

Supplementary information: Boundedness and exponential stabilization
for time-space fractional parabolic-elliptic Keller-Segel model
in higher dimensions

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Supplementary Note : Supplementary Explanation of Proof of Boundedness of Solutions

We will consider an initial boundary value problem for a time-space fractional parabolic-elliptic Keller-Segel (KS) model

$$\begin{cases} {}^C_0D_t^\beta u = -(-\Delta)^{\alpha/2}(\rho(v)u), & (t, x) \in (0, T] \times \Omega, \\ (-\Delta)^{\alpha/2} v + v = u, & (t, x) \in (0, T] \times \Omega, \\ \partial_\nu u = \partial_\nu v = 0, & (t, x) \in (0, T] \times \partial\Omega, \\ u(0, x) = u_0(x), & x \in \Omega, \end{cases} \quad (1)$$

where $\beta \in (0, 1)$, $\alpha \in (1, 2)$, $\Omega \subset \mathbb{R}^n$, $n \geq 3$.

Lemma 4.3 Assume that $\rho(v) = v^{-k}$, $k > 0$ and that $n \geq 1$. Let $(u(t, x), v(t, x))$ be a unique non-negative classical solution to (1) on $[0, T_{\max})$, $T_{\max} \in (0, \infty]$. Then, there exists a positive constant $C = C(n, \Omega, k, u_0, \alpha)$ satisfying

$$\|u(t, \cdot)\|_{L^\infty(\Omega)} \leq C, \quad \text{for all } t \in (0, T_{\max}), \quad (2)$$

$$\|v(t, \cdot)\|_{W^{1,\infty}(\Omega)} \leq C, \quad \text{for all } t \in (0, T_{\max}). \quad (3)$$

Proof We have demonstrated that L^{p+1} integrability of u for some $p+1 > n/2$. In the following, the method of proving Lemma 4.3 in Ref. [1] is referred to in order to prove the Lemma. Let $p+1 > n$. Multiplying the first equation of (1) by u^p and integrating over Ω , by [[2],p579,(3.15)] and [[3],p1248,(3.1)] we obtain

$$\begin{aligned} & \frac{1}{p+1} {}^C_0D_t^\beta \int_\Omega u^{p+1} dx + \int_\Omega D^{\frac{\alpha}{2}} u^p D^{\frac{\alpha}{2}} (v^{-k} u) dx \\ & \leq \int_\Omega u^p {}^C_0D_t^\beta u dx + \int_\Omega u^p (-\Delta)^{\alpha/2} (\rho(v)u) dx = 0. \end{aligned} \quad (4)$$

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Using further Eilertsen equality and "Cordoba-Cordoba" type inequalities from [2] to deflate $\int_{\Omega} D^{\frac{\alpha}{2}} u^p D^{\frac{\alpha}{2}} (v^{-k} u) dx$, for $\rho(v)$ and u with the same monotonicity, we get

$$\begin{aligned} \int_{\Omega} D^{\frac{\alpha}{2}} u^p D^{\frac{\alpha}{2}} (v^{-k} u) dx &\leq \int_{\Omega} p u^{p-1} D^{\alpha/2} u \left[v^{-k} D^{\alpha/2} u + u D^{\alpha/2} v^{-k} \right] dx \\ &\quad - \int_{\Omega} p u^{p-1} D^{\alpha/2} u A_{\alpha/2} \int_{\Omega} \frac{(v^{-k}(x) - v^{-k}(y))(u(x) - u(y))}{|x - y|^{n+\alpha}} dx \\ &\leq \int_{\Omega} p u^{p-1} D^{\alpha/2} u \left[v^{-k} D^{\alpha/2} u + u D^{\alpha/2} v^{-k} \right] dx \\ &\leq \int_{\Omega} p v^{-k} u^{p-1} |D^{\alpha/2} u|^2 dx - p k \int_{\Omega} \frac{1}{v^{k+1}} u^p D^{\alpha/2} u D^{\alpha/2} v dx. \end{aligned} \quad (5)$$

Combining (4) and (5), we obtain the following inequality

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \int_{\Omega} \frac{1}{v^k} u^{p-1} |D^{\alpha/2} u|^2 dx \leq k \int_{\Omega} \frac{1}{v^{k+1}} u^p D^{\alpha/2} u D^{\alpha/2} v dx. \quad (6)$$

By applying Young's inequality to (6), we derive

$$\begin{aligned} &\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \int_{\Omega} \frac{1}{v^k} u^{p-1} |D^{\alpha/2} u|^2 dx \\ &\leq \frac{1}{2} \int_{\Omega} \frac{1}{v^k} u^{p-1} |D^{\alpha/2} u|^2 dx + \frac{k^2}{2} \int_{\Omega} \frac{1}{v^{k+2}} u^{p+1} |D^{\alpha/2} v|^2 dx. \end{aligned}$$

Due to $v_* \leq v \leq v^*$ and [[2], Theorem 3.2], there exists $C > 0$ such that

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \frac{1}{C} \int_{\Omega} \left| D^{\alpha/2} u \frac{v^{p+1}}{2} \right|^2 dx \leq C \int_{\Omega} u^{p+1} |D^{\alpha/2} v|^2 dx. \quad (7)$$

Using Young's inequality again for the right-hand side of equation (7) yields

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \frac{1}{C} \int_{\Omega} \left| D^{\alpha/2} u \frac{v^{p+1}}{2} \right|^2 dx \leq C \int_{\Omega} u^{p+2} dx + C \int_{\Omega} |D^{\alpha/2} v|^{2(p+2)} dx.$$

Furthermore, applying Nash-Gagliardo-Nirenberg-type inequality (see [3], Lemma 5.3), standard elliptic regularity theory and $\|v\|_{L^\infty(\Omega)} \leq C$, we obtain

$$\left\| D^{\alpha/2} v \right\|_{L^{2(p+2)}(\Omega)}^{2(p+2)} \leq C \|v(t, \cdot)\|_{L^\infty(\Omega)}^{p+2} \|v(t, \cdot)\|_{W^{2,p+2}(\Omega)}^{p+2} \leq C \|u(t, \cdot)\|_{L^{p+2}(\Omega)}^{p+2}.$$

Therefore, we have

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \frac{1}{C} \int_{\Omega} \left| D^{\alpha/2} u \frac{v^{p+1}}{2} \right|^2 dx \leq C \int_{\Omega} u^{p+2} dx. \quad (8)$$

Using Nash-Gagliardo-Nirenberg-type inequality again, we notice that

$$\int_{\Omega} u^{p+2} dx = \left\| u^{\frac{p+1}{2}} \right\|_{L^{\frac{2(p+2)}{p+1}}(\Omega)}^2 \leq C \left\| u^{\frac{p+1}{2}} \right\|_{L^1(\Omega)}^2 \left\| D^{\alpha/2} u^{\frac{p+1}{2}} \right\|_{L^2(\Omega)}^{\frac{2(p+2)}{p+1} \lambda^*}, \quad (9)$$

with $\lambda^* = \frac{n+1}{n+3} \frac{p+3}{p+2} \in (0,1)$. By L^{p+1} integrability of u for some $p+1 > \frac{n}{2}$, we

choose $p+1 > n$ satisfying $\|u(t, \cdot)\|_{L^{\frac{p+1}{2}}(\Omega)} \leq C$, for all $t \in (0, T_{max})$. Thus, the

inequality (9) can be expressed as

$$C \int_{\Omega} u^{p+2} dx \leq C \left\| D^{\alpha/2} u^{\frac{p+1}{2}} \right\|_{L^2(\Omega)}^{\frac{2(p+2)}{p+1} \lambda^*}. \quad (10)$$

Due to $\frac{2(p+2)}{p+1} \lambda^* < 2$, by applying Young's inequality to (10), then from (8) we

deduce that

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \frac{1}{2C} \int_{\Omega} \left| D^{\alpha/2} u^{\frac{p+1}{2}} \right|^2 dx \leq C.$$

Analogous to (9), Nash-Gagliardo-Nirenberg-type inequality gives us that there exist a number $C > 0$, such that

$$\frac{1}{p(p+1)} {}^C D_t^\beta \int_{\Omega} u^{p+1} dx + \int_{\Omega} u^{p+1} dx \leq C.$$

Let $w(t) = \int_{\Omega} u^{p+1} dx$, and with reference to [[4], Corollary 2.1], we obtain for some

$p+1 > n$ that

$$w(t) = \|u(t, \cdot)\|_{L^{p+1}(\Omega)}^{p+1} \leq w(0) E_\beta(-t^\beta) - C(E_\beta(-t^\beta) - 1) \leq C, \quad (11)$$

for any $t \in (0, T_{max})$. The final deflation of the above inequality makes use of the fact

that there exists $t_1 > 0$, $\varepsilon > 0$, such that Mittag-Leffler function $w(0) E_\beta(-t^\beta) < \varepsilon$ for

any $t > t_1$.

Finally, (3) is obtained by using standard elliptic regularity theory (see, e.g. [5]) for (11). And (2) can be deduced by standard iteration argument (see, e.g. [6]). This completes the proof.

Reference

- [1] Ahn J, Yoon C. Global well-posedness and stability of constant equilibria in parabolic–elliptic chemotaxis systems without gradient sensing. *Nonlinearity*, 2019, 32(4): 1327.
- [2] Alsaedi A, Ahmad B, Kirane M. A survey of useful inequalities in fractional calculus. *Fractional Calculus and Applied Analysis*, 2017, 20(3): 574-594.
- [3] de Pablo A, Quirós F, Rodríguez A, et al. A general fractional porous medium equation. *Communications on Pure and Applied Mathematics*, 2012, 65(9): 1242-1284.
- [4] Li L, Liu J G, Wang L. Cauchy problems for Keller–Segel type time–space fractional diffusion equation. *Journal of Differential Equations*, 2018, 265(3): 1044-1096.
- [5] Brézis H. *Analyse fonctionnelle, Théorie et applications*. Collection Mathématiques Appliquées pour la Maîtrise. Paris: Masson, 1983.
- [6] Alikakos N D. L^p bounds of solutions of reaction-diffusion equations. *Communications in Partial Differential Equations*, 1979, 4(8): 827-868.