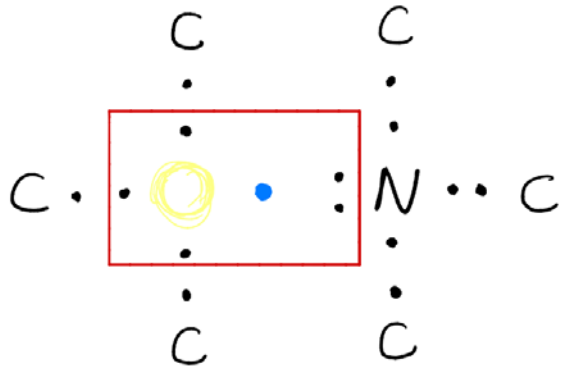
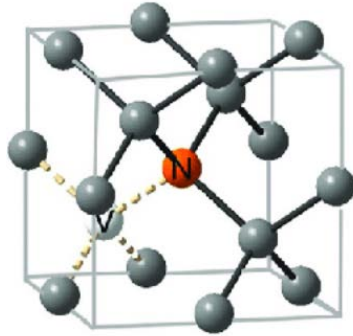


NV center electronic state

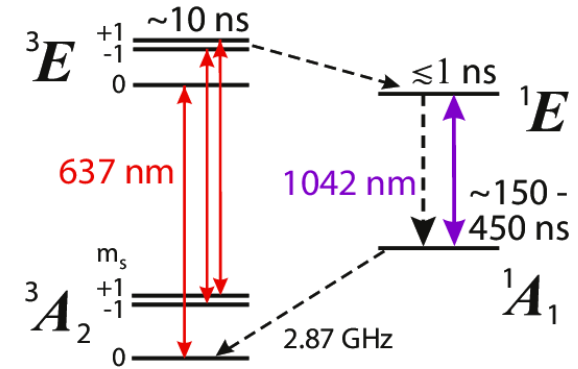
NV center
(sensor)



Structure of NV^-

○ vacancy

• additional electron



Energy level diagram of NV^-

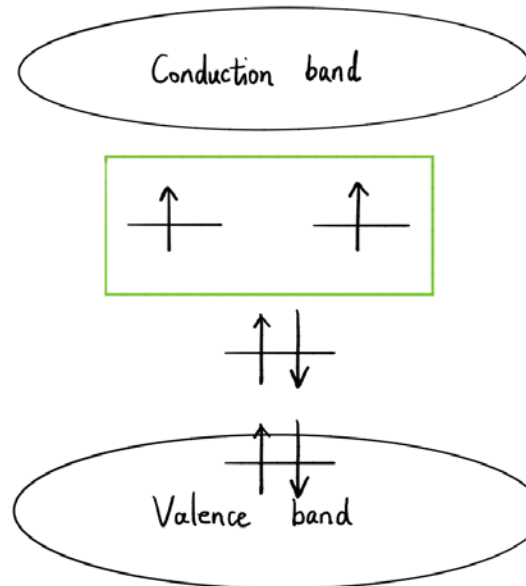
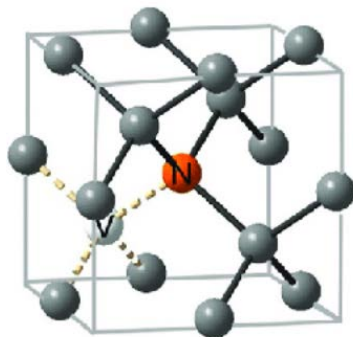
negative NV^- : $S = 1$

neutral NV^0 : $S = \frac{1}{2}$

positive NV^+ : $S = 0$

NV center electronic state

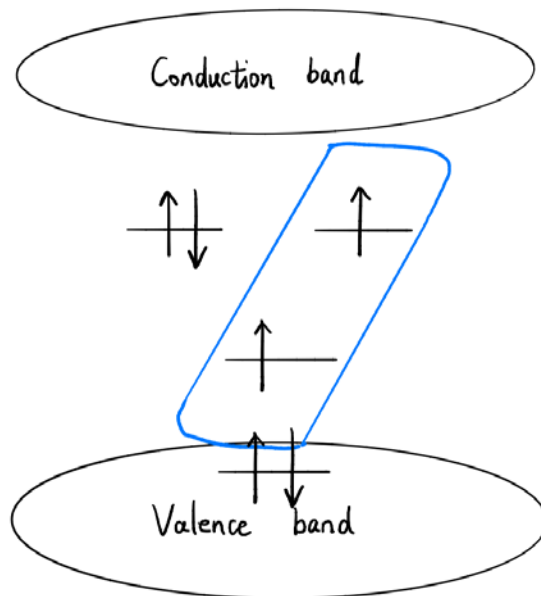
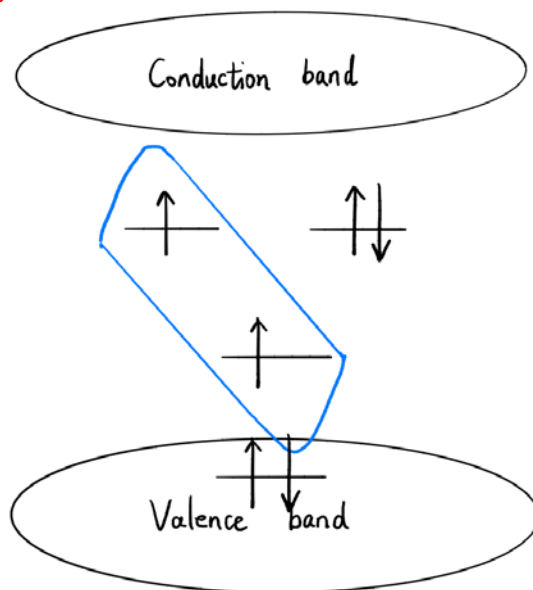
NV center
(sensor)



Ground state

Triplet $\begin{cases} \uparrow\uparrow = |1,1\rangle \\ \frac{\uparrow\downarrow + \downarrow\uparrow}{\sqrt{2}} = |1,0\rangle \\ \downarrow\downarrow = |1,-1\rangle \end{cases}$

Singlet $\frac{\uparrow\downarrow - \downarrow\uparrow}{\sqrt{2}} = |0,0\rangle$



Excited state

NV center electronic state

Zero-filed H: SOC & spin-spin interaction

$$H_g = \frac{1}{\hbar} (\vec{S} \cdot \vec{D} \cdot \vec{S}) \quad \text{Zero field splitting tensor} \quad \vec{D} = \begin{pmatrix} D_{xx} & D_{xy} & D_{xz} \\ & D_{yy} & \\ & & D_{zz} \end{pmatrix}$$

$$= \frac{1}{\hbar} (S_x D_{xx} S_x + S_y D_{yy} S_y + S_z D_{zz} S_z)$$

$$= \frac{1}{\hbar} (D_x S_x^2 + D_y S_y^2 + D_z S_z^2)$$

$$S_z^2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad S_x^2 = \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \end{pmatrix} \quad S_y^2 = \begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & \frac{1}{2} \end{pmatrix}$$

$$= \hbar \begin{pmatrix} \frac{D}{3} & E \\ E & -\frac{2}{3}D \end{pmatrix} = \hbar \left\{ D \left[\hat{S}_z^2 - \frac{S(S+1)}{3} \right] + E \left(\hat{S}_+^2 + \hat{S}_-^2 \right) \right\}$$

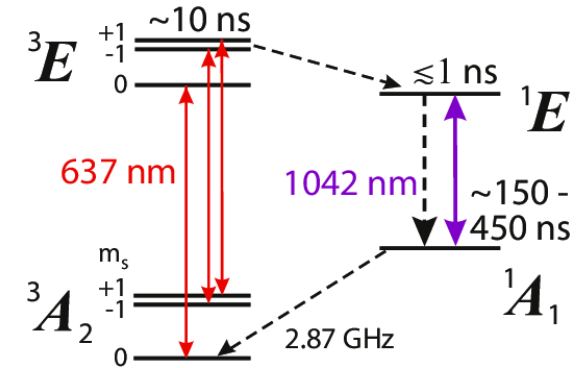
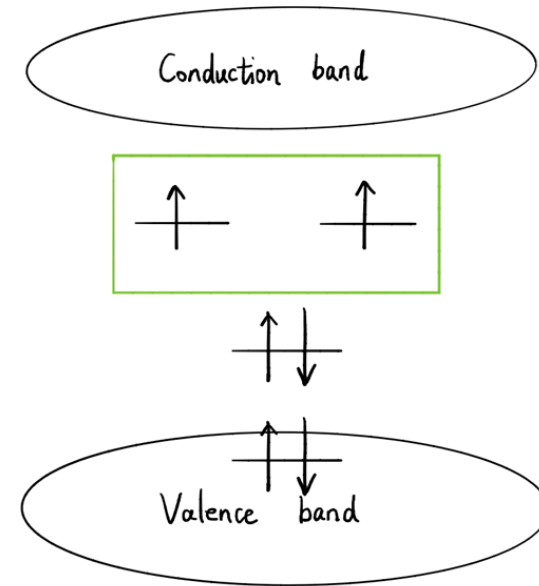
$$2(S_x^2 - S_y^2)$$

Ground state longitudinal and transverse zero-field splittings D_{GS} and E_{GS}

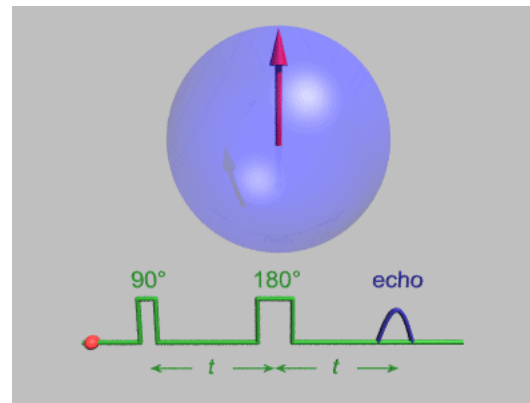
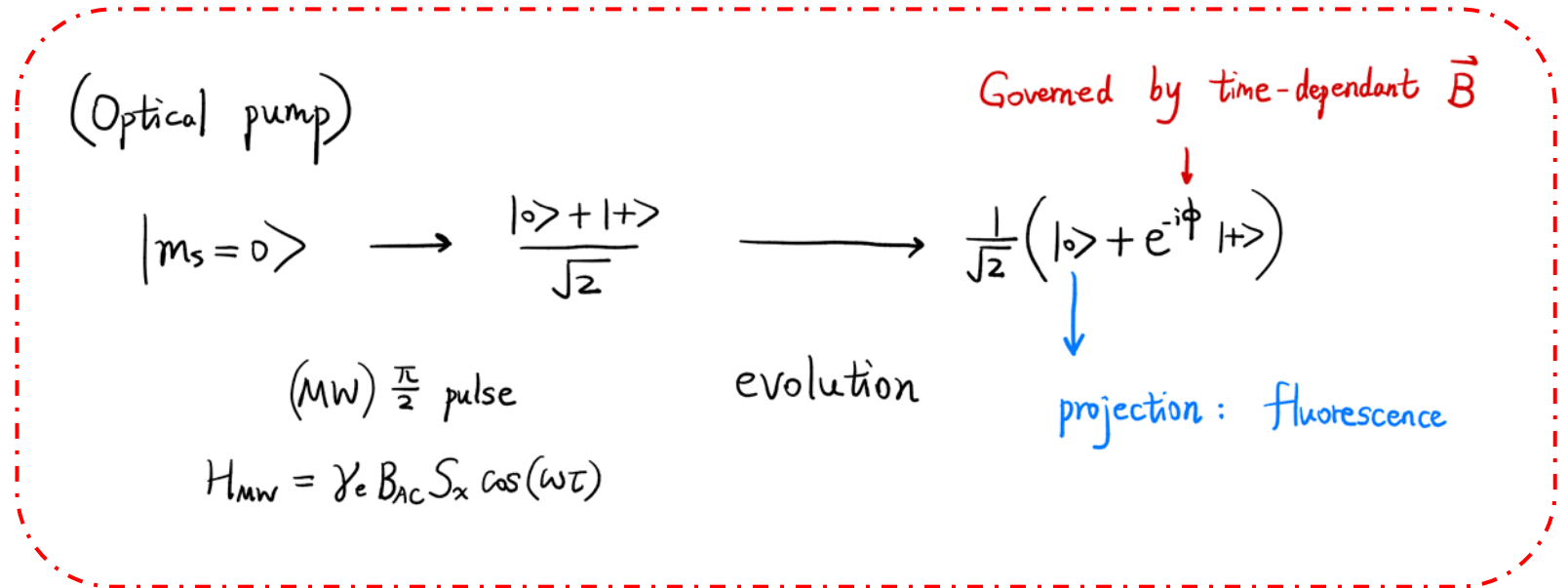
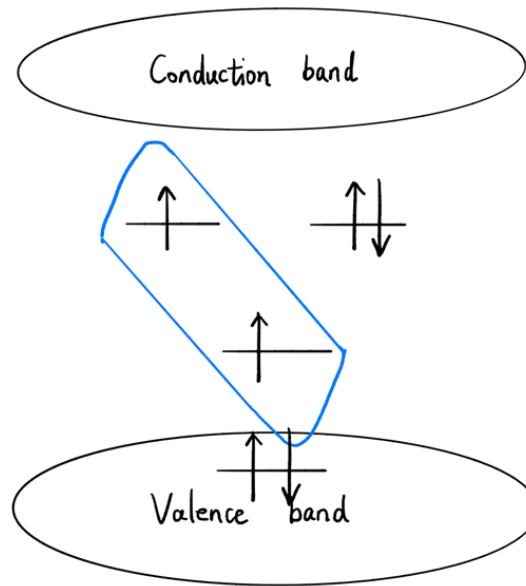
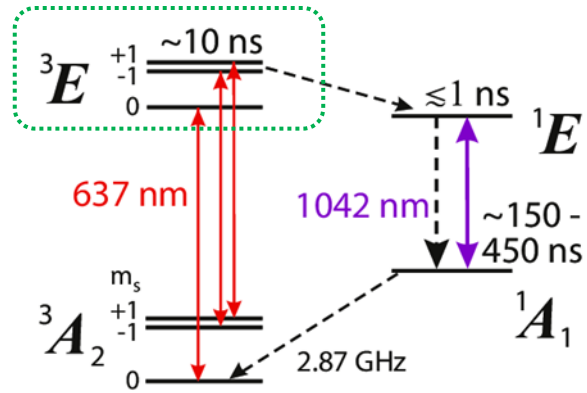
With energy & states:

$$-\frac{2}{3}D\hbar, \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}; \quad \left(\frac{D}{3}-E\right)\hbar, \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}; \quad \left(\frac{D}{3}+E\right)\hbar, \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix};$$

$|m_s = 0\rangle \qquad \qquad \qquad |-\rangle \qquad \qquad \qquad |+\rangle$

Energy level diagram of NV²⁻

Coherent manipulation of NV center



simplest form of dynamical decoupling

coherence is mapped by another $\pi/2$ pulse into a population difference

Echo amplitude \sim polarization difference

Traditional Nuclear Quadrupole Resonance (NQR)



non-spherical nuclear charge
distributions for $I \geq 1$ nuclei

^{14}N , ^{17}O , ^{35}Cl , ^{63}Cu ...

Interact with EFG:
electric field gradient
(crystal field?)

Transition frequency for axial symmetry

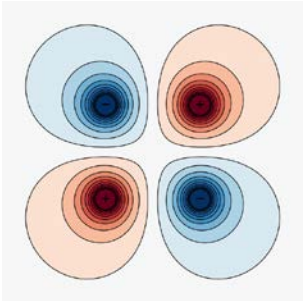
Specific frequency
for a given system

$$\nu_Q = \frac{3e^2 q Q}{4I(2I-1)\hbar}$$

Q : nuclear quadrupole moment

$q \sim$ largest principal component of the
EFG tensor at the nucleus

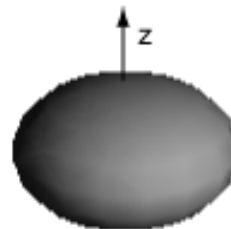
I : nuclear spin



Charge = 0
Dipole moment = 0
 $Q \neq 0$



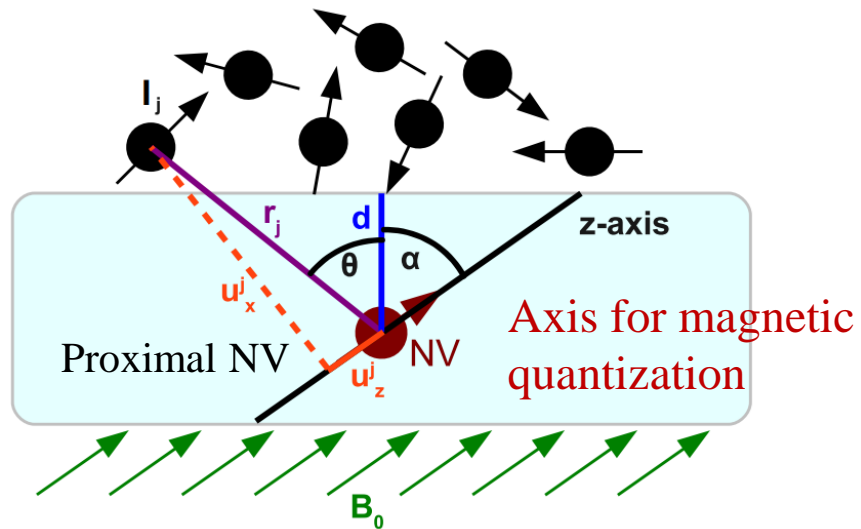
$Q > 0$
Prolate



$Q < 0$
Oblate

*B. H. Suits, Handbook of Applied
Solid State Spectroscopy
(Springer, New York, USA, 2006).*

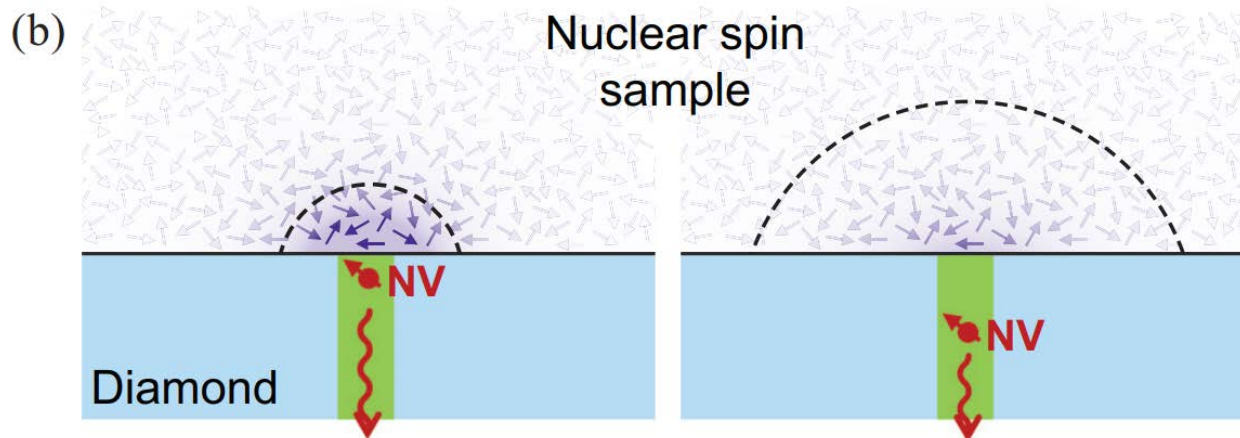
NV center as a nanoscale NQR spectrometer



Measure statistical fluctuations of the spin polarization $\sim \sqrt{N}$
(independent of the applied field)

N : number of nuclear spins
in the sensing volume

High sensitivity and small sensing volume:
ability to probe atomically thin layers

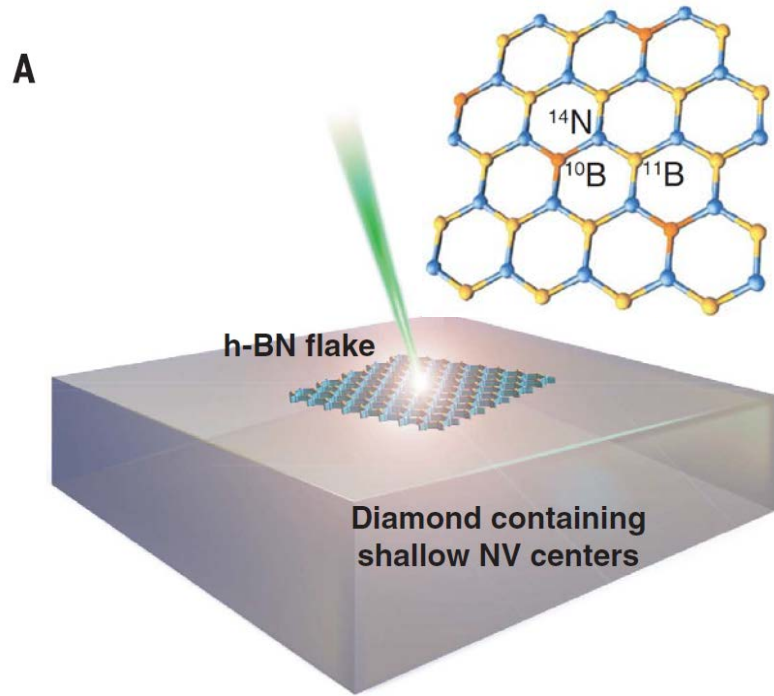


Depth: 6.8 ± 0.1 nm

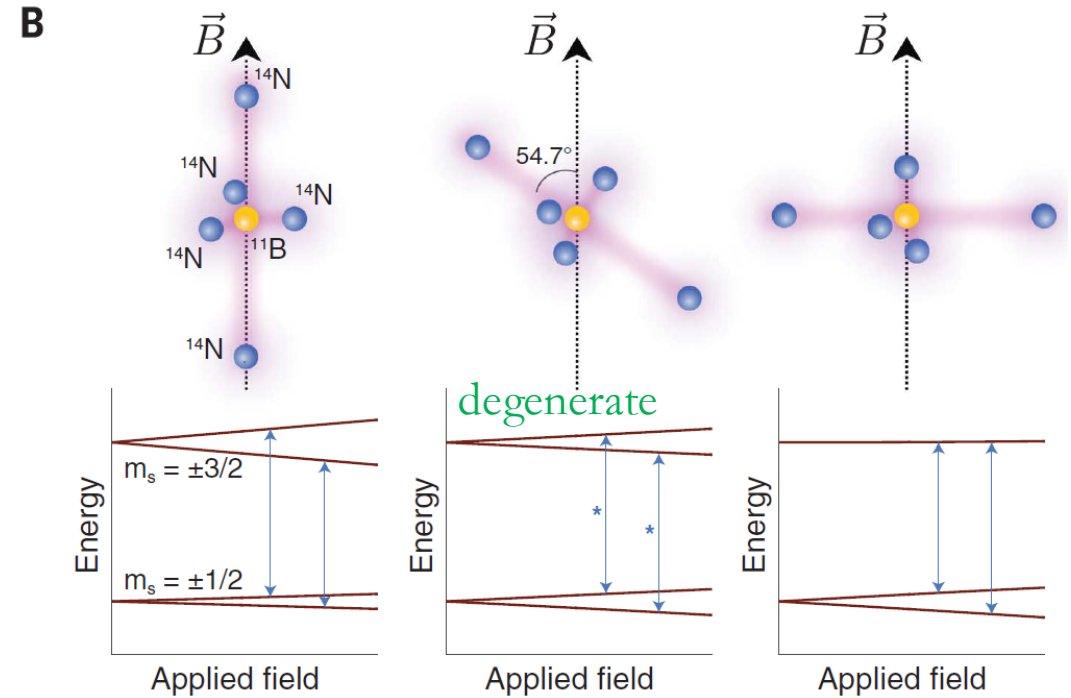
Due to dipolar coupling, a shallow NV center experiences a significantly stronger magnetic field from a smaller nuclear spin sample volume than a deep NV center experiences.

NMR technique for determining the depth of shallow nitrogen-vacancy centers in diamond
Linh M. Pham, Stephen J. DeVience et al

Special case: h-BN flake



$I = 3/2$ $I = 3$
 (80% ^{11}B , 20% ^{10}B)
 (close to 100% ^{14}N) $I = 1$
 honeycomb layered structure



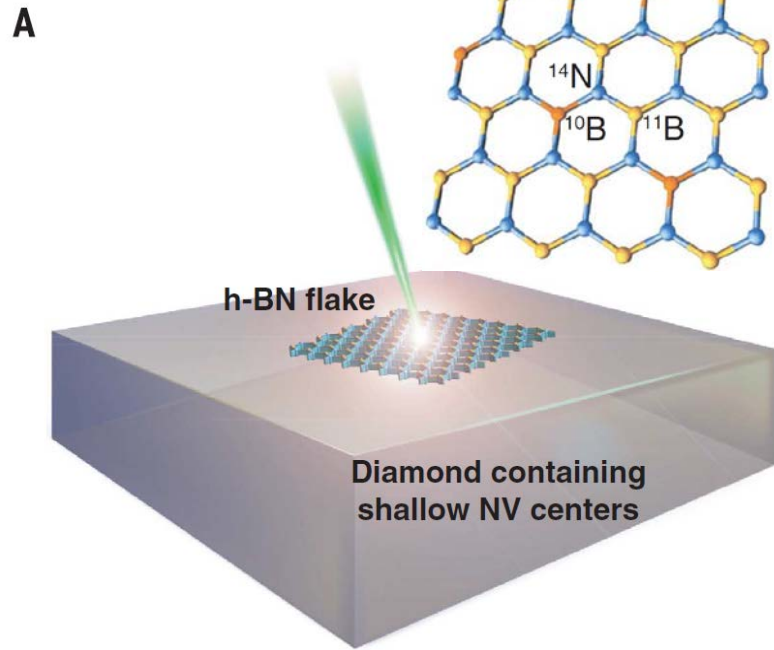
$$m_I = -I, -I+1, \dots, I-1, I$$

$$E_{\pm 3/2} = h\nu_Q \left(1 + \frac{\eta^2}{3} \right)^{1/2}$$

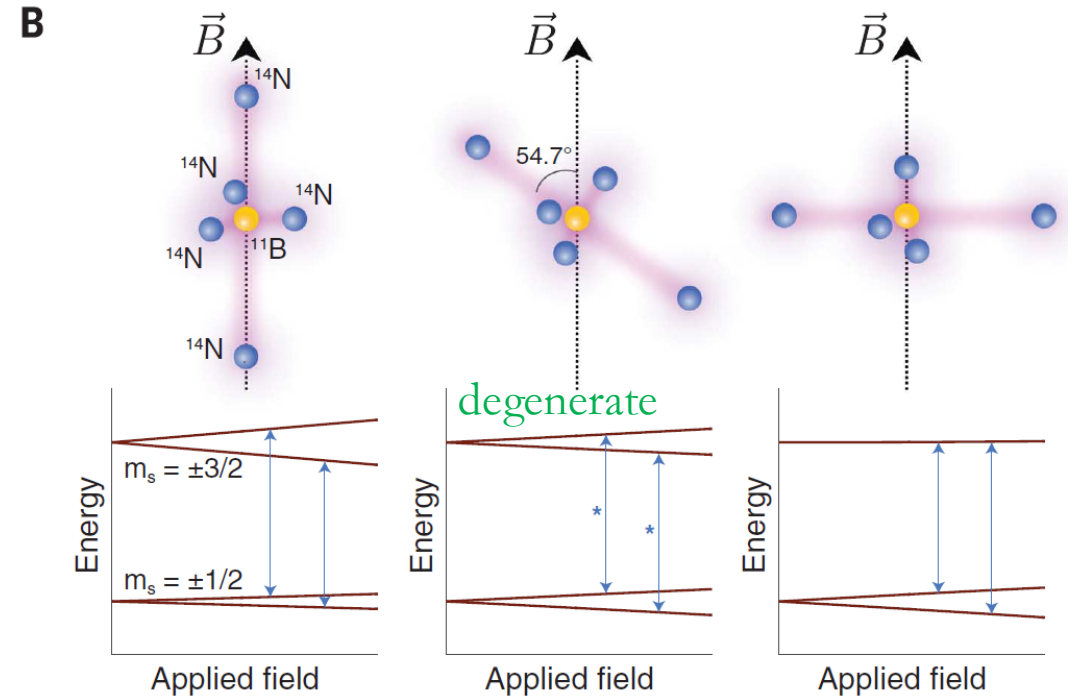
$$E_{\pm 1/2} = -h\nu_Q \left(1 + \frac{\eta^2}{3} \right)^{1/2}$$

B. H. Suits, Handbook of Applied Solid State Spectroscopy (Springer, New York, USA, 2006).

Special case: h-BN flake



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$$m_I = -I, -I+1, \dots, I-1, I$$

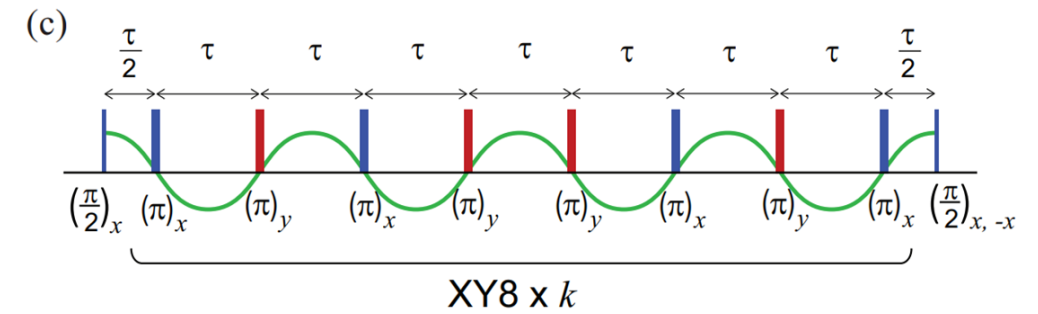
$$E_{\pm 3/2} = \frac{\hbar \nu_Q \rho}{2} \pm \frac{\hbar \nu_0}{2\rho} \left[(\rho - 1 + \eta)^2 c_x^2 + (\rho - 1 - \eta)^2 c_y^2 + (2 + \rho)^2 c_z^2 \right]^{1/2}$$

$$E_{\pm 1/2} = -\frac{\hbar \nu_Q \rho}{2} \pm \frac{\hbar \nu_0}{2\rho} \left[(\rho + 1 - \eta)^2 c_x^2 + (\rho + 1 + \eta)^2 c_y^2 + (2 - \rho)^2 c_z^2 \right]^{1/2}$$

*B. H. Suits, Handbook of Applied Solid State Spectroscopy
 (Springer, New York, USA, 2006).*

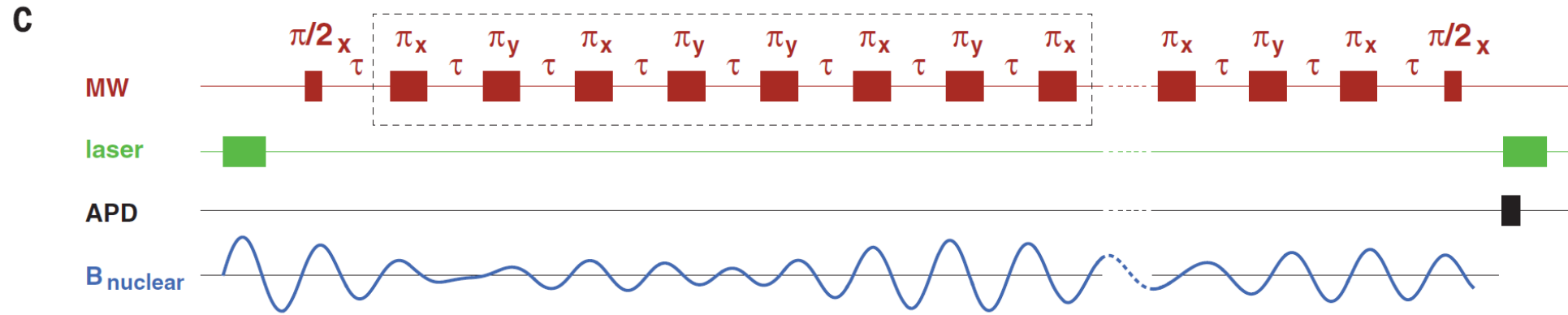
Pulse sequence:

measure individual Fourier components of NMR



Modified: separated by identical
free evolution intervals

modified XY8-371, 251...



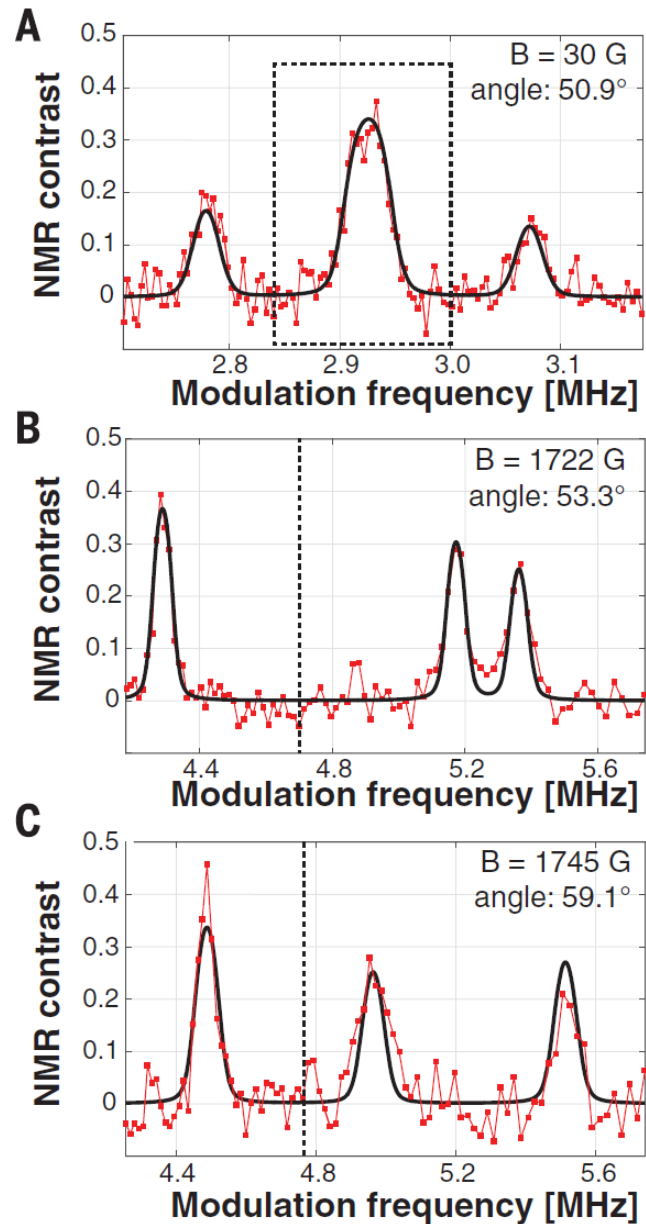
detected phase signal stored on the
populations of NV dressed states

effective B_{ac} detected in a
frequency-selective manner



Repetition of this sequence at different modulation
frequencies yields the NQR spectrum

NQR spectroscopy of h-BN



Satellite peaks: dipole-forbidden transitions become weakly allowed due to mixing by the **B** component perpendicular to the h-BN principal axis

