1 Problem 5

5.3
$$\frac{2\lambda}{d} \times 50 \text{cm} = 1 \text{cm} \Rightarrow d = 63 \mu \text{m}.$$

5.4
$$D = 2 \times \frac{1.22\lambda}{d} \times 3.76 \times 10^8 \text{m}, \quad \lambda = 632.8 \text{nm}.$$

Table 1: Diameters

on Earth d	on Moon D
2mm	290.3km
2m	290.3m
5m	116.1m

- **5.6** $\theta_0 = 1.22 \lambda/D = 1.342 \times 10^{-7} \text{rad} = 0.02768''$. The angular limit of resolution of a human eye is 1', 1'/0.02768'' = 2167.
- 5.9 (1) $\delta y = 0.61 \lambda / 0.85 = 287 \text{nm}$. (2) $\delta y = 0.61 \lambda / 1.45 = 168 \text{nm}$, the limit of resolution of a human eye is $\delta y' = 25 \text{cm} \times 1' = 72.7 \mu \text{m}$. $\delta y' / \delta y = 432$.
- 5.13 $\theta = 2\lambda/d$, thus $\delta\theta = 2\delta\lambda/d = 5 \times 10^{-6} \text{rad} = 1.03''$. While half the linewidth is $\Delta\theta \approx \frac{\lambda}{Nd} = 5 \times 10^{-6} \text{rad} \sim \delta\theta$. 恰能分辨.

5.16 Let

$$s(x) = \sum_{n=-N}^{N-1} \delta(x - 6a(n+1/2)), \quad t(x) = \begin{cases} 1, & x \le a/2, \\ 0, & \text{otherwise.} \end{cases}$$
 (I)

Their Fourier transformations

$$\mathcal{F}\{s\} := \int_{-\infty}^{\infty} s(x) e^{ikx \sin \theta} dx = \sum_{n=-N}^{N-1} e^{i6ak \sin \theta (n+1/2)} = \frac{\sin 6Nak \sin \theta}{\sin 3ak \sin \theta}, \quad (2)$$

$$\mathcal{F}\{t\} := \int_{-\infty}^{\infty} t(x) e^{ikx \sin \theta} dx = \left. \frac{e^{ikx \sin \theta}}{ik \sin \theta} \right|_{-a/2}^{a/2} = a \frac{\sin \frac{1}{2} ka \sin \theta}{\frac{1}{2} ka \sin \theta}.$$
 (3)

If we define convolution $*: \mathbb{R}^{\mathbb{R}} \times \mathbb{R}^{\mathbb{R}} \to \mathbb{R}^{\mathbb{R}}$ as follows:

$$(f * g)(x) = \int_{-\infty}^{\infty} f(t)g(x - t)dt,$$
(4)

it is well-know that \mathcal{F} is an isomorphism: $\mathcal{F}\{f * g\} = \mathcal{F}\{f\} \cdot \mathcal{F}\{g\}$.

For (\mathbf{I}) even numbered slits are blocked, and ($\mathbf{2}$) odd numbered slits are blocked, both of their aperature functions are s*t. Then the field distribution is

$$E(\theta) = \frac{E_0}{a} \mathcal{F}\{s * t\} = \frac{E_0}{a} \mathcal{F}\{s\} \cdot \mathcal{F}\{t\} = E_0 \left(\frac{\sin 12N\alpha}{\sin 6\alpha}\right) \left(\frac{\sin \alpha}{\alpha}\right),\tag{5}$$

where $\alpha = \frac{1}{2}ka\sin\theta$. Intensity distribution

$$I(\theta) = \frac{I_0}{E_0^2} |E(\theta)|^2 = I_0 \left(\frac{\sin 12N\alpha}{\sin 6N\alpha}\right)^2 \left(\frac{\sin \alpha}{\alpha}\right)^2.$$
 (6)

(3) The aperature function is (s(x) + s(x - 2a)) * t. By calculation:

$$\mathcal{F}\{s(x) + s(x - 2a)\} = 2e^{2i\alpha}\cos 2\alpha \cdot \mathcal{F}\{s\}. \tag{7}$$

Therefore the field distribution is

$$E(\theta) = \frac{E_0}{a} \mathcal{F}\{(s(x) + s(x - 2a)) * t\} = 2E_0 e^{2i\alpha} \cos 2\alpha \left(\frac{\sin 12N\alpha}{\sin 6\alpha}\right) \left(\frac{\sin \alpha}{\alpha}\right), \quad (8)$$

and intensity distribution

$$I(\theta) = \frac{I_0}{E_0^2} |E(\theta)|^2 = 4I_0 \cos^2 2\alpha \left(\frac{\sin 12N\alpha}{\sin 6N\alpha}\right)^2 \left(\frac{\sin \alpha}{\alpha}\right)^2.$$
 (9)

5.18 According to Fresnel-Kirchhoff formula,

$$E = -\frac{iE_0}{\lambda} \int_{L}^{\sqrt{L^2 + (D/2)^2}} 2\pi r dr \frac{e^{ikr}}{r} \frac{1 + L/r}{2} = -i\pi \frac{E_0}{\lambda} \int_{L}^{\sqrt{L^2 + (D/2)^2}} \left(1 + \frac{L}{r}\right) e^{ikr} dr, \text{ (10)}$$

where L=1.5m, D=2.6mm and $\lambda=589$ nm. Since $D\ll L$, we have $L/r\approx 1$, and thus

$$E \approx -ikE_0 \int_{L}^{\sqrt{L^2 + (D/2)^2}} e^{ikr} dr = 2E_0 e^{i\delta} \sin \frac{\pi}{\lambda} \left(\sqrt{L^2 + (D/2)^2} - L \right),$$
 (11)

where $\delta = -\frac{\pi}{2} + \frac{k}{2} (L + \sqrt{L^2 + (D/2)^2})$. That is

$$I = \frac{I_0}{E_0^2} |E|^2 \approx 4I_0 \sin^2 \frac{\pi D^2}{8\lambda L} = 4I_0 \times 0.14.$$
 (12)

Then center point is dark.

To make the center point bright, let $\frac{\pi D^2}{8\lambda L} = (2n+1)\frac{\pi}{2}$, only for small $n \in \mathbb{N}$,

$$D = \sqrt{4\lambda L(2n+1)} = \sqrt{2n+1} \times 1.88 \text{mm}$$

= 1.88mm, 3.26mm, 4.20mm, 4.97mm, ...

5.22 Using the result obtained in 5.18,

$$\begin{split} E &\approx -E_0 \left[\mathrm{e}^{\mathrm{i}k\sqrt{L^2 + r_2^2}} - \mathrm{e}^{\mathrm{i}kL} + \frac{1}{4} \mathrm{e}^{\mathrm{i}k\sqrt{L^2 + r_1^2}} - \frac{1}{4} \mathrm{e}^{\mathrm{i}k\sqrt{L^2 + r_2^2}} \right] \\ &\approx -E_0 \mathrm{e}^{\mathrm{i}kL} \left[\left(\mathrm{e}^{\mathrm{i}kr_2^2/(2L)} - 1 \right) + \frac{1}{4} \left(\mathrm{e}^{\mathrm{i}kr_1^2/(2L)} - \mathrm{e}^{\mathrm{i}kr_2^2/(2L)} \right) \right] \\ &= -2E_0 \left(1 + \frac{1}{4} \mathrm{e}^{\mathrm{i}\delta} \right) \mathrm{e}^{\mathrm{i}(\delta/2 + kL)} \sin \frac{\delta}{2}, \quad \delta = \frac{kr_2^2}{2L} = \pi, \quad \text{using } r_1^2 = 2r_1^2, \\ &= -\mathrm{i}\frac{3}{2} E_0 \mathrm{e}^{\mathrm{i}kL} = -\frac{3\mathrm{i}}{2} E_0, \end{split}$$

where E_0 is the amplitude on the aperature. Then the intensity is

$$I = \frac{I_0}{E_0^2} |E|^2 = \frac{9}{4} I_0. \tag{15}$$

5.23 (1)
$$\frac{2\pi}{\lambda} \left(\sqrt{L^2 + r_4^2} - L \right) = \pi + 3 \times 2\pi \quad \Rightarrow \quad r_4 \approx \sqrt{7\lambda L} = 5.916 \text{mm}.$$
 (2) $\frac{1}{10\text{m}} = \frac{(2n-1)\lambda}{r_n^2} = \frac{1}{2\text{m}} + \frac{1}{s} \quad \Rightarrow \quad s = -2.5 \text{m}$, that is, S is imaged 2.5m from the Fresnel zone plate, and on the same side of S .

2 Simulation

Applying the same method in **Problem 5.16**, we have

$$I_{\text{Diffraction}}(\theta) = \left(\frac{\sin \alpha}{\alpha}\right)^2,$$
 (16)

$$I_{\text{Interference}}(\theta) = \left(\frac{1}{N} \frac{\sin N\beta}{\sin \beta}\right)^2,\tag{17}$$

$$I(\theta) = I_0 \cdot I_{\text{Diffraction}}(\theta) \cdot I_{\text{Interference}}(\theta),$$
 (18)

where $\alpha = \frac{1}{2}ka\sin\theta$ and $\beta = \frac{1}{2}kd\sin\theta$. $I_{\text{Diffraction}}(\theta)$ and $I_{\text{Interference}}(\theta)$ are dimensionless and normalized so that their maxima are 1.

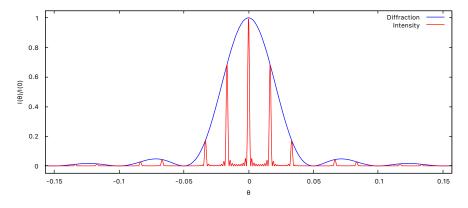
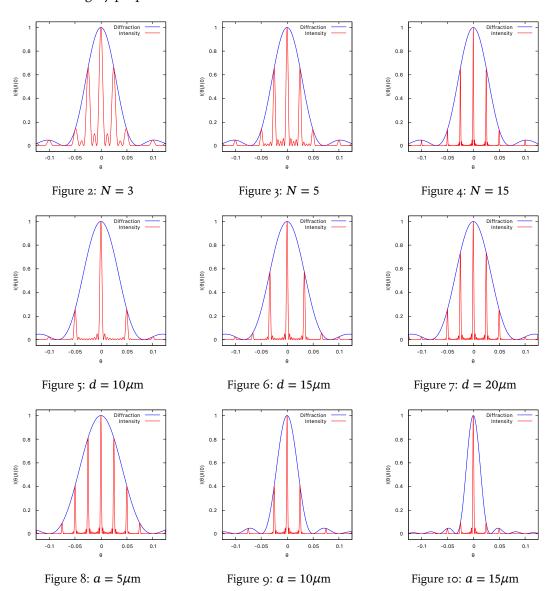


Figure 1: $\lambda = 500$ nm, N = 10, $d = 30 \mu$ m, $a = 10 \mu$ m

- (1) We can see in Figure 1 that the intensity $I(\theta)$ is enveloped by the diffraction factor $I_{\text{Diffraction}}(\theta)$. Thus the position of the principal maxima is determined by the interference factor $I_{\text{Interference}}(\theta)$, while their intensities are controlled by the diffraction factor.
- (2) Diffraction figures of different N are plotted in Figure 2-4 ($\lambda = 500$ nm, $d = 20\mu$ m, $a = 7\mu$ m). As N increases, the width of the principal maxima narrows.

- (3) Diffraction figures of different d are plotted in Figure 5-7 (λ = 500nm, N = 10, a = 7 μ m). As d increases, the distance between two nearest principal maxima narrows and roughly proportional to 1/d.
- (4) Diffraction figures of different a are plotted in Figure 8-10 ($\lambda = 500$ nm, N = 10, $d = 20\mu$ m). As d increases, the distance between two nearest interference minima narrows and roughly proportional to 1/a.



(5) A missing order means an overlap of the maximum of interference factor and the minimum of diffraction factor. The missing orders are $\pm \left(d \cdot \operatorname{lcm}\left(\frac{1}{a}, \frac{1}{d}\right)\right) \mathbb{N}^+$, where lcm denotes the least common multiple.

