

# Circuits

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## Sinusoidal Analysis



Spring 2022

# Context

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## DC sources

Up to now, we have only considered loadings from DC sources.

Another classical way is to load circuits with alternating currents (AC).  
For example, sinusoidal sources are AC sources.

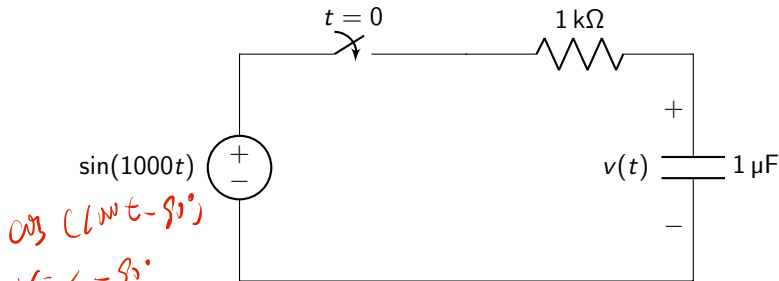
## Does it change anything compared to DC?

No! Most of the laws/theorems we have seen so far still apply.

But, the response of a circuit can be divided into 2 components:

- a **steady-state response**, the part that remains when  $t \rightarrow \infty$
- a **transient response**, the part that vanishes when  $t \rightarrow \infty$

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NYU SHANGHAIInitial condition:  $v(0) = 0$ 

$$V = 1\text{ k}\Omega \cdot I + \frac{1}{j\omega C} I \rightarrow I = \frac{V}{1\text{ k}\Omega + \frac{1}{j\omega C}} = \frac{1}{1000(1 - j)}$$

$$V_C = \frac{1}{j\omega C} I = \frac{V}{j\omega C} \cdot \frac{1}{1000(1 - j)} = \frac{V}{1000(1 - j)} = \frac{4}{1 - j}$$

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## Equation for $t \geq 0$

By applying KVL:

$$v(t) + \tau \frac{dv(t)}{dt} = \sin(1000t)$$

with  $\tau = RC = 1 \text{ ms}$

*Why the phase shift of  $V_{ss}$  is  $45^\circ$ ?  
What's the equivalent DC value of a sin signal?*

## Solution?

The **steady-state solution** is:

$$\blacksquare v_{ss}(t) = \frac{1}{\sqrt{2}} \sin(1000t - 45^\circ)$$

*half of the max abs of  
sin signal*

The **transient solution** is:

$$\blacksquare v_{tr}(t) = \frac{1}{2} e^{-t/\tau}$$

*How do we find the  $V_0$  here?*

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when timing is infinit. the first part is

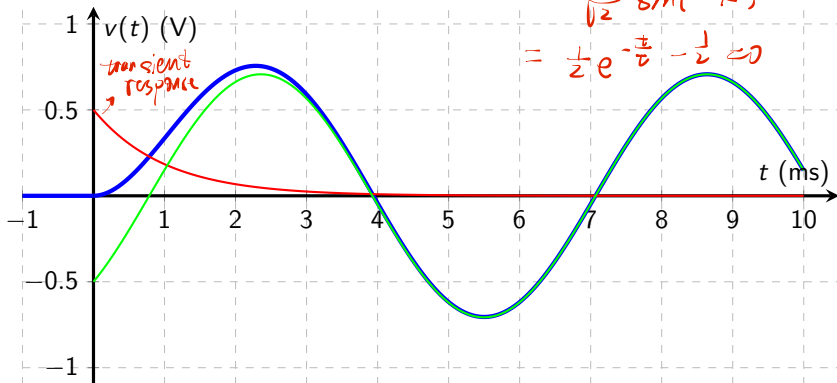
Total solution

almost 0 if  $t$  is very short, the steady state is

$$v(t) = \frac{1}{2}e^{-t/\tau} + \frac{1}{\sqrt{2}}\sin(1000t - 45^\circ)$$

$\frac{1}{\sqrt{2}}\sin(-45^\circ)$

$= \frac{1}{2}e^{-\frac{t}{\tau}} - \frac{1}{2} \approx 0$



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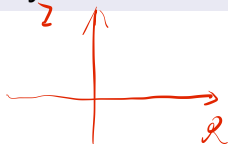
$$a + jb = \sqrt{a^2 + b^2} \angle \arctan \frac{b}{a}$$

## How to determine the solution?

In this lecture, we will only focus on **methods** to determine the **steady-state solution**.

The **transient solution** can be retrieved from **DC analysis** and proper initial conditions.

In this course, when dealing with AC circuits, we will **only consider the steady-state solution**.



# Complex numbers

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## Complex representation

Since we are gonna deal with **sinusoidal loadings and responses**, it will be very convenient (mathematically) to represent **every signal** by a **complex number**.

## Phasor

**Phasor**: a complex number that represents the amplitude  $A$  and the phase  $\phi$  of a sinusoid

$$A/\underline{\phi}$$

# Complex numbers


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## Equivalence

$$z = a + jb$$

$$z = A \angle \phi$$

$$z = Ae^{j\phi}$$

where  $A = \sqrt{a^2 + b^2}$

and  $\phi = \text{angle}(a, b) = \arctan\left(\frac{b}{a}\right)$

$$= A (\cos \phi + j \sin \phi)$$

$$= A \cos \phi + j A \sin \phi$$

$$a = A \cos \phi \quad b = A \sin \phi$$

## Euler's identity

$$e^{j\phi} = \cos(\phi) + j \sin(\phi)$$

$$e^{-j\phi} = \cos(-\phi)$$

$$+ j \sin(-\phi)$$

$$\Rightarrow \cos(\phi) = \text{Re}(e^{j\phi}) = \frac{e^{j\phi} + e^{-j\phi}}{2}$$

$$\Rightarrow \sin(\phi) = \text{Im}(e^{j\phi}) = \frac{e^{j\phi} - e^{-j\phi}}{2j}$$

$$= \cos \phi - j \sin \phi$$



# Complex numbers

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## Phasor transform

Suppose we have a sinusoidal voltage:

$$v(t) = A \cos(\omega t + \phi)$$

*same  $\omega$  in the RLC circuit  
resonant frequency  
oscillating speed*

We can represent that voltage by the phasor  $\mathbf{V}$ :

$$\mathbf{V} = A \angle \phi$$

## Explanations

$$v(t) = A \cos(\omega t + \phi) = \operatorname{Re}(A e^{j(\omega t + \phi)}) = \operatorname{Re}(A e^{j\phi} \cdot e^{j\omega t}) = \operatorname{Re}(\mathbf{V} \cdot e^{j\omega t})$$

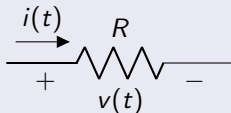
As we are gonna see in the next slides,  $e^{j\omega t}$  will be a common factor for every signal.

So, in the phasor representation, we only consider the amplitude  $A$  and the phase  $\phi$

# Resistor



## Resistor and phasors



Suppose we have a current  $i(t) = A \cos(\omega t + \phi)$   
Then,  $v(t) = Ri(t) = RA \cos(\omega t + \phi)$

When representing with phasors, we have:

$$\mathbf{I} = A/\underline{\phi} \qquad \mathbf{V} = RA/\underline{\phi} = R\mathbf{I}$$

## Ohm's law

Ohm's law is the same with phasors:

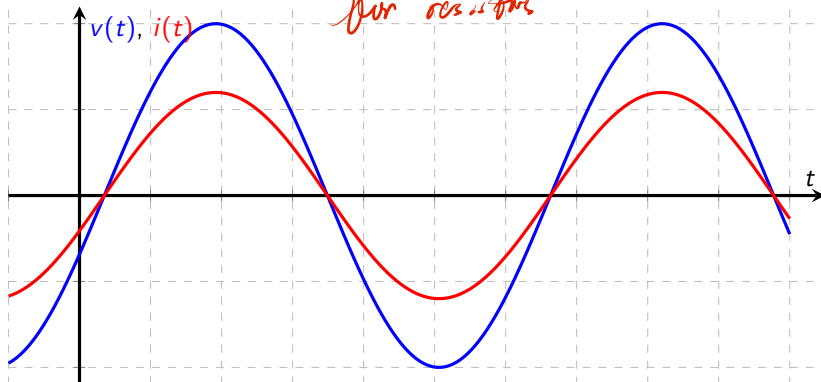
$$\mathbf{V} = R\mathbf{I}$$

# Resistor

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$v(t)$  and  $i(t)$  are in phase

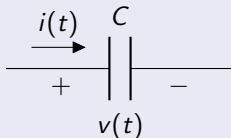
*no phase difference  
between  $v(t)$  &  $i(t)$   
for resistor*



# Capacitor


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## Capacitor and phasors



Suppose we have a voltage  $v(t) = A \cos(\omega t + \phi)$

$$\text{Then, } i(t) = C \frac{dv(t)}{dt} = -CA\omega \sin(\omega t + \phi) = -CA\omega \cos(\omega t + \phi - 90^\circ)$$

$$= -CA\omega \cos(\omega t + \phi + 270^\circ)$$

When representing with phasors, we have:

$$\mathbf{V} = A \angle \phi \qquad \mathbf{I} = -CA\omega e^{j(\phi - 90^\circ)} = -CA\omega e^{j\phi} e^{-j90^\circ}$$

## Phasor relationship

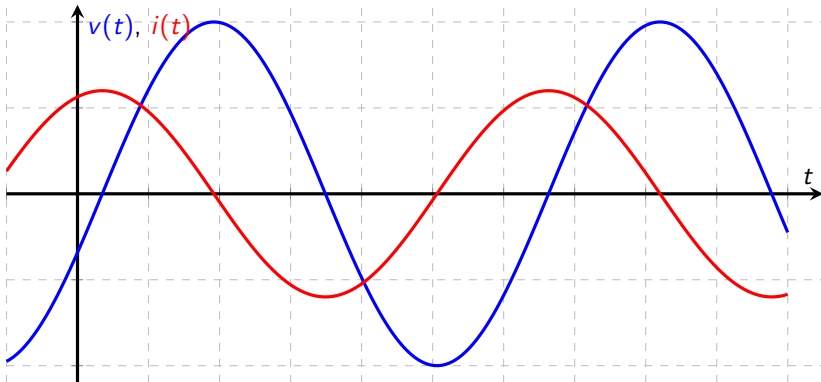
Kind of Ohm's law for phasors:

$$\mathbf{I} = jC\omega \mathbf{V} \qquad \text{or} \qquad \mathbf{V} = \frac{1}{jC\omega} \mathbf{I}$$

# Capacitor

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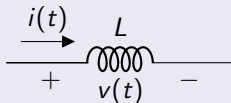
$v(t)$  and  $i(t)$  are out of phase  
 $i(t)$  is leading  $v(t)$  by  $90^\circ$   
 $v(t)$  is lagging  $i(t)$  by  $90^\circ$



# Inductor



## Inductor and phasors



Suppose we have a current  $i(t) = A \cos(\omega t + \phi)$

Then,  $v(t) = L \frac{di(t)}{dt} = -LA\omega \sin(\omega t + \phi) = -LA\omega \cos(\omega t + \phi - 90^\circ)$

When representing with phasors, we have:

$$\mathbf{I} = A \angle \phi \quad \mathbf{V} = -LA\omega e^{j(\phi - 90^\circ)} = -LA\omega e^{j\phi} e^{-j90^\circ}$$

## Phasor relationship

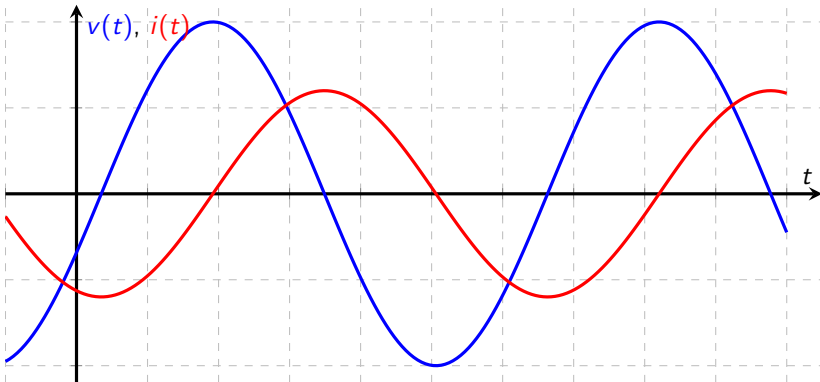
*Kind* of Ohm's law for phasors:

$$\mathbf{V} = jL\omega \mathbf{I} \quad \text{or} \quad \mathbf{I} = \frac{1}{jL\omega} \mathbf{V}$$

# Inductor

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$v(t)$  and  $i(t)$  are out of phase  
 $v(t)$  is leading  $i(t)$  by  $90^\circ$   
 $i(t)$  is lagging  $v(t)$  by  $90^\circ$

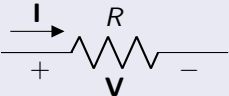
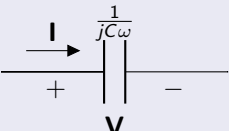



# Generalization



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## Generalization of Ohm's law

Element	Equation	Phasor equation	Circuit
Resistor	$v(t) = Ri(t)$	$\mathbf{V} = R\mathbf{I}$	
Capacitor	$i(t) = C \frac{dv(t)}{dt}$	$\mathbf{V} = \frac{1}{jC\omega} \mathbf{I}$	
Inductor	$v(t) = L \frac{di(t)}{dt}$	$\mathbf{V} = jL\omega \mathbf{I}$	



# Impedance/Admittance

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## Impedance

In the previous table, the phasor equation was written in the form:

$$\mathbf{V} = \mathbf{Z}\mathbf{I}$$

$\mathbf{Z}$  is called the **impedance** (unit: ohms -  $\Omega$ )

## Impedance

Sometimes, it can be useful to use another phasor equation:

$$\mathbf{I} = \mathbf{Y}\mathbf{V}$$

$\mathbf{Y}$  is called the **admittance** (unit: siemens - S)

## Relationship

Obviously, we have:

$$\mathbf{Z} = \frac{1}{\mathbf{Y}}$$

# Impedance/Admittance

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## Real and imaginary parts

In general, impedances or admittances are complex numbers.

$$\mathbf{Z} = R + jX \text{ and } \mathbf{Y} = G + jB$$

- $R$  is called the **resistance** ( $\Omega$ )
- $X$  is called the **reactance** ( $\Omega$ )
- $G$  is called the **conductance** (S)
- $B$  is called the **susceptance** (S)

# Circuits rules



## KVL

Sum of voltages around a closed path is equal to zero.

- $v_1(t) + v_2(t) + \dots + v_N(t) = 0$
- $A_1 \cos(\omega t + \phi_1) + A_2 \cos(\omega t + \phi_2) + \dots + A_N \cos(\omega t + \phi_N) = 0$
- $\text{Re}(A_1 e^{j(\omega t + \phi_1)} + A_2 e^{j(\omega t + \phi_2)} + \dots + A_N e^{j(\omega t + \phi_N)}) = 0$
- $\text{Re}((A_1 e^{j\phi_1} + A_2 e^{j\phi_2} + \dots + A_N e^{j\phi_N}) e^{j\omega t}) = 0$
- $\text{Re}(\mathbf{V}_1 + \mathbf{V}_2 + \dots + \mathbf{V}_N) e^{j\omega t} = 0$

Since  $e^{j\omega t} \neq 0$ , we have:

$$\mathbf{V}_1 + \mathbf{V}_2 + \dots + \mathbf{V}_N = 0$$

It means that **KVL still applies to phasors.**

## KCL

Similarly, it can be shown that **KCL still applies to phasors.**

# Circuit rules



## Impedances in series

When several impedances  $\mathbf{Z}_i$  are connected in series, it can be replaced by an equivalent impedance  $\mathbf{Z}_{eq}$ :

$$\mathbf{Z}_{eq} = \sum_{i=1}^N \mathbf{Z}_i$$

## Impedances in parallel

When several impedances  $\mathbf{Z}_i$  (or admittances  $\mathbf{Y}_i$ ) are connected in parallel, it can be replaced by an equivalent impedance  $\mathbf{Z}_{eq}$  (admittance  $\mathbf{Y}_{eq}$ ):

$$\mathbf{Z}_{eq} = \frac{1}{\sum_{i=1}^N \frac{1}{\mathbf{Z}_i}} \quad \text{or} \quad \mathbf{Y}_{eq} = \sum_{i=1}^N \mathbf{Y}_i$$

# Circuit rules



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## Source transformation

Sources can be transformed in phasor representation

## Thevenin/Norton equivalences

Thevenin and Norton equivalences can be given in phasor representation

## Superposition

Superposition can be applied in the phasor domain

# Circuit analysis

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## Analysis

Since we are only considering the **sinusoidal steady-state response**, the analysis of the circuits will be equivalent to the study of a network of resistors (no differential equation to solve!!!).

The only difference with DC resistor circuits is that we are going to manipulate **phasors** (complex voltages and currents) and **complex impedances**.

# Circuit analysis

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## Different steps

- ① Transform independent sources to phasors\*
- ② Determine the impedance of every passive element
- ③ Apply the analysis tools learned this semester
- ④ Find the phasors of interest and transform them back to a time-domain expression

\* Be careful if sources have different frequencies!!! Everything detailed earlier only applies for signals with the same frequency.

If different frequencies, you can still apply **superposition**. (Check exercise in recitation)