}}$ during 4 years. These stars apparently exhibit relatively stable velocities on timescales of a decade and represent candidates for radial velocity standard stars. However, their integrity as velocity standard stars requires future observations to verify their stability. We expect that some of these 782 "stable" stars may reveal slow drifts in radial velocity on timescales longer than $10 \, \mathrm{yr}$.

The accuracy of the present velocities offers an opportunity to detect such slow drifts by future measurements made with comparable accuracy. Such velocity variations may prove useful in identifying unseen companions at large orbital distances, i.e., over 10 AU. We found that 107 stars exhibited velocity variations of over 100 m s⁻¹ (rms). For these stars, we have provided the rms velocity, and also either an orbital solution or a description of the linear trends. We intend these measurements to provide dynamical constraints on the nature of the companions.

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