

## Linear algebra

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For unitary matrix  $U$ , (normalised) eigenvectors  $|v_i\rangle$  and eigenvalues  $\lambda_i$ :  $U|v_i\rangle = \lambda_i|v_i\rangle$ .

For diagonalisable matrix, spectral decomposition  $U = \sum_{i=1}^n \lambda_i |v_i\rangle \langle v_i|$ .

Unitary  $A^\dagger A = I$ , Hermitian  $A = A^\dagger \subseteq$  normal matrices  $A^\dagger A = AA^\dagger$ .

## Postulates of quantum mechanics

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Superposition, interference

Entanglement: non-separability

## Concepts in quantum mechanics

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Measurement and the Helstrom-Holevo bound  $p \leq \frac{1+\sin\theta}{2}$ , where  $|\langle\psi_a|\psi_b\rangle| = \cos\theta$ .

The no-signalling principle: after measurement, the entanglement is collapsed, thus not possible to transmit information.

The no-cloning principle: impossible to copy an unknown quantum state.  $\nexists U. U(|\psi\rangle|0\rangle) = |\psi\rangle|\psi\rangle$ .

The no-deleting principle: impossible to delete one of the unknown quantum state copies.

$\nexists \tilde{U}. \tilde{U}(|\psi\rangle|\psi\rangle) = |\psi\rangle|0\rangle$ .

## Quantum circuits

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Universal gate set:  $\{H, T, CNOT\}$ , where  $\pi/8$  gate is  $T$ .  $\pi/4$  gate is  $S$ . (not self-invertible)

Phase gate  $S = T^2$ ,  $Z = S^2$ ,  $X = HZH$ ,  $Y = iXZ = SXSZ$ .

- proof for  $Z = HXH$  (L8. search)
  - either by matrix multiplication.
  - or geometric interpretation ( $X/Z$ : rotate 180 degree about x/z-axis,  $H$ : swap x and z axis).

by self-inverse,  $CNOT = CX = (I \otimes H) \times CZ \times (I \otimes H)$ .

SWAP can be decomposed into 3 CNOTs.

Entanglement circuits via Hadamard-CNOT combination  $CNOT(H \otimes I)|00\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$

## Quantum information applications

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Teleportation (send a qubit via two bits)

Superdense coding (send **two bits** via one qubit)  $\{I, X, Z, XZ\} \rightarrow \text{CNOT} + \text{Hadamard}$ .

## Deutsch-Jozsa algorithm

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$f : \{0, 1\}^n \rightarrow \{0, 1\}$ , which is either constant or balanced.

$$H^{\otimes n}|x\rangle = \frac{1}{\sqrt{2^n}} \sum_{z \in \{0,1\}^n} (-1)^{x \cdot z} |z\rangle$$

Proof: as  $|x\rangle = |x_1 \dots x_n\rangle$ , where  $x_i \in \{0, 1\}$  and

$$\begin{aligned} H|x_i\rangle &= \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x_i}|1\rangle) \\ &= \frac{1}{\sqrt{2}}(|z_1 = 0\rangle + (-1)^{x_i}|z_j = 1\rangle) \\ &= \frac{1}{\sqrt{2}}((-1)^{x_i \times 0}|z_1 = 0\rangle + (-1)^{x_i \times 1}|z_2 = 1\rangle) \\ &= \frac{1}{\sqrt{2}}((-1)^{x_i \times z_1}|z_1 = 0\rangle + (-1)^{x_i \times z_2}|z_2 = 1\rangle) \\ &= \frac{1}{\sqrt{2}} \sum_{z_j \in \{0,1\}} (-1)^{x_i \times z_j} |z_j\rangle \end{aligned}$$

$H^{\otimes n}|x_1 \dots x_n\rangle = \otimes_i (H|x_i\rangle)$ , and the power of the function is  $\sum_i x_i \times z_i = x \cdot z$ , we are done.

## Quantum Search

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Grover's algorithm

## QFT & QPE

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### Quantum Fourier Transform (QFT)

$|x\rangle \rightarrow |y\rangle: \sum_{j=0}^{N-1} x_j |j\rangle \rightarrow \sum_{k=0}^{N-1} y_k |k\rangle$ , where  $y_k = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} w^{jk} x_j$  and  $w^{jk} = e^{i \frac{2\pi}{N} jk}$ .

In the matrix form, we have the following transformation,

$$\begin{bmatrix} y_0 \\ y_1 \\ y_2 \\ \dots \\ y_N \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & \dots & 1 \\ 1 & \omega & \omega^2 & \dots & \omega^{N-1} \\ 1 & \omega^2 & \omega^4 & \dots & \omega^{2(N-1)} \\ 1 & \dots & \dots & \dots & \dots \\ 1 & \omega^{N-1} & \omega^{2(N-1)} & \dots & \omega^{(N-1)(N-1)} \end{bmatrix} \cdot \begin{bmatrix} x_0 \\ x_1 \\ x_2 \\ \dots \\ x_N \end{bmatrix}, \text{ where } \omega = e^{i \frac{2\pi}{N}}.$$

The dimension of Hilbert space for  $n$  qubits  $N = 2^n$ . The sinusoid's frequency  $f = \frac{k}{N}$ , i.e.,  $k$  cycles per  $N$  samples.

## inverse QFT (iQFT)

$$|y\rangle \rightarrow |x\rangle: \sum_{k=0}^{N-1} y_k |k\rangle \rightarrow \sum_{j=0}^{N-1} x_j |j\rangle, \text{ where } x_j = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} w^{-jk} y_k \text{ and } w^{-jk} = e^{-i\frac{2\pi}{N}jk}.$$

Note that the normalizing terms should be a product of  $\frac{1}{N}$ , where the above satisfies unitary. The exponential term is negated in one of the two.

## Quantum Phase Estimation (QPE)

precision up to  $t$  bits. If given the eigenvector  $|u\rangle$  of  $U$  and eigenvalue  $e^{i2\pi\phi}$  with **phase**  $\phi \in [0, 1)$ , we have  $U|u\rangle = e^{i2\pi\phi}|u\rangle$ .

- preparation
  - $1^{st}$  register:  $H^{\otimes t}|0\rangle^{\otimes t} = \frac{1}{\sqrt{2^t}} \sum_{x \in \{0,1\}^t} |x\rangle$  (superposition)
  - $2^{nd}$  register: the (superposition of) given eigenvector(s)  $|u\rangle$  with eigenvalue  $e^{i2\pi\phi}$ ,
- oracle  $U^j$  on the  $1^{st}$  register (Entanglement)
  - $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \rightarrow \frac{1}{\sqrt{2}}(|0\rangle + (e^{i2\pi\phi})^j|1\rangle)$
  - $\frac{1}{\sqrt{2^t}} \sum_{x=0}^{2^t-1} |x\rangle \rightarrow \frac{1}{\sqrt{2^t}} \sum_{j=0}^{2^t-1} (e^{i2\pi\phi})^j |j\rangle$
  - $2^{nd}$  register: respective  $|u\rangle$  with eigenvalue  $e^{i2\pi\phi}$  and phase  $\phi$ .
- iQFT (Interference)
- measurement
  - $1^{st}$  register:  $t$  bits approximation of  $|\tilde{\phi}\rangle$
  - $2^{nd}$  register:  $|u\rangle$  with phase  $\phi$ .

## Application: factoring

**order finding:** for coprime  $x$  and  $N$ , find  $x^r \equiv 1 \pmod{N}$ , where  $r$  is the least positive integer.

$U|r\rangle = |(x \cdot r) \pmod{N}\rangle \implies$  For eigenstates  $s \in [0, r-1]$ , we have eigenvectors  $|u_s\rangle = \frac{1}{\sqrt{r}} \sum_{j=0}^{r-1} e^{-i2\pi\frac{s}{r}j} |x^j \pmod{N}\rangle$  with **phase**  $\phi = \frac{s}{r}$ .

Use QPE,  $2^{nd}$  register prepared with equal superposition of unknown eigenvectors  $\frac{1}{\sqrt{r}} \sum_{j=0}^{r-1} |u_j\rangle = |1\rangle$  (shallow-depth quantum circuit  $X$ ).

**factoring:** for composite integer  $N$ ,  $N = p \cdot q$ , where  $p$  and  $q$  are prime numbers.

Shor's algorithm

## Application: quantum chemistry

Trotter formula:  $U = e^{-i(H_1+H_2)t} = U_1 U_2 = e^{-iH_1 t} e^{-iH_2 t} + O(t^2)$ , where  $U_1$  and  $U_2$  don't commute.

Projective measurement with (normalized) eigenvectors

**Ground state energy estimation**  $|e_0\rangle$  of a  $H$  with eigenvalue  $\lambda_0 = E_0$ .

Use QPE,  $2^{nd}$  register should be prepared as close to the eigenvector such that it's sufficiently dominated by the ground state  $|e_0\rangle$  (L15. adiabatic state preparation).

## Fault tolerance

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bit-flip, phase-flip, Shor code, Steane code

Fault tolerance threshold  $p_{th} = \frac{1}{c}$ , for suppressed error rate  $p = cp_e^2 + O(p_e^3)$ . Per-gate error rate  $\frac{(cp_e)^{2^k}}{c}$  after  $k$  concatenation.