# MESH PARTITIONING FOR HPX PARALLEL COMPUTING FOR NONLOCAL COMPUTATIONAL MODELS

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In this document, we consider a general class of nonlocal computational model seen in various fields, such as Peridynamics [21, 4, 14, 1, 12, 8, 19, 16, 17], discrete element method [6, 7, 13, 20, 9], Peridynamics plus discrete element method for granular media [26, 3], nonlocal cellcell adhesion in computational biology [2, 11, 22], nonlocal heat equation [5, 10]. References above are only minuscule fraction of what is available. The application of nonlocal modeling method for understanding of complex spatially multiscale phenomenon is seen in various new fields such as fluid mechanics, particulate media, directed self-assembly of Block-Copolymers, tumor modeling, etc. What is more interesting and also convenient is that underlying algorithm which is used to numerically solve these nonlocal models is more or less same and therefore the efficient computational method available for one type of nonlocal model can be easily applied to the nonlocal model in other fields.

It is our understanding that efficient computational method to solve the nonlocal model is very important. It is widely known that the nonlocal models are much more computationally demanding compared to their coun terpart, namely local models (such as heat equation, wave equation, etc). This is because the length of interaction in nonlocal model is typically 3 to 10 times larger than the size of discretization and therefore, the usual local assembly of matrix and vector considered in solution of pdes (partial differential equation) is not feasible. The efficient parallel implementation because of long-range interaction is difficult due to dependence over much larger length scale.

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In this document we present this problem using simple example of nonlocal diffusion equation, see [5] and references therein for more information. Our main goal is to highlight key difficulties in designing massively parallel scheme for nonlocal models while keeping the model related complexity to minimum. For this purpose nonlocal heat equation is suitable. Further, to fix the ideas we only consider two dimensional setting.

In this work we use parallel API called HPX. HPX is an open source asynchronous many task run time system that focuses on high performance computing [15, 23, 18]. HPX provides wait-free asynchronous execution and futurization for synchronization. It also features parallel execution policies utilizing a task scheduler, which enables a fine-grained load balancing parallelization and synchronization due to work stealing. HPX is in strict adherence to the C++ 11 [24] and C++ 17 standard definitions [25].

## 1. Nonlocal diffusion equation

In this section we give brief overview of nonlocal diffusion equation for temperature field u over an square domain  $D = [0,1]^2$  subjected to zero temperature condition on its boundary and subjected to given heat source distribution within the domain. It is a first order transient equation in time. For simplification, we only consider explicit time integration scheme, namely forward Euler scheme. For spatial discretization, we consider uniform mesh over domain D and consider finite difference approximation (also commonly referred to as particle discretization). The resulting final equations are very simple and serial implementation requires only few lines of code, see 1. With HPX, thread-level parallel implementation and shared-memory parallel implementation is almost as simple as serial implementation, see 2.

Our goal is to fully parallelize the solver, i.e. starting from mesh partition to computation of fields at mesh nodes. We proceed step wise starting from serial implementation, to shared-memory parallel implementation, and to fully parallel implementation and discuss the challenges going from one step to next.

Let material domain is  $D = [0, 1]^2$  and time domain I = [0, T]. We fix size of horizon (length over one point interacts with another, also called nonlocal length scale)  $\epsilon > 0$ . Let  $D_c = (-\epsilon, 1 + \epsilon)^2 - D$  be the nonlocal boundary of thickness  $\epsilon$  surrounding D. See Figure 1. We define  $H_r(x)$  as two-dimensional open ball of radius r centered at  $x \in \mathbf{R}^2$ .

Let  $u: I \times D \cup D_c \to \mathbf{R}$  is a temperature field. Let  $J: \mathbf{R} \to \mathbf{R}$  be positive function such that  $0 \le J(r) \le M$ 

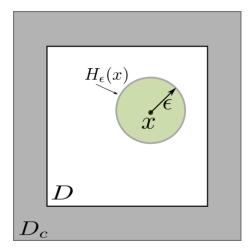


FIGURE 1. Material domain D and non-local boundary  $D_c$ . Figure shows typical material point  $x \in D$  and ball  $H_{\epsilon}(x)$ .

for  $r \in [0,1]$  and J(r) = 0 for  $r \notin [0,1]$ . We refer to J as influence function.

u(t,x) satisfies following nonlocal diffusion equation,  $\forall x \in D$  and  $\forall t \in I$ ,

$$\frac{\partial}{\partial t}u(t,x) = c\int_{H_{\epsilon}(x)}J(|y-x|/\epsilon)(u(t,y)-u(t,x))dy.$$

Initial condition is given by

$$(1.2) u(0,x) = u_0(x) \forall x \in D$$

and boundary condition is given by

$$(1.3) u(t,x) = 0 \forall x \in D_c \text{ and } \forall t \in I.$$

Constant c is chosen such that in the limit  $\epsilon \to 0$  the nonlocal operator in right-hand side of Equation 1.1. goes to  $k\nabla \cdot \nabla u$ . Local diffusion equation is given by

(1.4) 
$$\frac{\partial}{\partial t}u(t,x) = k\nabla \cdot \nabla u(t,x),$$

where k is the diffusivity constant. Boundary condition is u = 0 on  $\partial D$  (consistent with boundary condition in Equation 1.3 and initial condition is  $u(0, x) = u_0(x)$ . Constant c is taken as

$$(1.5) \quad c:=\begin{cases} \frac{k}{\epsilon^3 M_2}, & \text{when dimension } d=1\\ \frac{k}{\pi \epsilon^4 M_3}, & \text{when dimension } d=2, \end{cases}$$

where

$$(1.6) M_i = \int_0^1 J(r)r^i dr.$$

For standard influence functions,  $M_i$  can be computed easily. For example, when J = 1,  $M_i = 1/(i+1)$ , when J = 1 - r,  $M_2 = 1/12$ ,  $M_3 = 1/20$ , etc.

With choice of constant c in Equation 1.5, it can be shown formally that

(1.7) 
$$c \int_{H_{\epsilon}(x)} J(|y-x|/\epsilon)(u(t,y)-u(t,x))dy \to k\nabla \cdot \nabla u(t,x).$$

# Algorithm 1 Serial implementation

```
1: % Create neighbor list
2: for each integer i \in K do
      if |x_i - x_i| \le \epsilon then
         Add j to neighborList[i]
4:
      end if
5:
6: end for
7:
8: \% U is the vector of temperatures of
9: % all mesh nodes.
10:
11: % integrate in time
12: for each integer 0 \le k \le T/\Delta t do
      \% time step k
13:
      for each integer i \in K do
14:
         % Loop over neighbors of i
15:
16:
         val_{-}i = 0
         for each integer j \in \text{neighborList}[i] do
17:
           val_{-}i = val_{-}i +
18:
              cJ(|x_j - x_i|/\epsilon)(U[j] - U[i])V_j
19:
20:
         end for
         % Update temperature
21:
         U[i] = U[i] + \Delta t \times val_{-}i
22:
23:
      end for
      % Output temperature at time t^k = k\Delta t
24:
      Output(U)
25:
26: end for
```

#### 2. Finite difference approximation

Consider uniform mesh  $D_h = (D \cup D_c) \cap (h\mathbf{Z})^2$  of mesh size h > 0, see Figure 2.  $\mathbf{Z}$  is the set of positive and negative integers. Let  $\Delta t$  is size of time step and  $[0,T] \cap (\Delta t\mathbf{Z})$  be the discretization of time domain. We assume  $\epsilon = mh$  where  $m \geq 1$  is some integer.

We consider index set  $K \subset \mathbf{Z}^2$  such that for  $i \in K$ ,  $x_i = hi \in D$ . Similarly, consider index set  $K_c \subset \mathbf{Z}^2$  such that for  $i \in K_c$ ,  $x_i = hi \in D_c$ .

Let  $\hat{u}_i^k$  be the solution of forward Euler discretization in time. It satisfies,  $\forall 1 \leq k \leq T/\Delta t, \forall i \in K$ ,

$$(2.1) \frac{\hat{u}_{i}^{k+1} - \hat{u}_{i}^{k}}{\Delta t} = c \sum_{\substack{j \in K \cup K_{c}, \\ |x_{i} - x_{i}| < \epsilon}} J(|x_{j} - x_{i}|/\epsilon)(\hat{u}_{j}^{k} - \hat{u}_{i}^{k})V_{j}$$

and

(2.2) 
$$\hat{u}_i^k = 0 \quad \forall 0 \le k \le T/\Delta t, \forall i \in K_c.$$

At k = 0, we have  $\hat{u}_i^0 = u_0(x_i)$  for all  $i \in K$ .  $V_j$  is the volume occupied by node j and in our case we have  $V_j = h^2$ .

2.1. Serial and semi-parallel implementation of forward euler scheme. Algorithm for serial implementation is given in Algorithm 1. For parallel implementation, we will use parallel for loop (HPX utility).

28: **end for** 

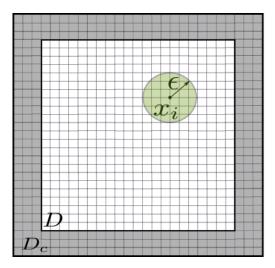


FIGURE 2. Uniform mesh of size h. Shaded area corresponds to nonlocal boundary  $D_c$ . We specify zero temperature on all the mesh nodes in the shaded area. For mesh node  $x_i$  in D, we consider interaction of  $x_i$  with all the mesh nodes inside the green shaded ball.

See Algorithm 2. In Algorithm 2, we observe that temperature data of all mesh nodes is stored in single data variable, "U", and similarly, neighbor list of all the mesh nodes are stored in single data variable, "neighbor List". We refer to this as semi-parallel as we have only parallelized for-loop and not the data. In fully parallel implementation, data will also be divided in number of computational nodes and each computational node will own data corresponding to it. This requires mesh partition. We discuss this in next section.

## 3. PARALLELIZATION AND MESH PARTITION

Given N number of computational nodes, mesh will be partitioned in N parts such that each computational node will store the information corresponding to the mesh partition it owns. For example, consider 4 computers connected by network. We would like to run the problem on all 4 computers in parallel. We will partition the mesh in four parts, see Figure 3.

In Figure 3, the mesh is colored to show the partition. Consider part of mesh colored as green. Let us suppose the id of computer who owns the green, blue, yellow, and red are 1, 2, 3, and 4 respectively. We now focus on one computational node, say green, and present two possible situations.

Case 1: Consider mesh node  $x_i$  which belongs to computer 1's partition. Let us suppose that  $x_i$  is within the dashed line, see Figure 4. For such  $x_i$ , all mesh nodes  $x_j \in H_{\epsilon}(x_i)$  will be within the green partition. Since computer 1 owns the mesh data of green partition, calculation of right hand side of Equation 2.1, i.e.

# Algorithm 2 Semi-parallel implementation

```
1: % Create neighbor list using
2: % hpx parallel for loop
3: hpx::parallel::for_loop each integer i \in K do
      if |x_i - x_j| \le \epsilon then
5:
         Add j to neighborList[i]
      end if
6:
7: end parallel for
8:
9: \% is the vector of temperatures of
   % all mesh nodes.
12: % integrate in time
13: for each integer 1 \le k \le T/\Delta t do
      \% time step k
14:
      % process mesh nodes in parallel
15:
      hpx::parallel::for_loop each integer i \in K
16:
         % Loop over neighbors of i
17:
         val_{-}i = 0
18:
         for each integer j \in \text{neighborList}[i] do
19:
           val_{-}i = val_{-}i +
20:
              +cJ(|x_i - x_i|/\epsilon)(U[j] - U[i])V_i
21:
         end for
22:
         \% Update temperature
23:
         U[i] = U[i] + \Delta t \times val_{-i}
24:
      end parallel for
25:
      % Output temperature at time t^k = k\Delta t
26:
      Output(U)
27:
```

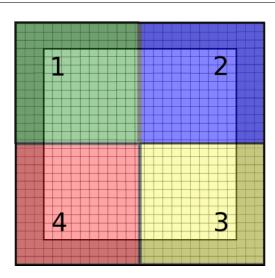


FIGURE 3. Typical mesh partition. Mesh nodes under each color are owned by the respective computer.

 $cJ(|x_j - x_i|/\epsilon)(\hat{u}_j^k - \hat{u}_i^k)V_j$ , can be carried out without needing to interact with other computers.

Case 2: Now consider another mesh node  $x_i$  in green partition. This mesh node is close to the boundary of green partition, and we see that there are mesh nodes  $x_j$  which are in ball  $H_{\epsilon}(x_i)$ , but which are not part of green

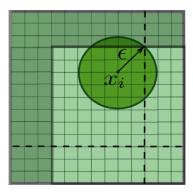


FIGURE 4. Mesh nodes of green partition which are within dashed line do not interact with partition owned by other computers.

partition. For example, we consider  $x_j \in H_{\epsilon}(x_i)$  in Figure 5. Mesh node  $x_j$  is in Blue partition. Blue partition is owned by computer 2. Therefore, to compute the right hand side term in Equation 2.1, computer 1 has to request the information associated to mesh node  $x_j$  from computer 2. We see that for all the mesh nodes in Green partition which are outside dash line, computer 1 will have to communicate with other computers for the information.

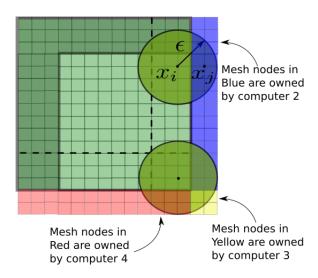


FIGURE 5. In first example, mesh node  $x_i$  of Green partition interacts with few mesh nodes of Blue partition owned by computer 2. In second example, we see that  $x_i$  of Green mesh interacts with mesh nodes owned by computer 2, 3, and 4.

3.1. Algorithm for fully parallel code. Consider Algorithm 3 which outlines the steps needed to run the problem in fully parallel framework. Following are the list of steps which will require effort in implementing Algorithm 3.

- 1. Partitioning of mesh: Libraries are available which can partition the mesh into given number of computational nodes.
- 2. List of interacting mesh nodes owned by other computational node: If global mesh data is available to each computational node then we can create a list of interacting neighboring mesh nodes which are owned by other computational node. This list will have global id, which is unique, of mesh nodes.
- 3. Sharing of information: For given computational node, we need to implement the method which shares information to other computational node and which receives the information from other computational node. For this we need a list of information to be shared and communication method available in HPX framework.

# 4. Constructing exact solution

In this section we show how to construct exact solution to test the numerical results. Recall that nonlocal diffusion equation is given by

$$\frac{\partial}{\partial t}u(t,x) = c \int_{H_{\epsilon}(x)} J(|y-x|/\epsilon)(u(t,y) - u(t,x))dy,$$

where u is the temperature field, c is a constant given by Equation 1.5. u satisfies boundary condition Equation 1.3 and initial condition Equation 1.2. We modify above equation to include the external source b:

$$\frac{\partial}{\partial t}u(t,x) = c \int_{H_{\epsilon}(x)} J(|y-x|/\epsilon)(u(t,y) - u(t,x))dy + b(t,x).$$

## 4.1. **One dimension.** Let

$$w(t,x) = \cos(2\pi t)\sin(2\pi x)$$

when  $x \in [0,1]$  and w(t,x) = 0 when  $x \notin [0,1]$ . We want solution of Equation 4.1 u to be equal to w, i.e., u = w. To achieve this we substitute

$$b(t,x) = \frac{\partial}{\partial t}w(t,x) - c\int_{H_{\epsilon}(x)} J(|y-x|/\epsilon)(w(t,y) - w(t,x))dy$$

in Equation 4.1 and prescribe  $u(0,x) = w(0,x) = \sin(2\pi x)$  as initial condition on u and u(t,x) = 0 on  $x \in D_c$  as boundary condition.

#### 4.2. Two dimension. Let

$$w(t, x) = \cos(2\pi t)\sin(2\pi x_1)\sin(2\pi x_2)$$

when  $x \in [0,1]^2$  and w(t,x) = 0 when  $x \notin [0,1]^2$ . We want solution of Equation 4.1 u to be equal to w, i.e., u = w. To achieve this we substitute (4.3)

$$b(t,x) = \frac{\partial}{\partial t}w(t,x) - c\int_{H_{\epsilon}(x)} J(|y-x|/\epsilon)(w(t,y) - w(t,x))dy$$

in Equation 4.1 and prescribe

$$u(0,x) = w(0,x) = \sin(2\pi x_1)\sin(2\pi x_2)$$

#### Algorithm 3 Fully parallel implementation 1: % mesh\_file is the file containing mesh data 2: % N is the number of computational nodes 3: 4: % get the id of this computational node $5: \text{my\_id} = \text{get\_id}()$ 6: 7: % read mesh data and create mesh partition 8: myMeshNodes = create\_mesh\_partition(mesh\_file) 9: % create neighbor list. 10: neighborList = create\_neighbor\_list(mesh\_file) 13: % next task is to create a list of mesh 14: % nodes, associated to other computational 15: % nodes, which interact with mesh nodes 16: % owned by this computational node 17: for each integer $0 \le i \le N$ do if $i == mv_i d$ then 18: % skip this i 19: 20: else % create a list of interacting nodes 21: % owned by comp. node i 22: neighborsOutside[i] 23: = create\_list\_interacting\_nodes(my\_id, i) 24: end if 25: end for 26: 27: % U is the vector of temperatures of 28: % mesh nodes owned by this comp. node 29: 30: 31: % integrate in time 32: for each integer $0 \le k \le T/\Delta t$ do % time step k 33: 34: % get data from other computer 35: % before integrating in time 36: for each integer $0 \le i \le N$ do 37: if $i == mv_i d$ then 38: % skip this i 39: 40: else % request data from comp. node i 41: getData[i] = request\_data(my\_id, i) 42: 43: % share data to comp. node i 44: % as comp. node i may also have 45: % mesh nodes which have neighbors 46: % in this comp. node 47: share\_data(my\_id, i) 48: end if 49: end for 50:

```
as initial condition on u and u(t,x) = 0 on x \in D_c as boundary condition.
```

```
% we now have all the relevant data so
51:
      % we can solve for temperature
52:
      % at next time step
53:
      hpx::parallel::for\_loop i \in myMeshNodes do
54:
         % Loop over neighbors of i
55:
         val_{-}i = 0
56:
         for each integer j \in \text{neighborList}[i] do
57:
           % check if j is in myMeshNode
58:
           if j \in \text{myMeshNode then}
59:
              % this mesh node is within dashed line
60:
61:
              % compute contribution
62:
              val_{-}i = val_{-}i
63:
                 +cJ(|x_j-x_i|/\epsilon)
64:
                 (U[j] - U[i])V_i
65:
           else
66:
              % this mesh node is outside dashed line
67:
              % search and find which comp. node
68:
              % owns mesh node j
69:
              get\_comp\_id =
70:
               search_for_mesh_node(my_id, j)
71:
72:
              % Look for j in getData[get_comp_id]
73:
              \% and get the temperature of mesh node j
74:
              t_{-j} = find_{-in}(j, getData[get\_comp\_id])
75:
76:
              % compute contribution
77:
              val_{-}i = val_{-}i
78:
                +cJ(|x_i-x_i|/\epsilon)
79:
                (t_{-}j - U[i])V_i
80:
           end if
81:
         end for
82:
         \% Update temperature
83:
         U[i] = U[i] + \Delta t \times val_{-}i
84:
      end parallel for
85:
      % Output temperature at time t^k = k\Delta t
86:
      Output(U)
87:
88: end for
```

#### References

- [1] Abigail Agwai, Ibrahim Guven, and Erdogan Madenci. Predicting crack propagation with peridynamics: a comparative study. *International journal of fracture*, 171(1):65–78, 2011. (document)
- [2] Nicola J Armstrong, Kevin J Painter, and Jonathan A Sherratt. A continuum approach to modelling cell-cell adhesion. *Journal of theoretical biology*, 243(1):98-113, 2006. (document)
- [3] Masoud Behzadinasab, Tracy J Vogler, Amanda M Peterson, Rezwanur Rahman, and John T Foster. Peridynamics modeling of a shock wave perturbation decay experiment in granular materials with intra-granular fracture. *Journal of Dynamic Behavior of Materials*, 4(4):529–542, 2018. (document)
- [4] Florin Bobaru and Wenke Hu. The meaning, selection, and use of the peridynamic horizon and its relation to crack

- branching in brittle materials. *International journal of fracture*, 176(2):215–222, 2012. (document)
- [5] Nathanial Burch and Richard Lehoucq. Classical, nonlocal, and fractional diffusion equations on bounded domains. *Inter*national Journal for Multiscale Computational Engineering, 9(6), 2011. (document)
- [6] Prathamesh Desai. Tribosurface interactions involving particulate media with dem-calibrated properties: Experiments and modeling. 2017. (document)
- [7] Prathamesh S Desai, Akash Mehta, Patrick SM Dougherty, and Fred C Higgs. A rheometry based calibration of a firstorder dem model to generate virtual avatars of metal additive manufacturing (am) powders. *Powder technology*, 342:441– 456, 2019. (document)
- [8] P Diehl, R Lipton, and MA Schweitzer. Numerical verification of a bond-based softening peridynamic model for small displacements: Deducing material parameters from classical linear theory. 2016. (document)
- [9] Maksym Dosta, Steven Dale, Sergiy Antonyuk, Carl Wassgren, Stefan Heinrich, and James D Litster. Numerical and experimental analysis of influence of granule microstructure on its compression breakage. *Powder technology*, 299:87–97, 2016. (document)
- [10] Qiang Du, Max Gunzburger, Richard B Lehoucq, and Kun Zhou. Analysis and approximation of nonlocal diffusion problems with volume constraints. SIAM review, 54(4):667–696, 2012. (document)
- [11] Christian Engwer, Christian Stinner, and Christina Surulescu. On a structured multiscale model for acid-mediated tumor invasion: the effects of adhesion and proliferation. *Mathematical Models and Methods in Applied Sciences*, 27(07):1355–1390, 2017. (document)
- [12] M Ghajari, L Iannucci, and P Curtis. A peridynamic material model for the analysis of dynamic crack propagation in orthotropic media. *Computer Methods in Applied Mechanics and Engineering*, 276:431–452, 2014. (document)
- [13] Anton Gladkyy and Meinhard Kuna. Dem simulation of polyhedral particle cracking using a combined mohr–coulomb—weibull failure criterion. *Granular Matter*, 19(3):41, 2017. (document)
- [14] Youn Doh Ha and Florin Bobaru. Studies of dynamic crack propagation and crack branching with peridynamics. *Inter*national Journal of Fracture, 162(1-2):229–244, 2010. (document)
- [15] Thomas Heller, Patrick Diehl, Zachary Byerly, John Biddiscombe, and Hartmut Kaiser. HPX An open source C++ Standard Library for Parallelism and Concurrency. In Proceedings of OpenSuCo 2017, Denver, Colorado USA, November 2017 (OpenSuCo'17), page 5, 2017. (document)
- [16] Prashant K. Jha and Robert Lipton. Numerical convergence of nonlinear nonlocal continuum models to local elastodynamics. *International Journal for Numerical Methods in En*gineering, 114(13):1389–1410, 2018. (document)
- [17] Prashant K Jha and Robert Lipton. Numerical convergence of finite difference approximations for state based peridynamic fracture models. Computer Methods in Applied Mechanics and Engineering, 351:184–225, July 2019. (document)
- [18] Hartmut Kaiser, Thomas Heller, Bryce Adelstein-Lelbach, Adrian Serio, and Dietmar Fey. Hpx: A task based programming model in a global address space. In *Proceedings of the* 8th International Conference on Partitioned Global Address Space Programming Models, page 6. ACM, 2014. (document)
- [19] Robert Lipton, Stewart Silling, and Richard Lehoucq. Complex fracture nucleation and evolution with nonlocal elastodynamics. arXiv preprint arXiv:1602.00247, 2016. (document)

- [20] Sebastian Lobo-Guerrero and Luis E Vallejo. Discrete element method evaluation of granular crushing under direct shear test conditions. *Journal of Geotechnical and Geoenvironmental Engineering*, 131(10):1295–1300, 2005. (document)
- [21] SA Silling, O Weckner, E Askari, and Florin Bobaru. Crack nucleation in a peridynamic solid. *International Journal of Fracture*, 162(1-2):219–227, 2010. (document)
- [22] Christian Stinner, Christina Surulescu, and Michael Winkler. Global weak solutions in a pde-ode system modeling multiscale cancer cell invasion. SIAM Journal on Mathematical Analysis, 46(3):1969–2007, 2014. (document)
- [23] Alexandre Tabbal, Matthew Anderson, Maciej Brodowicz, Hartmut Kaiser, and Thomas Sterling. Preliminary design examination of the parallex system from a software and hardware perspective. ACM SIGMETRICS Performance Evaluation Review, 38(4):81–87, 2011. (document)
- [24] The C++ Standards Committee. ISO International Standard ISO/IEC 14882:2011, Programming Language C++. Technical report, Geneva, Switzerland: International Organization for Standardization (ISO)., 2011. http://www.open-std.org/jtc1/sc22/wg21. (document)
- [25] The C++ Standards Committee. ISO International Standard ISO/IEC 14882:2017, Programming Language C++. Technical report, Geneva, Switzerland: International Organization for Standardization (ISO)., 2017. http://www.open-std.org/jtc1/sc22/wg21. (document)
- [26] F Zhu and J Zhao. A peridynamic investigation on crushing of sand particles. Géotechnique, 69(6):526-540, 2018. (document)