

1 Relationship between I/I_{TOF} and $S(\mathbf{Q})$

The expression for the scattered intensity in the direction \mathbf{k}' by a collection of atoms is given by

$$I = \left(\frac{\hbar c k \Gamma}{r_D^2} \frac{9}{24\pi} \Lambda \right) \frac{I_p/I_{\text{sat}}}{4\Delta^2 + 2I_p/I_{\text{sat}}} \left[\frac{e^{-2W} 4\Delta^2}{(4\Delta^2 + 2I_p/I_{\text{sat}})} \left(4 \sum_{mn} e^{i\mathbf{Q}(\mathbf{R}_m - \mathbf{R}_n)} S_{zm} S_{zn} - N \right) + N \right] \quad (1)$$

Here, the momentum transfer is $\mathbf{Q} = \mathbf{k}' - \mathbf{k}$, where \mathbf{k} is the wave vector of the probe light. The probe light has intensity I_p , and I_{sat} is the saturation intensity of the optical transition. This expression is obtained for a mixture of atoms in states $|1\rangle$ and $|2\rangle$, for light with a detuning $\Delta_1 = -\Delta_2$, in the limit $|\Delta_{1,2}| \gg 1$ (the detuning is in units of the linewidth of the transition). No approximations have been made regarding the relative magnitude of $|\Delta_{1,2}|$ and I_p/I_{sat} . The Debye-Waller factor is e^{-2W} , and the letter Λ denotes a sum over output polarizations

$$\Lambda = \sum_{\boldsymbol{\varepsilon}'} |(\boldsymbol{\varepsilon}_+ \cdot \boldsymbol{\varepsilon}')(\boldsymbol{\varepsilon}_- \cdot \boldsymbol{\varepsilon}')|^2 \quad (2)$$

The two states satisfy

$$\hat{S}_z|1\rangle = +\frac{1}{2}|1\rangle \quad \hat{S}_z|2\rangle = -\frac{1}{2}|2\rangle \quad (3)$$

The indices m, n denote different atoms, we identify the sum with the spin structure factor at momentum \mathbf{Q} :

$$S(\mathbf{Q}) = \frac{4}{N} \sum_{mn} e^{i\mathbf{Q}(\mathbf{R}_m - \mathbf{R}_n)} S_{zm} S_{zn} \quad (4)$$

We consider a measurement of the intensity after some time-of-flight (TOF). After TOF, the Debye-Waller factor goes to zero due to the expanding size of the atomic wavefunction, so we have

$$I_{\text{TOF}} = \left(\frac{\hbar c k \Gamma}{r_D^2} \frac{9}{24\pi} \Lambda \right) \frac{I_p/I_{\text{sat}}}{4\Delta^2 + 2I_p/I_{\text{sat}}} N \quad (5)$$

$$\frac{I}{I_{\text{TOF}}} = \frac{e^{-2W} 4\Delta^2}{(4\Delta^2 + 2I_p/I_{\text{sat}})} (S(\mathbf{Q}) - 1) + 1 \quad (6)$$

We note that doublons do not contribute to $S(\mathbf{Q})$. This can be easily seen from

$$S(\mathbf{Q}) = \frac{4}{N} \sum_m e^{i\mathbf{Q}\mathbf{R}_m} S_{zm} \sum_n e^{-i\mathbf{Q}\mathbf{R}_n} S_{zn} \quad (7)$$

A doublon pair contributes $e^{i\mathbf{Q}\mathbf{R}_D} (+1/2 - 1/2) = 0$ to each sum. If we think about a sample with only doublons, according to the equation above for I/I_{TOF} we have

$$\frac{I_D}{I_{\text{TOF}}} = 1 - \frac{e^{-2W} 4\Delta^2}{(4\Delta^2 + 2I_p/I_{\text{sat}})} \quad (8)$$

Lately we have been performing scattered intensity measurements where we magneto-associate doublons into molecules. The molecules do not scatter the probe light, so they do not contribute at all to the scattered intensity. Under this protocol, for a sample with only doublons we would have

$$\frac{I_D}{I_{\text{TOF}}} = 0 \quad (9)$$

which is consistent with scattering in the limit of an infinitely deep lattice ($e^{-2W} \rightarrow 1$) and for a weak probe ($\Delta^2 \gg I_p/I_{\text{sat}}$), where we have the simple equivalence

$$\frac{I}{I_{\text{TOF}}} = S(\mathbf{Q}) \quad (10)$$

The magneto-association protocol thus makes the doublons contribute to the scattering in an ideal manner, whereas the atoms in singly occupied sites still suffer from Debye-Waller and saturation effects. This suggests an expression like the following

$$\frac{I}{I_{\text{TOF}}} = [c(1 - D) + D] (S(\mathbf{Q}) - 1) + 1 \quad (11)$$

Where D is the fraction of atoms in doubly occupied sites, and the correction factor (Debye-Waller plus saturation corrections) for a sample with only singly occupied sites is

$$c = \frac{e^{-2W} 4\Delta^2}{4\Delta^2 + 2I_p/I_{\text{sat}}} \quad (12)$$

The fraction of atoms in doubly occupied sites can be obtain from TOF measurements with and without magneto-association as

$$D = 1 - I_{\text{TOFAssoc}}/I_{\text{TOF}} \quad (13)$$

2 Latest data at $5.5 E_R$

Lately we have been taking data with three shots:

1. In-situ with magneto-association, I
2. TOF without magneto-association, I_{TOF}
3. TOF with magneto-association, I_{TOFAssoc}

The data is shown at the end of this document. For this data we have calculated $S(\mathbf{Q})$ using the two different TOF measurements:

$$S_{\text{Assoc}}(\mathbf{Q}) = \frac{1}{c} \left(\frac{I}{I_{\text{TOFAssoc}}} - 1 \right) + 1 \quad (14)$$

$$S(\mathbf{Q}) = \frac{1}{c} \left(\frac{I}{I_{\text{TOF}}} - 1 \right) + 1 \quad (15)$$

where, again

$$c = \frac{e^{-2W} 4\Delta^2}{4\Delta^2 + 2I_p/I_{\text{sat}}} \quad (16)$$

We perform in-situ measurements in a $50E_R$ lattice and for $I_p/I_{\text{sat}} \approx 15$. In units of the transition linewidth, $\Delta = 6.5$ for our case. With this values we have

$$\begin{aligned} e^{-2W} &= 0.81 \\ 1 + \frac{I_p/I_{\text{sat}}}{2\Delta^2} &= 1.18 \end{aligned} \quad (17)$$

As usual the data has measurements at two values of \mathbf{Q} , the one labeled 2 corresponds to the $(1/2, 1/2, 1/2)$ direction, and the other one, labeled 1, corresponds to the intensity measured with our other camera.

We still have not applied a correction such as the one suggested in Eq. 11 because we have not yet reached consensus on the validity of this equation. In fact I am just suggesting it in this pdf to see what everybody's opinions are on this matter.

An interesting feature of this data is that, with the magneto-association, we get better agreement with $S(Q_1)$ as calculated by Thereza. In the data that we took at $7E_R$ we did not have much agreement with $S(Q_1)$.

The corrected value of $S(Q_2)$ is 1.4, which would be consistent with a temperature T/t of around 1.1.

We have also plotted $S(Q_1)/S(Q_2)$ for this data. In the absence of antiferromagnetic correlations this should go to 1.

