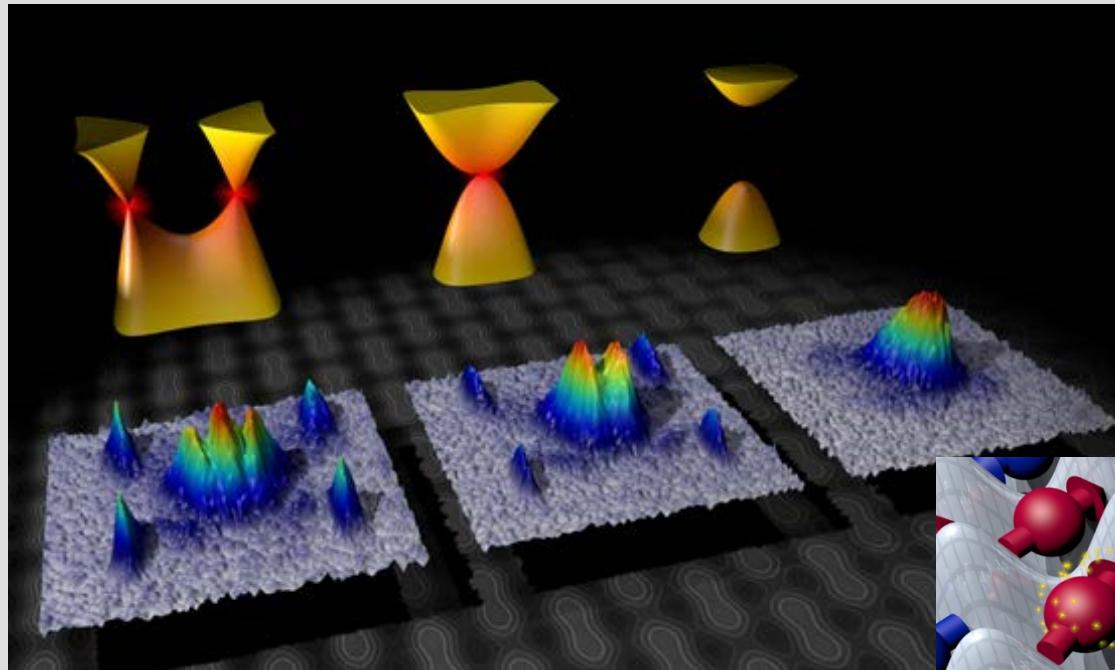
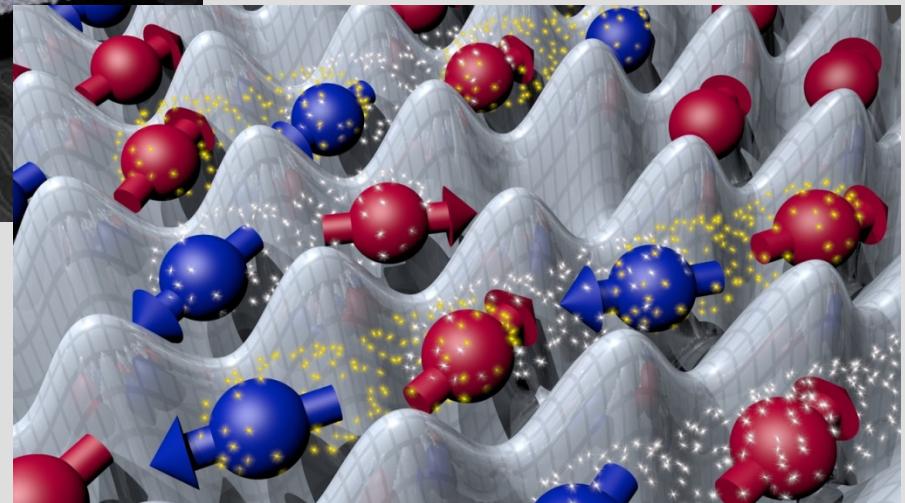


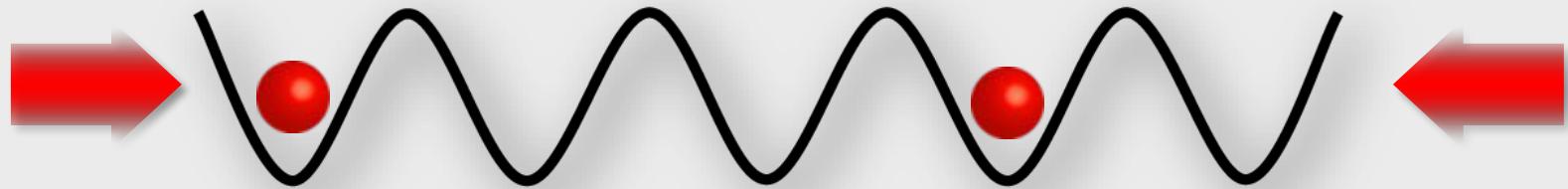
Ultracold fermions in tunable-geometry optical lattices



Leticia Tarruell



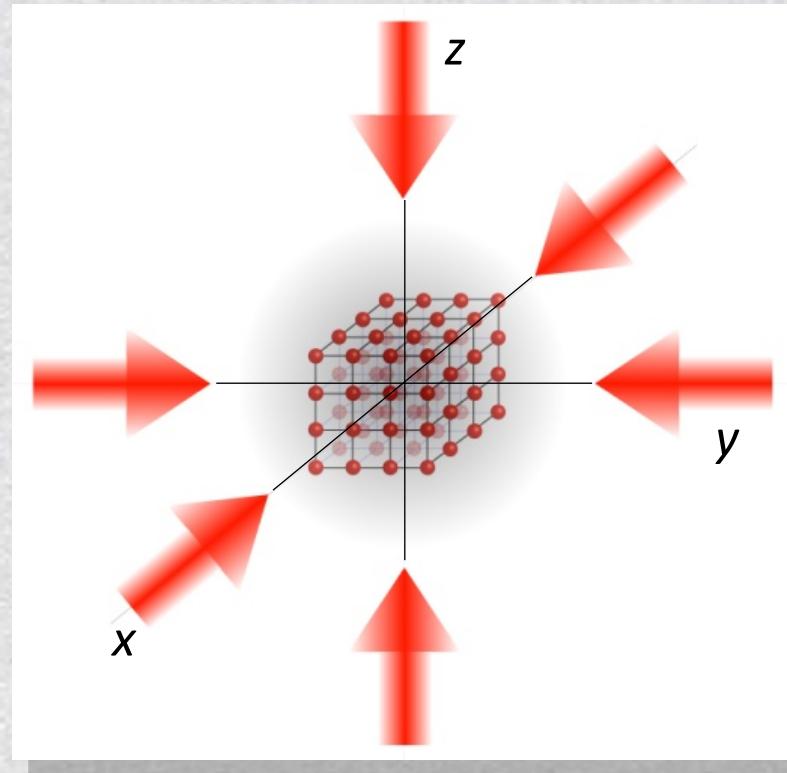
Ultracold fermions in optical lattices

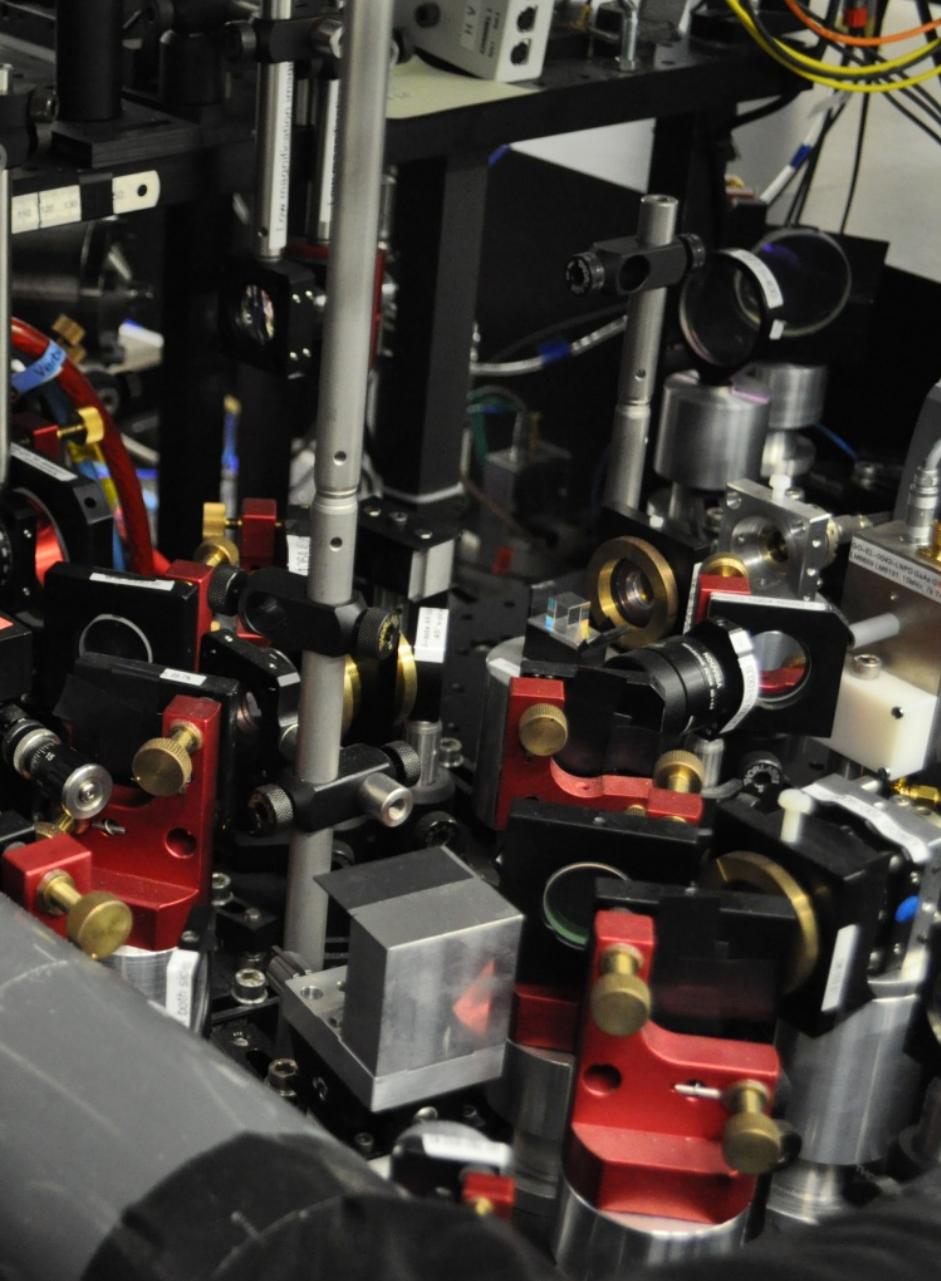
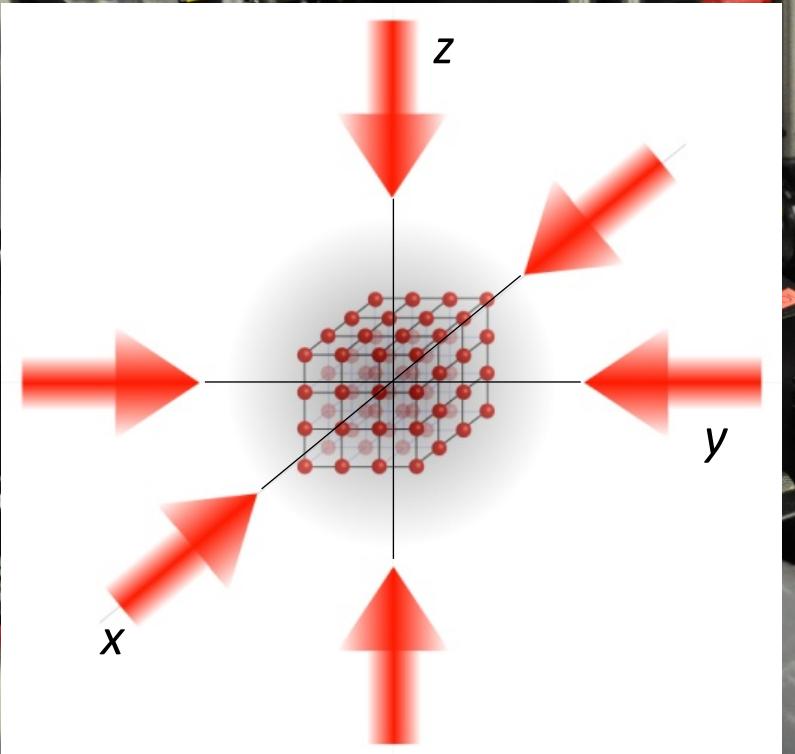


- Ultracold fermionics atoms (^{40}K)
- Interfering laser beams (optical lattice)

*Fermions in optical lattices
as condensed matter model systems*

Ultracold fermions in optical lattices





Also fermions in 3D optical lattices: Munich, Boston, Hamburg, Kyoto, Houston, Florence, Amsterdam, Bonn...

Condensed matter model systems

*Ultracold fermions
in optical lattices*

Strongly correlated materials

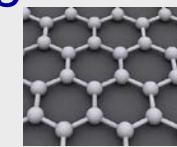
(Mott insulators,
high- T_c superconductors)



Spin systems
(quantum phase transitions,
frustration, spin-liquids)



Novel materials
(graphene,
topological insulators)



ETH approach: tunable-geometry optical lattice

Outline

An optical lattice of tunable geometry

Engineering Dirac points in a honeycomb lattice

Short-range quantum magnetism

Outline

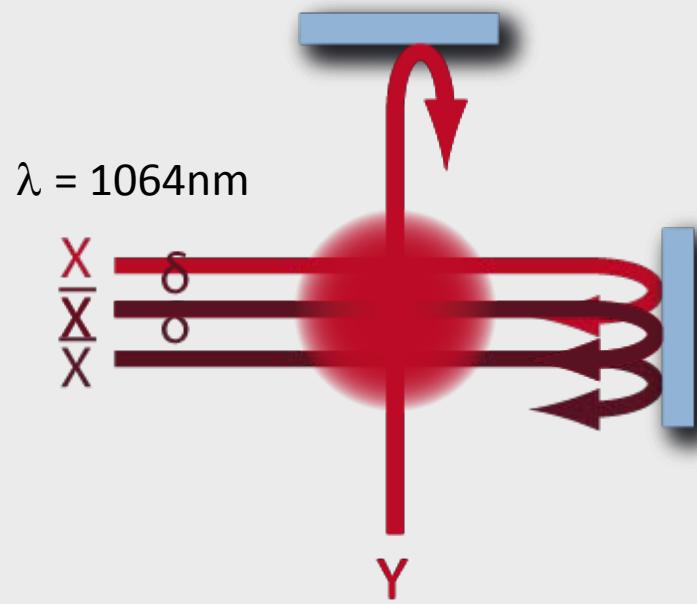
An optical lattice of tunable geometry

Engineering Dirac points in a honeycomb lattice

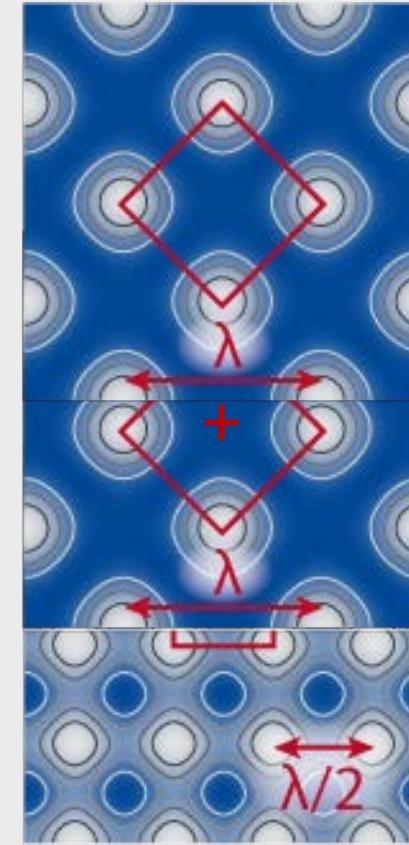
Short-range quantum magnetism

The tunable-geometry optical lattice

Setup



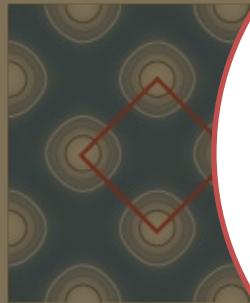
Optical potential



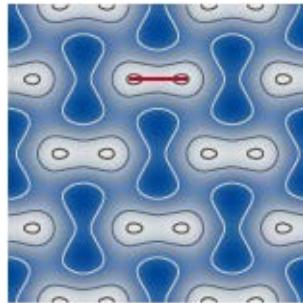
$$V(x, y) = V_{\bar{X}} \cos^2(kx + \theta/2) + V_X \cos^2(kx) + V_Y \cos^2(ky) + 2\alpha \sqrt{V_X V_Y} \cos(kx) \cos(ky)$$

The tunable-geometry optical lattice

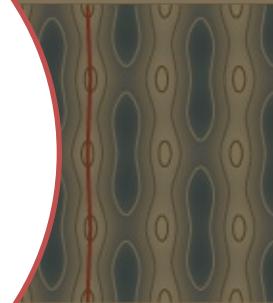
Chequerboard



Dimer



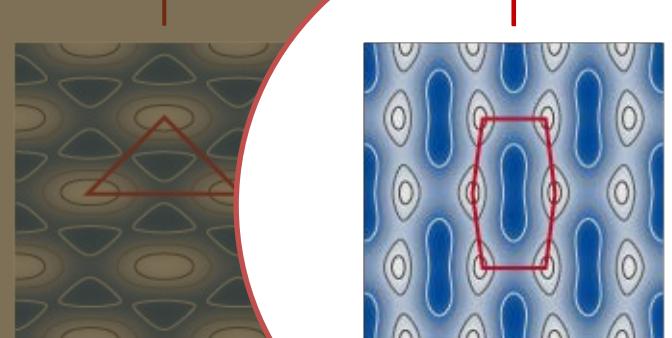
1D chains



$V_{\bar{X}} [E_R]$



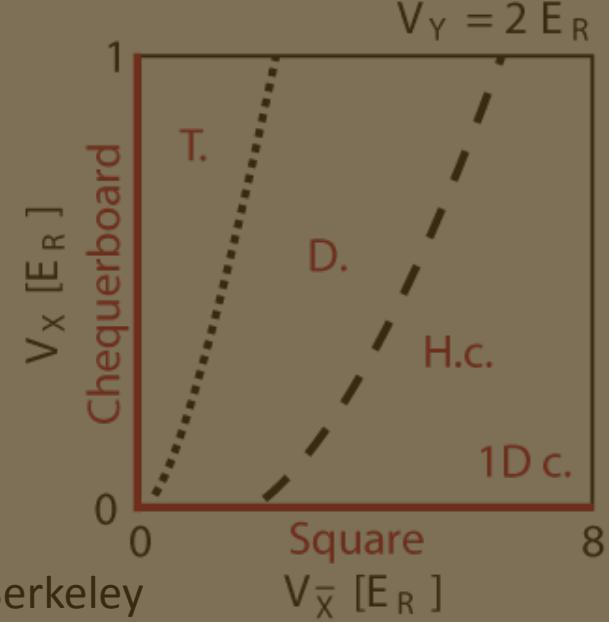
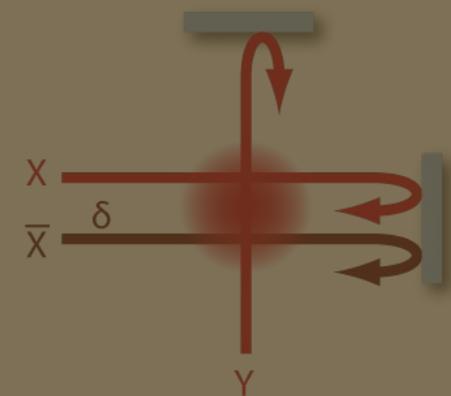
Triangular



Honeycomb

$V_x = 0$

Square



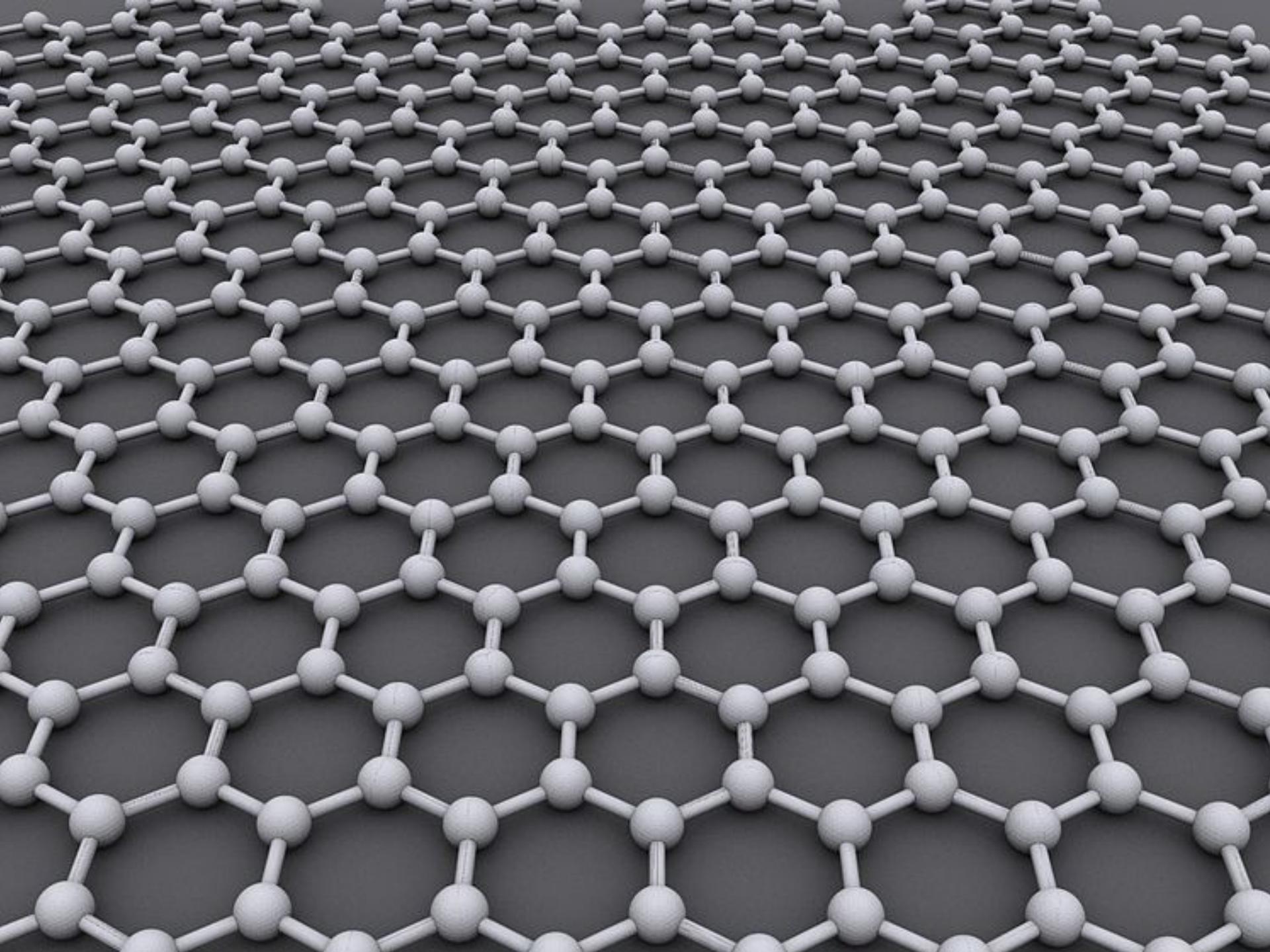
Other non-standard lattices: Niels Bohr, Aachen, Bonn, Hamburg, Berkeley

Outline

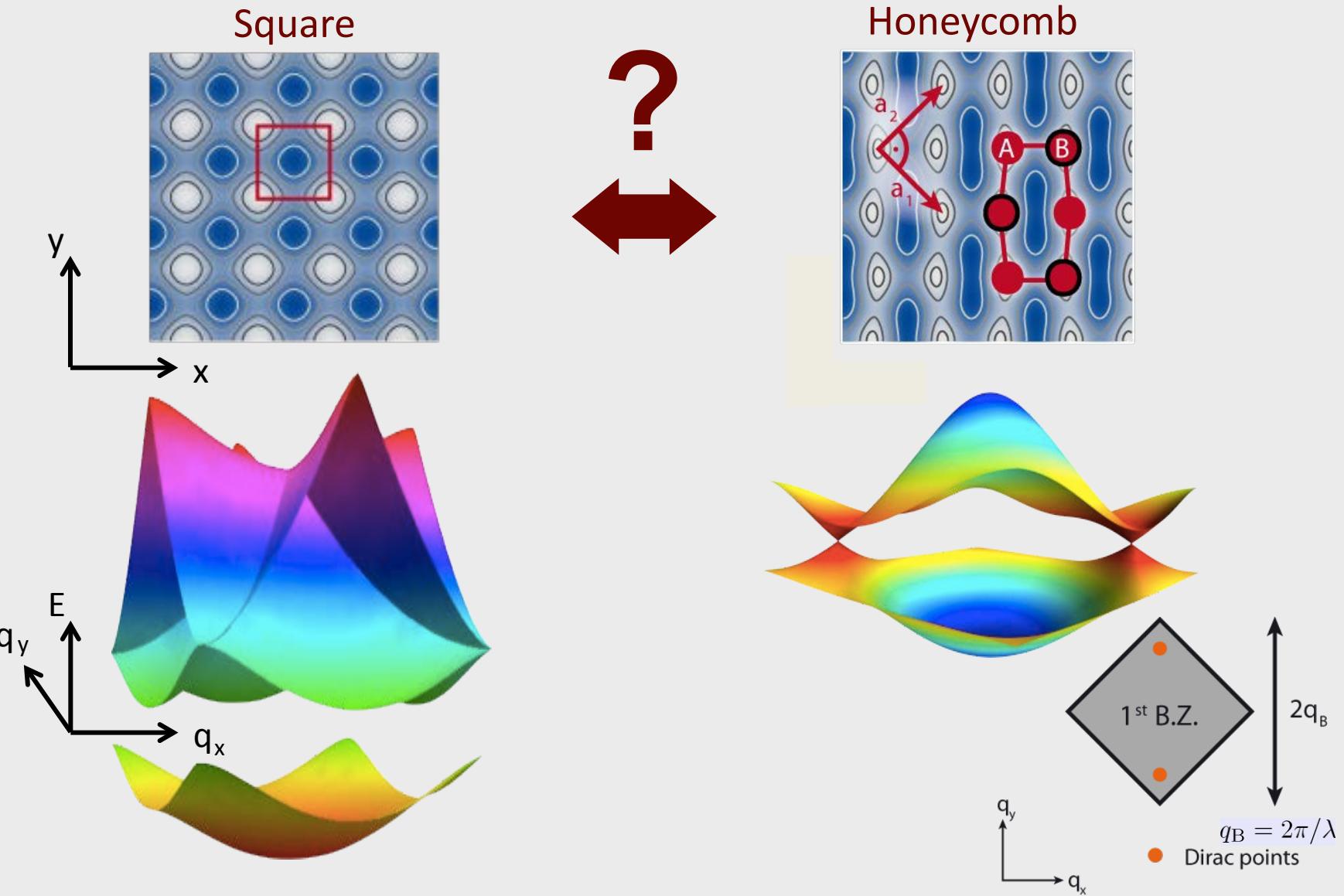
An optical lattice of tunable geometry

Engineering Dirac points in a honeycomb lattice

Short-range quantum magnetism



Deforming the band structure



Probing the Dirac points

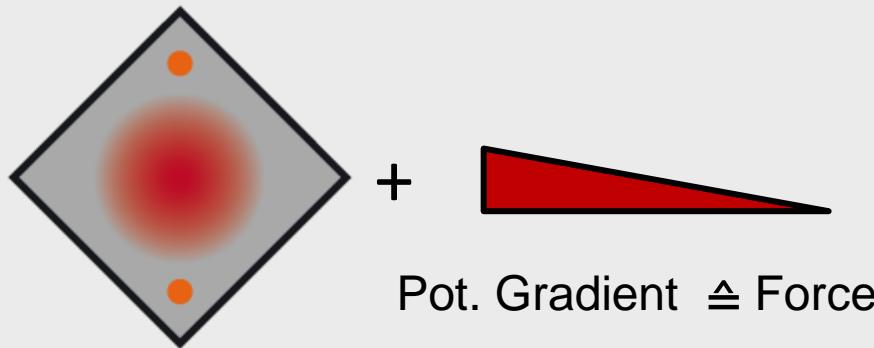
Challenges: vanishing density of states
small energy scales



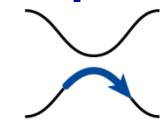
Probe energy splitting of
the bands **dynamically**

T. Salger *et al.*, Phys. Rev. Lett. 99, 190405 (2007)

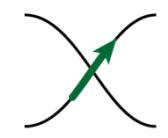
Bloch oscillations + interband transitions



Passing away from Dirac point:
stay in lowest band



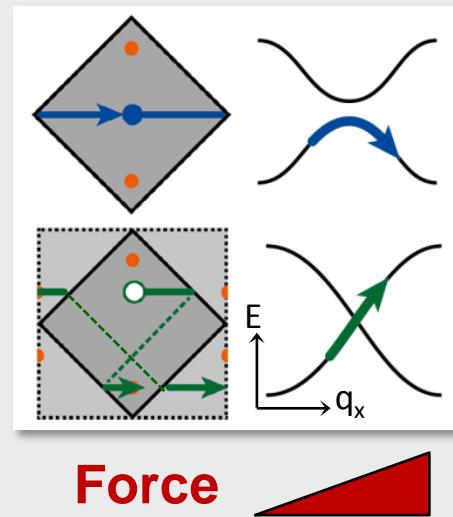
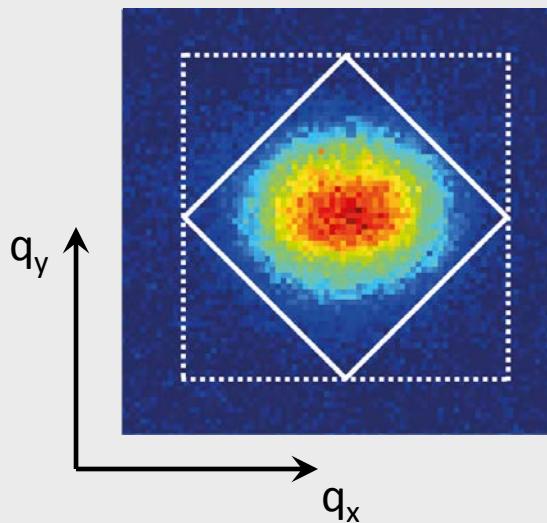
Passing through Dirac point:
transfer to 2nd band



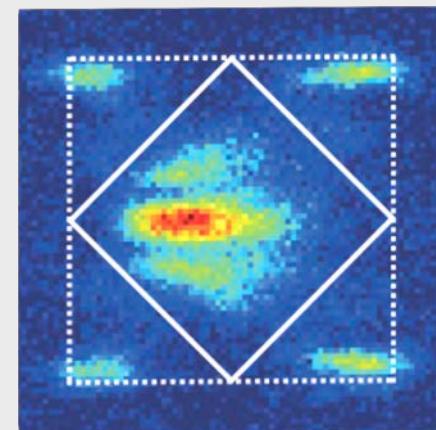
Observable: quasi-momentum distribution

Interband transitions

Spin-polarized ^{40}K gas



After a Bloch cycle: $t=T_B$



Transfer to 2nd band at the position of the Dirac points

Quantitative measurements

Higher band fraction

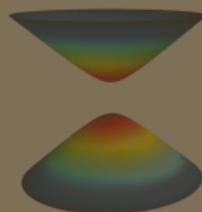
A diagram showing a diamond-shaped region divided into four quadrants. The top-right quadrant is shaded green, representing the occupied band. The other three quadrants are shaded grey, representing the empty band. The text $\xi = \frac{N(\text{green})}{N(\text{grey}) + N(\text{green})}$ is displayed, where N represents the number of atoms in each quadrant.

$$\xi = \frac{N(\text{green})}{N(\text{grey}) + N(\text{green})}$$

Engineering Dirac points

Tunability

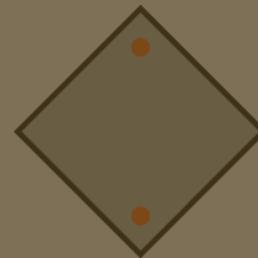
Gap



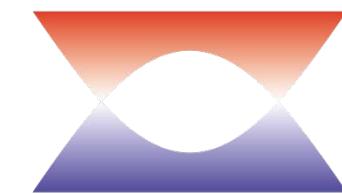
Cone shape



Position

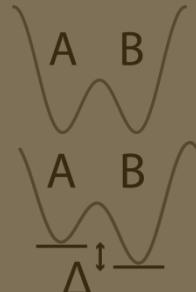
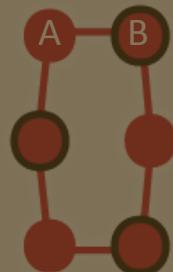


Merging

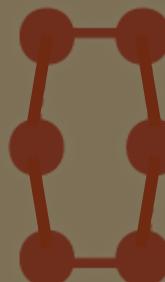


Tools

Inversion symmetry



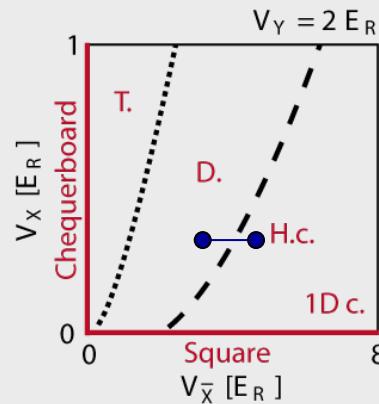
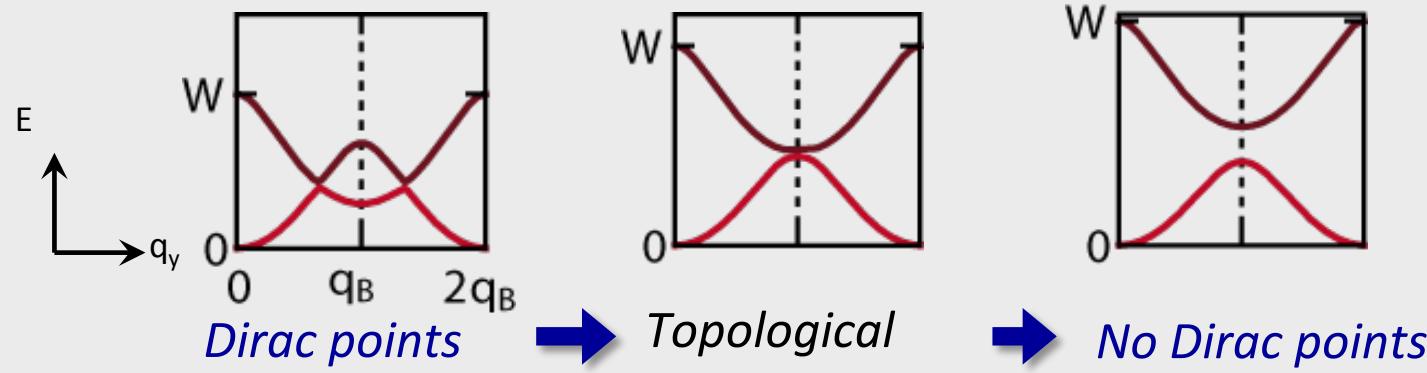
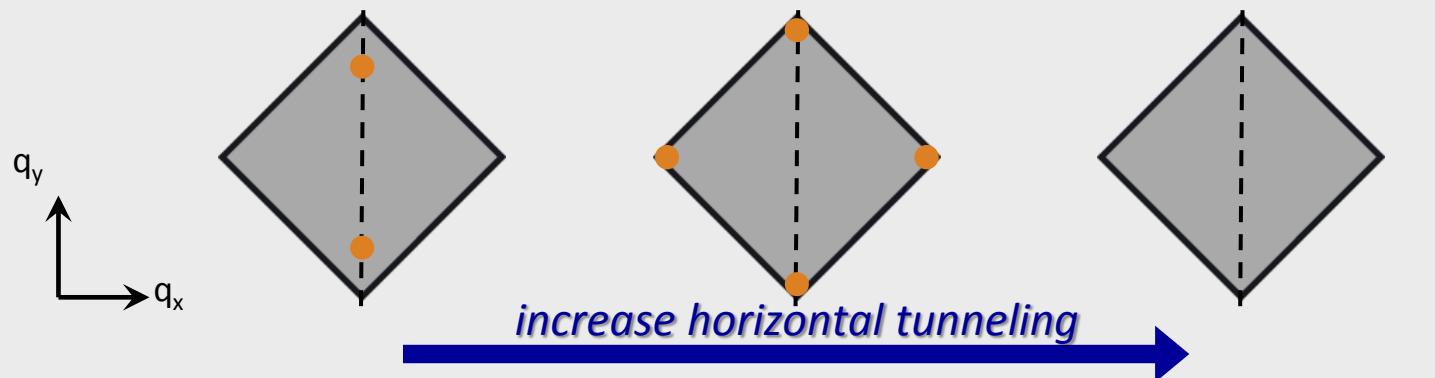
Tunneling imbalance



vs.

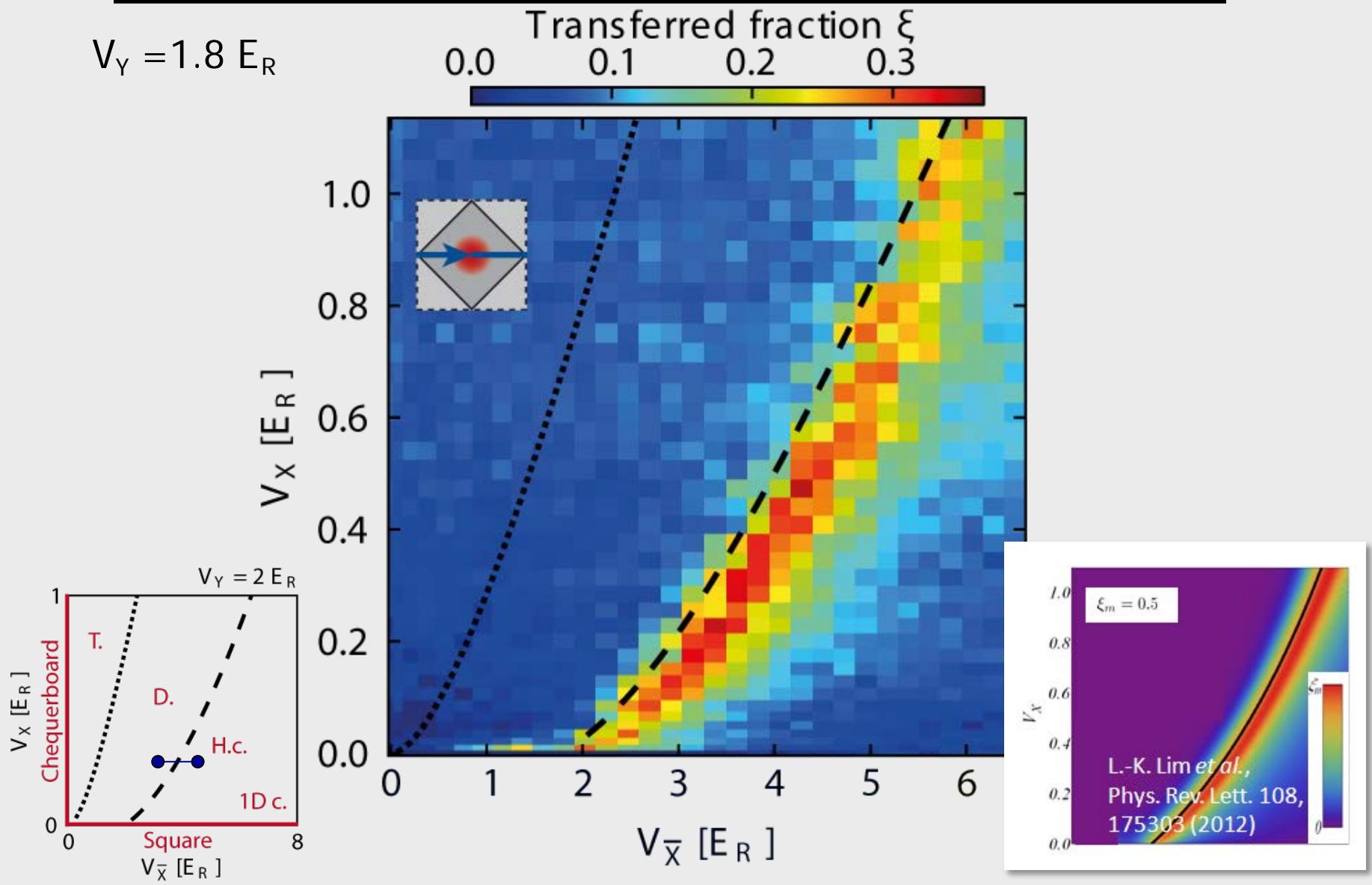


Merging Dirac points

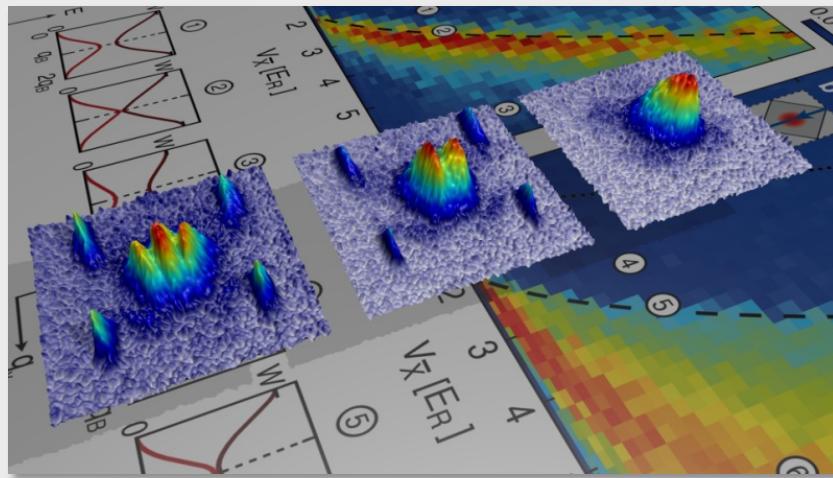


Lifshitz transition, Sov. Phys. JETP **11**, 1130 (1960)

The topological transition



Engineering Dirac points



L. Tarruell, D. Greif, T. Uehlinger, G. Jotzu, and T. Esslinger, *Nature* **483**, 302 (2012)

T. Uehlinger, D. Greif, G. Jotzu, L. Tarruell, T. Esslinger, L. Wang, and M. Troyer, *EPJ ST* **217**, 121 (2013)

Combination of honeycomb lattice and interactions (Mott insulator)

T. Uehlinger, G. Jotzu, M. Messer, D. Greif, U. Bissbort, W. Hoffstetter, and T. Esslinger,
Phys. Rev. Lett. **111**, 185307 (2013).

Lattice shaking to engineer the Haldane model

Outline

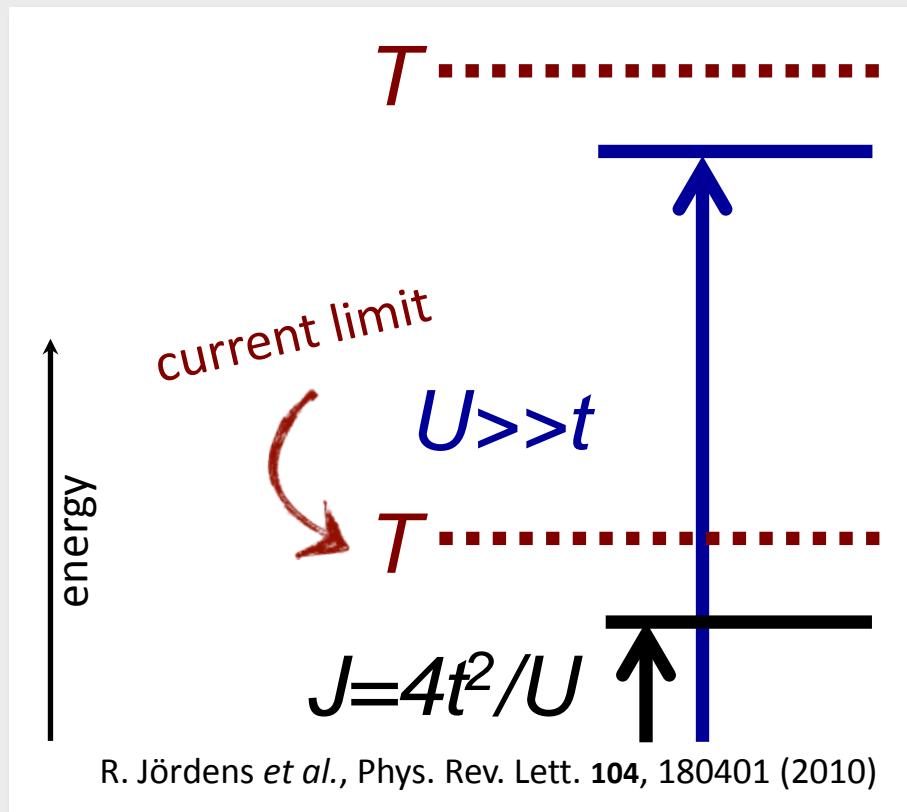
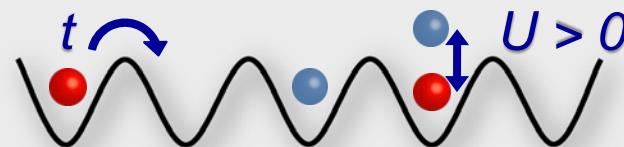
An optical lattice of tunable geometry

Engineering Dirac points in a honeycomb lattice

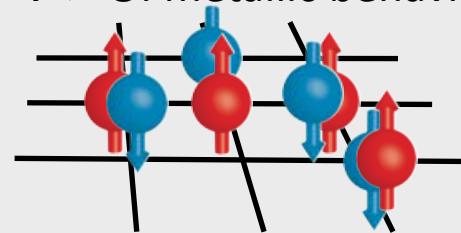
Short-range quantum magnetism

Magnetism: a temperature challenge

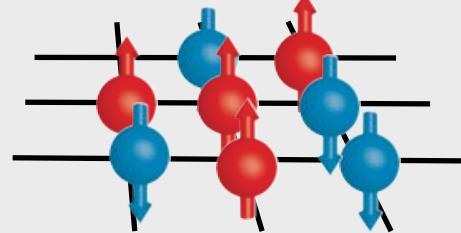
Fermi-Hubbard model



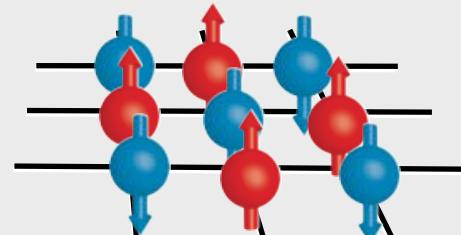
$T > U$: metallic behaviour



$T < U$: Mott insulator

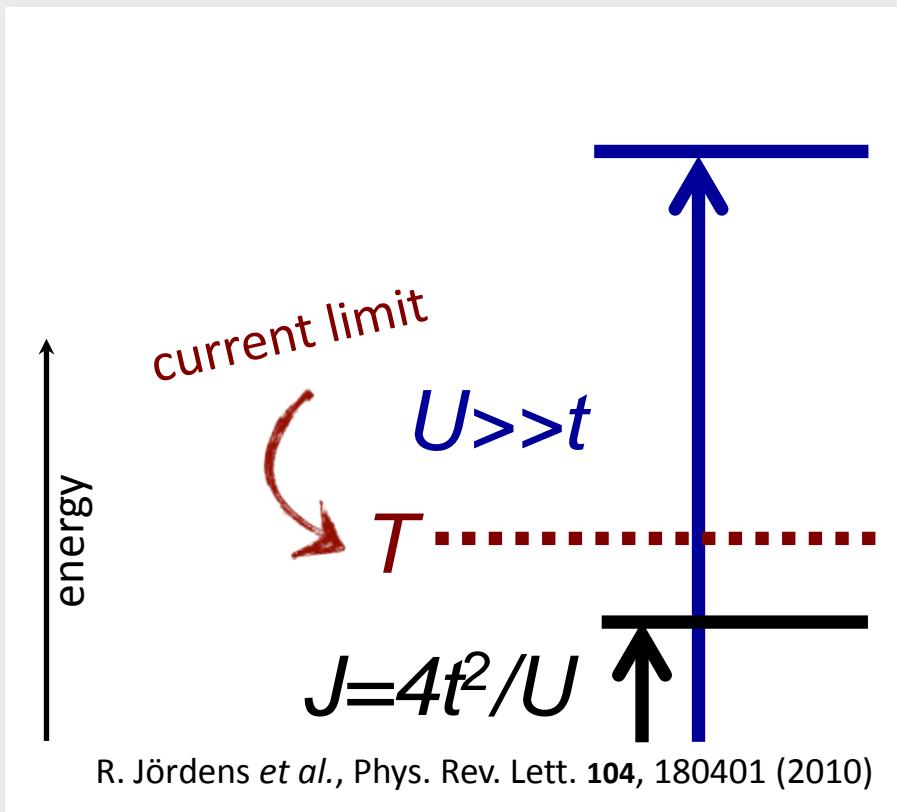
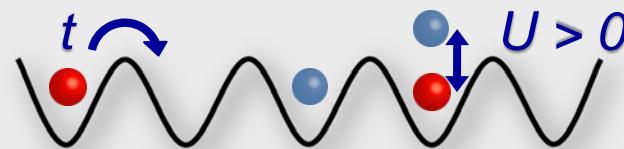


$T < J$: spin ordering



Magnetism: a temperature challenge

Fermi-Hubbard model



State of the art

Isolated double-wells or plaquettes (Munich)

S. Trotzky *et al.*, Science **319**, 295-299 (2008)
S. Nascimbène *et al.*, Phys. Rev. Lett. **108**, 205301 (2012)

Mappings

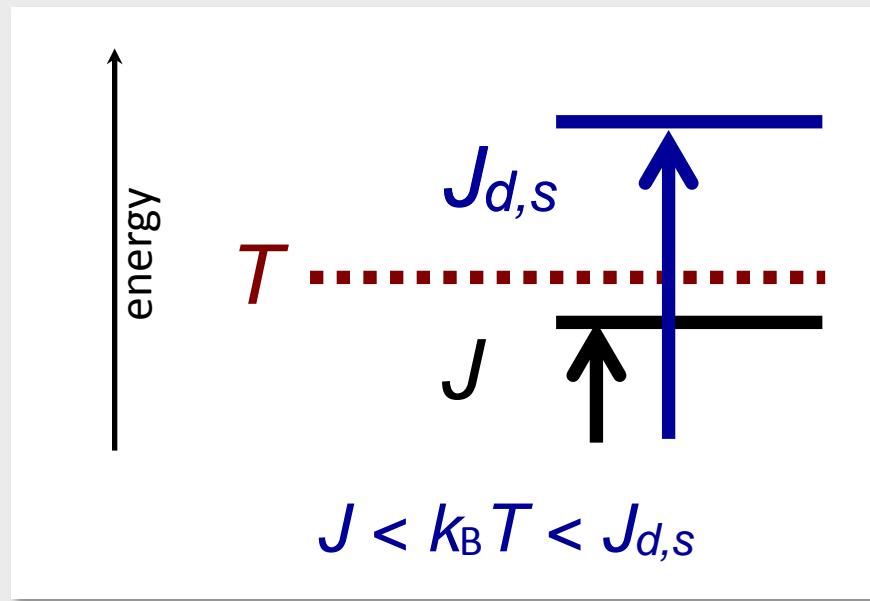
Ising spin chain (Harvard)
Classical magnets (Hamburg)
J. Simon *et al.*, Nature **472**, 307-312 (2011)
J. Struck *et al.*, Science **333**, 996-999 (2011)

Dipolar interactions (JILA, Paris)

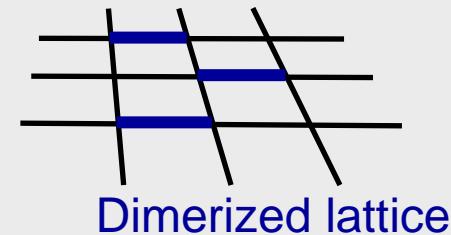
B. Yan *et al.*, Nature **501**, 521-525 (2013)
A. de Paz *et al.*, Phys. Rev. Lett. **111**, 185305 (2013)

Reaching magnetism: an energy trick

Reaching magnetic correlations:

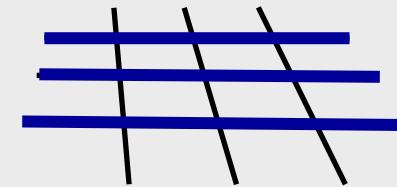


$$t_d \rightarrow J_d > J \quad t \rightarrow J$$



Dimerized lattice

$$t_s \rightarrow J_s > J$$



Anisotropic lattice

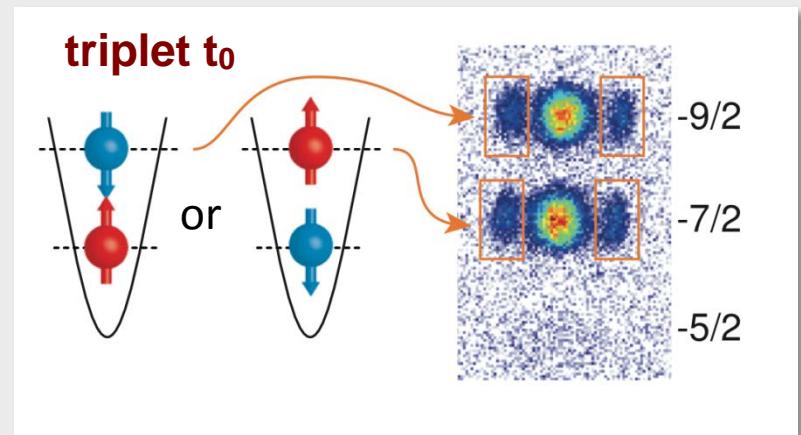
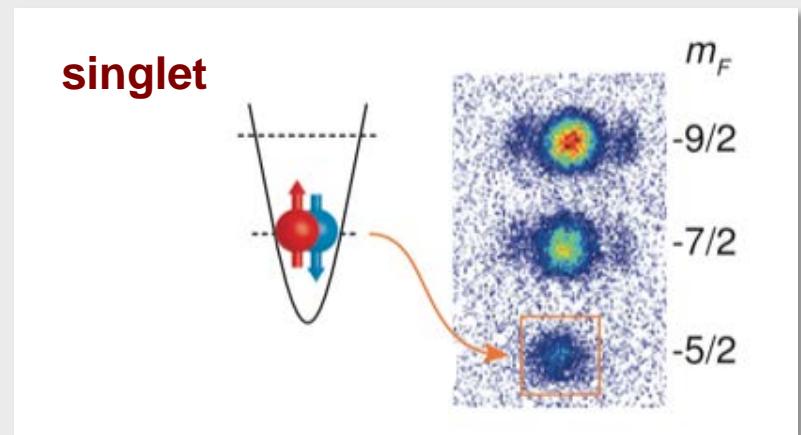
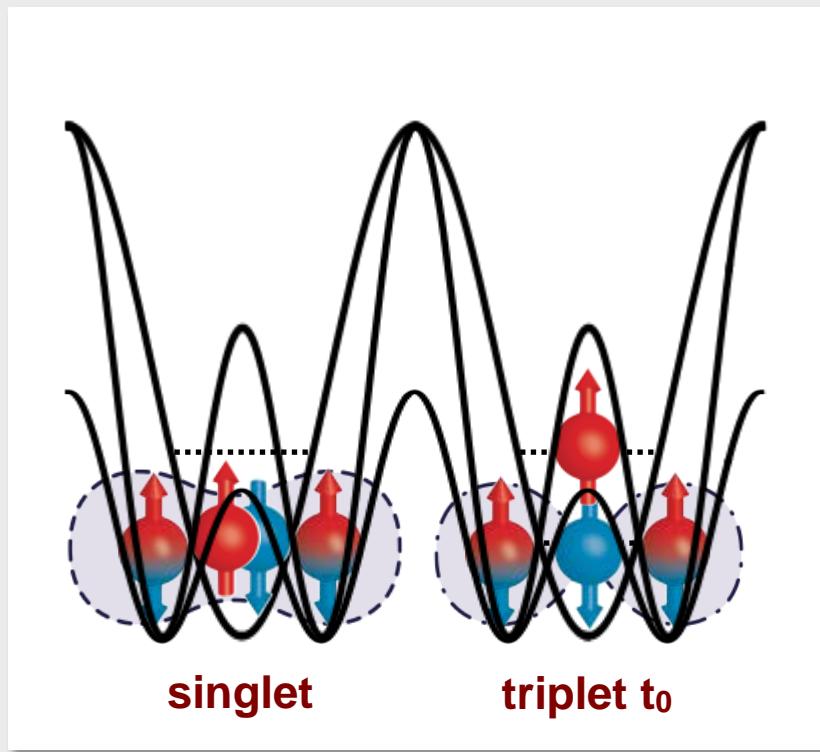
Probing: neighboring sites

The diagram shows a vertical double-headed arrow between two horizontal black lines. The top line is labeled "triplet" and the bottom line is labeled "singlet". To the right, the energy difference is given as $J_{d,s}$.

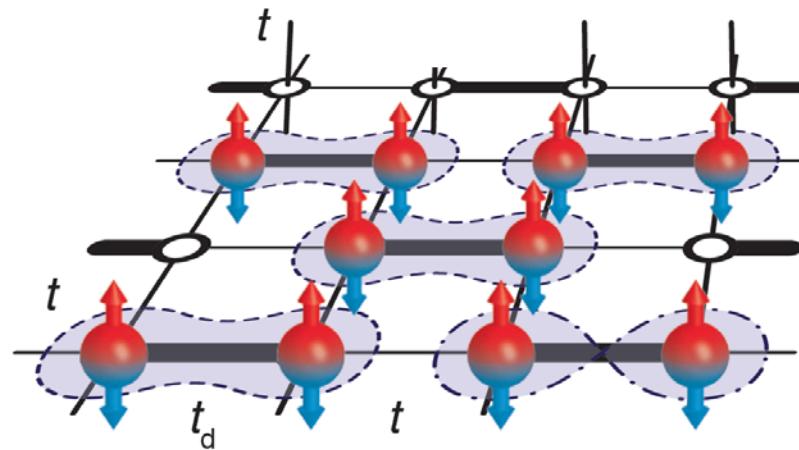
$$\text{triplet} \quad |\uparrow\uparrow\rangle, |\downarrow\downarrow\rangle, \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$
$$\text{singlet} \quad \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$k_B T < J_{d,s} : N_S > N_T$$

Detecting magnetic correlations



Dimerized lattice



Dimerized cubic lattice



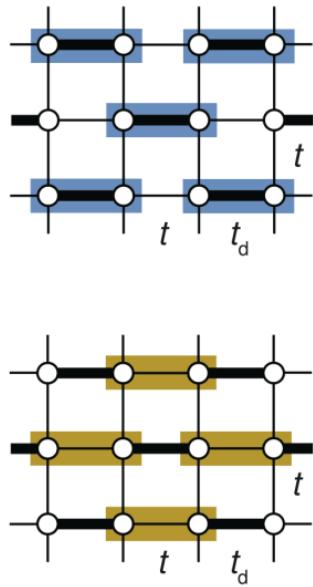
Singlet



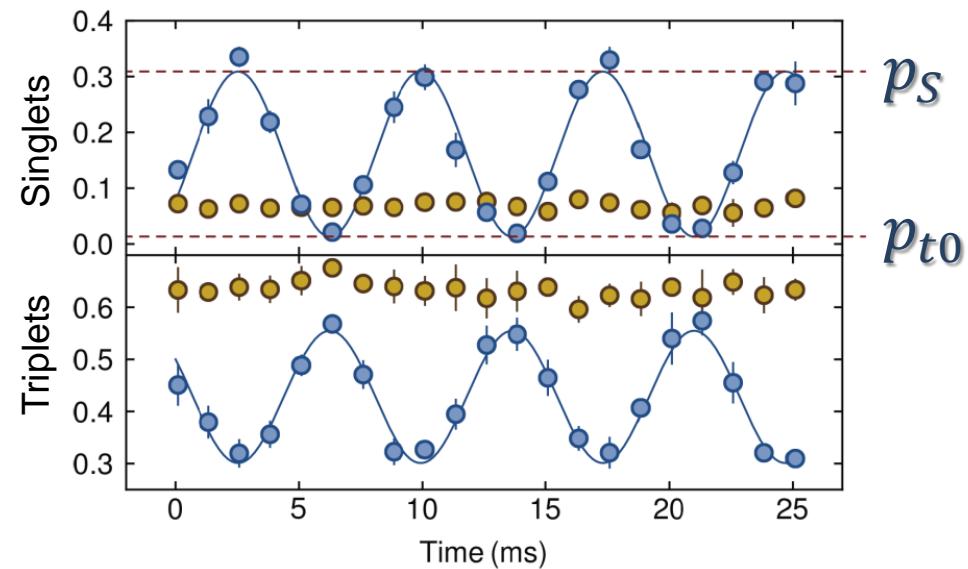
Triplet

Measuring singlets and triplets

Merging neighboring sites



Singlet-triplet oscillations

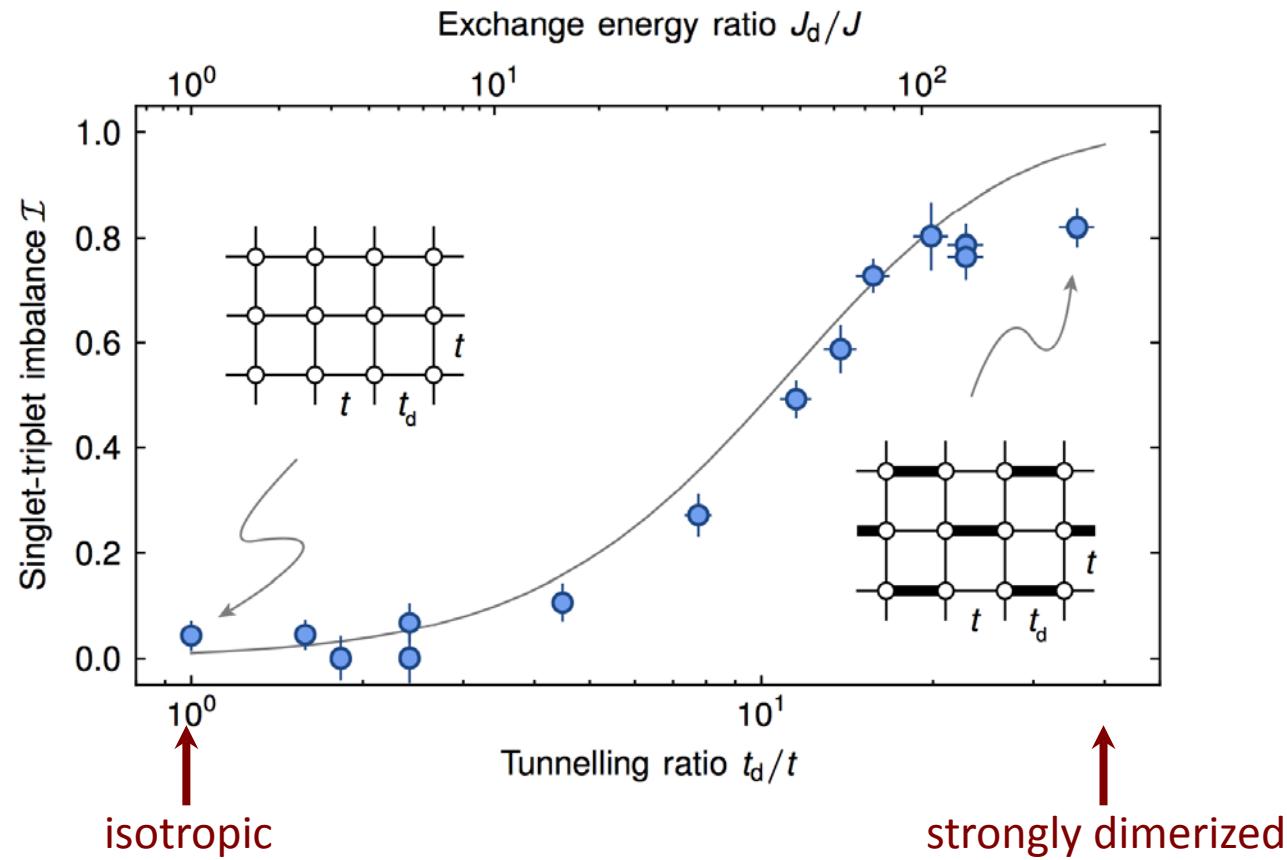


Singlet-Triplet Imbalance

$$\mathcal{I} = \frac{p_s - p_{t0}}{p_s + p_{t0}}$$

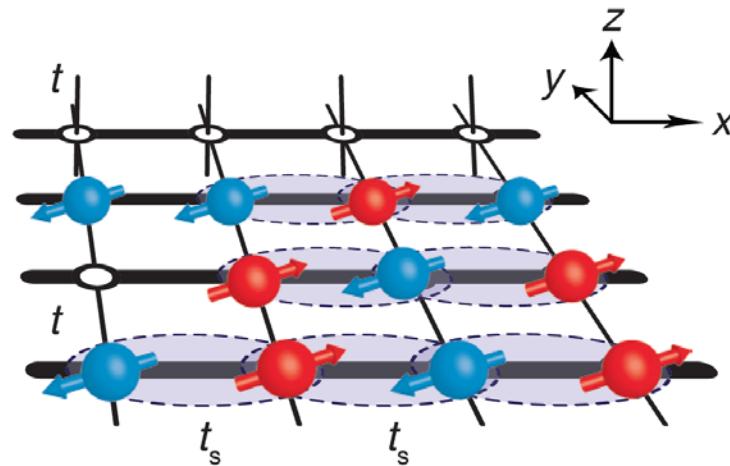
Singlet-Triplet Oscillations: S. Trotzky *et al.*, Phys. Rev. Lett. **105**, 265303 (2010)

Dimerization dependence



**Theory: second order high-temperature series expansion
of coupled dimers $s=1.7 k_B$**

Anisotropic cubic lattice



Anisotropic cubic lattice



Antiferromagnetic
spin correlation

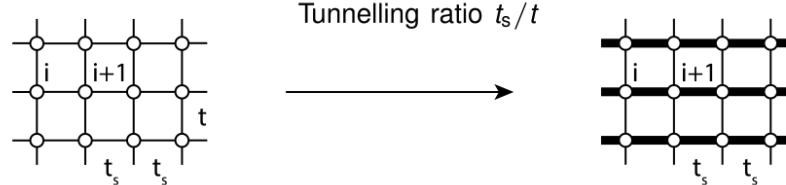
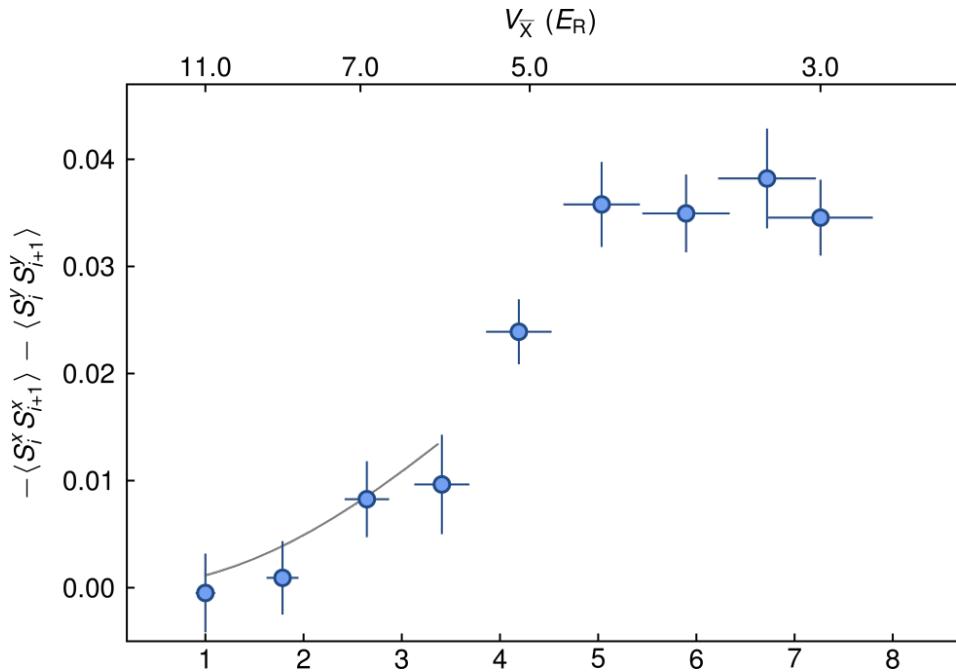
AFM correlations along x

$$\langle \mathbf{S}_i \cdot \mathbf{S}_{i+1} \rangle \neq 0$$

transverse spin correlator \Leftrightarrow population difference

$$-\langle S_i^x S_{i+1}^x \rangle - \langle S_i^y S_{i+1}^y \rangle = (p_s - p_{t_0}) / 2$$

Dependence on geometry

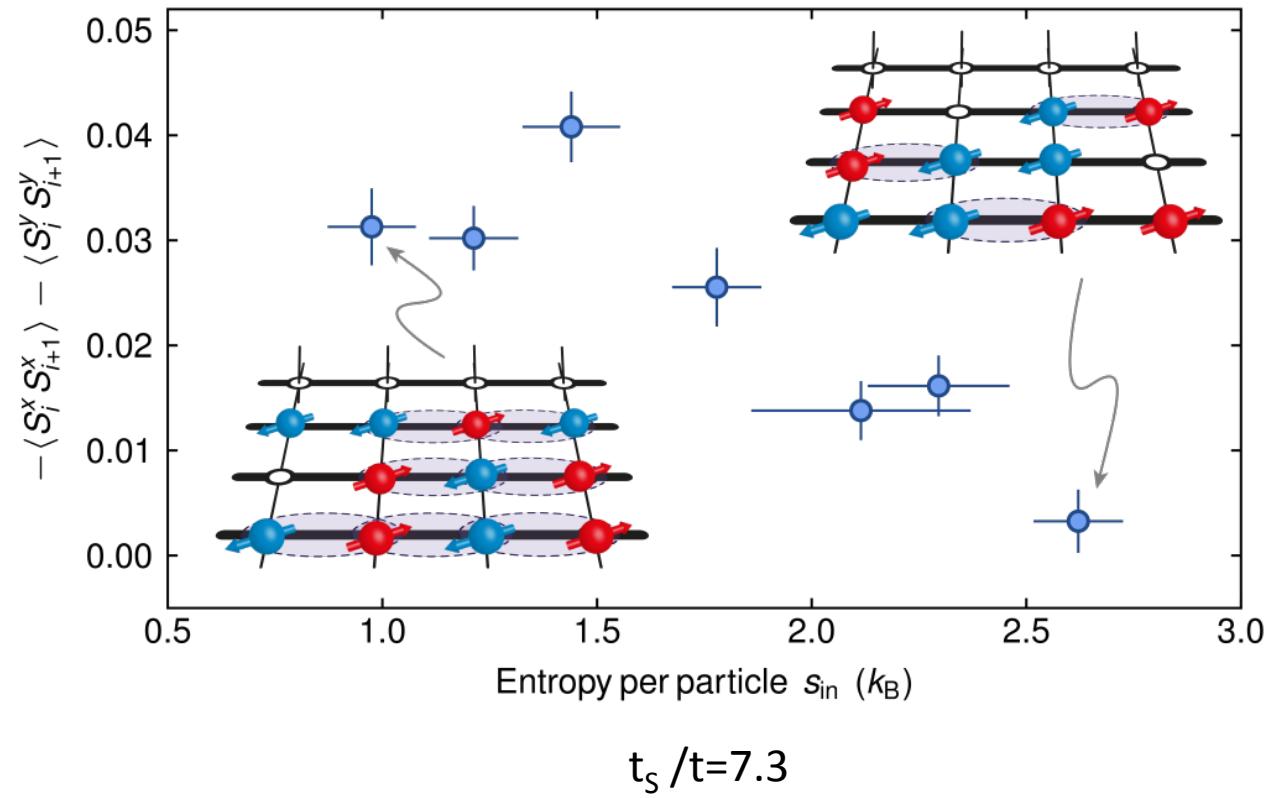


↑
isotropic

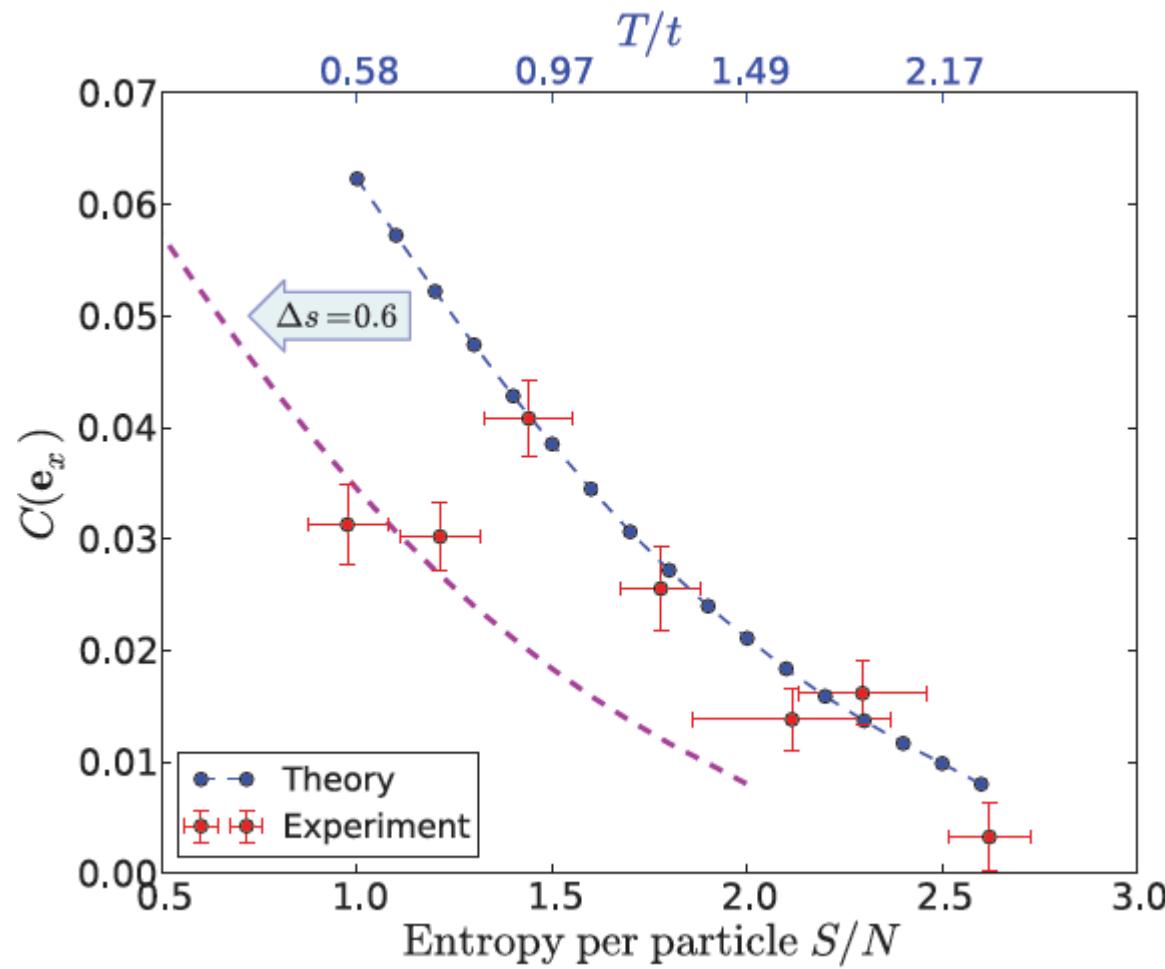
↑
strongly anisotropic

$$V_{Y,Z} = 11.0(3) E_R$$
$$s = 1.8 k_B$$

Dependence on entropy

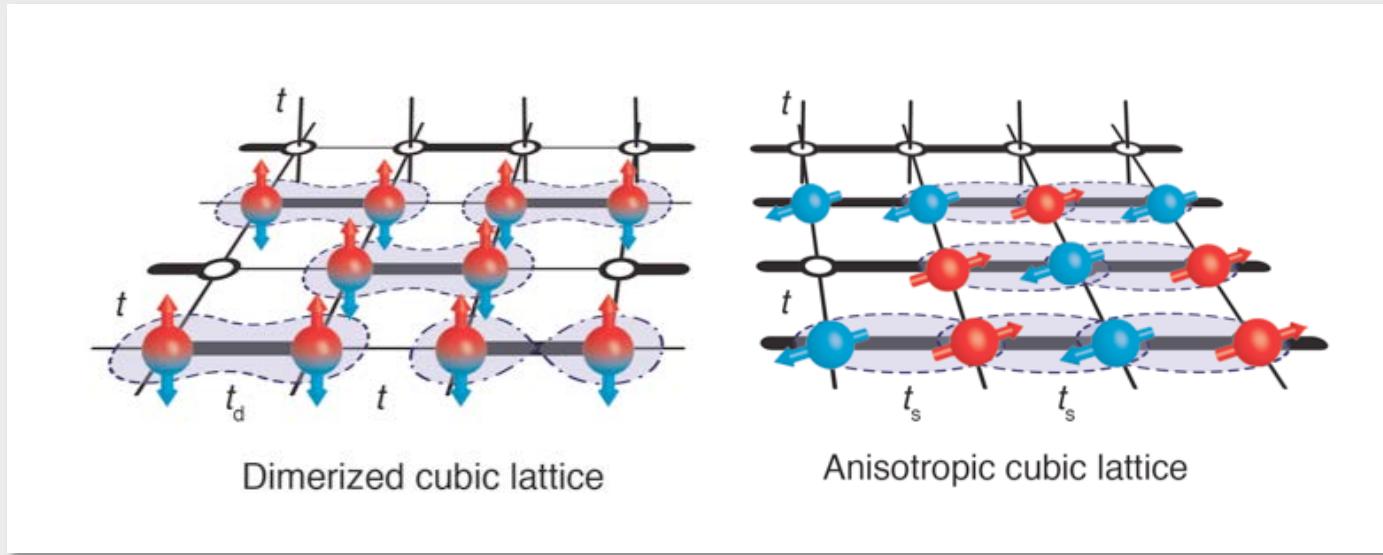


Comparison with theory



Theory (DCA+LDA): J. Imriška, M. Iazzi, L. Wang, E. Gull, and M. Troyer

Quantum magnetism



Nearest-neighbor magnetic correlations in thermalized ensembles

D. Greif, T. Uehlinger, G. Jotzu, L. Tarruell, and T. Esslinger, *Science* **340**, 1307 (2013)

Quantitative comparison with theory

J. Imriška, M. Iazzi, L. Wang, E. Gull, D. Greif, T. Uehlinger, G. Jotzu, L. Tarruell, T. Esslinger, and M. Troyer, *Phys. Rev. Lett.* **112**, 115301 (2014)

The ETH lattice team (fall 2012)



Magnetism theory: J. Imriška, M. Iazzi, L. Wang, E. Gull, M. Troyer

The ICFO Ultracold Quantum Gases group

