

# Examples of Basic Analysis Techniques for Neutral Meson and Photon Data from the PHENIX Detector

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## Abstract

This writeup is a part of an Analysis Preservation and educational package with specific focus on the Electromagnetic Calorimeter (EMCal) of the PHENIX detector. It aims to present a few basic techniques of identifying  $\pi^0$ -s and photons using that device. The data component of this package is provided in the form of ROOT files containing small samples of reconstructed  $p$ +Au data taken with the PHENIX detector in 2015 at  $\sqrt{s_{NN}}=200$  GeV. Both minimum bias and EMCal-triggered data are included. *These samples were created for educational and training purposes only and are not at all usable for deriving any real physics results.*

## 1 Introduction

Data samples described in this writeup contain information on *clusters* – small, contiguous regions of energy deposit in the EMCal – which are usually considered as *energy deposited by one single particle*. While this is not always true, this is our first working hypothesis, but we should always be aware that a cluster might contain energy from more than one particle. Also, not all clusters correspond to photons; many are from hadrons, especially at low and very high  $p_T$  (transverse momentum)<sup>1</sup>. In fact, it’s going to be your job in this exercise to find selection criteria (“cuts”) which eliminate hadrons (the *contamination*) but still preserves most of the photons (*signal*) in the sample.

As for  $\pi^0$ , recall that most of them decay via  $\pi^0 \rightarrow \gamma\gamma$ , so they can be reconstructed from photon candidate *pairs* using the invariant mass

$$m_{\gamma\gamma} = \sqrt{2E_1E_2(1 - \cos(\theta))}$$

where  $E_1, E_2$  are the energies of the two photon candidates and  $\theta$  is the *opening angle* between them. Note that the  $p_T$  of the  $\pi^0$  is simply the vector sum of the  $p_T$ -s of the two photons paired.

Please note that – as per PHENIX standards – time is always given in nanoseconds, distance in centimeters and energy in GeV.

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<sup>1</sup> This may sound surprising with regards to an electromagnetic calorimeter, but it is true: many high  $p_T$  single clusters are actually from two overlapping photons from a  $\pi^0$  decay, i.e. they are hadrons.

## 2 Contents of the Package

The package described here contains

- This writeup (file: "pi0ntup\_v2.pdf")
- Four ROOT-formatted data files as described in Sec.3
- Two ROOT macros designed to manipulate these data, as described in Sec.4

The ROOT macros and this document (pi0ntup\_v2.pdf) can be located on GitHub:

[https://github.com/PhenixCollaboration/opendata/tree/master/p0\\_photon\\_analysis](https://github.com/PhenixCollaboration/opendata/tree/master/p0_photon_analysis)

**Caveat:** the ROOT macros were developed in the standard PHENIX analysis environment which is based on ROOT version 5.34/36. You will need to make adjustments to the code to make it run under a more modern version e.g. version 6. Alternatively, you may obtain a Docker image of ROOT version 5.34 from the PHENIX Collaboration Organization on Docker Hub:

```
docker pull phenixcollaboration/tools:s17_root5
```

The data (e.g. the ROOT-formatted Ntuples) are available on the CERN OpenData portal and also can be obtained by request from the PHENIX Collaboration. Please visit <https://www.phenix.bnl.gov> for more detail.

## 3 The Data: *gnt* and *ggntuple* Ntuples

### 3.1 General comments about the data

There are a total of four files which are a part of the package in which this document is included:

- *MBntup\_ggntuple.root*
- *ERTntup\_ggntuple.root*
- *MBntup\_gnt.root*
- *ERTntup\_gnt.root*

The "*MB*" files correspond to the minimum bias data while the "*ERT*" files are based on the EMCal-triggered data. The *ggntuple* files contain information on the EMCal *cluster pairs* while *gnt* files describe *single clusters*.

These files were produced using a small fraction of the  $\sqrt{s_{NN}}=200$  GeV  $p$ +Au data collected by PHENIX in 2015. They are published for educational purposes only. You will not be able to derive any publishable physics results from them. Their purpose is to help you learn some basic techniques of particle identification (PID), practice making different cuts and estimate their effect, observe how the signal to background ratio can be improved at the expense of some loss in the signal and how to find an optimum between these two trends.

You may also learn that PID itself is also not an unambiguous procedure: you can apply much looser cuts on the "photonness" of the clusters when reconstructing  $\pi^0$ , because you identify the  $\pi^0$  by the *correlation* of the two "photons", namely, whether their invariant mass is in the expected range (around 0.135 GeV) or not. Only photon pairs from the decay of the same  $\pi^0$  are

truly correlated, random pairs of photons, photon-hadron or hadron-hadron pairs only rarely give invariant masses in the “ $\pi^0$  window”. Therefore, you can allow yourself to have more contamination in the “photon” sample – and gain significant statistics in reconstructed  $\pi^0$  (why?). On the other hand for reconstructing (inclusive) photons you don’t have such help: if you decide that a cluster is a photon, that’s it. In order to keep the contamination from misclassified hadrons you might want to make your photon PID cuts stricter, in order to increase *purity* (in other terminology decrease *contamination*) at the expense of *efficiency*.

### 3.2 The Gamma-Gamma Ntuples (*ggntuple*)

The first Ntuple type considered here is called *ggntuple* (gamma-gamma Ntuple). It contains information on EMCal cluster *pairs* and the two clusters in the pair (see Table 1). In each event all clusters in a sector are paired with all other clusters, and the pair variables (*pt*, *costheta*, *phi*, *mass*, *asym*) calculated, even if they are obviously not photons from a  $\pi^0$  decay (for instance because their invariant mass is, say, 0.4 GeV). These “obviously wrong” random pairs are called the *combinatorial background* and they help you to estimate the background in the  $\pi^0$  peak when you’ll try to extract the number of  $\pi^0$  in the sample. The momenta attributed to a cluster (particle candidate) is calculated from the vector connecting the collision point (*vtxZ*) with the impact point on the EMCal, and under the assumption that the particle is a photon (i.e.  $p = E$ , the measured energy). From these, the pair momentum *pt* is calculated as the vector sum of the transverse components; *costheta* is the cosine of the opening angle between the two momentum vectors, *phi* is the azimuth of the pair momentum (i.e. of the parent  $\pi^0$ ), *mass* is the invariant mass calculated with the formula above, and *asym* is the energy asymmetry  $\alpha$  of the two clusters

$$\alpha = \frac{|E_1 - E_2|}{E_1 + E_2}$$

The rest of the variables describe the two clusters (1,2), their meaning is spelled out below.

| Variable name | Description  |
|---------------|--|
| cent          | Event centrality   |
| vtxZ          | $z$ -vertex of the event   |
| pt            | Transverse momentum of cluster <i>pair</i> ( $\pi^0$ candidate)          |
| costheta      | $\cos$ of the opening angle between the two clusters                     |
| phi           | Azimuthal angle of the direction of the pair momentum (assumed $\pi^0$ ) |
| mass          | Invariant mass calculated from the two clusters (energy and position)    |
| asym          | Energy asymmetry ( $ E_1 - E_2 /(E_1 + E_2)$ ) of the two clusters       |
| sec1          | EMCal sector where the first cluster is                                  |
| Ecore1        | “Core” energy of the first cluster                                       |
| tof1          | Time-of-flight of the first cluster                                      |
| twrhit1       | Number of towers in the first cluster                                    |
| prob1         | Probability that the first cluster is a photon (based on $\chi^2$ )      |
| chisq1        | $\chi^2$ from expected photon shape for the first cluster                |
| stoch1        | Combined variable to describe “photonness” of first cluster              |
| sec2          | ... Same quantities as above for the second cluster of the pair          |

Table 1: Variables in the  $\gamma$ -pair Ntuple (*ggntuple* type).

### 3.3 The Single Cluster Ntuples (*gnt*)

The second Ntuple type is called *gnt*. It contains information on individual clusters (again, many of which are *not* photons!). Please see Table 2 for description of individual variables. Here *pt* is the transverse momentum of the (single) particle, calculated under the assumption that it is a photon (straight path from the collision vertex and  $p = E$ ), *costheta* is the cosine of the polar angle, *phi* is the azimuth, and *sec* is the sector number (0-5 are lead scintillator sectors, PbSc, 6-7 are lead glass, PbGl). For both detectors the smallest, individually read out units are called *towers*.

The variable *ecore* is an estimate of the particle energy (using various corrections to account for detector effects), *ecent* is the energy in the central (highest energy) tower of the cluster, *tof* is the time-of-flight measured in the central tower and corrected with the flight-path  $s/c$  such that for photons *tof* is a distribution around zero. Based on test beam data we have a model (parameterization) of electromagnetic showers depending on energy, impact point and angle of the photon; the *chisq* is calculated as the deviation of the actual shower (energy deposit pattern in the towers) from the “model” shower of the same energy, impact point and angle. The *prob* is calculated from this  $\chi^2$ , usually showers with *prob* > 0.05 are very likely to be photons; *disp* is the 2D dispersion, *twrhit* is the number of towers hit (i.e. included in the cluster). Finally, *stoch* is a combination of various probability metrics to characterize “photonness”, higher values mean higher probability that the cluster is indeed a photon. The calculated impact point of the particle is given as  $x, y, z$  in the PHENIX absolute coordinate system ( $x > 0$  is the West Arm, negative  $z$  is South).

The *MB\*.root* files are produced from minimum bias data (although with a lower limit  $p_T > 2.0 \text{ GeV}/c$  in *gnt* and pair  $p_T$  in *ggntuple*), whereas in *ERT\*.root* files the threshold for single cluster  $p_T$  in *gnt* is 8 GeV, and the threshold for pair  $p_T$  in *ggntuple* is also 8 GeV. Note that since here we restrict only the pair  $p_T$ , the energy of the individual clusters can be (and often is) significantly lower.

**Important:** data in both *MBntup\_ggntuple.root* and *ERTntup\_ggntuple.root* are subject to the condition that the asymmetry is less than 0.8.

| Variable name | Description  |
|---------------|--|
| cent          | Event centrality   |
| vtxZ          | $z$ -vertex of the event   |
| pt            | Transverse momentum of the cluster ( $\gamma$ -candidate)                        |
| costheta      | Polar angle of the cluster ( $\gamma$ -candidate)                                |
| phi           | Azimuthal angle of the cluster ( $\gamma$ -candidate)                            |
| sec           | EMCal sector of the cluster ( $\gamma$ -candidate)                               |
| ecore         | “Core” energy of the cluster ( $\gamma$ -candidate)                              |
| ecent         | Energy in the central tower of the cluster ( $\gamma$ -candidate)                |
| tof           | Time-of-flight in the central tower of the cluster ( $\gamma$ -candidate)        |
| prob          | Probability that the cluster is a photon (based on $\chi^2$ )                    |
| disp          | Dispersion of the cluster ( $\gamma$ -candidate)                                 |
| chisq         | $\chi^2$ from expected photon shape of the cluster ( $\gamma$ -candidate)        |
| twrhit        | Number of towers in the cluster ( $\gamma$ -candidate)                           |
| stoch         | Combined variable to describe “photonness” of the cluster ( $\gamma$ -candidate) |
| x             | $x$ -position of impact point on the EMCal surface                               |
| y             | $y$ -position of impact point on the EMCal surface                               |
| z             | $z$ -position of impact point on the EMCal surface                               |

Table 2: Variables in the single-cluster ( $\gamma$ ) Ntuple (the *gnt* type).

## 4 Analyzing the Data

### 4.1 Access to the Code

The two ROOT macros named *gamma\_gamma.C* and *single\_cluster.C* can be downloaded from this GitHub location:

[https://github.com/PhenixCollaboration/opendata/tree/master/p0\\_photon\\_analysis](https://github.com/PhenixCollaboration/opendata/tree/master/p0_photon_analysis)

Please see Sec.2 for an importance notice regarding the ROOT version compatibility. The macros were designed to be self-documenting and to produce a visually appealing result. In the following, we are using *paraphrased* fragments of the macros as the basis for commentary (i.e. they may differ a little for illustration purposes). Feel free to run modified snippets of the code to get a feel of how it works with the data.

### 4.2 Cluster Pairs (macro: *gamma\_gamma.C*)

#### 4.2.1 Effect of the $\chi^2$ cut on the invariant mass spectra

Below is a code sample demonstrating the effect of the  $\chi^2$  cut.

```
int W=900;
int H=400;

TFile* mb = new TFile("MBntup_ggntuple.root");
TFile* ert = new TFile("ERTntup_ggntuple.root");

TCanvas *canvas1 = new TCanvas("canvas1", "Effect_of_the_chi2_cut", W, H*3);
canvas1->Divide(1,3);

canvas1_1->Divide(2,1);
canvas1_2->Divide(2,1);

mb->cd();
canvas1_1_1->cd();
ggntuple->Draw("mass", "mass<1.0");

ert->cd();
canvas1_1_2->cd();
ggntuple->Draw("mass", "mass<1.0");

mb->cd();
canvas1_2_1->cd();
ggntuple->Draw("mass>>htemp1", "mass<1.0");
canvas1_2_2->cd();
ggntuple->Draw("mass>>htemp2", "mass<1.0&&chisq1<3.0&&chisq2<3.0");

canvas1_3->SetLogy();
canvas1_3->cd();

htemp2->SetLineColor(2);
htemp1->Draw();
htemp2->Draw("same");
```

For characteristic results see Fig. 1. The left panel shows the  $m_{\gamma\gamma}$  distribution in the MB data (no cut on  $p_T$  or any other observable). While the true  $\pi^0$  peak dominates (0.135 GeV), there

is a pronounced peak at 0.03 GeV (primarily due to “minimum ionization” hadrons), also, the combinatorial background under the true  $\pi^0$  peak is substantial. The middle panel is the same  $m_{\gamma\gamma}$  distribution from the ERT data. The fake peak at 0.03 GeV/c is gone, the combinatorial background is much smaller, and now you clearly see the peak at the  $\eta$  meson mass as well. The right panel is once again MB data; the blue histogram is the (uncut)  $m_{\gamma\gamma}$  (same as on the left panel), the red is the same  $m_{\gamma\gamma}$  if we require that both clusters of the pair satisfy the  $\chi^2 < 3$  condition (i.e. are consistent with the photon cluster shape). Note the log scale. As you can see, the cut on “photonness” leaves the  $\pi^0$  peak area nearly unchanged, but reduces the combinatorial background drastically.

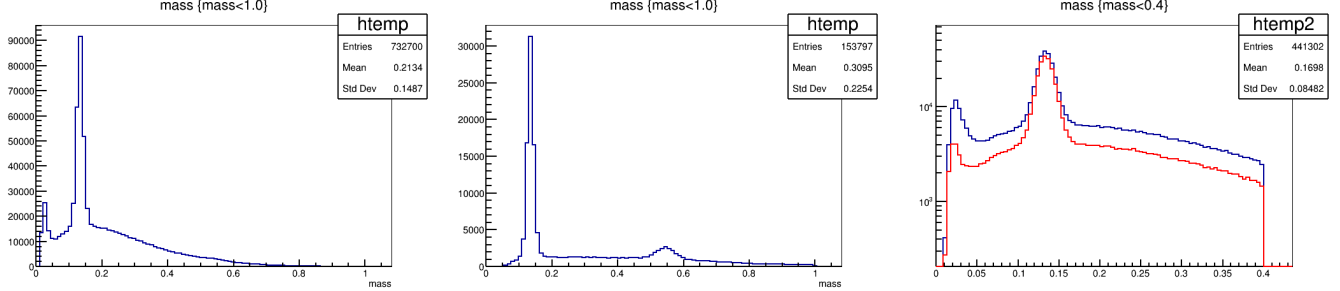


Figure 1: Left panel: Invariant mass from the MB data in the 0-1 GeV region. You can see a strong  $\pi^0$  peak and maybe a hint of  $\eta$ . Middle panel:  $\pi^0$  peak for pairs from the ERT data (i.e. the pair  $p_T$  is greater than 8 GeV/c. The combinatorial background is much smaller and the  $\eta$  peak is now clearly seen. Right panel: the effect of a strong  $\chi^2$  cut on the photon pairs in the MB data. The blue histogram is the same as in the left panel, the red is with the  $\chi^2$  cut on both clusters. While you barely lose any signal (due to the small inefficiency of the cut), the signal to background ratio improves dramatically.

## 4.2.2 Asymmetry

The (energy) asymmetry of the two  $\pi^0$  decay photons is a flat distribution by kinematics (at least in  $4\pi$  solid angle). This offers another opportunity to improve the signal/background ratio when extracting the  $\pi^0$ -s. In the left panel of Fig. 2 the asymmetry distribution of all  $m_{\gamma\gamma}$  pairs is shown (MB data). The enhancement at high asymmetries comes from the abundant low  $p_T$  photons and hadrons depositing small energy - i.e. non- $\pi^0$ , combinatorial background pairs. By applying a modest, say  $\alpha < 0.6$  asymmetry cut the S/B ratio can be significantly improved at a small cost of lost true  $\pi^0$ -s, as demonstrated in the right panel of Fig. 2. The ROOT commands are here:

```
ggntuple->Draw("asym","mass<1.0");
ggntuple->Draw("mass>>htemp1","mass<1.0");
ggntuple->Draw("mass>>htemp2","mass<1.0&&asym<0.6");
htemp1->Draw();
htemp2->Draw("same");
```

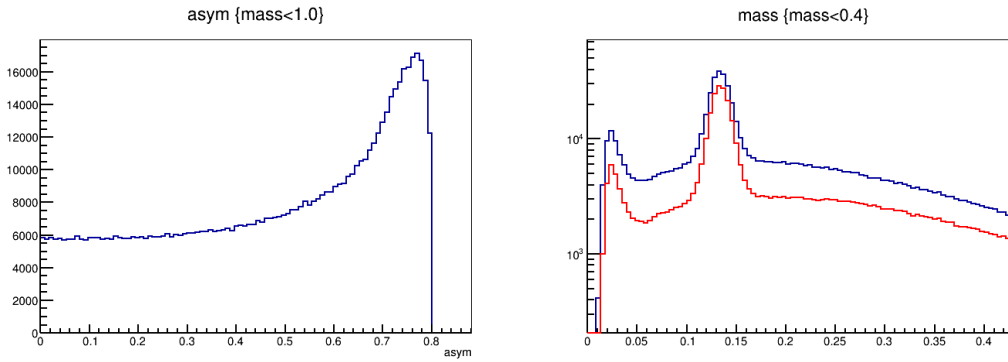


Figure 2: MB data, plots from the pair ntuple. Left panel: energy asymmetry distribution for all pairs. Right panel: invariant mass distribution without asymmetry cut (i.e. with the default  $\alpha < 0.8$ ) and with a tighter  $\alpha < 0.6$  cut (red histogram). The signal to combinatorial background ratio is significantly improved.

### 4.2.3 $p_T$

Let's make a crude estimate of the *raw*  $\pi^0$  and  $\eta$  spectrum, i.e. the pair  $p_T$  distribution of those pairs which give an invariant mass in the  $\pi^0$  or  $\eta$  “mass window”. The estimate will be crude because we don't attempt to account for the combinatorial background. However, we can try to clean up the spectra by making a cut on the “photonness” of both clusters. This is illustrated in Fig. 3 for  $\pi^0$  in both MB and ERT data, and  $\eta$  (ERT only). The blue histograms are the  $p_T$  distribution of pairs in the mass window, the red histoes the same but with an additional cut on the *chisq* variables. The principal commands:

```
ggntuple->Draw("pt","abs(mass-0.135)<0.015");
ggntuple->Draw("pt","abs(mass-0.135)<0.015&&chisq1<3.0&&chisq2<3.0");
ggntuple->Draw("pt","abs(mass-0.548)<0.1");
ggntuple->Draw("pt","abs(mass-0.548)<0.1&&chisq1<3.0&&chisq2<3.0");
```

Note that the mass window for  $\pi^0$  is much narrower than for  $\eta$ . In the MB  $\pi^0$  spectra at low  $p_T$  the *chisq* cut eliminates some 20-30% of  $\pi^0$  counts (red vs blue histogram), in line with what the left plot in Fig. 1 suggests (the estimated magnitude of the combinatorial background). In ERT there is very little change, again in line with the small combinatorial on the middle panel in Fig. 1. Finally, for  $\eta$  there is a significant drop in the yields for the entire (high  $p_T$ ) spectrum, in line with the large combinatorial (as compared to the small peak) in Fig 1. Also, the spurious “flattening” at very high  $p_T$  in the blue histogram is completely gone.

### 4.2.4 TOF

As stated earlier, the time-of-flight (*tof*) for clusters is readjusted such that it will be zero for true photons. In the left panel of Fig. 4 for cluster pairs which have an  $m_{\gamma\gamma}$  in the  $\pi^0$  window we plot *tof* values vs each other. True  $\pi^0$ -s are in the big peak at (0,0). If we now look at the horizontal band to the right of the peak, the first cluster is still likely to be a photon (small *tof*), but the second is more likely to be a hadron. Is this reflected in other observables characterizing the clusters, too? The answer is yes: on the right panel in Fig. 4 we plot the “photonness” variables *chisq1* in blue, *chisq2* in red. The second cluster, presumed hadronic, has on average higher *chisq2* value.

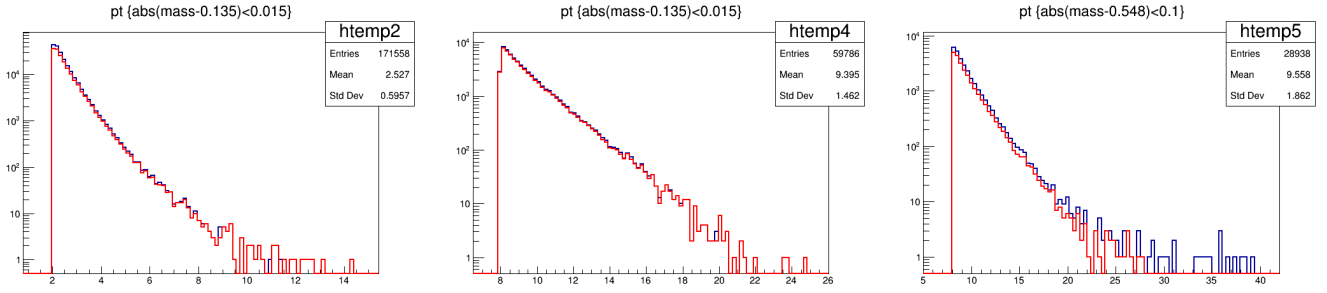


Figure 3: Transverse momentum ( $p_T$ ) distributions of photon candidate pairs in the  $\pi^0$  window (top panels) and  $\eta$  window (bottom panel). Blue histograms:  $p_T$  for all pairs (no identification of the cluster as photon). Red histograms:  $p_T$  for pairs where both clusters satisfy the *chisq* “photonness” criterion. Left: MB data, center and right: ERT data.

There is one more important lesson to be learned here. Real-life analyses are complicated. There is no “silver bullet”, single cut, that, by miracle, selects *all* photons and nothing else. Cuts are characterized by their *efficiency* (how much of the desired signal they preserve) and *contamination* (how much of the undesired signal passes). Tighter cuts may reduce contamination, but always at the cost of efficiency. The art of analysis is to optimize the signal/background, without sacrificing too much signal. This is usually achieved by combining different cuts with AND or even OR operations. In the last decades probabilistic approaches of combining cuts are also gaining ground.

```
ggntuple->Draw("tof1:tof2","abs(mass-0.135)<0.015&&abs(tof1)<20.0&&abs(tof2)<20.0","colz");
TH1D *hsmall = new TH1D("hsmall","hsmall",60,0.0,30.0);
TH1D *hlarge = new TH1D("hlarge","hlarge",60,0.0,30.0);
ggntuple->Draw("chisq1>>hsmall","abs(mass-0.135)<0.015&&abs(tof1)<3.0&&abs(tof2)>5.0");
ggntuple->Draw("chisq2>>hlarge","abs(mass-0.135)<0.015&&abs(tof1)<3.0&&abs(tof2)>5.0");
hlarge->SetLineColor(2);
hlarge->SetLineWidth(2);
hsmall->SetLineWidth(2);
hsmall->Draw();
hlarge->Draw("same");
```

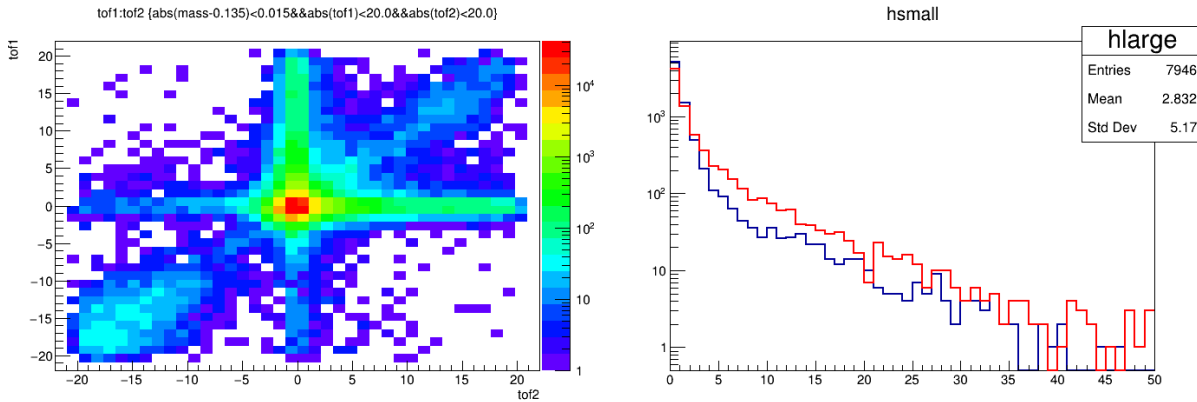


Figure 4: Left panel: time-of-flight values of two clusters which produce an invariant mass in the  $\pi^0$  window. TOF for true photons is per construction around zero so the true (non-combinatorial) pairs are in the sharp peak at (0,0). Right panel: *chisq* (“photonness”) of the clusters in the horizontal band above the main peak at (0,0). Blue histogram is *chisq1* of the presumed true photon, the red histogram is the *chisq2* of the presumed hadron. See text for more details.



## 4.3 Single Clusters (macro: *single\_cluster.C*)

### 4.3.1 Cluster Position

Let's turn to single clusters (photon candidates, *gnt* Ntuple), and look first at the position of the clusters (assumed impact point) in a particular sector. In Fig. 5 the  $y, z$  positions of the clusters are shown from MB and ERT data. The depletion at the edges is due to the algorithm (clusters too close to the edge are not reconstructed because energy may leak out and the resulting *ecore* will be incorrect). The white areas represent “dead” channels that never fire, and the blue areas have only few hits, although our expectation is that the clusters should be approximately uniformly distributed over the sector. Such dead or low-sensitivity areas should be taken account in a real analysis when the acceptance and efficiency of the detector are calculated.

```
TFile* mb = new TFile("MBntup_gnt.root");
TFile* ert = new TFile("ERTntup_gnt.root");

int W=1200;
int H=400;

TCanvas *canvas1 = new TCanvas("canvas1", "Cluster_position", W, H);
canvas1->Divide(2,1);

mb->cd();
canvas1_1->cd();
gnt->Draw("y:z", "sec==1", "colz");

ert->cd();
canvas1_2->cd();
gnt->Draw("y:z", "sec==1", "colz");
```

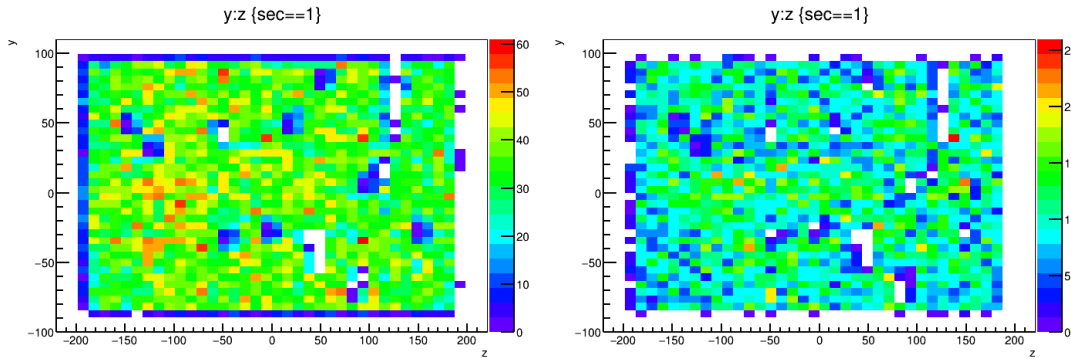


Figure 5: Position of cluster centers (calculated impact point) in sector 1 (PbSc West). Left panel: MB data. Right panel: ERT data.

### 4.3.2 Selecting photons with the $\chi^2$ cut

Unfortunately not all clusters in *gnt* are photons. The *chisq* variable is an estimate of the “photon-ness” based on a parametrized shower profile, but its detailed description is beyond the scope of this write-up. However, it is interesting to look at the compactness of the shower. Electromagnetic

showers are narrow, most of the energy is deposited in a single tower, while hadronic showers are fairly wide. Besides the total energy ( $ecore$ ) the Ntuple also stores the energy in the “central” (most energetic) tower ( $ecent$ ). The ratio  $ecent/ecore$  is expected to be high for photons, lower for hadrons. In Fig. 6 this ratio is plotted versus the total energy for MB data, first without any other cuts, then cutting on the  $chisq$  variable, first deliberately selecting “photon-like” showers, then “un-photon-like” showers. Without cuts the  $ecent/ecore$  has two characteristic peaks, which then get almost perfectly resolved by the “photon - not photon” cut. Of course that also means that  $ecent/ecore$  in its own right can be a powerful cut to select photons.

```
mb->cd();
TCanvas *canvas2 = new TCanvas("canvas2", "ecent/ecore:ecore", 3*W/2, H);
canvas2->Divide(3,1);

canvas2_1->cd();
gnt->Draw("ecent/ecore:ecore", "ecore < 4.0", "colz");

canvas2_2->cd();
gnt->Draw("ecent/ecore:ecore", "ecore < 4.0 && chisq < 3.0", "colz");

canvas2_3->cd();
gnt->Draw("ecent/ecore:ecore", "ecore < 4.0 && chisq > 5.0", "colz");
```

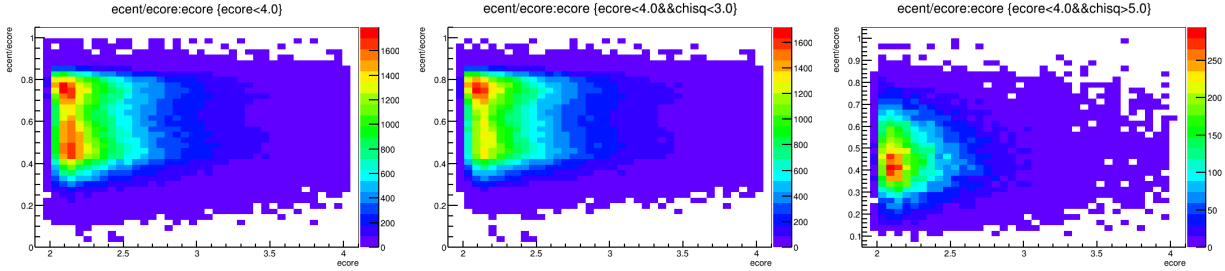


Figure 6: Shower compactness ( $ecent/ecore$ ) vs shower energy in MB data. Left panel: no cuts. There is a characteristic double-peak structure. Using the  $chisq$  (“photonness”) variable on the middle panel we select showers that are likely photons (small  $chisq$ ), on the right panel clusters that are most likely not photons (large  $chisq$ ). The compactness variable is quite powerful distinguishing between them (i.e. for photon identification).

### 4.3.3 The effect of the $\chi^2$ cut on the $p_T$ spectra

As mentioned before, the  $chisq$  variable is a powerful tool to enhance photons and reject hadrons in the sample. In Fig. 7 we plot the  $p_T$  distribution of the clusters in the ERT data sample, first with no cut (blue histogram), then with a cut on  $chisq$  (red histogram). Without the cut, the spectrum “flattens out” at high  $p_T$ , which is non-physical. The effect of the cut becomes visible early on, and will be quite dramatic above 15 GeV/c. Up to about 25 GeV/c the red histogram shows the expected power-law shape, above that it also flattens out. You may try to cut also on other  $gnt$  variables to get rid of the nonphysical flat part of the red histogram.

```
TH1D *h1 = new TH1D("h1", "h1", 100, 0.0, 50.0);
TH1D *h2 = new TH1D("h2", "h2", 100, 0.0, 50.0);
```

```

h1->SetLineWidth(2.0);
h2->SetLineWidth(2.0);
h2->SetLineColor(2);
gnt->Draw("pt>>h1","");
gnt->Draw("pt>>h2","chisq<3.0");
h1->Draw();
h2->Draw("same");

```

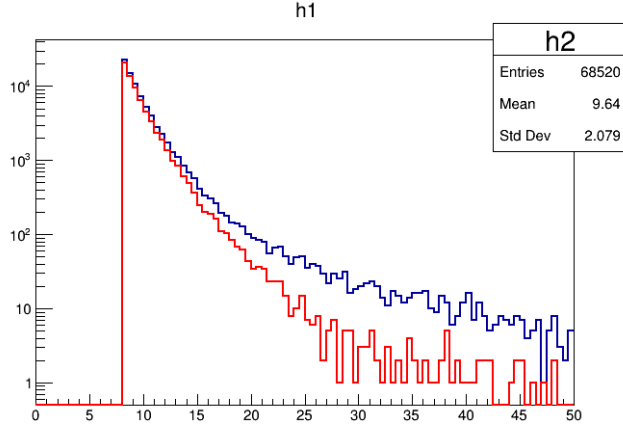


Figure 7: Blue histogram:  $p_T$  spectrum from the ERT data,  $gnt$  ntuple. Red spectrum: same, but with a cut on “photonness” ( $chisq$ ). Much - albeit not all - of the nonphysical flattening at high  $p_T$  is gone.

## 4.4 Final words

Please continue to play with the data! The variables in the ntuples are there for a reason! Try plotting them against each other, making them part of cuts when you plot a spectrum. Make invariant mass plots in bins of pair  $p_T$ , fit the combinatorial background with some simple function (first, second order polynomial) and try to get a “clean” raw  $\pi^0$  yield as a function of  $p_T$ . Try to imagine in advance, what effect a cut would have, then actually make it and compare the result with your expectations. If you are open-minded and play creatively with them, *the data will start speaking to you and telling you how to analyze them.*

**Have fun!**