Tight Binding method for 1D atomic chain

- LCAO for a periodic atomic chain: the Tight Binding method
- s- and p-orbitals overlap integral: relation with band curvature
- overlap integral and bandwidth
- overlap integral and interatomic distance

1 Semiconductor Nanostructures

03/11/2

$$V_{\text{cut}}(x)$$
 $(N=\emptyset)$

$$\Delta U = 2$$
 Vot $n \neq 0$

$$H = T + V_{ad}(r) + \Delta V$$
 V_{c}

$$(T+Vor(r))$$

$$\psi = \sum_{n=-\infty}^{+\infty} \phi_{S}(n-n)$$

$$\frac{e^{i h r} l(r)}{\sqrt{(r+h)}} = e^{i h r} \sqrt{(r)}$$

$$\psi = \phi_S(n) + \sum_{n \neq 0} i h n \phi_S(n-n)$$

$$(T + V_{OM}(R) + DV)(\phi_{S}(R) + \sum_{i=1}^{n} \phi_{S}(R))$$

$$E(h) \left(\phi_{S}(n) + \sum_{n \neq 0} e^{ihn} \phi_{S}(n-h) \right)$$

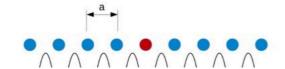
$$\begin{cases}
\phi_{s}^{(n)}(r + V_{oA} + DV) & (\phi_{s}(r) + \sum_{n \neq 0} 1^{n} V_{o}^{(n)} + \sum_{n \neq 0} 1^{n} V_{o}^{(n)} + \sum_{n \neq 0} 1^{n} V_{o}^{(n)} \\
& (\phi_{s}^{(n)}(r) + \sum_{n \neq 0} 1^{n} V_{o}^{(n)} + \sum_{n \neq 0} 1^{n} V_{o}^{(n)} \\
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& (\phi_{s}^{(n)}(r) + \sum_{n \neq 0$$

$$\sum_{R\neq 0} ihn \left\{ \phi_{S}(R) \otimes V \right\} \phi_{S}(R-R) dR = \frac{1}{R} \frac{1}{R} \frac{1}{R} \left(R \right)$$

$$= -\beta - \sum_{R\neq 0} ihR \qquad \beta(R) = E(H)$$

$$= E(H) = E - \beta - \sum_{R\neq 0} ihR \qquad R = \frac{1}{R} \frac{$$

Tight-Binding approximation



$$\begin{split} E(k) &= \varepsilon_s - \beta - 2\gamma(a)\cos(ka) & \text{here on } n \to \Gamma \\ &\approx \varepsilon_s - \beta - 2\gamma(a)\left(1 - \frac{1}{2}(ka)^2\right) & \text{cos}(ha) & \text{$\widehat{\wedge}$ $1-1$ (Ka)$} \\ &= \varepsilon_s - \beta - 2\gamma(a) + \gamma(a)(ka)^2 & \text{$\widehat{\wedge}$ 0} & \text{$S-$STATES$} \\ &&\text{$\widehat{\wedge}$ 0} & \text{$P-$STATES$} \end{split}$$

2 Semiconductor Nanostructures

03/11/2

Comparison with "real" bandstructures

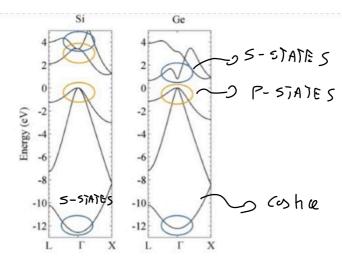


Figure 2.7: Bandstructure of Si and Ge taken from P. Moontragoon et al. J. Appl. Phys 112, 073106 (2012). The curvature of the bands around Γ (*i. e.* k=0) are highlighted in blue (orange) for bands originating from s-type (p-type) orbitals.

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03/11/2

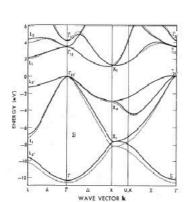
Comparison with "real" bandstructures

Relationship between the overlap integral γ and the energy width of the band

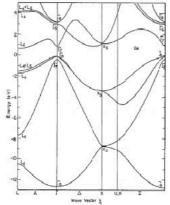
$$E(k) = \varepsilon_s - \beta - 2\gamma(a)\cos(ka)$$

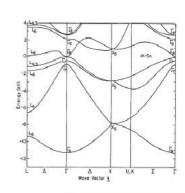
$$\approx \varepsilon_s - \beta - 2\gamma(a) \left(1 - \frac{1}{2} (ka)^2 \right)$$
$$= \varepsilon_s - \beta - 2\gamma(a) + \gamma(a)(ka)^2$$

Si









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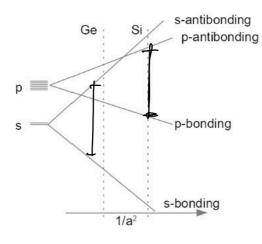
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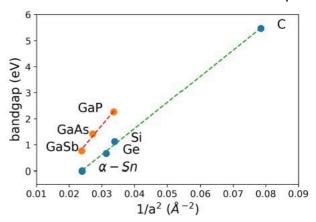
Overlap integral and effective mass

The dispersion around the Γ point (k \approx 0) of the Brillouin Zone can be described by means of the effective mass m*

$$\gamma = \frac{\hbar^2}{2m^*a^2}.$$

 $y \propto \frac{1}{a^2} \qquad \text{E(h)} \qquad \int h = 0$ $= \int h = 0$



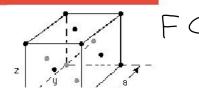


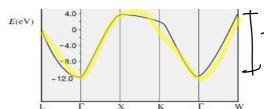
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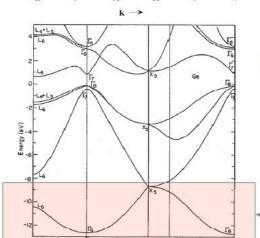
03/11/2

Comparison with "real" bandstructures

$\gamma = 1.0 \text{ eV}$ $E_s + \beta = 0$					
8.0-					







 $E(\mathbf{k}) = E_s - \beta_s - \sum_{\mathbf{R}} \gamma(\mathbf{R}) \; e^{i\mathbf{k}\cdot\mathbf{R}}$ For $\, {f R} \,$ spanning the first 12 near neighbours in a simple FCC lattice

$$\frac{a}{2}(\pm 1, \pm 1, 0); \frac{a}{2}(\pm 1, 0, \pm 1); \frac{a}{2}(0, \pm 1, \pm 1)$$

$$E(\mathbf{k}) = E_s - \beta_s - \gamma \left[e^{i(k_x + k_y)a/2} + e^{i(k_x - k_y)a/2} + e^{i(-k_x + k_y)a/2} + e^{i(k_x - k_y)a/2} + \dots \right]$$

$$= E_s - \beta_s - 4\gamma \left[\cos \frac{k_x a}{2} \cos \frac{k_y a}{2} + \cos \frac{k_y a}{2} \cos \frac{k_z a}{2} \right]$$

S-like bonding band

6

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03/11/2