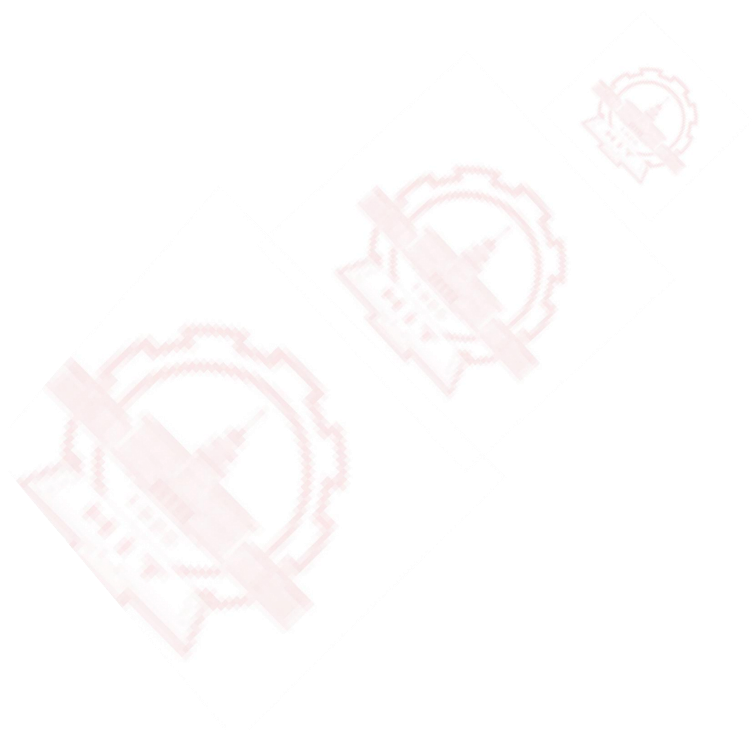


# 机器人学导论

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**2023春季**



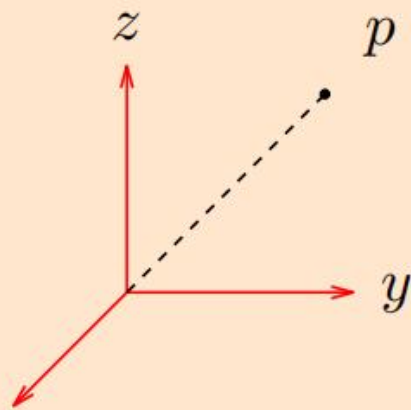
# [3] 刚体运动

HIT Rigid Body Motion

## Chapter 3 Rigid Body Motion

- **Notations**
- **Rigid Body Transformations**
- **Coordinate Frame Transformations in  $\mathbb{R}^3$**
- **Rigid Motion in  $\mathbb{R}^3$**
- **Exponential Coordinates and Screw Theory**
- **Other Parametrizations of  $SO(3)$**

# 3.1 Notations



$$p = \begin{bmatrix} p_x \\ p_y \\ p_z \end{bmatrix} \text{ or } p = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$$

For  $p \in \mathbb{R}^n, n = 2, 3$  (2 for planar, 3 for spatial)

$$\text{Point: } p = \begin{bmatrix} p_1 \\ p_2 \\ \vdots \\ p_n \end{bmatrix}, \|p\| = \sqrt{p_1^2 + \cdots + p_n^2}$$

$$\text{Vector: } v = p - q = \begin{bmatrix} p_1 - q_1 \\ p_2 - q_2 \\ \vdots \\ p_n - q_n \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ \vdots \\ v_n \end{bmatrix}, \|v\| = \sqrt{v_1^2 + \cdots + v_n^2}$$

$$\text{Matrix: } A \in \mathbb{R}^{n \times m}, A = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1m} \\ \vdots & & & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nm} \end{bmatrix}$$

## 3.2 Rigid Body Transformations (Description of point-mass motion)



$$p(0) = \begin{bmatrix} x(0) \\ y(0) \\ z(0) \end{bmatrix} : \text{initial position}$$

$$p(t) = \begin{bmatrix} x(t) \\ y(t) \\ z(t) \end{bmatrix}, t \in (-\varepsilon, \varepsilon)$$

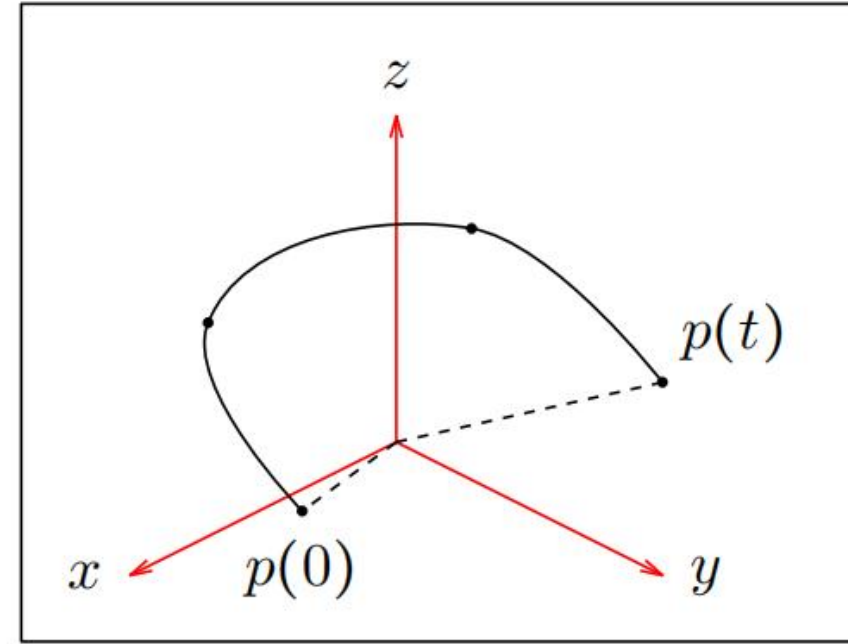


Figure 3.1

### Definition: Trajectory

A **trajectory** is a curve  $p : (-\varepsilon, \varepsilon) \mapsto \mathbb{R}^3, p(t) = \begin{bmatrix} x(t) \\ y(t) \\ z(t) \end{bmatrix}$

## 3.2 Rigid Body Transformations (Description of rigid body motion)

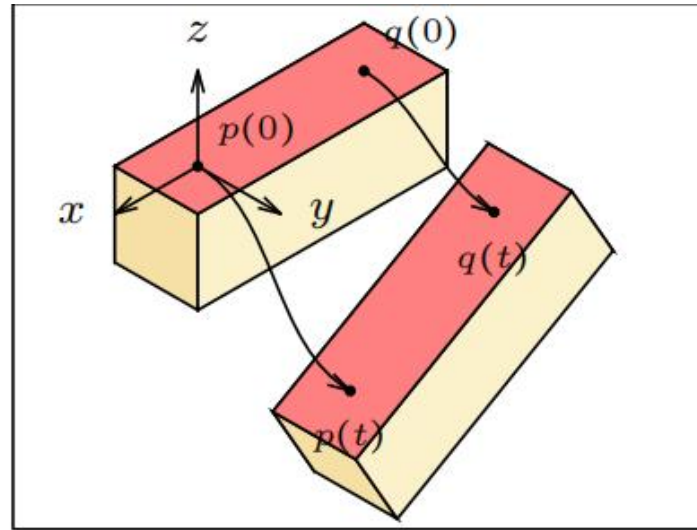


Figure 3.2

$$\|p(t) - q(t)\| = \|p(0) - q(0)\| = \text{constant}$$

**Definition: Rigid body transformation**

$$g : \mathbb{R}^3 \mapsto \mathbb{R}^3$$

s.t.

- ① Length preserving:  $\|g(p) - g(q)\| = \|p - q\|$
- ② Orientation preserving:  $g_*(v \times \omega) = g_*(v) \times g_*(\omega)$

## 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$



**A rigid body transformation**



**A coordinate frame transformation**

- ① Choose a reference frame A  
(spatial frame)
- ② Attach a frame B to the body  
(body frame)

$x_{ab} \in \mathbb{R}^3$ : coordinates of  $x_b$  in frame A  
 $R_{ab} = [x_{ab} \ y_{ab} \ z_{ab}] \in \mathbb{R}^{3 \times 3}$ : Rotation (or orientation) matrix of B w.r.t. A

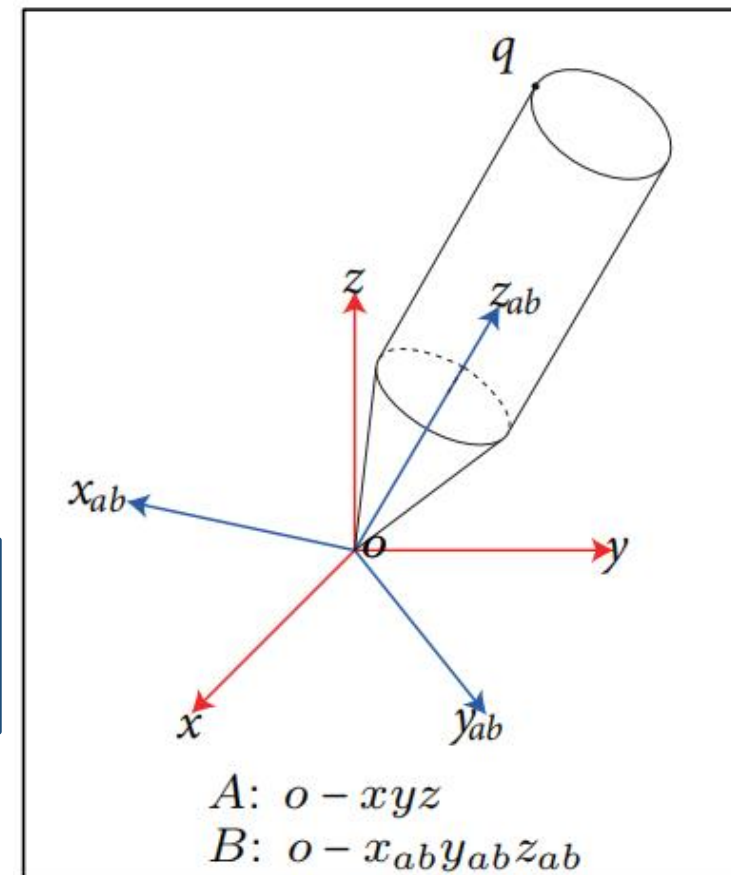


Figure 3.3

## 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$

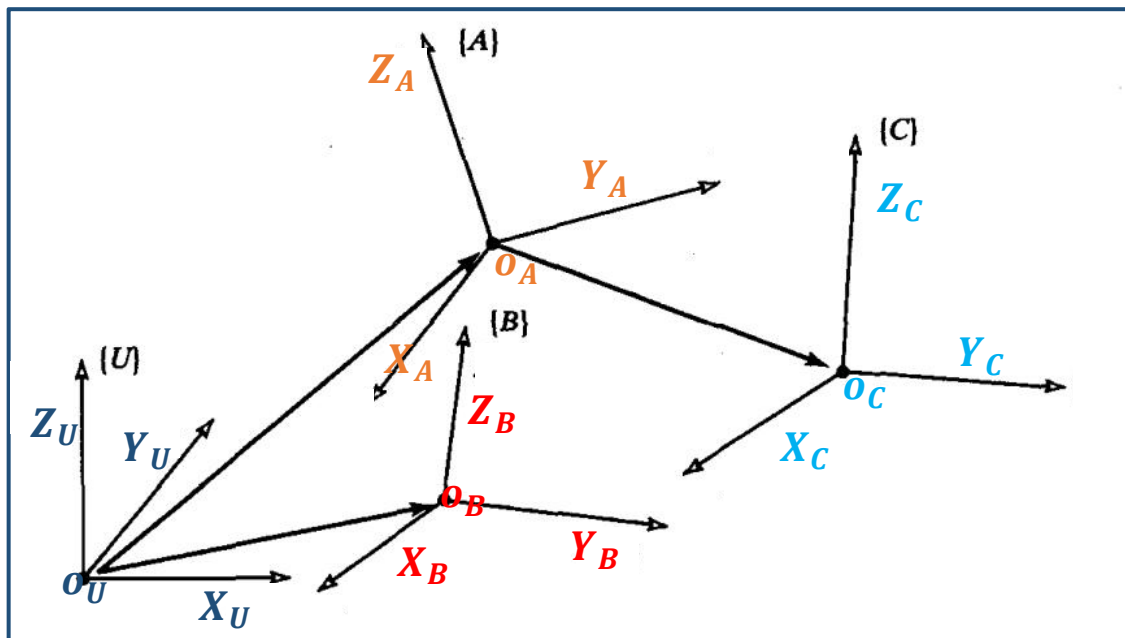


Figure 3.4

A coordinate frame  $\{A\}$  consists of two parts:

- ❑ Original Point  $o = [o_x \ o_y \ o_z]^T$
- ❑ Coordinat axes  $[A_x^T \ A_y^T \ A_z^T]$

Coordinate Frame A

Transform Matrix  
 $T \in \mathbb{R}^{4 \times 4}$

Coordinate Frame B

A coordinate frame transform matrix can be divided into two parts:

- ❑ Rotation Matrix  $R \in \mathbb{R}^{3 \times 3}$
- ❑ Ttranslation Vector  $p \in \mathbb{R}^{3 \times 1}$



### 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$ (Rotation Matrix)

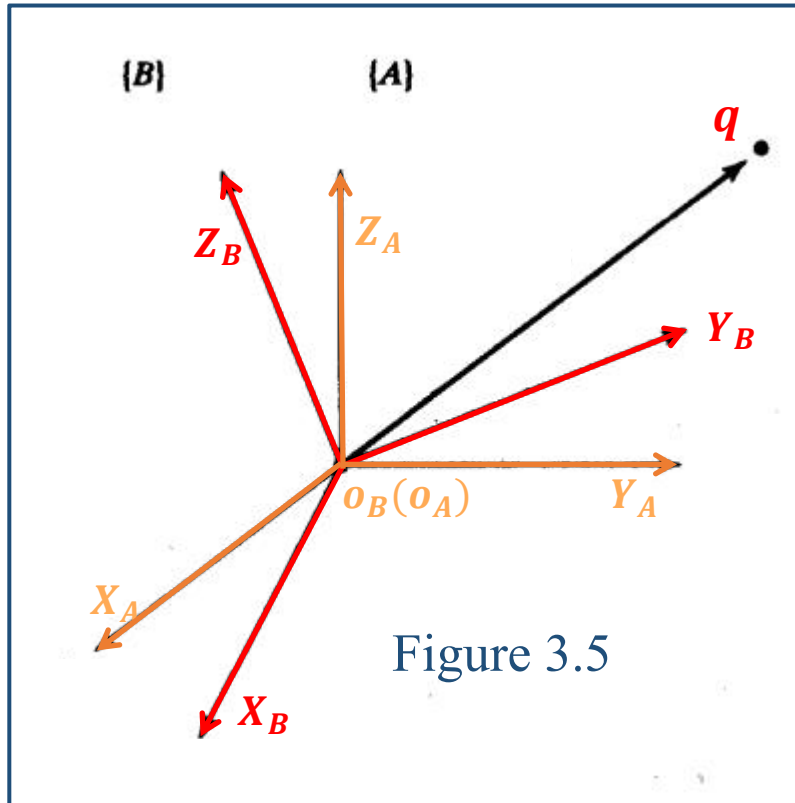


Figure 3.5

Let  $q_b = \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{B\}$

Let  $q_a = \begin{bmatrix} x_a \\ y_a \\ z_a \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{A\}$



$$\begin{aligned} q_a &= x_{ab} \cdot x_b + y_{ab} \cdot y_b + z_{ab} \cdot z_b \\ &= \begin{bmatrix} x_{ab} & y_{ab} & z_{ab} \end{bmatrix} \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} \\ &= R_{ab} \cdot q_b \end{aligned}$$

**Definition:**  $SO(3) = \{R \in \mathbb{R}^{3 \times 3} | R^T R = I, \det R = 1\}$

$R_{ab} \in SO(3)$  is the rotation matrix of frame  $\{B\}$  w.r.t frame  $\{A\}$

### 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$ (Property of Rotation Matrix)



Let  $R = [r_1 \ r_2 \ r_3]$  be a rotation matrix

$$\Rightarrow r_i^T \cdot r_j = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$$

or

$$R^T \cdot R = \begin{bmatrix} r_1^T \\ r_2^T \\ r_3^T \end{bmatrix} [r_1 \ r_2 \ r_3] = I$$

or  $R \cdot R^T = I$

We have:

$$\det(R^T R) = \det R^T \cdot \det R = (\det R)^2 = 1, \det R = \pm 1$$

$$\text{As } \det R = r_1^T (r_2 \times r_3) = 1 \Rightarrow \det R = 1$$

### 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$ (Property of Rotation Matrix)



**Property 2:**  $R_{ab}$  preserves distance between points and orientation.

①  $\|R_{ab} \cdot (p_b - q_b)\| = \|p_a - q_a\|$

②  $R(v \times \omega) = (Rv) \times R\omega$

**Proof :**

For  $a \in \mathbb{R}^3$ , let  $\hat{a} = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix}$

Note that  $\hat{a} \cdot b = a \times b$

① follows from 
$$\begin{aligned} \|R_{ab}(p_b - p_a)\|^2 &= (R_{ab}(p_b - p_a))^T R_{ab}(p_b - p_a) \\ &= (p_b - p_a)^T R_{ab}^T R_{ab} (p_b - p_a) \\ &= \|p_b - p_a\|^2 \end{aligned}$$

② follows from  $R\hat{v}R^T = (\hat{Rv})$  (prove it yourself) □

### 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$ (Translation Vector)

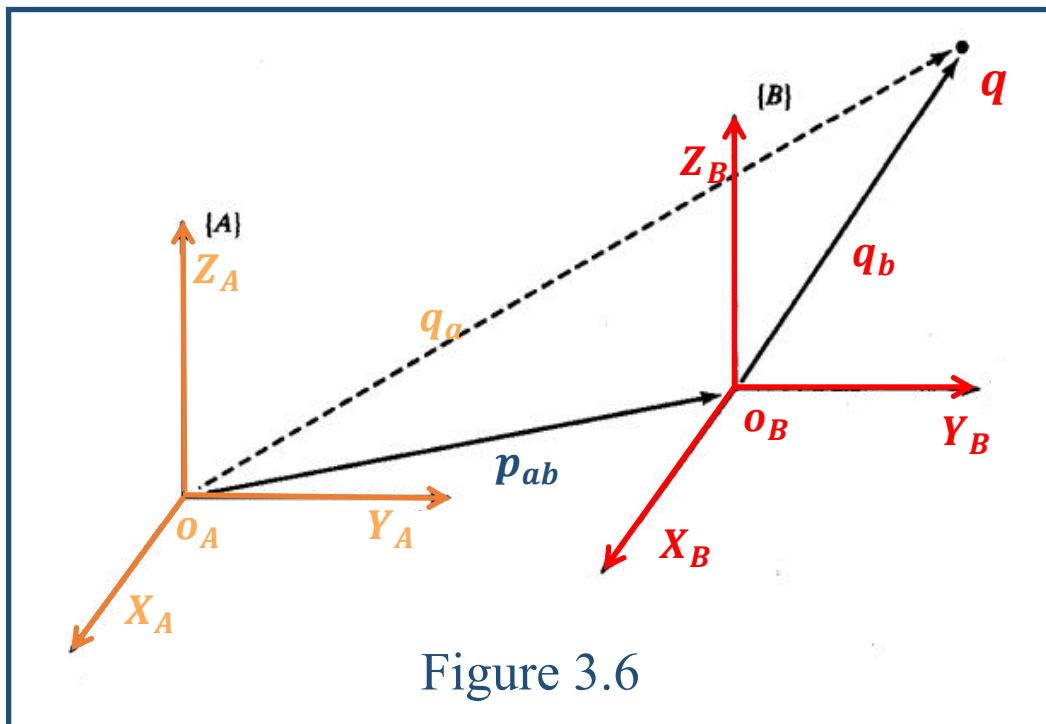


Figure 3.6

Let  $q_b = \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{B\}$

Let  $q_a = \begin{bmatrix} x_a \\ y_a \\ z_a \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{A\}$



$$q_a = q_b + p_{ab}$$

$$p_{ab} = \begin{bmatrix} O_{Bx} - O_{Ax} \\ O_{By} - O_{Ay} \\ O_{Bz} - O_{Az} \end{bmatrix}$$

$p_{ab} \in \mathbb{R}^3$  is the translation vector of frame  $\{B\}$  w.r.t frame  $\{A\}$

### 3.3 Coordinate Frame Transformations in $\mathbb{R}^3$ (General)

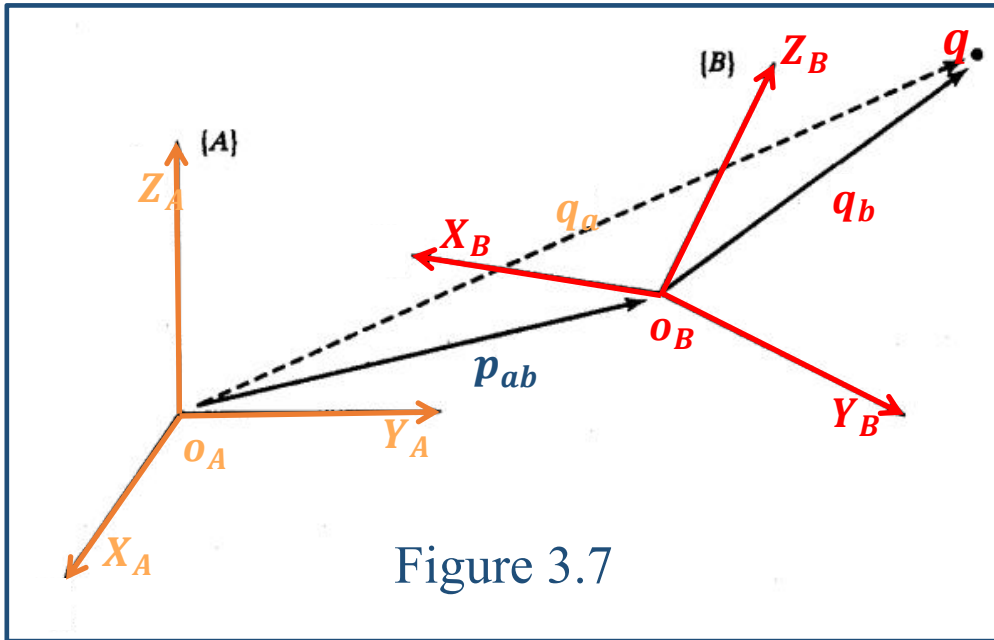


Figure 3.7

Let  $q_b = \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{B\}$

Let  $q_a = \begin{bmatrix} x_a \\ y_a \\ z_a \end{bmatrix} \in \mathbb{R}^3$ : coordinates of point  $q$  in frame  $\{A\}$

$$q_a = R_{ab} \cdot q_b + p_{ab}$$

A general coordinate frame transformation consists of two parts:

- Rotation Matrix  $R \in \mathbb{R}^{3 \times 3}$
- Translation Vector  $p \in \mathbb{R}^{3 \times 1}$

$$\begin{bmatrix} q_a \\ 1 \end{bmatrix} = \begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 0 & 0 & 1 \end{bmatrix} \cdot \begin{bmatrix} q_b \\ 1 \end{bmatrix}$$
$$\bar{q}_a = T_{ab} \cdot \bar{q}_b$$

$$T_{ab} = \begin{bmatrix} R_{ab} & p_{ab} \\ 0_{1 \times 3} & 1 \end{bmatrix} \text{ is the transformation matrix of } \{B\} \text{ w.r.t } \{A\}$$

## 3.4 Rigid Motion in $\mathbb{R}^3$

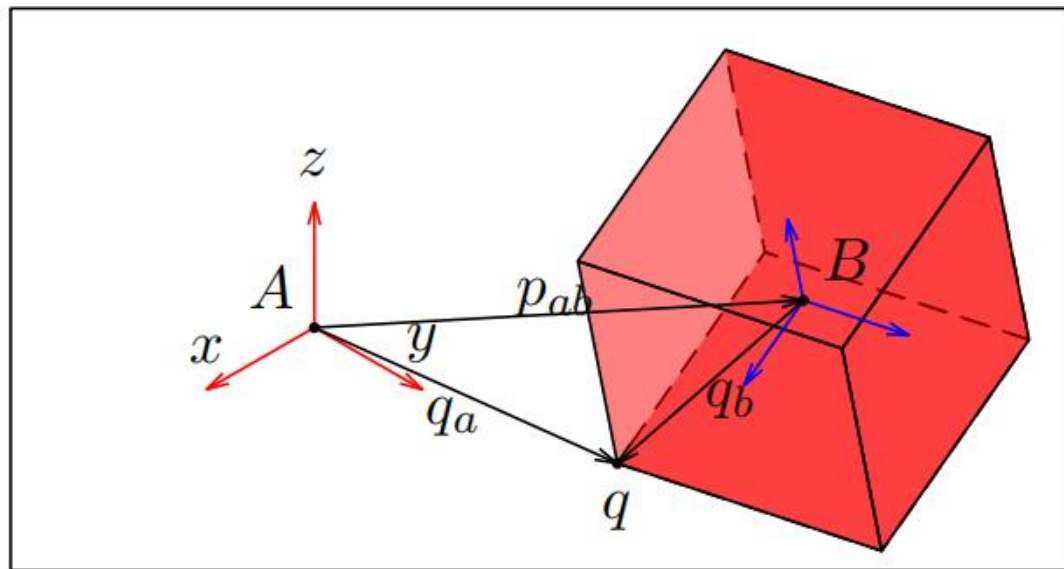


Figure 3.8

**Review:**  $SO(3) = \{R \in \mathbb{R}^{3 \times 3} | R^T R = I, \det R = 1\}$

$p_{ab} \in \mathbb{R}^3$ :	Coordinates of the origin of $B$
$R_{ab} \in SO(3)$ :	Orientation of $B$ relative to $A$
$SE(3) : \left\{ \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix} \middle  p \in \mathbb{R}^3, R \in SO(3) \right\}$ :	Orientation of $B$ relative to $A$

Or...as a transformation:

$$g_{ab} = (p_{ab}, R_{ab}) : \mathbb{R}^3 \mapsto \mathbb{R}^3$$

$$q_b \mapsto q_a = p_{ab} + R_{ab} \cdot q_b$$



## 3.4 Rigid Motion in $\mathbb{R}^3$ (Homogeneous Representation)



Points:

$$q = \begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} \in \mathbb{R}^3$$



$$\bar{q} = \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ 1 \end{bmatrix} \in \mathbb{R}^4$$

Vectors:

$$v = p - q = \begin{bmatrix} p_1 - q_1 \\ p_2 - q_2 \\ p_3 - q_3 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix}$$



$$\bar{v} = \bar{p} - \bar{q} = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \\ 1 \end{bmatrix} - \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ 1 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ 0 \end{bmatrix}$$

- ① Point-Point = Vector
- ② Vector+Point = Point
- ③ Vector+Vector = Vector
- ④ Point+Point: Meaningless

### 3.4 Rigid Motion in $\mathbb{R}^3$ (Homogeneous Representation)



$$q_a = p_{ab} + R_{ab} \cdot q_b$$

$$\begin{bmatrix} q_a \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix}}_{\bar{g}_{ab}} \begin{bmatrix} q_b \\ 1 \end{bmatrix}$$

$$\bar{q}_a = \bar{g}_{ab} \cdot \bar{q}_b$$

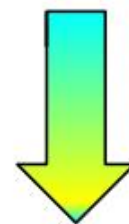
□ **Composition Rule:**

$$\bar{q}_b = \bar{g}_{bc} \cdot \bar{q}_c$$

$$\bar{q}_a = \bar{g}_{ab} \cdot \bar{q}_b = \underbrace{\bar{g}_{ab} \cdot \bar{g}_{bc}}_{\bar{g}_{ac}} \cdot \bar{q}_c$$

$$\bar{g}_{ac} = \bar{g}_{ab} \cdot \bar{g}_{bc} = \begin{bmatrix} R_{ab}R_{bc} & R_{ab}p_{bc} + p_{ab} \\ 0 & 1 \end{bmatrix}$$

$$g_{ab} = (p_{ab}, R_{ab})$$



$$\bar{g}_{ab} = \begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix}$$

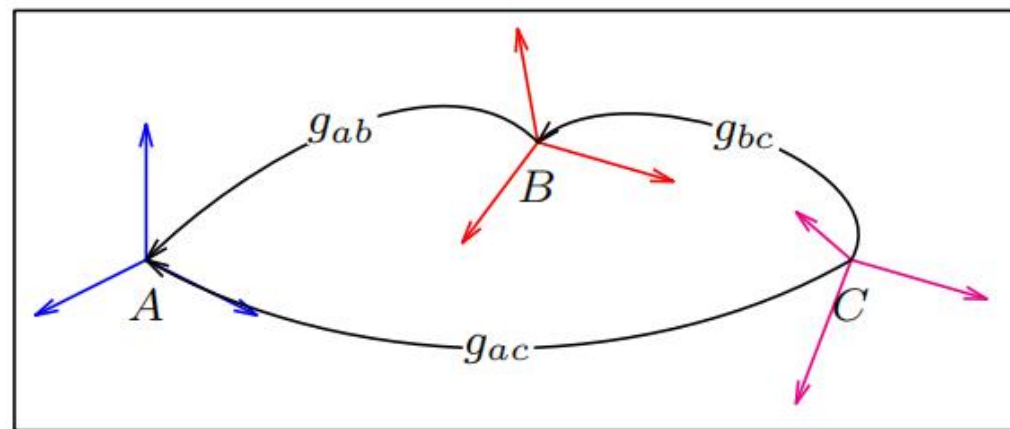


Figure 3.9



## 3.5 Exponential Coordinates and Screw Theory (Parametrization of Rotation Matrix)

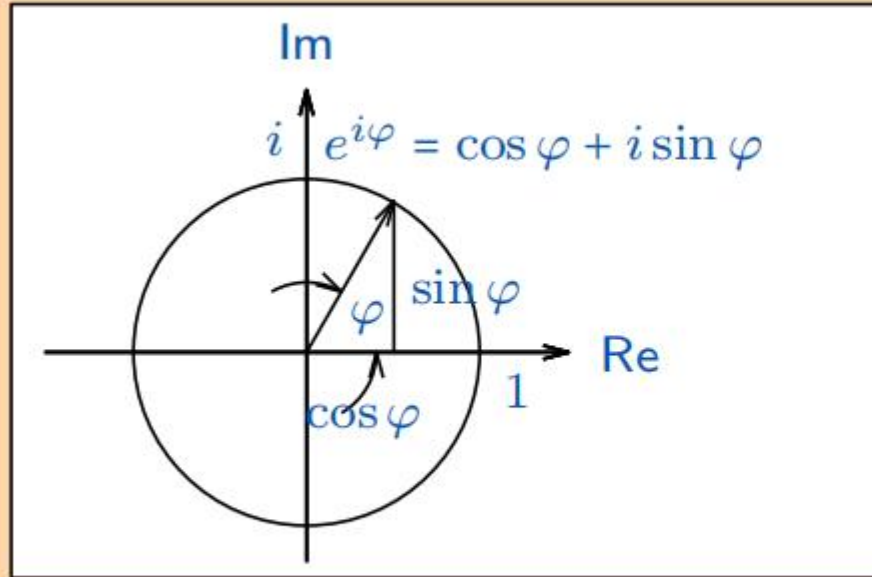


Figure 3.10

### Euler's Formula

“One of the most remarkable, almost astounding, formulas in all of mathematics.”

R. Feynman

### ◇ Review:

$$\begin{cases} \dot{x}(t) = ax(t) \\ x(0) = x_0 \end{cases} \Rightarrow x(t) = e^{at}x_0$$

## 3.5 Exponential Coordinates and Screw Theory (Parametrization of Rotation Matrix)



$$R \in SO(3), R = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix}$$
$$r_i \cdot r_j = \begin{cases} 0, & i \neq j \\ 1, & i = j \end{cases} \leftarrow 6 \text{ constraints}$$

$\Rightarrow 3$  independent parameters!

Consider motion of a point  $q$  on a rotating link

$$\begin{cases} \dot{q}(t) = \omega \times q(t) = \hat{\omega}q(t) \\ q(0): \text{Initial coordinates} \end{cases}$$

$$\Rightarrow q(t) = e^{\hat{\omega}t} q_0 \text{ where } e^{\hat{\omega}t} = I + \hat{\omega}t + \frac{(\hat{\omega}t)^2}{2!} + \frac{(\hat{\omega}t)^3}{3!} + \dots$$

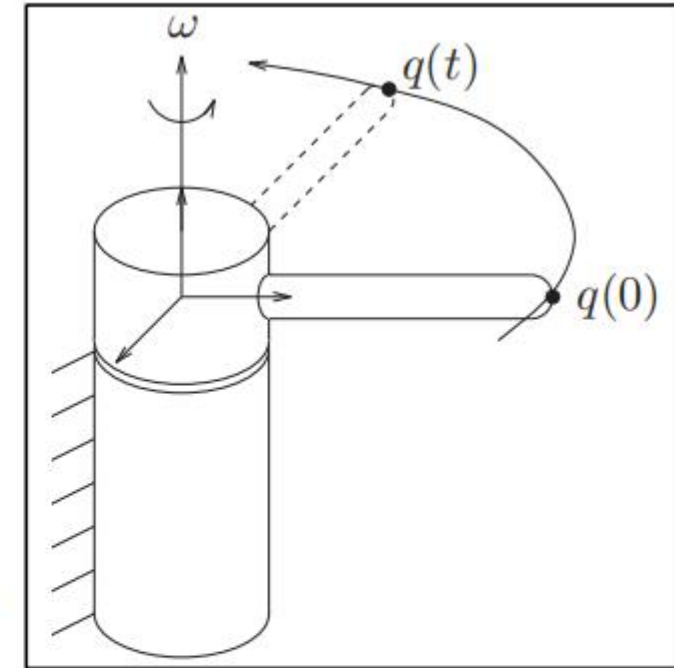


Figure 3.11

By the definition of rigid transformation,  $R(\omega, \theta) = e^{\hat{\omega}\theta}$

## 3.5 Exponential Coordinates and Screw Theory (Rodrigues formula)



**Rodrigues' formula** ( $\|\omega\| = 1$ ):

$$e^{\hat{\omega}\theta} = I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta)$$

**Proof :**

Let  $a \in \mathbb{R}^3$ , write

$$a = \omega\theta, \omega = \frac{a}{\|a\|} \text{ (or } \|\omega\| = 1), \text{ and } \theta = \|a\|$$

$$e^{\hat{\omega}\theta} = I + \hat{\omega}\theta + \frac{(\hat{\omega}\theta)^2}{2!} + \frac{(\hat{\omega}\theta)^3}{3!} + \dots$$

As

we have:

$$\hat{a}^2 = aa^T - \|a\|^2 I, \hat{a}^3 = -\|a\|^2 \hat{a}$$

$$\begin{aligned} e^{\hat{\omega}\theta} &= I + \left(\theta - \frac{\theta^3}{3!} + \frac{\theta^3}{5!} - \dots\right)\hat{\omega} + \left(\frac{\theta^2}{2!} - \frac{\theta^4}{4!} + \dots\right)\hat{\omega}^2 \\ &= I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta) \end{aligned}$$



## 3.5 Exponential Coordinates and Screw Theory (Rodrigues formula)



**Rodrigues' formula for  $\|\omega\| \neq 1$ :**

$$e^{\hat{\omega}\theta} = I + \frac{\hat{\omega}}{\|\omega\|} \sin \|\omega\|\theta + \frac{\hat{\omega}^2}{\|\omega\|^2} (1 - \cos \|\omega\|\theta)$$

**Proof for Property 3:**

Let  $R \triangleq e^{\hat{\omega}\theta}$ , then:

$$\begin{aligned} (e^{\hat{\omega}\theta})^{-1} &= e^{-\hat{\omega}\theta} = e^{\hat{\omega}^T \theta} = (e^{\hat{\omega}\theta})^T \\ \Rightarrow R^{-1} &= R^T \Rightarrow R^T R = I \Rightarrow \det R = \pm 1 \end{aligned}$$

From  $\det \exp(0) = 1$ , and the continuity of  $\det$  function w.r.t.  $\theta$ , we have  $\det e^{\hat{\omega}\theta} = 1, \forall \theta \in \mathbb{R}$  □



## 3.5 Exponential Coordinates and Screw Theory (Rodrigues formula)



□ Let  $\omega = \begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}$ , we have  $\hat{\omega} = \begin{bmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & \omega_1 \\ -\omega_2 & -\omega_1 & 0 \end{bmatrix}$  and  $\hat{\omega} = -\hat{\omega}^T$

□ Let  $v\theta = 1 - \cos\theta$ ,  $c\theta = \cos\theta$ ,  $s\theta = \sin\theta$ ,  $\|\omega\| = 1$ . By Rodrigues' formula,

$$e^{\hat{\omega}\theta} = \begin{bmatrix} \omega_1^2 v\theta + c\theta & \omega_1 \omega_2 v\theta - \omega_3 s\theta & \omega_1 \omega_3 v\theta + \omega_2 s\theta \\ \omega_1 \omega_2 v\theta + \omega_3 s\theta & \omega_2^2 v\theta + c\theta & \omega_2 \omega_3 v\theta - \omega_1 s\theta \\ \omega_1 \omega_3 v\theta - \omega_2 s\theta & \omega_2 \omega_3 v\theta + \omega_1 s\theta & \omega_3^2 v\theta + c\theta \end{bmatrix} = R = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix} \in SO(3)$$

□ Let  $\text{tr}(R) = r_{11} + r_{22} + r_{33} = \sum_{i=1}^3 \lambda_i$ , where  $\lambda_i$  is the eigenvalue of  $R$ ,  $i = 1, 2, 3$

Case 1:  $\text{tr}(R) = 3$  or  $R = I$ ,  $\theta = 0 \Rightarrow \omega\theta = 0$

Case 2:  $-1 < \text{tr}(R) < 3$ ,

$$\theta = \arccos \frac{\text{tr}(R) - 1}{2} \Rightarrow \omega = \frac{1}{2s\theta} \begin{bmatrix} r_{32} - r_{23} \\ r_{13} - r_{31} \\ r_{21} - r_{12} \end{bmatrix}$$

Case 3:  $\text{tr}(R) = -1 \Rightarrow \cos\theta = -1 \Rightarrow \theta = \pm\pi$

## 3.5 Exponential Coordinates and Screw Theory (Exponential coordinates of SE(3))



**For rotational motion:**

$$\dot{p}(t) = \omega \times (p(t) - q)$$

$$\begin{bmatrix} \dot{p} \\ 0 \end{bmatrix} = \begin{bmatrix} \hat{\omega} & -\omega \times q \\ 0 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix}$$

or  $\dot{\bar{p}} = \hat{\xi} \cdot \bar{p} \Rightarrow \bar{p}(t) = e^{\hat{\xi}t} \bar{p}(0)$

where  $e^{\hat{\xi}t} = I + \hat{\xi}t + \frac{(\hat{\xi}t)^2}{2!} + \dots$

**For translational motion:**

$$\dot{p}(t) = v$$

$$\begin{bmatrix} \dot{p}(t) \\ 0 \end{bmatrix} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix}$$

$$\dot{\bar{p}}(t) = \hat{\xi} \cdot \bar{p}(t) \Rightarrow \bar{p}(t) = e^{\hat{\xi}t} \bar{p}(0)$$

$$\hat{\xi} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix}$$

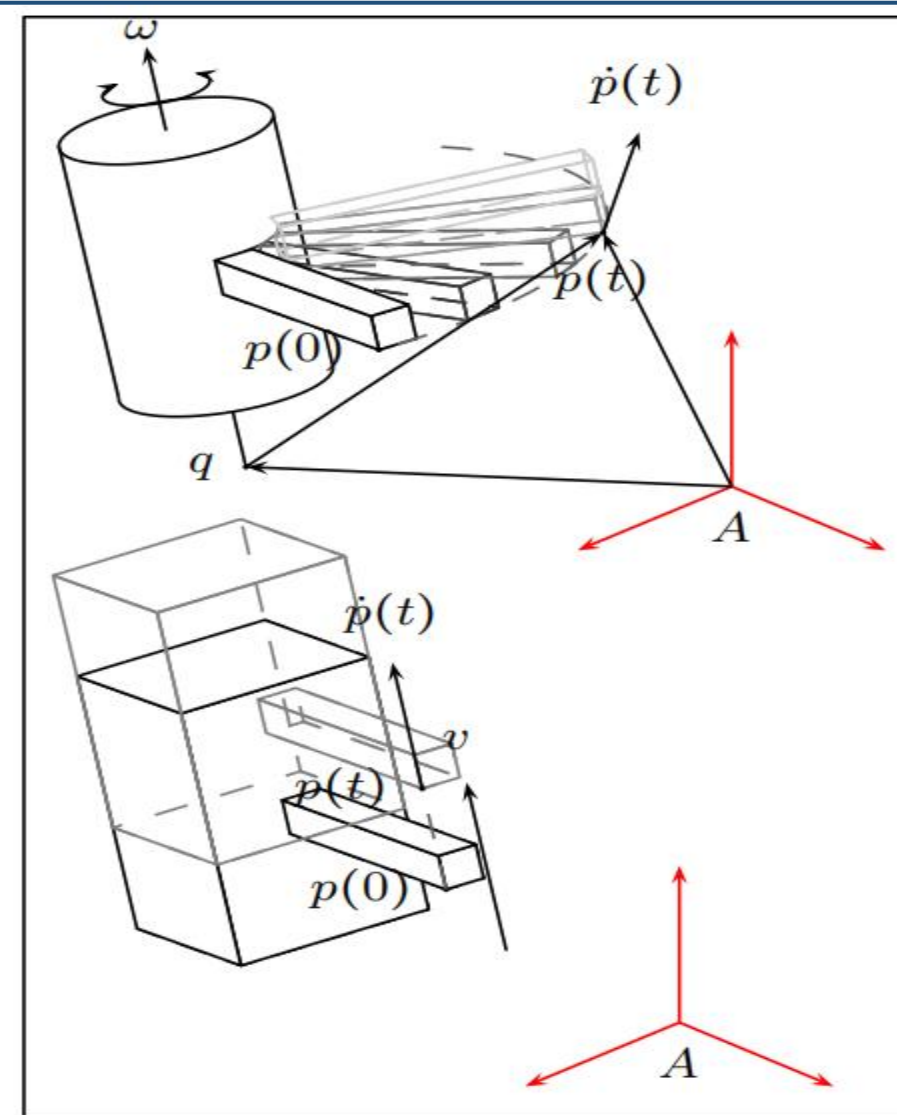


Figure 3.12

## 3.5 Exponential Coordinates and Screw Theory (Exponential coordinates of SE(3))



$$\xi := \begin{bmatrix} v \\ \omega \end{bmatrix} \mapsto \hat{\xi} = \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix}$$

- If  $\omega = 0$ , then  $\hat{\xi}^2 = \hat{\xi}^3 = \dots = 0$ ,  $e^{\hat{\xi}\theta} = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix} \in SE(3)$

- If  $\omega$  is not 0, assume  $\|\omega\| = 1$ .

Define:

$$g_0 = \begin{bmatrix} I & \omega \times v \\ 0 & 1 \end{bmatrix}, \hat{\xi}' = g_0^{-1} \cdot \hat{\xi} \cdot g_0 = \begin{bmatrix} \hat{\omega} & h\omega \\ 0 & 0 \end{bmatrix}$$

where  $h = \omega^T \cdot v$ .

$$e^{\hat{\xi}\theta} = e^{g_0 \cdot \hat{\xi}' \cdot g_0^{-1}} = g_0 \cdot e^{\hat{\xi}'\theta} \cdot g_0^{-1}$$

and as

$$\hat{\xi}'^2 = \begin{bmatrix} \hat{\omega}^2 & 0 \\ 0 & 0 \end{bmatrix}, \hat{\xi}'^3 = \begin{bmatrix} \hat{\omega}^3 & 0 \\ 0 & 0 \end{bmatrix}$$

we have

$$e^{\hat{\xi}'\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & h\omega\theta \\ 0 & 1 \end{bmatrix} \Rightarrow e^{\hat{\xi}\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})\hat{\omega}v + \omega\omega^T v\theta \\ 0 & 1 \end{bmatrix}$$

## 3.5 Exponential Coordinates and Screw Theory (Screw Theory)



Screw attributes

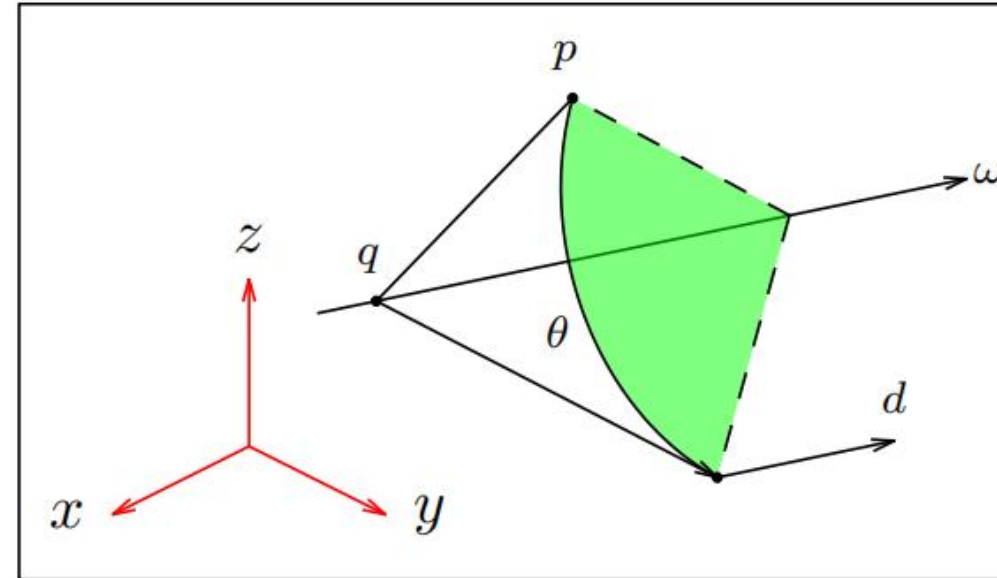


Figure 3.13

$$\begin{aligned} \text{Pitch:} \quad & h = \frac{d}{\theta} (\theta = 0, h = \infty), d = h \cdot \theta \\ \text{Axis:} \quad & l = \{q + \lambda \omega \mid \lambda \in \mathbb{R}\} \\ \text{Magnitude:} \quad & M = \theta \end{aligned}$$

### Definition:

A **screw**  $S$  consists of an axis  $l$ , pitch  $h$ , and magnitude  $M$ . A **screw motion** is a rotation by  $\theta = M$  about  $l$ , followed by translation by  $h\theta$ , parallel to  $l$ . If  $h = \infty$ , then, translation about  $v$  by  $\theta = M$



## 3.5 Exponential Coordinates and Screw Theory (Screw Theory)



Corresponding  $g \in SE(3)$ :

$$g \cdot p = q + e^{\hat{\omega}\theta}(p - q) + h\theta\omega$$

$$g \cdot \begin{bmatrix} p \\ 1 \end{bmatrix} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})q + h\theta\omega \\ 0 & 1 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix} \Rightarrow$$

$$g = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})q + h\theta\omega \\ 0 & 1 \end{bmatrix}$$

On the other hand...

$$e^{\hat{\xi}\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})\omega \times v + \omega\omega^T v\theta \\ 0 & 1 \end{bmatrix}$$

If we let  $v = -\omega \times q + h\omega$ , then

$$(I - e^{\hat{\omega}\theta})(-\hat{\omega}^2 q) = (I - e^{\hat{\omega}\theta})(-\omega\omega^T q + q) = (I - e^{\hat{\omega}\theta})q$$

Thus,  $e^{\hat{\xi}\theta} = g$

For pure rotation ( $h = 0$ ):  $\xi = (-\omega \times q, \omega)$

For pure translation:  $g = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix}$ ,  $\Rightarrow \xi = (v, 0)$ , and  $e^{\hat{\xi}\theta} = g$

## 3.5 Exponential Coordinates and Screw Theory (Screw Theory)



$$\xi = (v, \omega) \in \mathbb{R}^6$$

- ① Pitch:  $h = \begin{cases} \frac{\omega^T v}{\|\omega\|^2}, & \text{if } \omega \neq 0 \\ \infty, & \text{if } \omega = 0 \end{cases}$
- ② Axis:  $l = \begin{cases} \frac{\omega \times v}{\|\omega\|^2} + \lambda \omega, & \lambda \in \mathbb{R}, \text{ if } \omega \neq 0 \\ 0 + \lambda v & \lambda \in \mathbb{R}, \text{ if } \omega = 0 \end{cases}$
- ③ Magnitude:  $M = \begin{cases} \|\omega\|, & \text{if } \omega \neq 0 \\ \|v\|, & \text{if } \omega = 0 \end{cases}$

Screw	Twist: $\hat{\xi}\theta$
Case 1: Pitch: $h = \infty$ Axis: $l = \{q + \lambda v   \ v\  = 1, \lambda \in \mathbb{R}\}$ Magnitude: $M$	$\theta = M,$ $\hat{\xi} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix}$
Case 2: Pitch: $h \neq \infty$ Axis: $l = \{q + \lambda \omega   \ \omega\  = 1, \lambda \in \mathbb{R}\}$ Magnitude: $M$	$\theta = M,$ $\hat{\xi} = \begin{bmatrix} \hat{\omega} & -\hat{\omega}q + h\omega \\ 0 & 0 \end{bmatrix}$

### Special cases:

- ①  $h = \infty$ , Pure translation (prismatic joint)
- ②  $h = 0$ , Pure rotation (revolute joint)

### Definition: Screw Motion

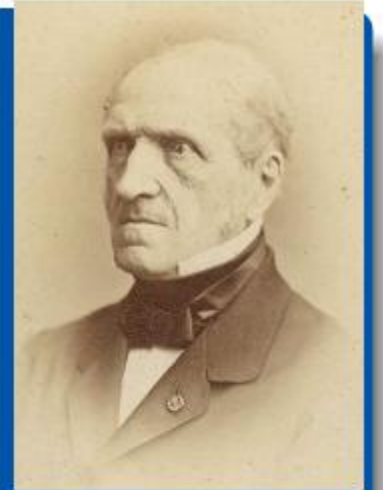
Rotation about an axis by  $\theta = M$ , followed by translation about the same axis by  $h\theta$

## 3.5 Exponential Coordinates and Screw Theory (Chasles Theorem)



### Theorem 2 (Chasles):

Every rigid body motion can be realized by a rotation about an axis combined with a translation parallel to that axis.



1793–1880

### Proof :

For  $\hat{\xi} \in se(3)$ :

$$\hat{\xi} = \hat{\xi}_1 + \hat{\xi}_2 = \begin{bmatrix} \hat{\omega} & -\omega \times q \\ 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & h\omega \\ 0 & 0 \end{bmatrix}$$
$$[\hat{\xi}_1, \hat{\xi}_2] = 0 \Rightarrow e^{\hat{\xi}\theta} = e^{\hat{\xi}_1\theta} e^{\hat{\xi}_2\theta}$$



### 3.6 Other Parametrizations of SO(3) (XYZ fixed angles)

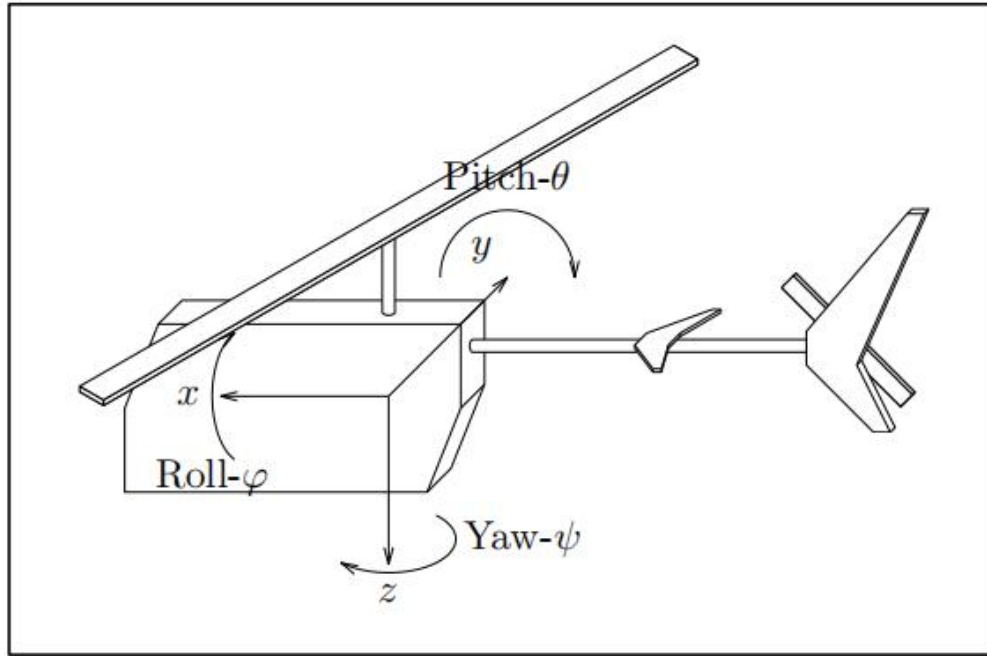


Figure 3.14

#### □ XYZ fixed angles (or Roll-Pitch-Yaw angle)

$$R_x(\varphi) := e^{\hat{x}\varphi} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \varphi & -\sin \varphi \\ 0 & \sin \varphi & \cos \varphi \end{bmatrix}$$
$$R_y(\theta) := e^{\hat{y}\theta} = \begin{bmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{bmatrix}$$
$$R_z(\psi) := e^{\hat{z}\psi} = \begin{bmatrix} \cos \psi & -\sin \psi & 0 \\ \sin \psi & \cos \psi & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

1. Rotate  $\varphi$  about x-axis
2. Rotate  $\theta$  about y-axis
3. Rotate  $\psi$  about z-axis



$$R_{ab} = R_z(\psi)R_y(\theta)R_x(\varphi)$$
$$R_{ab} = \begin{bmatrix} c\psi c\theta & -s\psi c\varphi + c\psi s\theta s\varphi & s\psi s\varphi + c\psi s\theta c\varphi \\ s\psi c\theta & c\psi c\varphi + s\psi s\theta s\varphi & -c\psi s\varphi + s\psi s\theta c\varphi \\ -s\theta & c\theta s\varphi & c\theta c\varphi \end{bmatrix}$$



## 3.6 Other Parametrizations of SO(3) (ZYX Euler angle)



### □ ZYX Euler angle

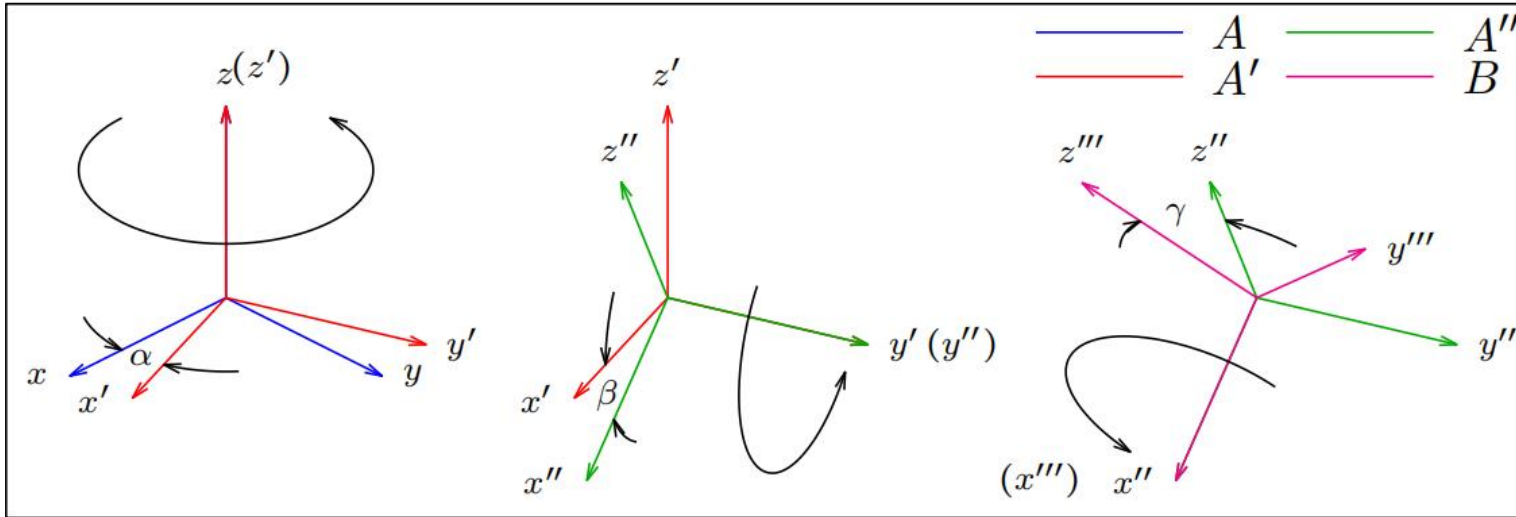


Figure 3.15

$$\beta = \text{atan2}(-r_{31}, \sqrt{r_{32}^2 + r_{33}^2})$$

$$\alpha = \text{atan2}(r_{21}/c_\beta, r_{11}/c_\beta)$$

$$\gamma = \text{atan2}(r_{32}/c_\beta, r_{33}/c_\beta)$$

1. Rotate  $\alpha$  about z-axis
2. Rotate  $\beta$  about  $y'$ -axis
3. Rotate  $\gamma$  about  $x''$ -axis



$$R_{aa'} = R_z(\alpha)$$

$$R_{a'a''} = R_y(\beta)$$

$$R_{a''b} = R_x(\gamma)$$

$$R_{ab} = R_z(\alpha)R_y(\beta)R_x(\gamma)$$

$$R_{ab}(\alpha, \beta, \gamma) = \begin{bmatrix} c_\alpha c_\beta & -s_\alpha c_\gamma + c_\alpha s_\beta s_\gamma & s_\alpha s_\gamma + c_\alpha s_\beta c_\gamma \\ s_\alpha c_\beta & c_\alpha c_\gamma + s_\alpha s_\beta s_\gamma & -c_\alpha s_\gamma + s_\alpha s_\beta c_\gamma \\ -s_\beta & c_\beta s_\gamma & c_\beta c_\gamma \end{bmatrix}$$

**Note:** When  $\beta = \frac{\pi}{2}$ ,  $\cos \beta = 0$ ,  $\alpha + \gamma = \text{const} \Rightarrow \text{singularity!}$

## 3.6 Other Parametrizations of $SO(3)$ (Quaternions)



### § Quaternions:

$$Q = q_0 + q_1i + q_2j + q_3k$$

$$\text{where } i^2 = j^2 = k^2 = -1, i \cdot j = k, j \cdot k = i, k \cdot i = j$$

**Property 1:** Define  $Q^* = (q_0, q)^* = (q_0, -q)$ ,  $q_0 \in \mathbb{R}, q \in \mathbb{R}^3$

$$\|Q\|^2 = QQ^* = q_0^2 + q_1^2 + q_2^2 + q_3^2$$

**Property 2:**  $Q = (q_0, q), P = (p_0, p)$

$$QP = (q_0p_0 - q \cdot p, q_0p + p_0q + q \times p)$$

**Property 3:** (a) The set of unit quaternions forms a group

(b) If  $R = e^{\hat{\omega}\theta}$ , then  $Q = (\cos \frac{\theta}{2}, \omega \sin \frac{\theta}{2})$

(c)  $Q$  acts on  $x \in \mathbb{R}^3$  by  $QXQ^*$ , where  $X = (0, x)$

## 3.6 Other Parametrizations of SO(3) (Quaternions)



$$q_0 = \cos \frac{\theta}{2}, q = \omega \sin \frac{\theta}{2}$$

and the Rodrigues' formula:

$$e^{\hat{\omega}\theta} = I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta)$$

then

$$\begin{aligned} R(Q) &= I + 2q_0\hat{q} + 2\hat{q}^2 \\ &= \begin{bmatrix} 1 - 2(q_2^2 + q_3^2) & -2q_0q_3 + 2q_1q_2 & 2q_0q_2 + 2q_1q_3 \\ 2q_0q_3 + 2q_1q_2 & 1 - 2(q_1^2 + q_3^2) & -2q_0q_1 + 2q_2q_3 \\ -2q_0q_2 + 2q_1q_3 & 2q_0q_1 + 2q_2q_3 & 1 - 2(q_1^2 + q_2^2) \end{bmatrix} \end{aligned}$$

where  $\|Q\| \triangleq q_0^2 + q_1^2 + q_2^2 + q_3^2 = 1$

■ Conversion from Roll-Pitch-Yaw angle to unit quaternions

$$Q = \left(\cos \frac{\varphi}{2}, x \sin \frac{\varphi}{2}\right) \left(\cos \frac{\theta}{2}, y \sin \frac{\theta}{2}\right) \left(\cos \frac{\psi}{2}, z \sin \frac{\psi}{2}\right)$$



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