

QFT I, Homework 4

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Scalar QED Consider the theory of a complex scalar field ϕ interacting with the electromagnetic field A^μ . The Lagrangian is

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_\mu\phi)^* D^\mu\phi - m^2\phi^*\phi. \quad (1)$$

where $D_\mu = \partial_\mu + ieA_\mu$ is the usual gauge covariant derivative.

(a) Show the Lagrangian is invariant under the gauge transformations

$$\phi(x) \rightarrow e^{-i\alpha(x)}\phi(x), \quad A_\mu(x) \rightarrow A_\mu(x) + \frac{1}{e}\partial_\mu\alpha(x). \quad (2)$$

(b) Derive the Feynman rules for the interaction between photons and scalar particles.

(c) Draw all the leading-order Feynman diagrams and compute the amplitude for the process $\gamma\gamma \rightarrow \phi\phi^*$.

(d) Compute the differential cross section $d\sigma/d\cos\theta$. You can take an average over all initial state polarizations. For simplicity, you can restrict your calculation in the limit $m = 0$.

(e) Draw all leading order Feynman diagrams, that contribute to the Compton scattering process $\gamma\phi \rightarrow \gamma\phi$ and compute the differential cross section $d\sigma/d\cos\theta$ with $m = 0$.

Solution

(a) Under the gauge transformation (2), we have

$$F_{\mu\nu} \rightarrow F'_{\mu\nu} = \partial_\mu A'_\nu - \partial_\nu A'_\mu = \partial_\mu \left(A_\nu + \frac{1}{e}\partial_\nu\alpha \right) - \partial_\nu \left(A_\mu + \frac{1}{e}\partial_\mu\alpha \right) = \partial_\mu A_\nu - \partial_\nu A_\mu = F_{\mu\nu},$$

so the first term in (1) remains the same. It is obvious that under (2)

$$\phi^*\phi \rightarrow \phi'^*\phi' = e^{i\alpha}\phi^*e^{-i\alpha}\phi = \phi^*\phi,$$

so the third term in (1) is also invariant. Also we have

$$\begin{aligned} D^\mu\phi &\rightarrow (\partial^\mu + ieA'^\mu)\phi' = (\partial^\mu + ieA^\mu + i\partial^\mu\alpha)e^{-i\alpha}\phi \\ &= e^{-i\alpha}(\partial^\mu - i\partial^\mu\alpha + ieA^\mu + i\partial^\mu\alpha)\phi \\ &= e^{-i\alpha}D^\mu\phi, \end{aligned}$$

and also

$$(D^\mu\phi)^* = e^{i\alpha}D^\mu\phi^*,$$

so $D^\mu\phi(D^\mu\phi)^*$ is also invariant. Therefore (1) is invariant under (2).

(b) Expanding (2) we have

$$\mathcal{L} = \mathcal{L}_{\text{scalar}} + \mathcal{L}_{\text{vector}} + \mathcal{L}_{\text{scalarQED}}, \quad (3)$$

where $\mathcal{L}_{\text{scalar}}$ and $\mathcal{L}_{\text{vector}}$ are Lagrangians of free scalar field and free massless vector field, and

$$\begin{aligned} \mathcal{L}_{\text{scalarQED}} &= (D_\mu\phi)^* D^\mu\phi - (\partial_\mu\phi)^* \partial^\mu\phi \\ &= e^2 A_\mu A^\mu \phi\phi^* - ieA_\mu\phi^* \partial^\mu\phi + ie\partial_\mu\phi^* A^\mu\phi. \end{aligned} \quad (4)$$

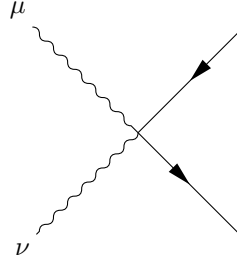
We make the following expansion of Fourier transformation. For the complex scalar field we have

$$\phi(x) = \int \frac{d^3\mathbf{p}}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} (a_{\mathbf{p}} e^{-ip \cdot x} + b_{\mathbf{p}}^\dagger e^{ip \cdot x}). \quad (5)$$

which was proved in (10) in [Homework 2](#). The vector field is expanded as

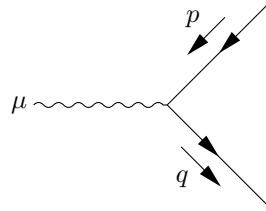
$$A_\mu(x) = \int \frac{d^3\mathbf{p}}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} \sum_{r=1}^2 \epsilon_\mu^r(\mathbf{p}) (a_{\mathbf{p},r}^\dagger e^{ip \cdot x} + a_{\mathbf{p},r} e^{-ip \cdot x}). \quad (6)$$

The first term gives the following (momentum space) diagram:

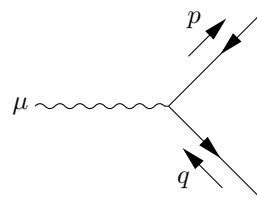


$$= ie^2 \delta_\mu^\nu. \quad (7)$$

The second and the third term gives



$$= -ie(p_\mu + q_\mu), \quad (8)$$



$$= ie(p_\mu + q_\mu). \quad (9)$$

Note

Here we follow the notation of Peskin, i.e. using the arrow *on* a particle line to show whether this line represents a particle or an antiparticle and using the *momentum* arrow to denote whether this line represents creation or annihilation. The direction and sign of a 4-momentum is *not* represented in any arrow.

(c)