

A Study On The Entropy Of Wave Function

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$$\int_{-\infty}^{+\infty} |f(x)|^2 \text{Ln}(|f(x)|^2) dx + \int_{-\infty}^{+\infty} |F(k)|^2 \text{Ln}(|F(k)|^2) dk = \text{Constant}$$

We use the common definition of Fourier Transform in physics :

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{ikx} dx$$

so that it satisfies the Plancherel theorem in the following form :

$$\int_{-\infty}^{+\infty} |F(k)|^2 dk = \int_{-\infty}^{+\infty} |f(x)|^2 dx$$

In quantum mechanics, if $f(x)$ represents the wave function in the position space, then its fourier transform $F(k)$ represents the wave function in the momentum space.

Meanwhile, $|f(x)|^2$ is the probability that, say, a particle exists at position x (i.e. $|f(x)|^2$ is the probability density function, p.d.f.). Thus, $|f(x)|^2$ must satisfy the normalization condition:

$$\int_{-\infty}^{+\infty} |f(x)|^2 dx = 1$$

Due to the Plancherel Theorem, $|F(k)|^2$ satisfies the normalization condition as well (Probably why physicists like to use the definition of Fourier transform with a constant $\frac{1}{\sqrt{2\pi}}$)

Now, if we define the "width" of the p.d.f as :

$$(\Delta x)^2 \equiv \frac{\int_{-\infty}^{\infty} x^2 |f(x)|^2 dx}{\int_{-\infty}^{\infty} |f(x)|^2 dx} \quad (\Delta k)^2 \equiv \frac{\int_{-\infty}^{\infty} k^2 |F(k)|^2 dk}{\int_{-\infty}^{\infty} |F(k)|^2 dk}$$

then it has been proven in the theory of Fourier Transform that $(\Delta x) (\Delta k) \geq \frac{1}{2}$.

Heisenberg used the "width" to describe the uncertainty of a wave function, and wrote down the famous principle of uncertainty:

$$(\Delta x) (\Delta p) \geq \frac{\hbar}{2}$$

where p represents momentum, and satisfies the de Broglie relation :

$$p = \hbar k$$

where \hbar is the reduced Planck constant :

$$\hbar = h / (2 \pi)$$

I got the idea to use the informational entropy from Shannon to describe the uncertainty (the natural logarithm is used here for simplicity of calculation).

$$H \equiv \int_{-\infty}^{+\infty} p(x) \text{Log}\left(\frac{1}{p(x)}\right) dx = - \int_{-\infty}^{+\infty} p(x) \text{Log}(p(x)) dx$$

Then I propose that the total entropy of $|f(x)|^2$ and $|F(k)|^2$ is a constant. It's actually quite intuitive and consistent with the principle of uncertainty that, if you know more information in one space, you lose information in another space.

I first made some calculation with "Gaussian wave function"

$$f(x) = \sqrt{\frac{2}{\pi}} * \frac{1}{\sqrt{a}} * e^{-\left(\frac{x}{a}\right)^2}$$

where a is a real number that's bigger than 0 and bigger a indicates more uncertainty. The constant is chosen to satisfy the normalization condition.

Its Fourier transform is (simply replace a with $\frac{2}{a}$):

$$F(k) = \sqrt{\frac{2}{\pi}} * \sqrt{\frac{a}{2}} * e^{-\left(\frac{ak}{2}\right)^2}$$

The total entropy can be calculated and indeed is independent of a , which is $1 + \text{Log}(\pi)$.

If the wave function is multi-dimensional, similarly we obtain:

$$\text{Entropy}_{\text{total}} = n(1 + \text{Log}(\pi))$$

Then I turned to Sech function:

$$f(x) = \sqrt{\frac{1}{2a}} * \text{Sech}\left(\frac{x}{a}\right)$$

Its Fourier transform is (simply replace a with $\frac{2}{a\pi}$):

$$f(x) = \sqrt{\frac{a\pi}{4}} * \text{Sech}\left(\frac{a\pi k}{2}\right)$$

Again, the total entropy is independent of a . Generally in n dimensions:

$$\text{Entropy}_{\text{total}} = n(4 - \text{Log}(2\pi))$$

The total entropy can be calculated easily by hand in these two situations, and in both cases the Fourier transform takes the same form as the original function.

Then I use Mathematica to calculate some different cases. For example:

$$\sqrt{\frac{1}{2a}} \quad (-a \leq x \leq a) \longleftrightarrow \frac{\sin(ak)}{\sqrt{a\pi} k}$$

$$\sqrt{\frac{3}{2a}} \left(1 - \frac{|x|}{a}\right) \quad (-a \leq x \leq a) \longleftrightarrow -\sqrt{\frac{3}{4\pi a^3}} * e^{-iak} * (e^{iak} - 1)^2$$

$$\sqrt{\frac{a}{\pi(1+(ax)^2)}} \longleftrightarrow \sqrt{\frac{2}{a}} \frac{\text{BesselK}(0, \frac{|k|}{a})}{\pi}$$

The total entropy is calculated numerically for different a (say 2, 1, $\frac{1}{2}$, $\frac{1}{4}$, etc) , and all results are (quite) the same when a changes.

So I got enough courage to make the conjecture that

$$\int_{-\infty}^{+\infty} |f(x)|^2 \ln(|f(x)|^2) dx + \int_{-\infty}^{+\infty} |F(k)|^2 \ln(|F(k)|^2) dk = C$$

where C is a constant that only depend on the “type” of $f(x)$, but not its “shape”.