Title

Ben

1 BPS states

Definition 1.1. BPS states or counts are invariants of 3CY categories.

A CY category is a category with a duality condition on associated algebras.

Examples of 3 CY categories:

- 1. Coherent sheaves on X where X is a smooth projective 3-fold over \mathbb{C} with trivial canonical bundle. Serre duality $E, F \in \text{Coh}(X)$ then $\text{Ext}^i(E, F) \cong \text{Ext}(F, E)^*$.
- 2. Local systems on a real 3 dimensional closed orientable manifold. IN this case $H^i*(M, E^*\otimes F)\cong H^{3-i}(M, F^*\otimes E^*)^*$.
- 3. M as above. Then Representations of $\mathbb{C}[\pi_1(M)]$.

Any one of these examples gives us the input of DT theory.

Let C be 3CY category. Introduce the notion of charge lattice. Target for some [[SOMETHING]] invariant of algebras of $\mathbb C$

Let $N = \mathbb{Z}$ - dimension of $\mathbb{C}[\pi(M)]$ - module

 \mathcal{M}_{γ} for $\gamma \in N$ stack of objects in C of class γ . This is manageably small.

Vanishing cycles: Say $\operatorname{Spec}(\mathbb{C}[x]/x^d)$ is a moduli space of objects in some category B. This naively has a only a single object of a single point. However a deformation of x^d will see the point split into d objects. Thus we would like $\#\operatorname{Spec}(\mathbb{C}[x]/x^d)$ to be d and not 1.

Consider $H^*(X, \mathbb{Q})$ [[CANT SEE]]

DT theory is the study of the cohomology $H(M(C), \mathcal{Q})$ of the stack M(C) over the category into sheaf \mathcal{Q} We use the charge lattice and mixed hodge structures to decompose this space $D(MHS_N)$.

More precisely $[H(M(C), \mathcal{Q})] \in K_0(MHS_N)$ where $K_0(MHS_N) = \mathbb{Z}[[L]]/([L'] + [L''] = [L])$ if $L' \to L \to L'' \to \Delta \in D(MHS_N)$. This is huge and unwieldly.

How to even write this down? Answer: via Plethystic exponential EXP, $f = \sum_{d \in \mathbb{Z}^m} a_d x^d$, EXP $(f) = \prod_{d \in \mathbb{Z}^m} (1 - x^d)^{-a_d}$.

Where does this come from? Origin: consider

$$K(\operatorname{Vect}_{\mathbb{Z}^m}) \xrightarrow{\chi} \mathbb{Z}[[\mathbb{Z}^m]]$$
 (1)

$$[v] \mapsto \sum_{d} \dim V_d x^d$$
 (2)

but this ill defined...

cf Grothendeick something. Commutative diagramme between characters and symmetric algebra maps. [[ERRRRRR]]

Question: How does on understand $[H(M(C), \mathcal{Q})]$? Its just too big.

Integrality conjecture (historically misnamed): This says that $[H(M(C), \mathcal{Q})] = \text{EXP}(\bigoplus_{\gamma} BPS_{\gamma} \otimes H_{\mathbb{C}^*}())$ where BPS_{γ} is a finite dimensional mixed hodge structures.

Many stupid things have appeared in this talk. We should try to address these...

Let M be an n dimensional orientable closed real manifold. $\mathbb{C}[\pi_1(M)]$ -modules. nCY categories, with $\mathrm{Ext}^i(E,F)\cong\mathrm{Ext}^{n-i}(F,E)^*$ we want n=2 not 3.

Theorem 1.2. (BEN) Let B be a 2CY category. Let C be a category of pairs (E, f) for $E \in B$, and $f: E \to E^*$. Then C is 3CY and $H(M(C), \mathcal{Q}) \cong H(M(B), \mathbb{Q})$

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Conclusion: [M(\mathbb{C}[\pi_1(\Sigma_g)])] = \text{EXP}([\oplus_n BPS_{n,g} \otimes H_{\mathbb{C}^*}(pts)])
Hausel Villegas : [BPS_{n,g}] = [H(\text{Rep}_n^1(\Sigma_g)]
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 $\operatorname{Rep}_n^1(\Sigma_g) = [A, B...]/\operatorname{PGL}_n$ and by NAHT \cong Higgs_{n,1}(C)

Categorified Integrability theorem. this is much much stronger.

What does this have to do with the P = W? Only conjectures to follow?

Conjecture: $BPS_{g,n} \cong H(\operatorname{Rep}_n^1(\Sigma_g))$

There's analogue of the conjecture on the Higgs side. There's a network of conjectures.

Conjecture: $H(\operatorname{Rep}(\mathbb{C}[\pi_1(\Sigma_g)])) \cong \operatorname{Sym}(\bigoplus H(\operatorname{Rep}_n^1(\Sigma_g) \otimes H_{\mathbb{C}}(pt)))$

 $H(\operatorname{Higgs}_0^{ss}) \cong \operatorname{Sym}(\bigoplus H(\operatorname{Higgs}(C) \oplus H_{\mathbb{C}^*}(pt)))$

If we believe these conjectures, then P = W becomes equivalent to the an isomorphism between the two LHS above (call it α). α is an isomorphism between 2 quantum groups. They contain lie algebra structures on respective BPSg and BPS-Higgs

Claim: P,W filtrations come from induced \mathfrak{sl}_2 on these Lie algebra.