

## Abstract

Expressing neutrino oscillation probabilities in matter can be done either approximately with very good precision, or exactly with complicated expressions. Exact solutions require solving for the eigenvalues and, in turn, the eigenvectors. While there is no shortcut for the exact expressions for the eigenvalues, given those the eigenvectors can be determined in a straightforward fashion.

# Exact Neutrino Oscillation Probabilities in Matter

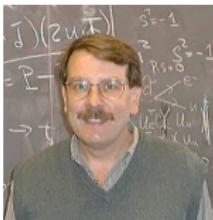
Peter B. Denton

TomFest

August 14, 2019



# Analytic Oscillation Probability Collaborators



Stephen Parke



Hisakazu Minakata



Gabriela Barenboim



Xining Zhang

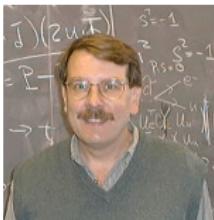


Christoph Ternes

[1604.08167](https://arxiv.org/abs/1604.08167), [1806.01277](https://arxiv.org/abs/1806.01277), [1808.09453](https://arxiv.org/abs/1808.09453),  
[1902.00517](https://arxiv.org/abs/1902.00517), [1902.07185](https://arxiv.org/abs/1902.07185), [1907.02534](https://arxiv.org/abs/1907.02534)

[github.com/PeterDenton/Nu-Pert](https://github.com/PeterDenton/Nu-Pert)  
[github.com/PeterDenton/Nu-Pert-Compare](https://github.com/PeterDenton/Nu-Pert-Compare)

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Terry Tao

[1604.08167](#), [1806.01277](#), [1808.09453](#),  
[1902.00517](#), [1902.07185](#), [1907.02534](#)  
[1908.03795](#)

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# The Billion Dollar Question

What is  $P(\nu_\mu \rightarrow \nu_e)$ ?

$$P({}^{\langle}\bar{\nu}_{\mu} \rightarrow {}^{\langle}\bar{\nu}_e) = |\mathcal{A}_{\mu e}|^2 \quad \mathcal{A}_{\mu e} = \mathcal{A}_{31} + e^{\pm i\Delta_{32}} \mathcal{A}_{21}$$

$$\mathcal{A}_{31} = 2s_{13}c_{13}s_{23} \sin \Delta_{31}$$

$$\mathcal{A}_{21} = 2s_{12}c_{13}(c_{12}c_{23}e^{i\delta} - s_{12}s_{13}s_{23}) \sin \Delta_{21}$$

$$\Delta_{ij} = \Delta m^2 {}_{ij} L / 4E$$

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...in matter?

Now: NOvA, T2K, MINOS, ...

Upcoming: DUNE, T2HK, ...

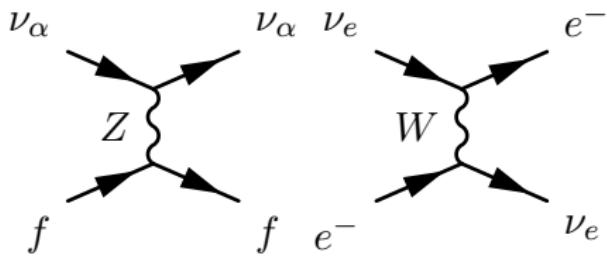
Second maximum: T2HKK? ESSnuSB? ...

# Matter Effects Matter

Call Schrödinger equation's eigenvalues  $m_i^2$  and eigenvectors  $U_i$ .

$$\mathcal{A}(\nu_\alpha \rightarrow \nu_\beta) = \sum_{i=1}^3 U_{\alpha i}^* U_{\beta i} e^{-im_i^2 L/2E} \quad P = |\mathcal{A}|^2$$

In matter  $\nu$ 's propagate in a new basis that depends on  $a \propto \rho E$ .



L. Wolfenstein, PRD 17 (1978)

Eigenvalues:  $m_i^2 \rightarrow \widehat{m^2}_i(a)$

Eigenvectors are given by  $\theta_{ij} \rightarrow \widehat{\theta}_{ij}(a)$        $\Leftarrow$       Unitarity

# Hamiltonian Dynamics

$$H = \frac{1}{2E} \left[ U \begin{pmatrix} 0 & & \\ & \Delta m_{21}^2 & \\ & & \Delta m_{31}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} a & 0 & \\ & 0 & \\ & & 0 \end{pmatrix} \right]$$
$$a = 2\sqrt{2}G_F N_e E$$

$$U = \begin{pmatrix} 1 & & \\ & c_{23} & s_{23} \\ & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & & s_{13}e^{-i\delta} \\ & 1 & \\ -s_{13}e^{i\delta} & & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & \\ -s_{12} & c_{12} & \\ & & 1 \end{pmatrix}$$

Find eigenvalues and eigenvectors:

$$H = \frac{1}{2E} \widehat{U} \begin{pmatrix} 0 & & \\ & \widehat{\Delta m_{21}^2} & \\ & & \widehat{\Delta m_{31}^2} \end{pmatrix} \widehat{U}^\dagger$$

Analytic expression?

# Eigenvalues Analytically: The Exact Solution

Solve the cubic characteristic equation

G. Cardano Ars Magna 1545

V. Barger, et al., PRD 22 (1980) 2718

H. Zaglauer, K. Schwarzer, Z.Phys. C40 (1988) 273

Then write down eigenvectors (mixing angles)

H. Zaglauer, K. Schwarzer, Z.Phys. C40 (1988) 273

K. Kimura, A. Takamura, H. Yokomakura [hep-ph/0205295](#)

PBD, S. Parke, X. Zhang [1907.02534](#)

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“Unfortunately, the algebra is rather impenetrable.”

V. Barger, et al.

# Eigenvalues Analytically: The Exact Solution

Solve the cubic characteristic equation.

$$\widehat{m^2}_1 = \frac{A}{3} - \frac{1}{3}\sqrt{A^2 - 3B}S - \frac{\sqrt{3}}{3}\sqrt{A^2 - 3B}\sqrt{1 - S^2}$$

$$\widehat{m^2}_2 = \frac{A}{3} - \frac{1}{3}\sqrt{A^2 - 3B}S + \frac{\sqrt{3}}{3}\sqrt{A^2 - 3B}\sqrt{1 - S^2}$$

$$\widehat{m^2}_3 = \frac{A}{3} + \frac{2}{3}\sqrt{A^2 - 3B}S$$

$$A = \Delta m_{21}^2 + \Delta m_{31}^2 + a$$

$$B = \Delta m_{21}^2 \Delta m_{31}^2 + a [c_{13}^2 \Delta m_{31}^2 + (c_{12}^2 c_{13}^2 + s_{13}^2) \Delta m_{21}^2]$$

$$C = a \Delta m_{21}^2 \Delta m_{31}^2 c_{12}^2 c_{13}^2$$

$$S = \cos \left\{ \frac{1}{3} \cos^{-1} \left[ \frac{2A^3 - 9AB + 27C}{2(A^2 - 3B)^{3/2}} \right] \right\}$$

# Exact Neutrino Oscillations in Matter: Mixing Angles

$$s_{\widehat{12}}^2 = \frac{- \left[ (\widehat{m^2}_2)^2 - \alpha \widehat{m^2}_2 + \beta \right] \Delta \widehat{m^2}_{31}}{\left[ (\widehat{m^2}_1)^2 - \alpha \widehat{m^2}_1 + \beta \right] \Delta \widehat{m^2}_{32} - \left[ (\widehat{m^2}_2)^2 - \alpha \widehat{m^2}_2 + \beta \right] \Delta \widehat{m^2}_{31}}$$

$$s_{\widehat{13}}^2 = \frac{(\widehat{m^2}_3)^2 - \alpha \widehat{m^2}_3 + \beta}{\Delta \widehat{m^2}_{31} \Delta \widehat{m^2}_{32}}$$

$$s_{\widehat{23}}^2 = \frac{s_{23}^2 E^2 + c_{23}^2 F^2 + 2c_{23}s_{23}c_\delta EF}{E^2 + F^2}$$

$$e^{-i\widehat{\delta}} = \frac{c_{23}s_{23} (e^{-i\delta}E^2 - e^{i\delta}F^2) + (c_{23}^2 - s_{23}^2) EF}{\sqrt{(s_{23}^2 E^2 + c_{23}^2 F^2 + 2EFc_{23}s_{23}c_\delta) (c_{23}^2 E^2 + s_{23}^2 F^2 - 2EFc_{23}s_{23}c_\delta)}}$$

$$\alpha = c_{13}^2 \Delta m_{31}^2 + (c_{12}^2 c_{13}^2 + s_{13}^2) \Delta m_{21}^2, \quad \beta = c_{12}^2 c_{13}^2 \Delta m_{21}^2 \Delta m_{31}^2$$

$$E = c_{13}s_{13} \left[ \left( \widehat{m^2}_3 - \Delta m_{21}^2 \right) \Delta m_{31}^2 - s_{12}^2 \left( \widehat{m^2}_3 - \Delta m_{31}^2 \right) \Delta m_{21}^2 \right]$$

$$F = c_{12}s_{12}c_{13} \left( \widehat{m^2}_3 - \Delta m_{31}^2 \right) \Delta m_{21}^2$$

H. Zaglauer, K. Schwarzer, [Z.Phys. C40 \(1988\) 273](#)

# Too Hard: Approximations

- ▶ Small matter potential:  $a/\Delta m^2$

Y. Li, Y. Wang, Z-z. Xing, [1605.00900](#)

- ▶  $s_{13}, s_{13}^2$

A. Cervera, et al., [hep-ph/0002108](#)

H. Minakata, [0910.5545](#)

K. Asano, H. Minakata, [1103.4387](#)

- ▶  $\Delta m_{21}^2/\Delta m_{31}^2 \sim 0.03$

J. Arafune, J. Sato, [hep-ph/9607437](#)

A. Cervera, et al., [hep-ph/0002108](#)

M. Freund, [hep-ph/0103300](#)

E. Akhmedov, et al., [hep-ph/0402175](#)

S. Agarwalla, Y. Kao, T. Takeuchi, [1302.6773](#)

M. Blennow, A. Smirnov, [1306.2903](#)

H. Minakata, S. Parke, [1505.01826](#)

**PBD**, H. Minakata, S. Parke, [1604.08167](#)

(See G. Barenboim, **PBD**, S. Parke, C. Ternes [1902.00517](#) for a review)



# Eigenvalues to Eigenvectors

KTY pushed calculating the eigenvectors from the eigenvalues.

K. Kimura, A. Takamura, H. Yokomakura [hep-ph/0205295](#)

Expressions for:

$$\hat{U}_{\alpha i} \hat{U}_{\beta i}^*$$

for  $\alpha \neq \beta$ .

Wanted to preserve phase information for  $\hat{\delta}$ .

# Eigenvalues: the Rosetta Stone

We realized:

$$|\hat{U}_{\alpha i}|^2 = \frac{(\widehat{m^2}_i - \xi_\alpha)(\widehat{m^2}_i - \chi_\alpha)}{\Delta\widehat{m^2}_{ij}\Delta\widehat{m^2}_{ik}}$$

where  $\xi_\alpha$  and  $\chi_\alpha$  are the submatrix eigenvalues of

$$H_\alpha \equiv \begin{pmatrix} H_{\beta\beta} & H_{\beta\gamma} \\ H_{\gamma\beta} & H_{\gamma\gamma} \end{pmatrix}$$

e.g.

$$\xi_e + \chi_e = \Delta m_{21}^2 + \Delta m_{ee}^2 c_{13}^2$$

$$\xi_e \chi_e = \Delta m_{21}^2 [\Delta m_{ee}^2 c_{13}^2 c_{12}^2 + \Delta m_{21}^2 (s_{12}^2 c_{12}^2 - s_{13}^2 s_{12}^2 c_{12}^2)]$$

# Eigenvalues: the Rosetta Stone

$$s_{\widehat{13}}^2 = |\widehat{U}_{e3}|^2 = \frac{(\widehat{m^2}_3 - \xi_e)(\widehat{m^2}_3 - \chi_e)}{\Delta \widehat{m^2}_{31} \Delta \widehat{m^2}_{32}}$$

$$s_{\widehat{12}}^2 c_{\widehat{13}}^2 = |\widehat{U}_{e2}|^2 = -\frac{(\widehat{m^2}_2 - \xi_e)(\widehat{m^2}_2 - \chi_e)}{\Delta \widehat{m^2}_{32} \Delta \widehat{m^2}_{21}}$$

$$s_{\widehat{23}}^2 c_{\widehat{13}}^2 = |\widehat{U}_{\mu 3}|^2 = \frac{(\widehat{m^2}_3 - \xi_\mu)(\widehat{m^2}_3 - \chi_\mu)}{\Delta \widehat{m^2}_{31} \Delta \widehat{m^2}_{32}}$$

What about  $\widehat{\delta}$ ?

# CPV From Rosetta

Toshev identity:

$$\sin \hat{\delta} = \frac{\sin 2\theta_{23}}{\sin 2\hat{\theta}_{23}} \sin \delta$$

S. Toshev [MPL A6 \(1991\) 455](#)

Get the sign of  $\cos \hat{\delta}$  from e.g.  $|U_{\mu 1}|^2$ .

## In General

Two flavor:

$$|\widehat{U}_{\alpha i}|^2 = \frac{\widehat{m^2}_i - \xi_\alpha}{\Delta \widehat{m^2}_{ij}}$$

leads to

$$\begin{aligned} \sin^2 \widehat{\theta} &= |\widehat{U}_{e2}|^2 = \frac{\widehat{m^2}_2 - \xi_e}{\widehat{m^2}_2 - \widehat{m^2}_1} \\ &= \frac{1}{2} \left( 1 - \frac{\Delta m^2 \cos 2\theta - a}{\sqrt{(\Delta m^2 \cos 2\theta - a)^2 + (\Delta m^2 \sin 2\theta)^2}} \right) \end{aligned}$$

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Numerically checked for  $N = 4, 5$ .

True for all  $N$ ?

We did it!

# EIGENVECTORS FROM EIGENVALUES

PETER B. DENTON, STEPHEN J. PARKE, TERENCE TAO, AND XINING ZHANG

ABSTRACT. We present a new method of succinctly determining eigenvectors from eigenvalues. Specifically, we relate the norm squared of the elements of eigenvectors to the eigenvalues and the submatrix eigenvalues.

$$|v_{i,j}|^2 = \frac{\prod_{k=1}^{n-1} (\lambda_i - \xi_{j,k})}{\prod_{k=1; k \neq i}^n (\lambda_i - \lambda_k)}$$

## Key Points

- ▶ Someone write a paper about the moon with Tom
- ▶ Tom gave me enough freedom to learn how to research
- ▶ Tom pushed me to go elsewhere
  
- ▶ All you need is Love Eigenvalues
- ▶ Terry is “cheery firehose”