

Precision measurements in the DUNE Near Detector Complex

Zahra Tabrizi

Virginia Tech/CNP

Based on:

- Ballett, Hostert, Pascoli, Perez-Gonzalez, Tabrizi and Funchal, JHEP 1901, 119 (2019)
- de Gouv  a, Machado, Perez-Gonzalez and Tabrizi, arXiv:1912.06658

Brookhaven Neutrino Theory Virtual Seminars

April 20th 2020

Outline

- Neutrino Trident production
- Neutrino-Electron scattering
- Measuring the Weak Mixing Angle at DUNE
- Conclusion

Open Questions:

- What is the order of neutrino masses?
- What is the value of the CP phase?
- Is neutrino its own anti particle?
- What is the origin of neutrino mass?
- What is the absolute neutrino masses?
- Are there more than three neutrinos?



We need next generation long baseline neutrino experiments!

Physics goals of near detectors:

Primary role of the ND is to study the systematic uncertainties.

High beam luminosity +
Large fiducial mass

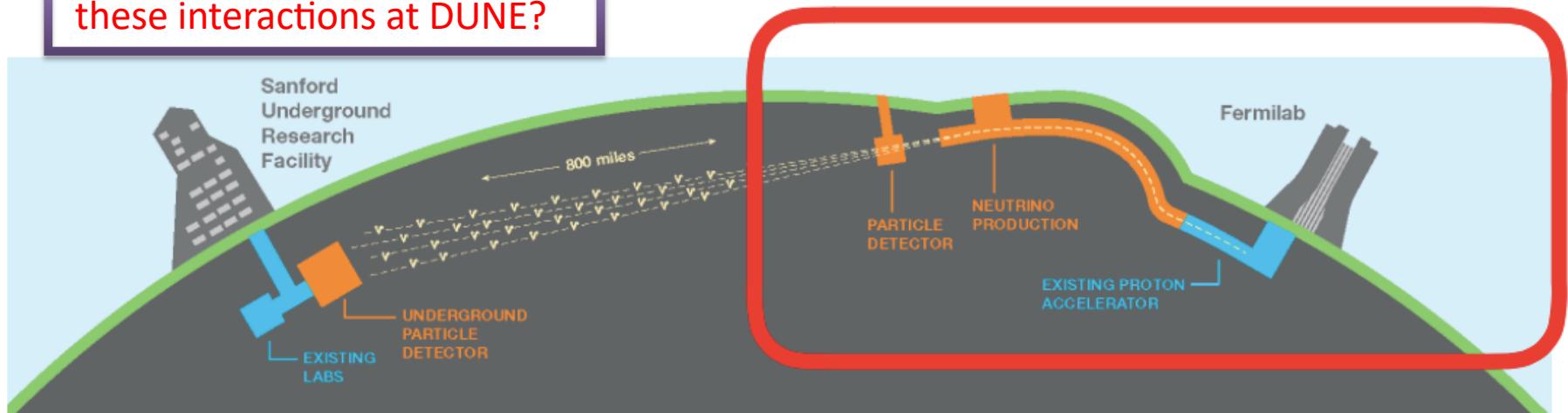
- ❖ neutrino-electron scattering
- ❖ neutrino trident production

Ideal to investigate (rare)
neutrino interactions

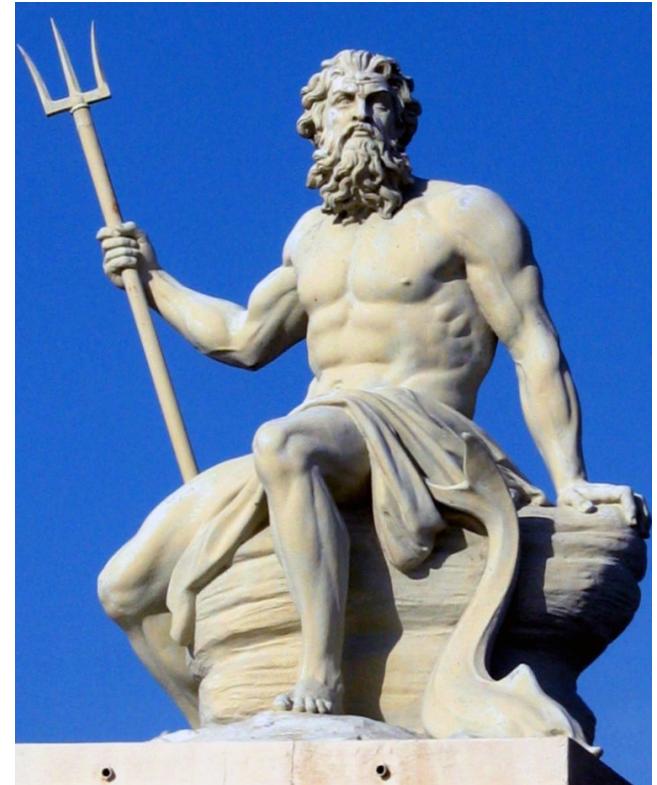
$$\sigma < 10^{-44} \text{ cm}^2$$

- Test SM predictions
- Search for BSM physics

What can we learn from
these interactions at DUNE?



Neutrino Trident Scattering at Near Detectors

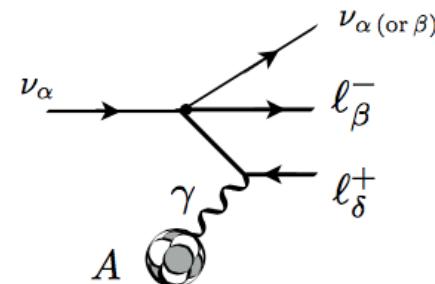


Ballett, Hostert, Pascoli, Perez, ZT and Funchal
JHEP **1901**, 119 (2019)

Neutrino Trident Scattering

Production of a **charged lepton pair**
in the scattering of a **neutrino**
in the Coulomb field of a **heavy nucleus/nucleon**

$$\nu_\alpha + \mathcal{N} \rightarrow \nu_\beta + \ell_\gamma^+ + \ell_\delta^- + \mathcal{N}$$



CHARM II

PLB 245 (1990) 271

$$\frac{\sigma_{\text{CHARM II}}}{\sigma_{\text{SM}}} = 1.58 \pm 0.57$$

CCFR

PRL 66 (1991) 3117

$$\frac{\sigma_{\text{CCFR}}}{\sigma_{\text{SM}}} = 0.82 \pm 0.28$$

NuTeV

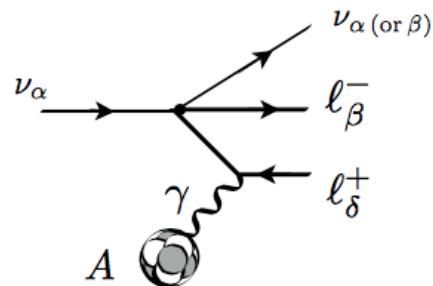
Vancouver 1998, High
energy physics, vol. 1

$$\frac{\sigma_{\text{NuTeV}}}{\sigma_{\text{SM}}} = 0.67 \pm 0.27$$

Neutrino Trident Scattering

Production of a **charged lepton pair**
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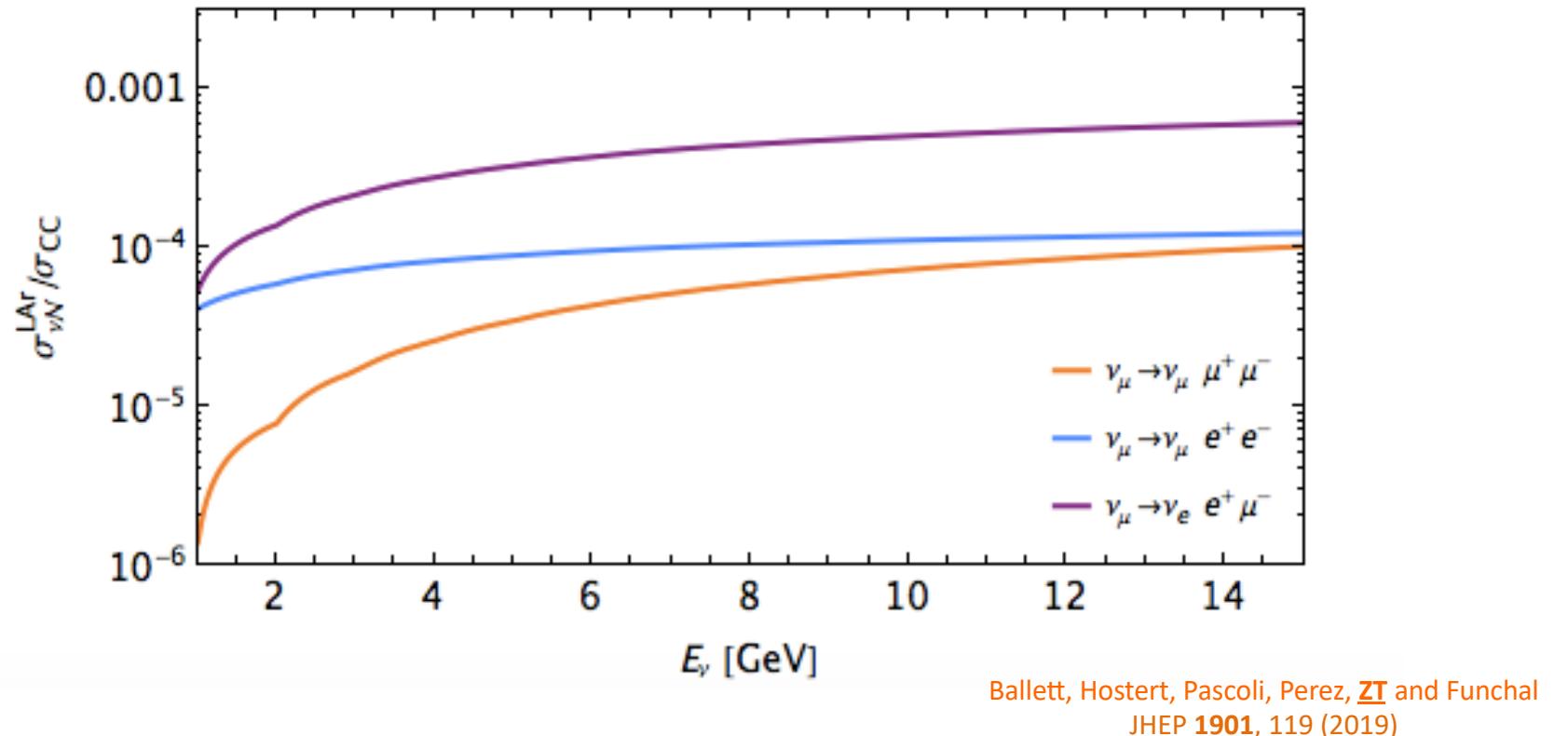
Neutrino	Antineutrino	SM Contributions
$\nu_\mu \rightarrow \nu_\mu \mu^+ \mu^-$	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu \mu^+ \mu^-$	CC, NC
$\nu_\mu \rightarrow \nu_e e^+ \mu^-$	$\bar{\nu}_\mu \rightarrow \bar{\nu}_e e^- \mu^+$	CC
$\nu_\mu \rightarrow \nu_\mu e^+ e^-$	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu e^+ e^-$	NC
$\nu_e \rightarrow \nu_e e^+ e^-$	$\bar{\nu}_e \rightarrow \bar{\nu}_e e^+ e^-$	CC, NC
$\nu_e \rightarrow \nu_\mu \mu^+ e^-$	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu e^+ \mu^-$	CC
$\nu_e \rightarrow \nu_e \mu^+ \mu^-$	$\bar{\nu}_e \rightarrow \bar{\nu}_e \mu^+ \mu^-$	NC

→ Observed

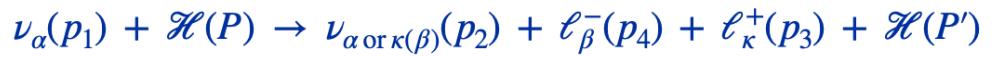
$$V_{\alpha\beta\kappa}(A_{\alpha\beta\kappa}) \equiv g_V^\beta(g_A^\beta)\delta_{\beta\kappa} + \delta_{\alpha\beta}$$

Measuring neutrino trident events give information on vector/axial couplings

How rare is it?



Trident cross section:



$$\frac{d^2\sigma_{\nu X}}{dQ^2 d\hat{s}} = \frac{1}{32\pi^2(s - M_{\mathcal{H}}^2)^2} \frac{H_X^{\mu\nu} L_{\mu\nu}}{Q^4}$$

hadronic tensor
(depends on scattering regime)

leptonic tensor

momentum transfer

Center-of-mass energy of the neutrino-photon system

Center-of-mass energy of the total system

*W. Czyz, G. C. Sheppley and J. D. Walecka,
Nuovo Cim. 34 (1964) 404.*

Coherent Regime: The neutrino interacts with the whole nucleus

Diffractive Regime: The neutrino interacts with the individual nucleons

We can separate the phc

$$\frac{d^2\sigma_{\nu X}}{dQ^2 d\hat{s}} = \frac{1}{32\pi^2} \frac{1}{\hat{s} Q^2} \left[h_X^T(Q^2, \hat{s}) \sigma_{\nu\gamma}^T(Q^2, \hat{s}) + h_X^L(Q^2, \hat{s}) \sigma_{\nu\gamma}^L(Q^2, \hat{s}) \right]$$

Equivalent Photon Approximation (EPA): The full cross section is related to the cross section of the neutrino scattering with a real photon, multiplied by the probability of creating a virtual photon.

EPA assumptions

- 1) Neglecting the L contribution ($h^L(q^2, \hat{s}) \sigma_{\nu\gamma}^L(q^2, \hat{s}) \approx 0$).
- 2) Taking the T contribution of the cross section to be on-shell ($\sigma_{\nu\gamma}^T(q^2, \hat{s}) \approx \sigma_{\nu\gamma}^T(0, \hat{s})$).

$$\sigma_t(P_i + C_s \rightarrow P_f + C_s) \approx \int dP(Q^2, \hat{s}) \sigma_\gamma(P_i + \gamma \rightarrow P_f; \hat{s}, Q^2 = 0)$$

EPA assumptions

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QED

$$\sigma_\gamma^{\text{QED}}(P_i + \gamma \rightarrow P_f; \hat{s}, 0) \propto \frac{1}{\hat{s}}$$

Decreases with
increasing transferred
four-momentum

On-shell $>>$ off-shell

Fermi Limit of the SM

$$\sigma_\gamma^{\text{FL}}(P_i + \gamma \rightarrow P_f; \hat{s}, 0) \propto G_F^2 \hat{s}$$

Increases with
increasing transferred
four-momentum

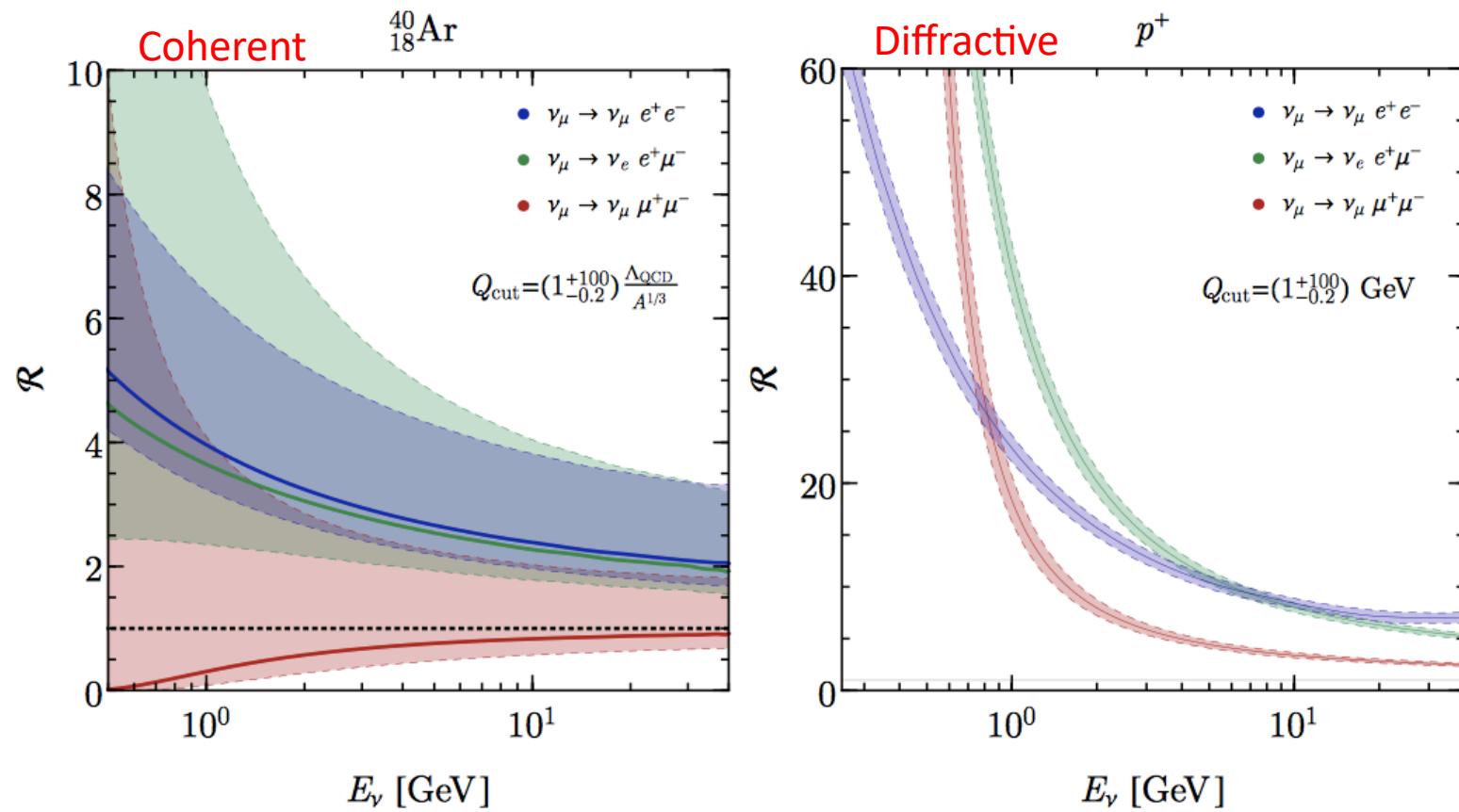
On-shell $<<$ off-shell

Ballett, Hostert, Pascoli, Perez, [ZT](#) and Funchal
JHEP **1901**, 119 (2019)

How bad is it?

$$\mathcal{R} = \frac{\sigma_{\text{EPA}}(E_\nu)|_{Q_{\max}}}{\sigma_{\text{4PS}}(E_\nu)}$$

Ballett, Hostert, Pascoli, Perez, ZT and Funchal
JHEP 1901, 119 (2019)



EPA approximation doesn't work. Full 4PS calculation must be done!!!

Trident rates at LAr Detectors

$$N = \text{time} \times \# \text{ of targets} \times \text{efficiency} \times \int_{E_i}^{E_f} dE_\nu \frac{d\phi(E_\nu)}{dE_\nu} \sigma(E_\nu)$$

Channel	SBND	μ BooNE	ICARUS	DUNE ND	ν STORM ND
Total $e^\pm \mu^\mp$	10	0.7	1	2993 (2307)	191
	2	0.1	0.2	692 (530)	41
Total $e^+ e^-$	6	0.4	0.7	1007 (800)	114
	0.7	0.0	0.1	143 (111)	14
Total $\mu^+ \mu^-$	0.4	0.0	0.0	286 (210)	11
	0.4	0.0	0.0	196 (147)	9

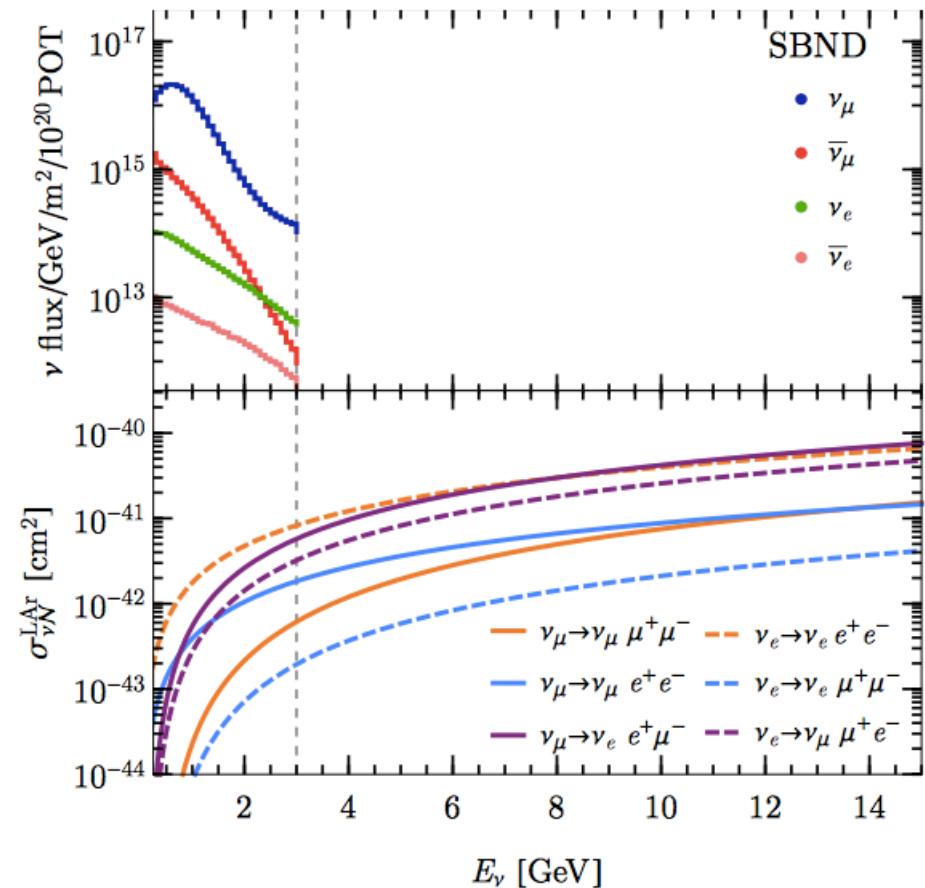
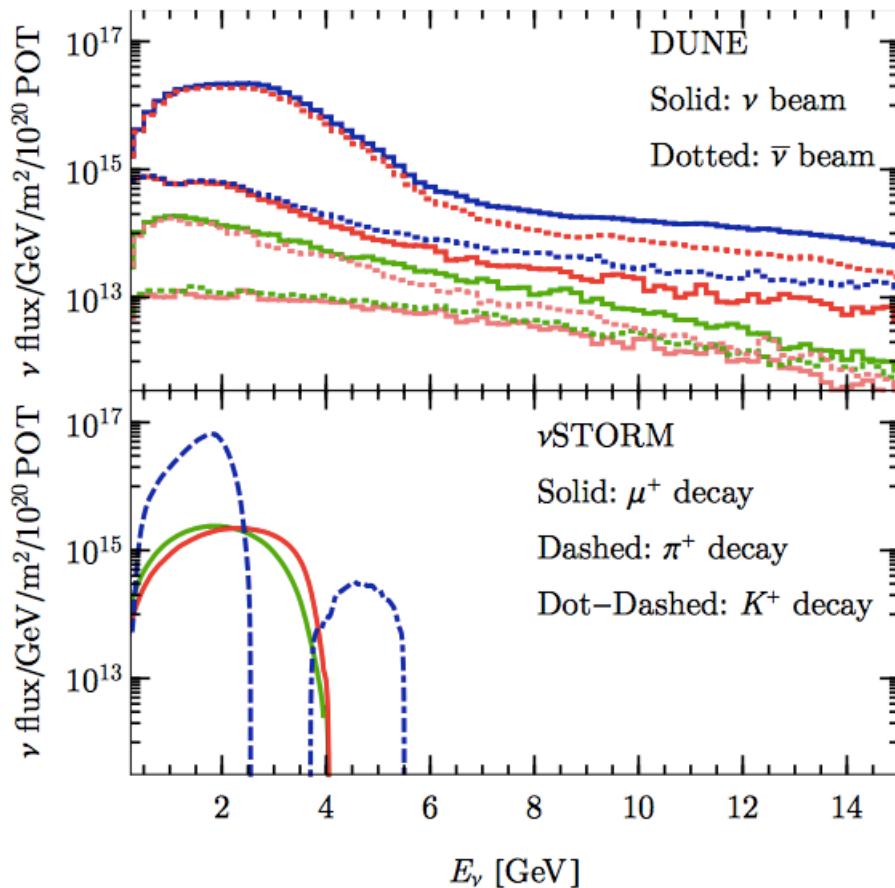
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Coherent (upper) and diffractive (lower) trident events for (anti)neutrino mode.



Trident rates at LAr Detectors

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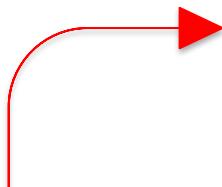
Trident background analysis

Genuine dilepton production is rare, but misID of particles is the problem.

misID	Channel	$N_B^{\text{misID}}/N_{\text{CC}}$	$N_B^{\text{had}}/N_{\text{CC}}$	$N_B^{\text{kin}}/N_{\text{CC}}$
γ as e^\pm	$e^\pm \mu^\mp$	$1.67 (1.62) \times 10^{-4}$	$2.68 (4.31) \times 10^{-5}$	$4.40 (3.17) \times 10^{-7}$
γ as $e^+ e^-$	$e^+ e^-$	$2.83 (4.19) \times 10^{-4}$	$1.30 (2.41) \times 10^{-4}$	$6.54 (14.1) \times 10^{-6}$
π^\pm as μ^\pm	$\mu^+ \mu^-$	$2.66 (2.73) \times 10^{-3}$	$10.4 (9.75) \times 10^{-4}$	$3.36 (3.10) \times 10^{-8}$

γ as $e^+ e^-$

π^\pm as μ^\pm

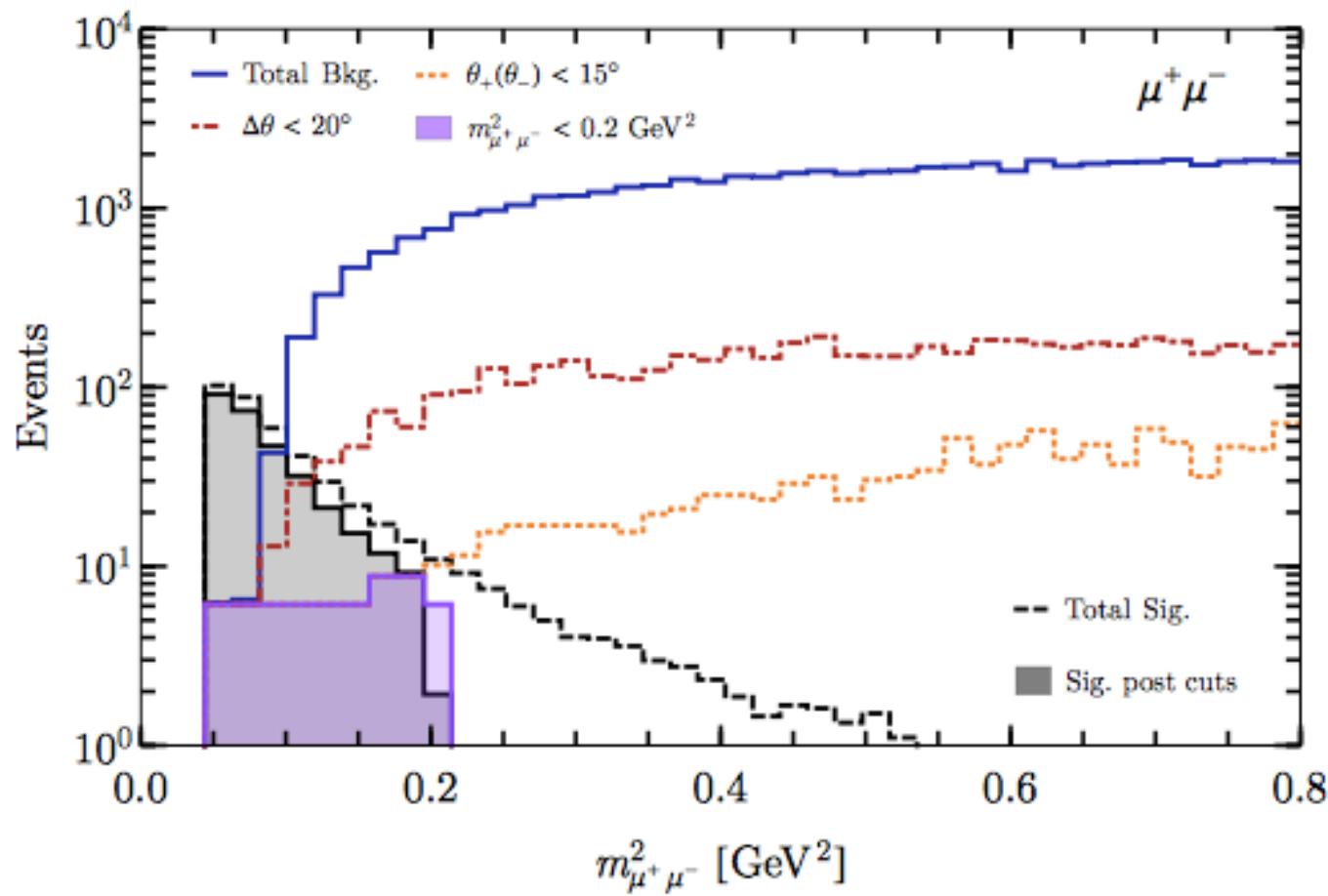


	$N_{\text{tot}}^{\text{CC}}$	$r_{\nu_\mu}^{\text{CC}}$	$r_{\bar{\nu}_\mu}^{\text{CC}}$	$r_{\nu_e}^{\text{CC}}$	$r_{\bar{\nu}_e}^{\text{CC}}$
ν -mode	4.25×10^8	0.964	0.028	0.007	0.001
$\bar{\nu}$ -mode	1.74×10^8	0.201	0.790	0.004	0.005

	$N_{\text{tot}}^{\text{NC}}$	$r_{\nu_\mu}^{\text{NC}}$	$r_{\bar{\nu}_\mu}^{\text{NC}}$	$r_{\nu_e}^{\text{NC}}$	$r_{\bar{\nu}_e}^{\text{NC}}$
ν -mode	1.48×10^8	0.956	0.037	0.006	0.001
$\bar{\nu}$ -mode	7.58×10^7	0.157	0.835	0.003	0.005

Reaching background rates of $O(10^{-6}-10^{-5})$ times the CC rate is necessary to observe trident events at DUNE ND, which is an attainable goal in a LAr detectors.

Ballett, Hostert, Pascoli, Perez, [ZT](#) and Funchal
JHEP **1901**, 119 (2019)



We apply consecutive cuts on the background, starting with cuts on the separation angle $\Delta\theta$ (red), both charged lepton angles to the beamline (θ_+ and θ_-) (orange) and the invariant mass.

Ballett, Hostert, Pascoli, Perez, [ZT](#) and Funchal
JHEP **1901**, 119 (2019)

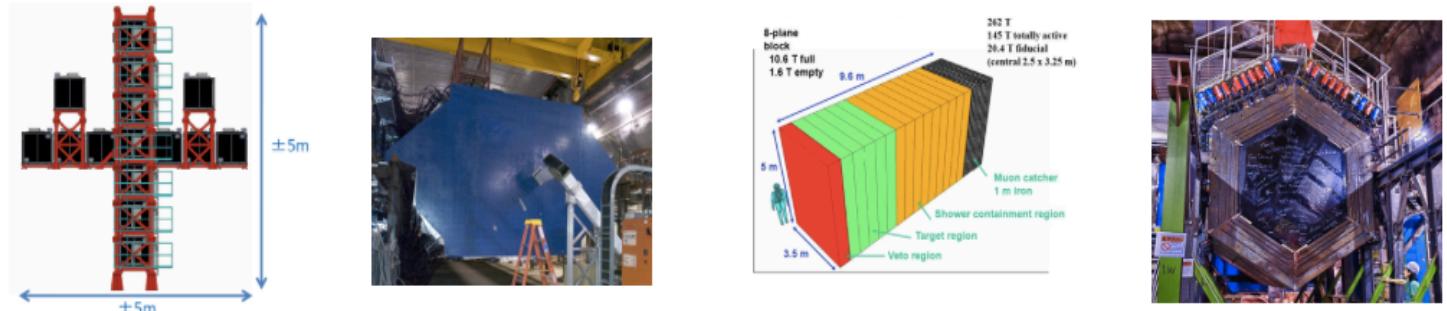
Trident rates at other Near Detectors

Experiment	Material	Baseline (m)	Exposure (POT)	Fiducial Mass (t)	E_ν (GeV)
INGRID	Fe	280	3.9×10^{21} [10^{22}] T2K-I [T2K-II]	99.4	0 – 4
MINOS[+]	Fe and C	1040	$10.56(3.36)[9.69] \times 10^{20}$	28.6	0 – 20
NO ν A	C ₂ H ₃ Cl and CH ₂	1000	$8.85(6.9) [36(36)] \times 10^{20}$ [NO ν A-II]	231	0 – 20
MINER ν A	CH, H ₂ O, Fe, Pb, C	1035	$12(12) \times 10^{20}$	7.98	0 – 20

All have finished data taking or are still running

Trident rates at other Near Detectors

INGRID



Channel	T2K-I	T2K-II	MINOS	MINOS+	NO ν A-I	NO ν A-II	MINER ν A
Total $e^\pm\mu^\mp$	563	1444	222 (56)	730	83 (72)	340 (374)	149 (102)
	96	246	46 (11)	151	25 (22)	102 (114)	56 (39)
Total e^+e^-	277	711	61 (15)	62	29 (22)	119 (114)	39 (27)
	24	62	9 (2)	8	4 (4)	16 (21)	10 (7)
Total $\mu^+\mu^-$	30	76	26 (6)	86	9 (9)	37 (47)	18 (13)
	21	54	15 (3)	49	8 (8)	34 (36)	18 (13)

Coherent (upper) and diffractive (lower) trident events for (anti)neutrino mode.

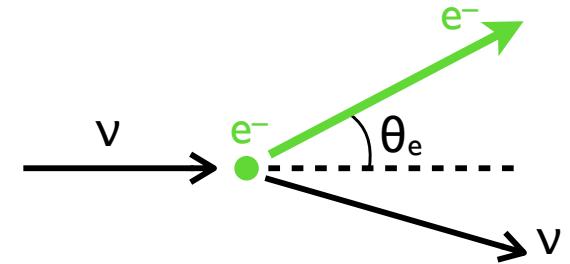
Ballett, Hostert, Pascoli, Perez, [ZT](#) and Funchal

JHEP 1901, 119 (2019)

Neutrino-Electron Scattering

$$\begin{aligned} \frac{d\sigma}{dE_R} &= \frac{2G_F^2 m_e}{\pi} \left\{ g_1^2 + g_2^2 \left(1 - \frac{E_R}{E_\nu}\right)^2 - g_1 g_2 \frac{m_e E_R}{E_\nu^2} \right\} \\ &\simeq 1.72 \times 10^{-41} \left\{ g_1^2 + g_2^2 \left(1 - \frac{E_R}{E_\nu}\right)^2 \right\} \frac{\text{cm}^2}{\text{GeV}} \end{aligned}$$

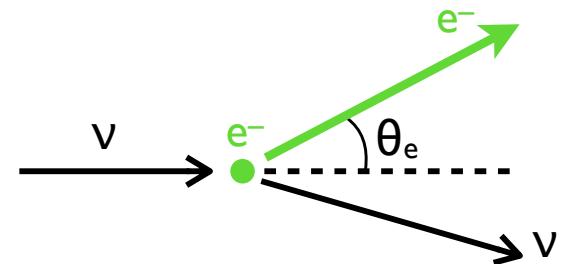
We translate neutrino-electron scattering measurements into a determination of the weak mixing angle at low scales.



ν_α	g_1	$g_1(\text{SM})$	g_2	$g_2(\text{SM})$
ν_e	$1 + (g_V + g_A)/2$	$1/2 + s_W^2$	$(g_V - g_A)/2$	s_W^2
$\nu_{\mu, \tau}$	$(g_V + g_A)/2$	$-1/2 + s_W^2$	$(g_V - g_A)/2$	s_W^2
$\bar{\nu}_e$	$(g_V - g_A)/2$	s_W^2	$1 + (g_V + g_A)/2$	$1/2 + s_W^2$
$\bar{\nu}_{\mu, \tau}$	$(g_V - g_A)/2$	s_W^2	$(g_V + g_A)/2$	$-1/2 + s_W^2$

Neutrino-Electron Scattering

Process	Total cross-section
$\nu_e + e^-$	$(G_F^2 s/4\pi) \left[(1 + 2 \sin^2 \vartheta_W)^2 + \frac{4}{3} \sin^4 \vartheta_W \right] \simeq 93 \text{ s/MeV}^2$
$\bar{\nu}_e + e^-$	$(G_F^2 s/4\pi) \left[\frac{1}{3} (1 + 2 \sin^2 \vartheta_W)^2 + 4 \sin^4 \vartheta_W \right] \simeq 39 \text{ s/MeV}^2$
$\nu_{\mu,\tau} + e^-$	$(G_F^2 s/4\pi) \left[(1 - 2 \sin^2 \vartheta_W)^2 + \frac{4}{3} \sin^4 \vartheta_W \right] \simeq 15 \text{ s/MeV}^2$
$\bar{\nu}_{\mu,\tau} + e^-$	$(G_F^2 s/4\pi) \left[\frac{1}{3} (1 - 2 \sin^2 \vartheta_W)^2 + 4 \sin^4 \vartheta_W \right] \simeq 13 \text{ s/MeV}^2$



$$s = 2 m_e E_\nu$$

The approximate ratios of the four cross-sections:

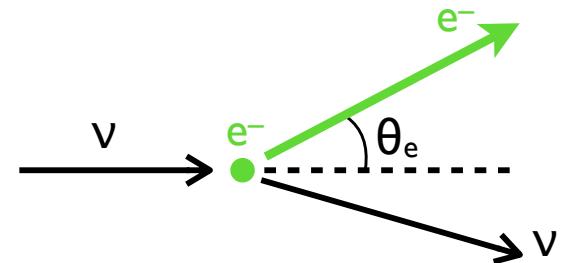
$$\sigma_{\nu_e} : \sigma_{\bar{\nu}_e} : \sigma_{\nu_{\mu,\tau}} : \sigma_{\bar{\nu}_{\mu,\tau}} \simeq 1 : 0.42 : 0.16 : 0.14$$

Neutrino-Electron Scattering

How many events at DUNE ND?

Channel	On axis
$\nu_\mu e \rightarrow \nu_\mu e$	27,705 2,888
$\bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$	2,066 21,926
$\nu_e e \rightarrow \nu_e e$	2,234 782
$\bar{\nu}_e e \rightarrow \bar{\nu}_e e$	226 705

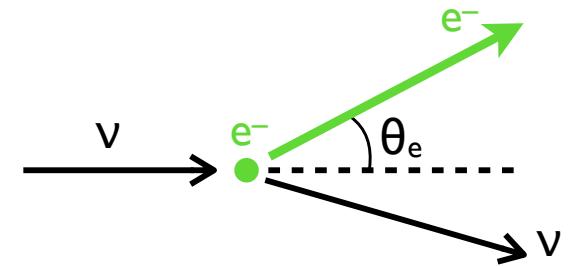
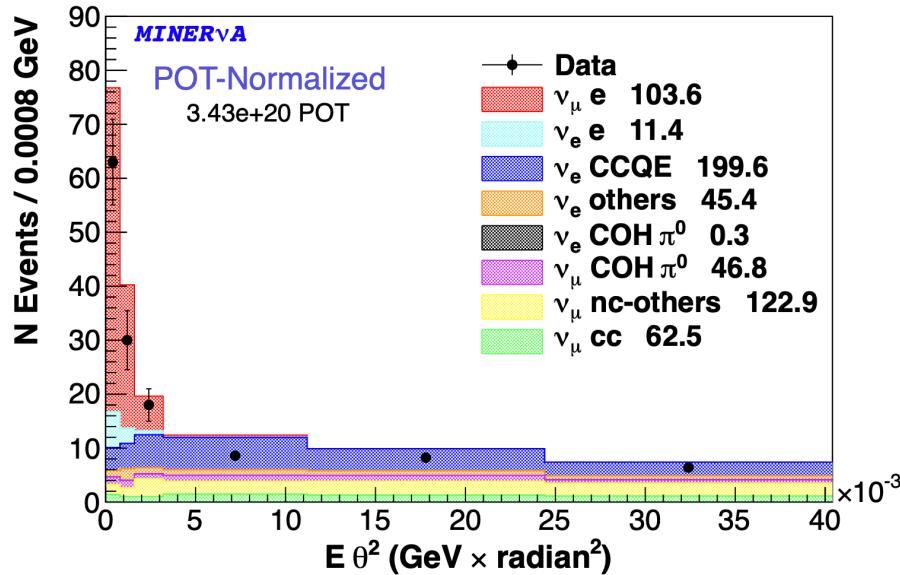
← neutrino beam
← anti-neutrino beam



Approximately 60,000 nu-e events for 75 tonnes-7 years!

Neutrino-Electron Scattering

Signature: forward going electrons!

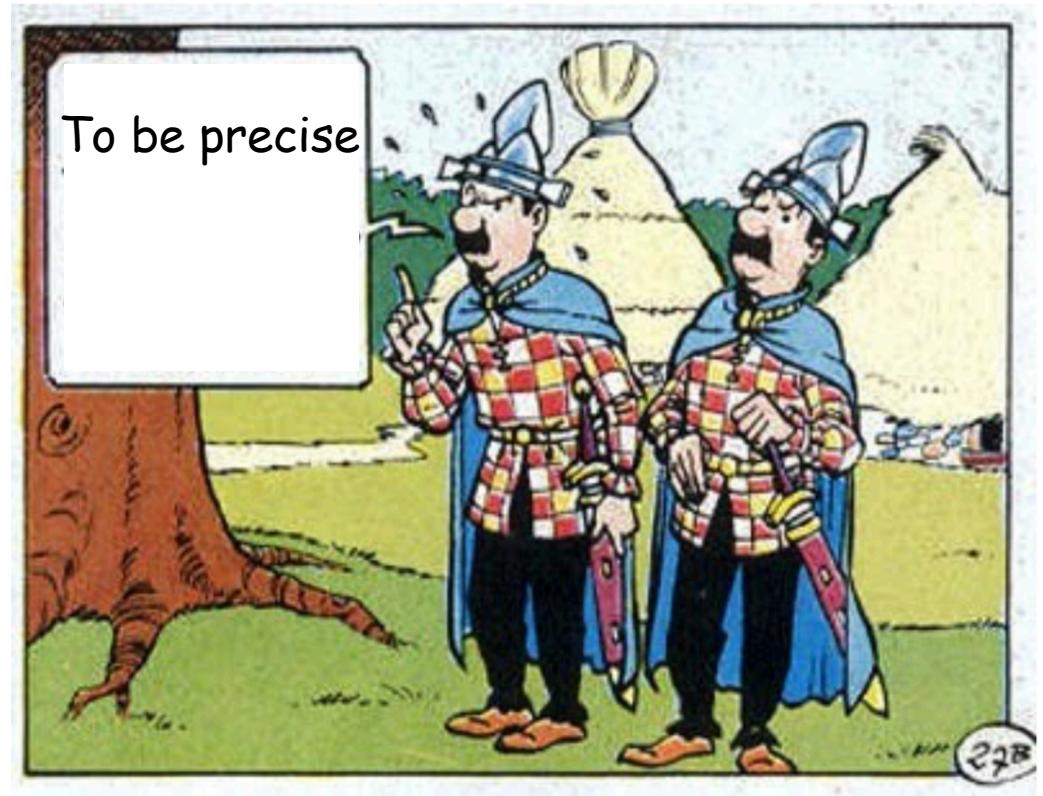


The main source of background is CC quasi-elastic (ν_e CCQE) ν_e scattering.

A cut on the angular distribution can suppress the background!

Jaewon Park, PhD thesis

Measuring the Weak Mixing Angle at DUNE-PRISM



Gouvêa, Machado, Perez-Gonzalez and ZT
arXiv:1912.06658

Measuring the Weak Mixing Angle at DUNE-PRISM

Gouvêa, Machado, Perez-Gonzalez and [ZT](#)
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$$\frac{d\sigma}{dE_R} \simeq 1.72 \times 10^{-41} \left\{ g_1^2 + g_2^2 \left(1 - \frac{E_R}{E_\nu} \right)^2 \right\} \frac{\text{cm}^2}{\text{GeV}}$$

ν_α	g_1	$g_1(\text{SM})$	g_2	$g_2(\text{SM})$
ν_e	$1+(g_V+g_A)/2$	$1/2+s_W^2$	$(g_V-g_A)/2$	s_W^2
$\nu_{\mu,\tau}$	$(g_V+g_A)/2$	$-1/2+s_W^2$	$(g_V-g_A)/2$	s_W^2
$\bar{\nu}_e$	$(g_V-g_A)/2$	s_W^2	$1+(g_V+g_A)/2$	$1/2+s_W^2$
$\bar{\nu}_{\mu,\tau}$	$(g_V-g_A)/2$	s_W^2	$(g_V+g_A)/2$	$-1/2+s_W^2$

There is an exact degeneracy in the differential cross section for $\nu_\mu - e$ scattering under the transformations:

$$(g_V, g_A) \rightarrow (g_A, g_V) \quad \text{and} \quad (g_V, g_A) \rightarrow (-g_V, -g_A)$$

Measuring the Weak Mixing Angle at DUNE-PRISM

Gouvêa, Machado, Perez-Gonzalez and [ZT](#)
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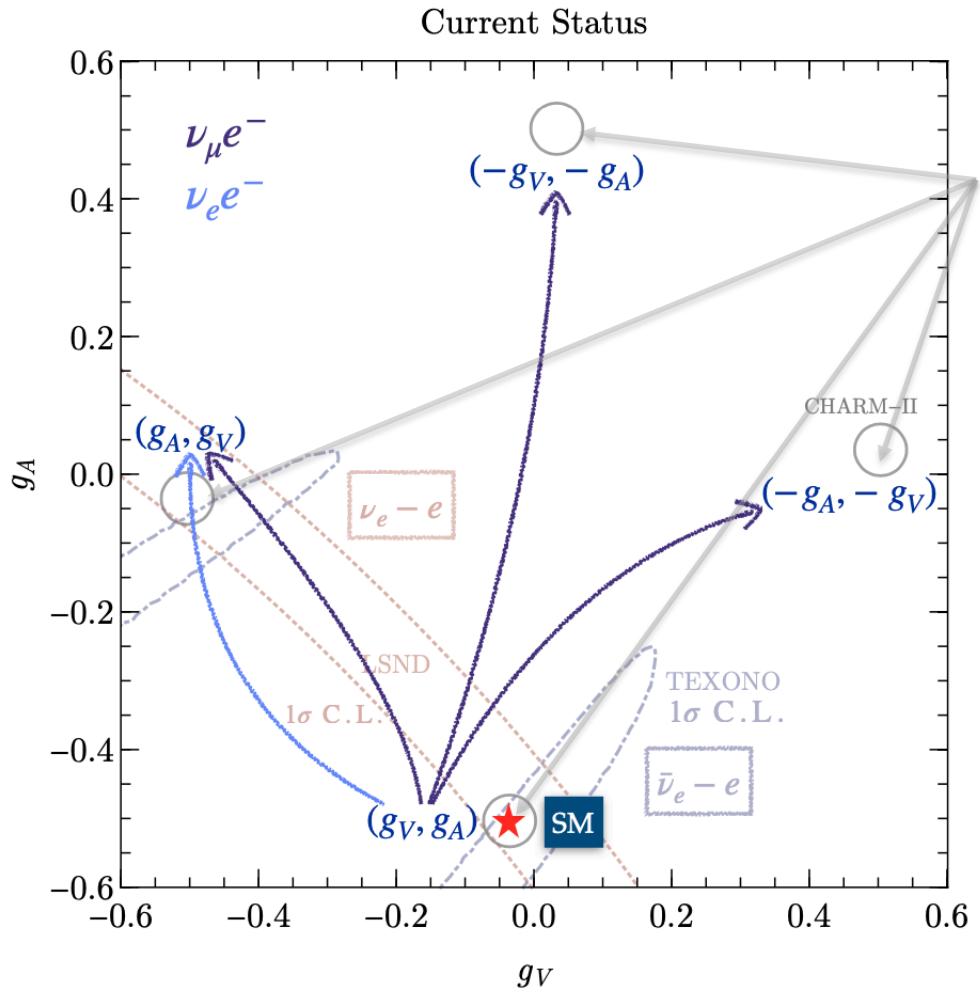
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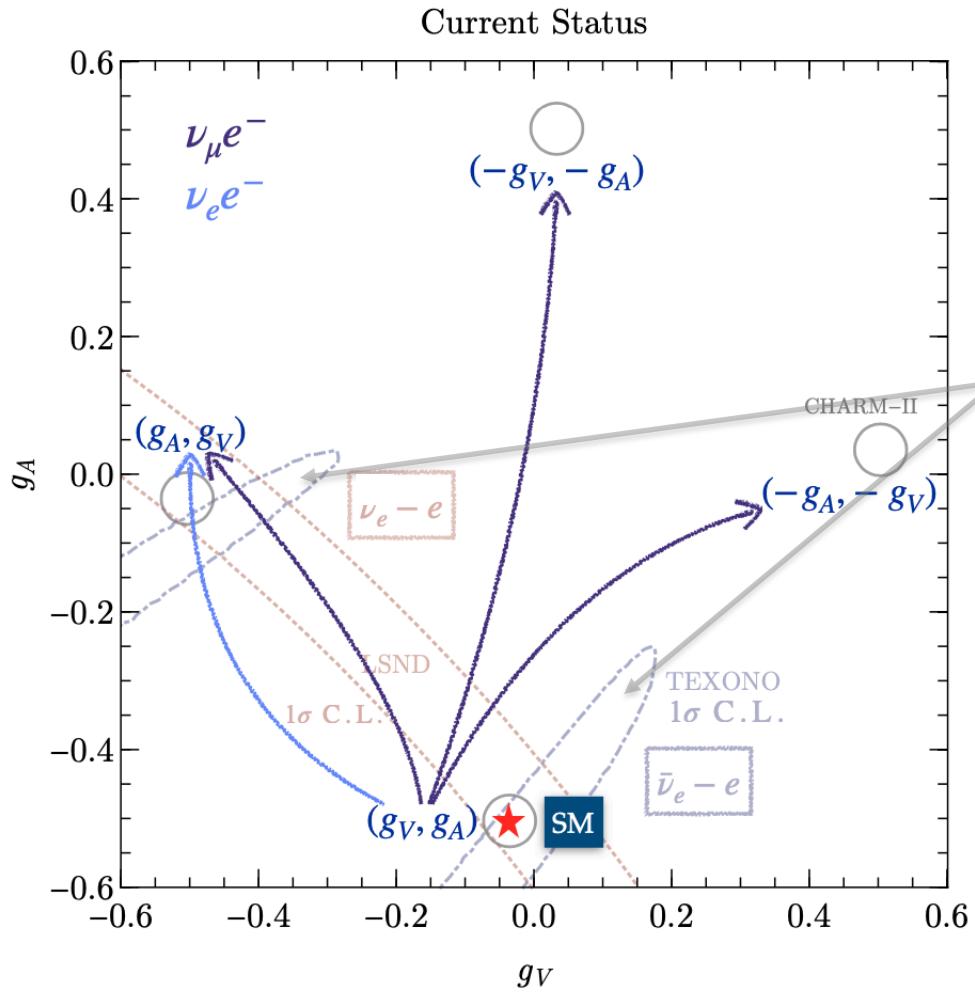
There are half solutions for $\nu_e - e$ scattering:

$$(g_V, g_A) \rightarrow (g_A, g_V)$$



CHARM-II measured the couplings from the scattering of muon-neutrinos on electrons

P. Vilain et al. (CHARM-II),
Phys. Lett. B335, 246 (1994).

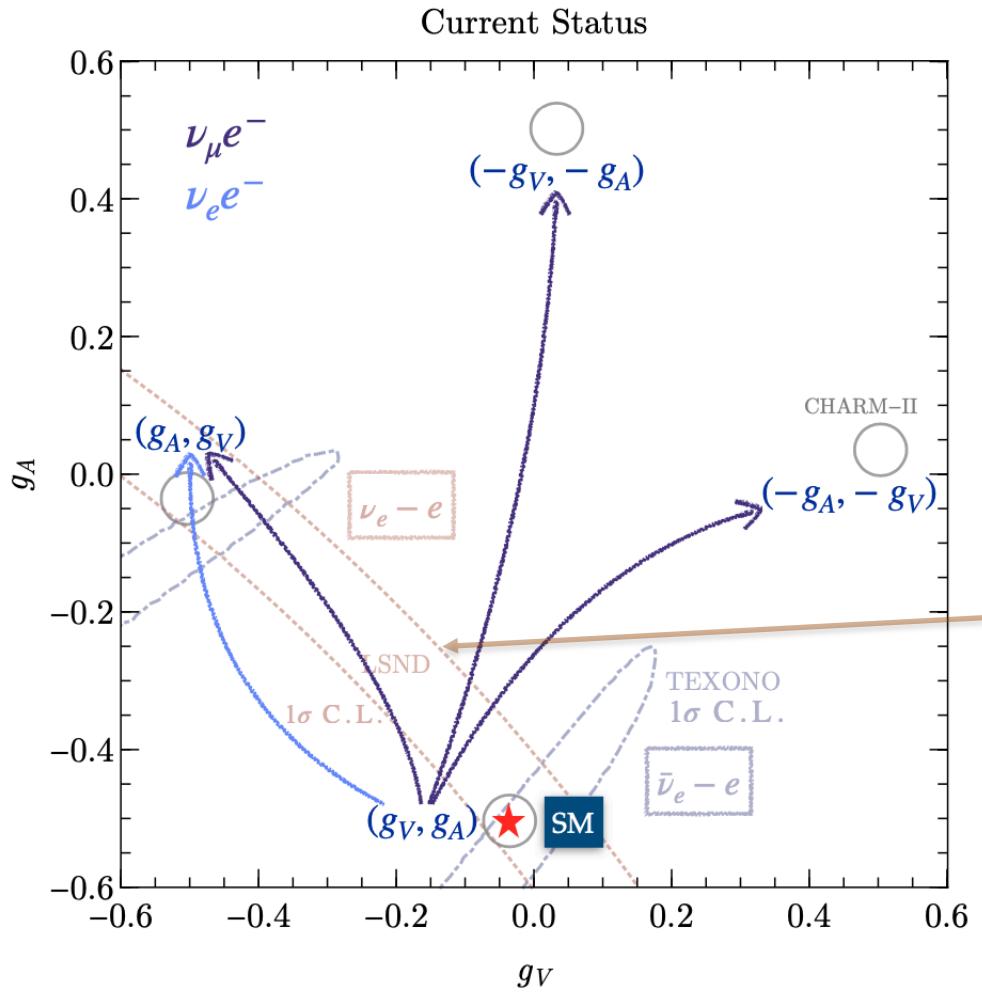


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TEXONO measured electron recoils from electron anti-neutrinos in a nuclear reactor.

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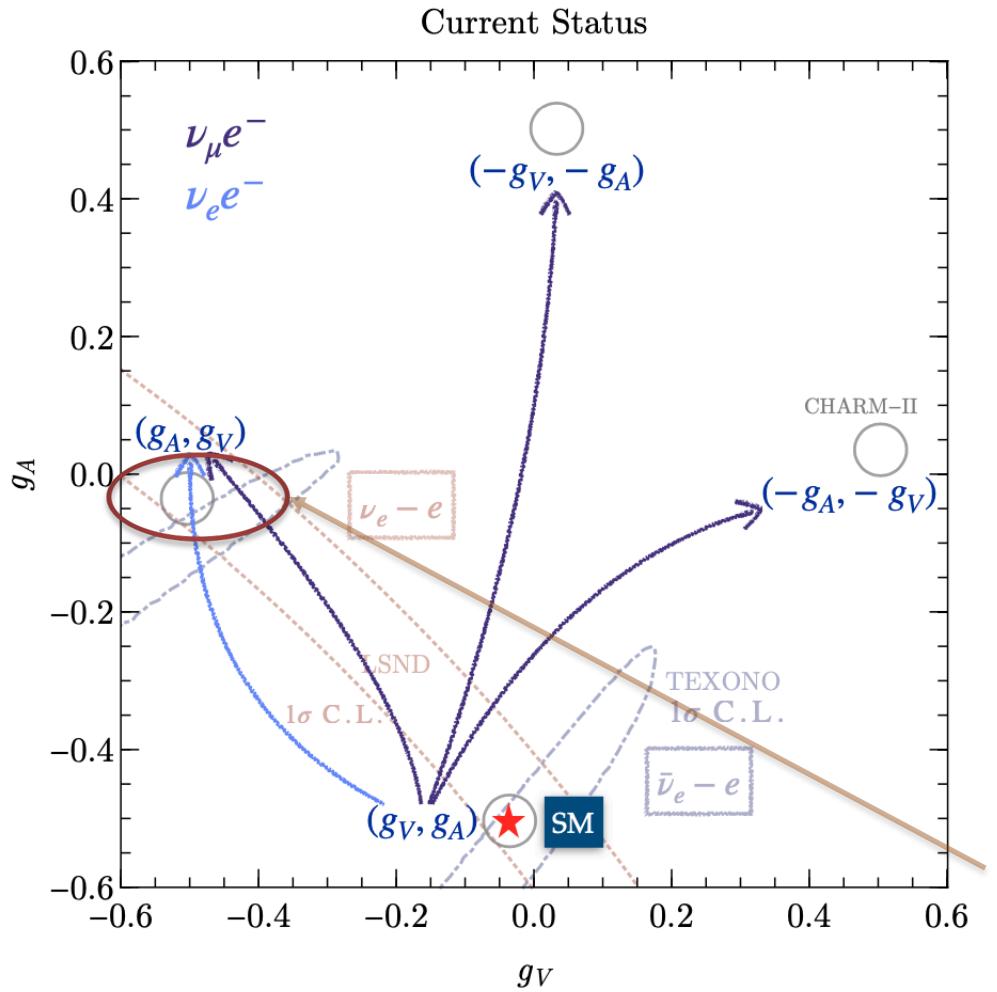
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LSND is similar, but with electron-neutrinos.

L. B. Auerbach et al. (LSND),
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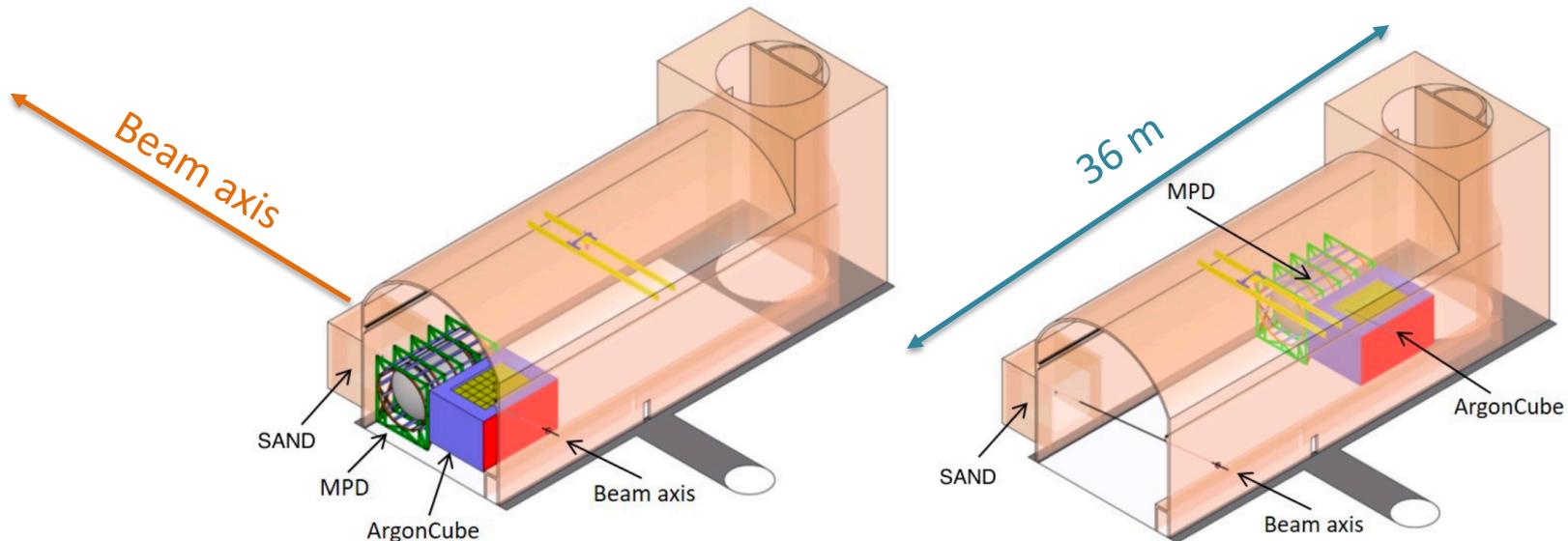
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Current data are not able to rule out very small g_A and $g_V \sim -0.5$.

Measuring the Weak Mixing Angle at DUNE-PRISM

DUNE-PRISM: a near detector that is capable of moving in the direction perpendicular to the neutrino-beam axis.

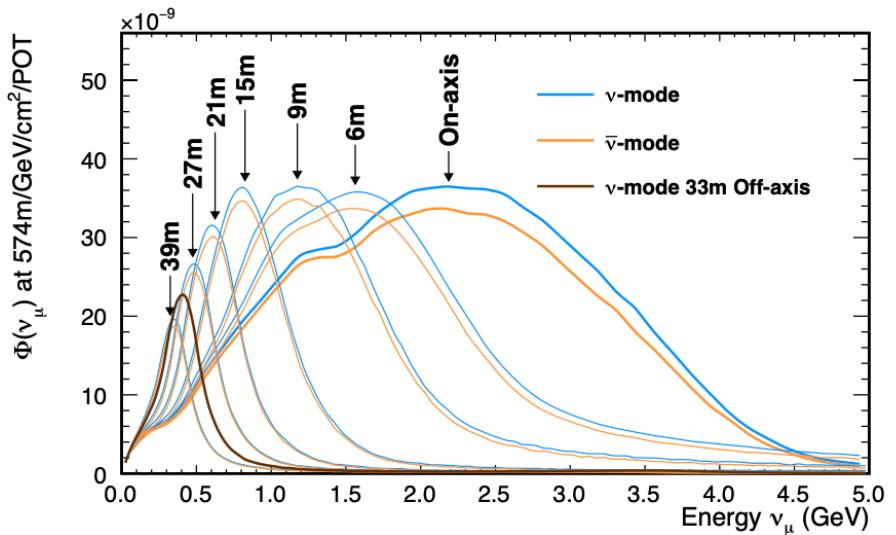


DUNE TDR, arXiv: 2002.03005

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Although the neutrino flux has prohibitively large uncertainties, the ratios of on-axis to off-axis fluxes are dictated only by meson-decay kinematics and thus are much better understood.



DUNE TDR, arXiv: 2002.03005

Measuring the Weak Mixing Angle at DUNE-PRISM

de Gouvêa, Machado, Perez-Gonzalez and [ZT](#)
arXiv:1912.06658 (In press)

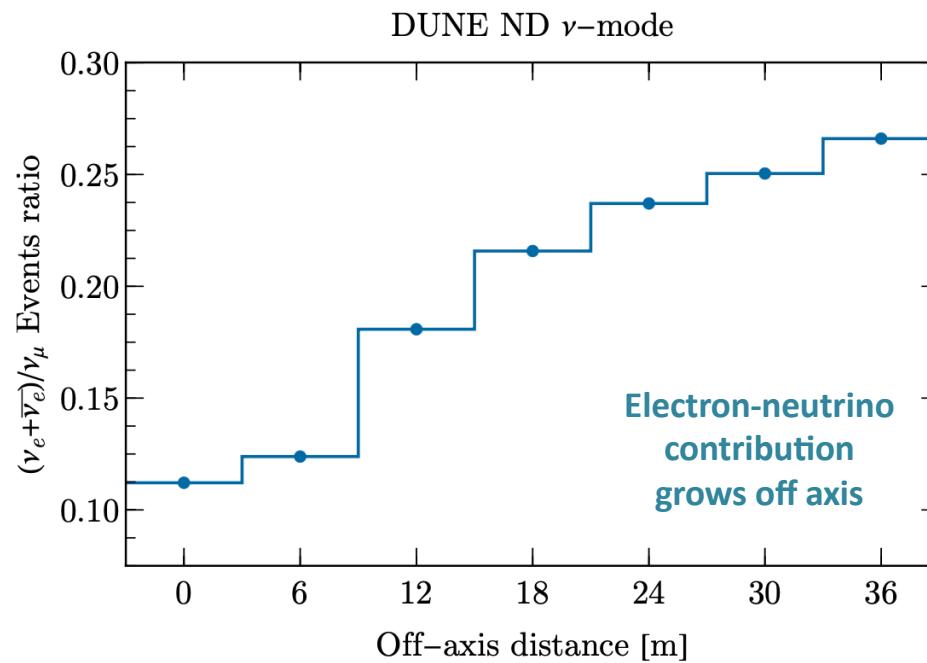
Less statistics in the off-axis locations:

Channel	0 m	6 m	12 m	18 m	24 m	30 m	36 m
$\nu_\mu e \rightarrow \nu_\mu e$	3,958 413	2,671 352	882 170	398 96	213 60	127 43	83 31
$\bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e$	295 3,132	232 2,068	114 679	64 284	40 151	26 90	20 59
$\nu_e e \rightarrow \nu_e e$	319 112	226 88	132 63	76 38	45 24	29 17	19 13
$\bar{\nu}_e e \rightarrow \bar{\nu}_e e$	32 101	31 78	20 41	13 24	9 14	6 9	4 6

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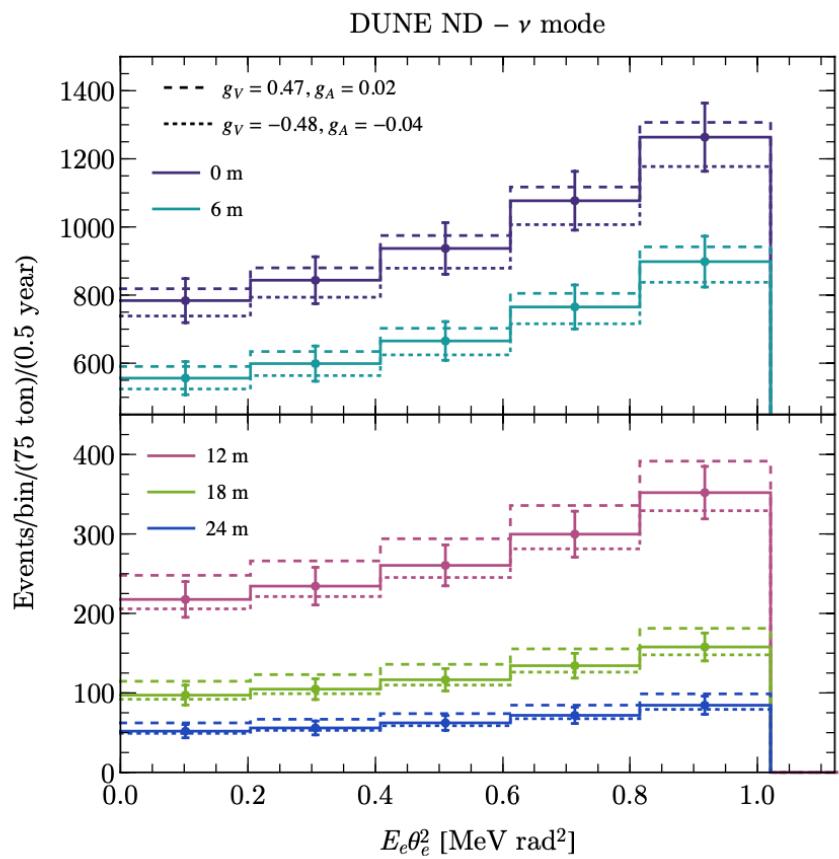
The relevance of the ν_e events grows significantly with the off-axis angle.



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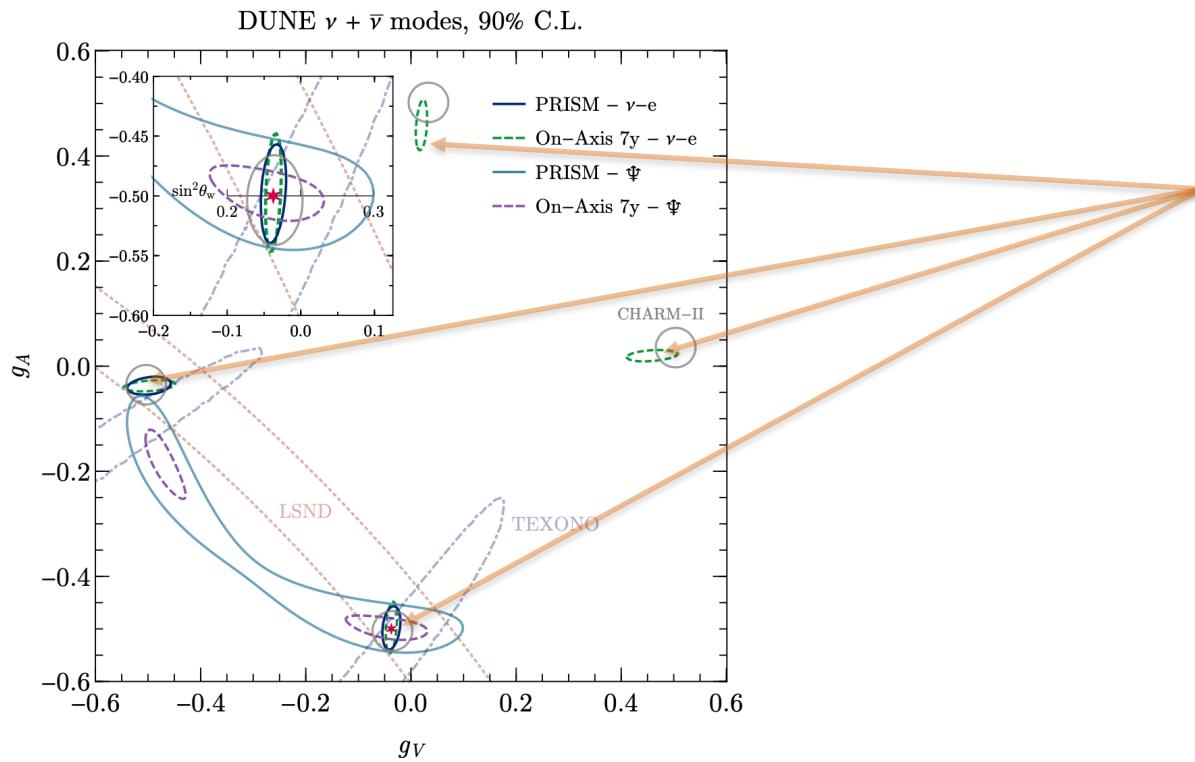
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Error bars illustrating the statistical and systematic errors, are included for the SM case.



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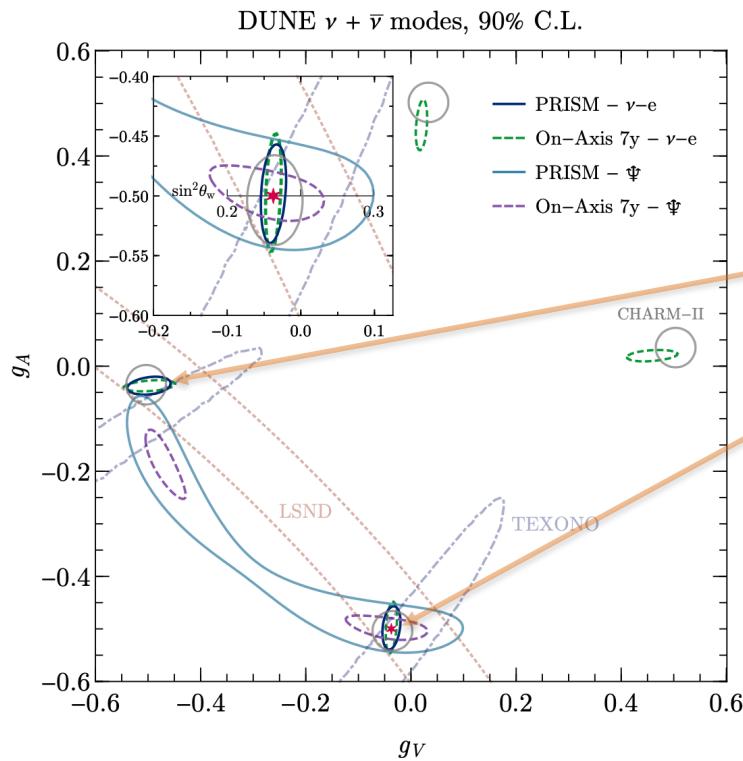
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Both DUNE on-axis and CHARM-II have almost pure ν_μ flux and suffer from a four-fold degeneracy.

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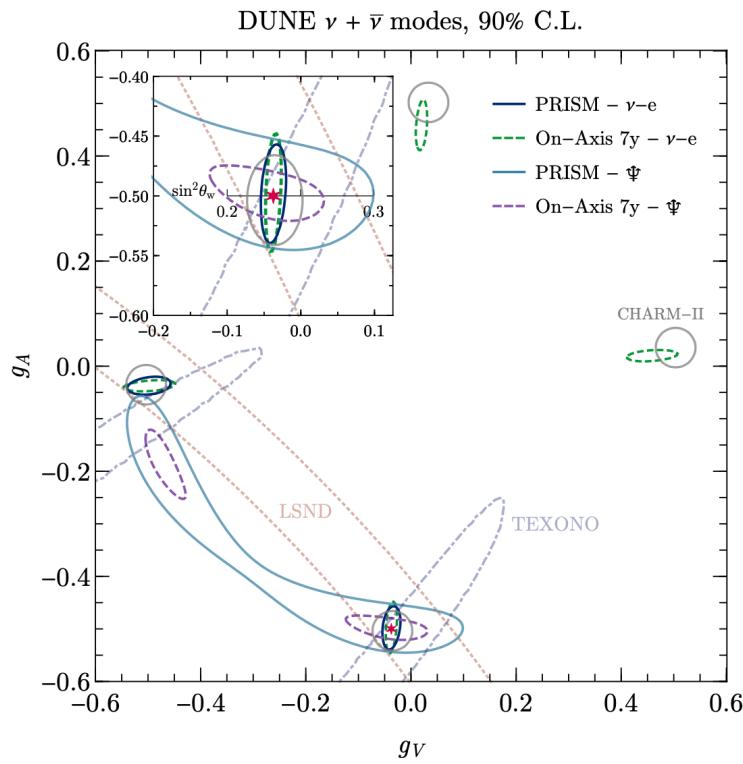
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The sub-dominant ν_e beam component in DUNE-PRISM can break half the degeneracy!

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[arXiv:1912.06658](#)



The trident cross section is more involved, in the limit where the muon mass vanishes, all cross sections are invariant under $g_V \leftrightarrow g_A$.

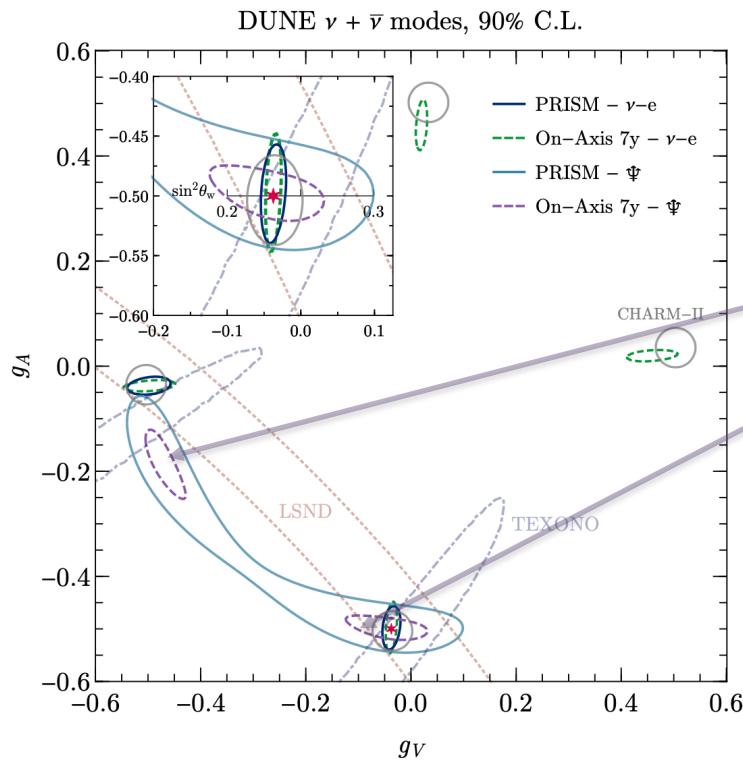
$$m_e, m_\mu \rightarrow 0$$

$$\sigma \sim g_V^2 + g_A^2 \quad (\nu_\mu e^+ e^-)$$

$$\sigma \sim (g_V + 1)^2 + (g_A + 1)^2 \quad (\nu_\mu \mu^+ \mu^-)$$

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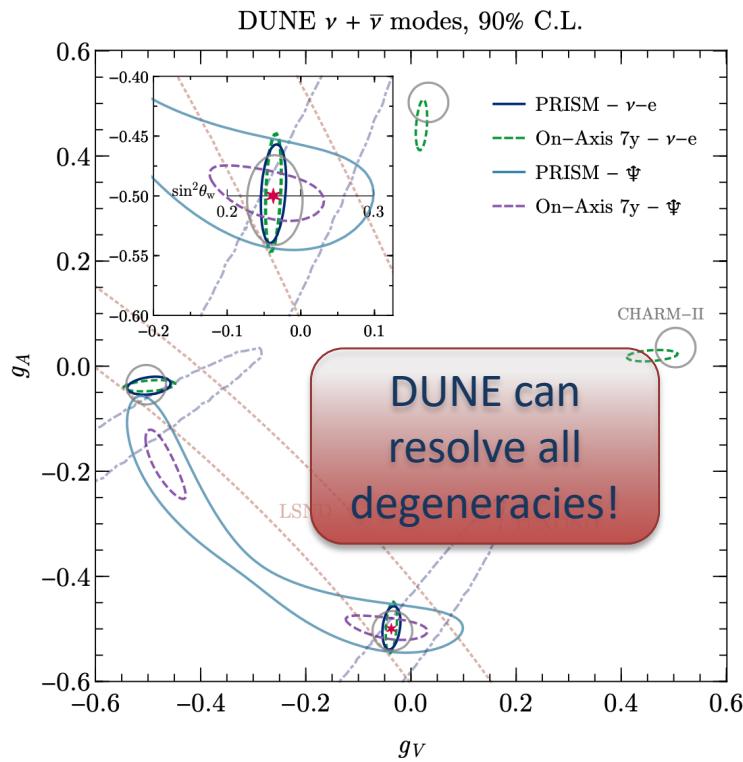
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Lepton masses break the final degeneracy!

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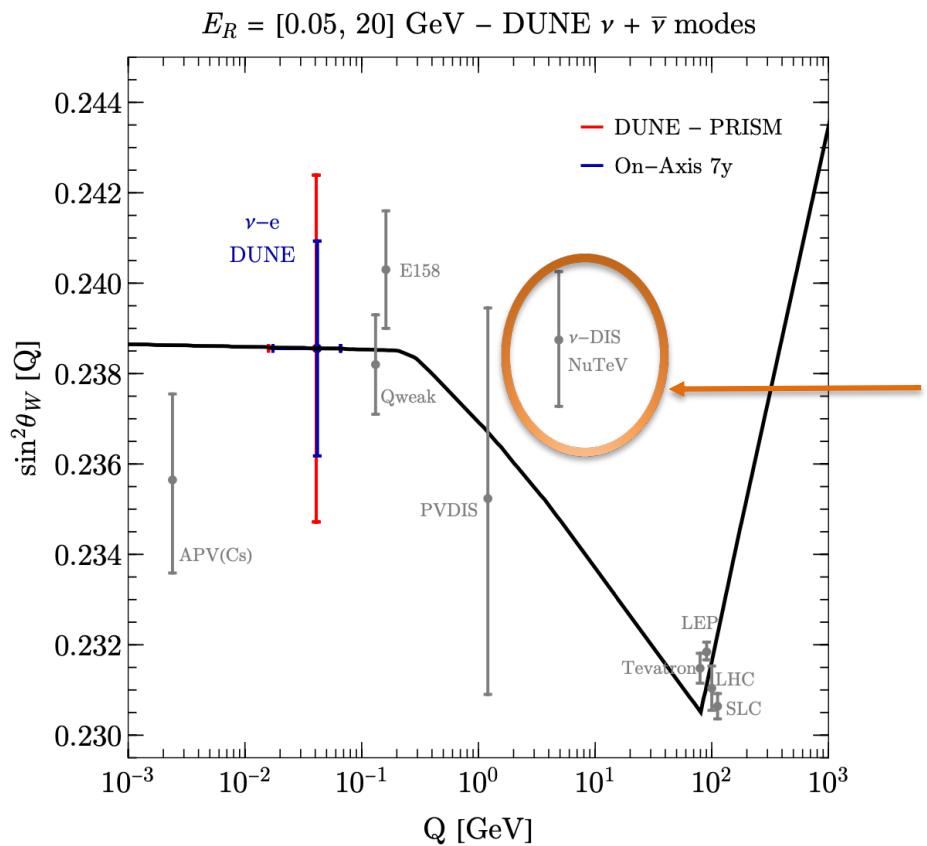
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We translate the neutrino–electron and trident scattering measurements into a determination of the vector and axial couplings of the electron to the Z-boson and the weak mixing angle at low scales.

in the modified minimal subtraction scheme:

$$\sin^2 \theta_W(\mu) \equiv \frac{g'^2(\mu)}{g^2(\mu) + g'^2(\mu)}$$

Measuring the Weak Mixing Angle at DUNE-PRISM



The most precise measurement of $\sin^2 \theta_W$ using neutrino scattering, at $\langle Q \rangle \simeq 4.5 \text{ GeV}$.

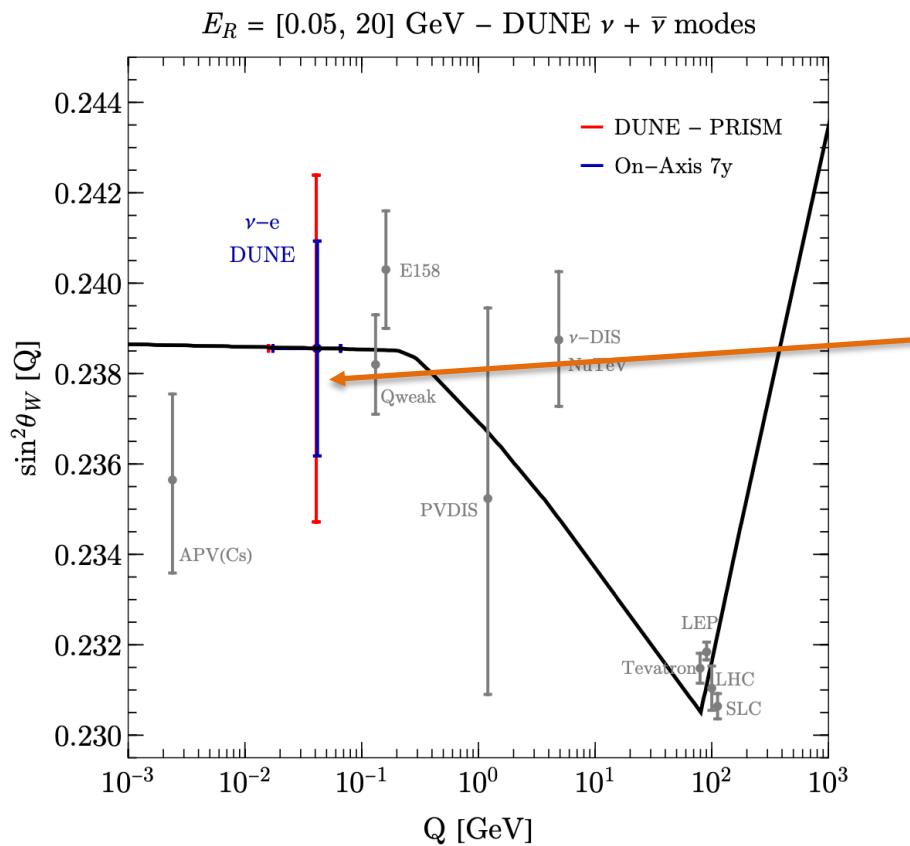
Deviates from the LEP measurement at 3σ level.

$$R^{\nu(\bar{\nu})} = \frac{\sigma(\nu(\bar{\nu})N \rightarrow \nu(\bar{\nu})X)}{\sigma(\nu(\bar{\nu})N \rightarrow \ell^{-(+)X})} \approx g_L^2 + 2g_R^2$$

$$\sin^2 \theta_W (\langle Q^2 \rangle = 20 \text{ GeV}^2) = 0.2277 \pm 0.0013 \pm 0.0009$$

G. P. Zeller et al. (NuTeV),
Phys. Rev. Lett. **88**, 091802 (2002)

Measuring the Weak Mixing Angle at DUNE-PRISM

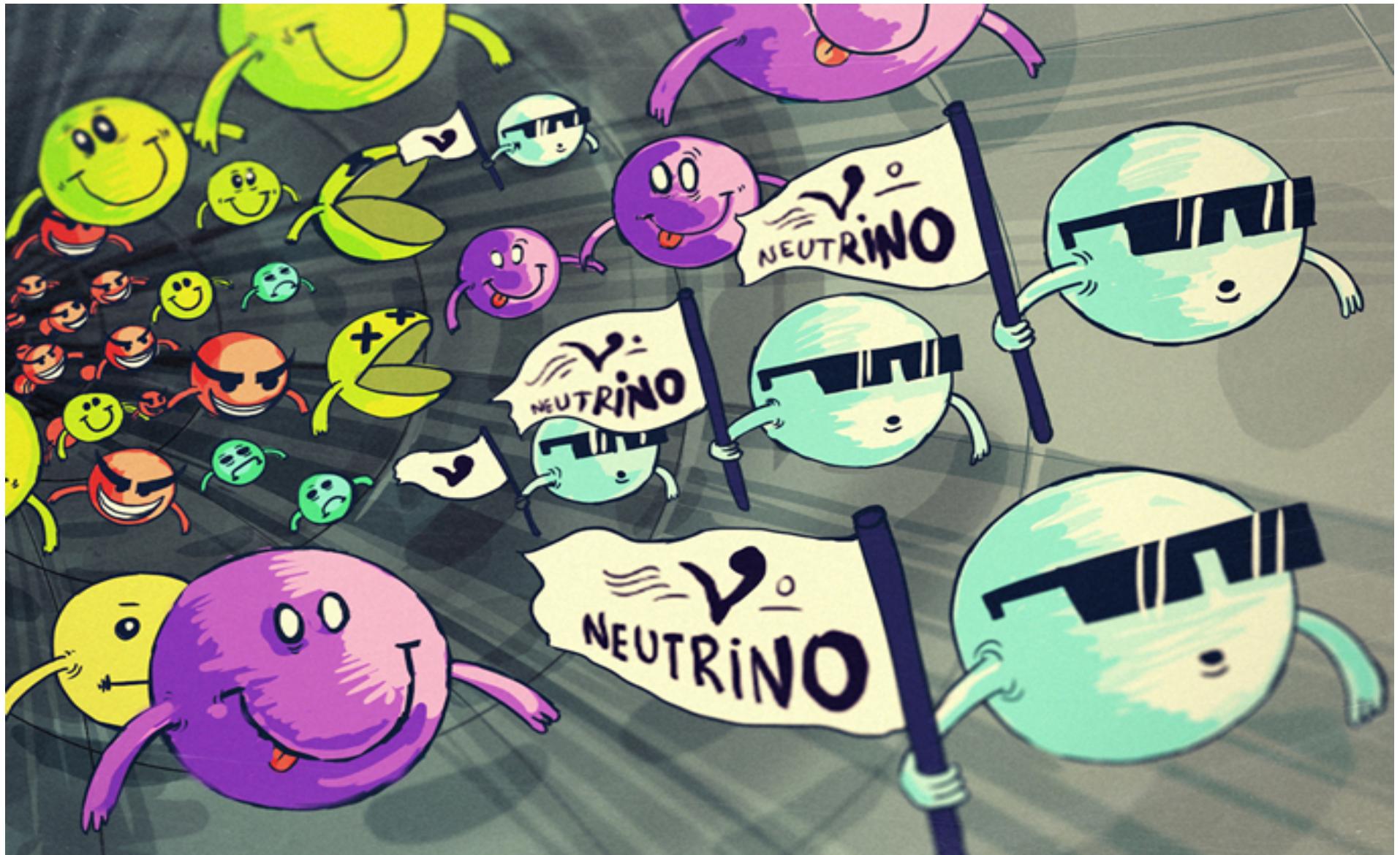


DUNE-ND can be used to measure $\sin^2 \theta_W$ with better than 2% precision, at $\langle Q^2 \rangle = (40 \text{ MeV})^2$

de Gouvêa, Machado, Perez-Gonzalez and [ZT](#)
arXiv:1912.06658 (In press)

Conclusion:

- The future DUNE experiment opens up the possibility to perform many measurements of rare neutrino processes at near detectors.
- We study Neutrino trident and Neutrino-Electron scatterings at DUNE.
- We investigate the sensitivity of DUNE-PRISM, and find that it will qualitatively impact our ability to constrain the weak couplings of the electron.
- Trident measurements could contribute to break the final degeneracy.
- DUNE can achieve competitive measurements and test (g_V, g_A) without the aid of external data.
- The DUNE near-detector can be used to measure $\sin^2\theta_W$ with better than 2% precision.



Thanks for your attention