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Chapter 1

Conceptual Overview

Newtonian physics serves as an excellent approximation at low speeds, accurately predicting the behavior of objects in our everyday world, but it gives way to different effects as we approach the speed of light. It turns out that an observer moving at high speed relative to another observer would notice lengths contract, time run slower, and events that were simultaneous to the other observer, no longer simultaneous to them. There is however one thing that remains constant, the speed that light travels relative to each observer is the same. To understand how all this works, we have to redefine our understanding of time and space, with special relativity.

Special relativity is vital for particle, nuclear, and astrophysics. It is needed for calculating the power output when it comes to nuclear fusion and fission or measuring the speed of galaxies and stars from the relativistic effects on the frequency of their emitted light. It is also needed for precision timekeeping for the likes of the fast-moving GPS satellites which need it for accurately calculating positions.

To understand how special relativity works we will need to first look at our classical understanding of the world, and using the fact that light somehow always moves at constant speed relative to everything, we will see how we must change our laws of physics for this to be true. Which in turn will lead to

many interesting consequences. We will start with an overview of all the main concepts, leaving the mathematical details and deeper insights for later chapters.

1.1 Classical Addition of Velocities

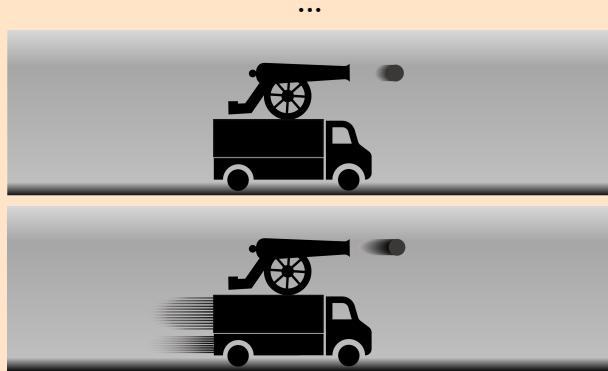


Figure 1.1: Diagram, showing the speed of cannonball from two different perspectives, (top) is from someone's perspective who is on the truck who sees the cannon at rest with the road moving backward, the cannonball is shot and is moving forward. (bottom) is from someone's perspective who is at rest on the road and sees the truck and cannon moving relative to them, when the cannonball is shot it looks like it is moving forward faster for them, as it seems to have the cannonball's velocity plus that of the truck it was shot from.

Let's imagine that a truck is moving forward on a road at a constant speed of 20 m/s (meters per second), on the top of that truck, is a cannon that fires a cannonball in the same direction the truck is moving. The ball travels at a speed of 300 m/s relative to the truck. Classically to find out how fast the ball is moving relative to the road, we add the speed of the truck and the speed of the ball together: $20 + 300 = 320$ m/s. So we find the ball is moving at a speed of 320 m/s relative to the road, and this is what we observe. So in classical physics, velocities directly add, the ball relative to the road would move at the speed of the ball relative to the truck plus the truck's speed relative to the road. But this turns out to just be a very good approximation for objects in our normal day life which we are used to seeing, which are moving much slower than the speed of light (which is roughly 300 000 000 m/s), but we will get to this later.

For now, we will explain more about the classical view, so that you will be able to see more clearly the differences when it comes to special relativity.

1.2 Inertial Reference Frame

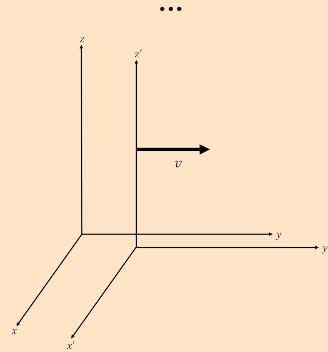


Figure 1.2: Diagram of a reference frame moving relative to another.

A reference frame can be thought of as an abstract coordinate system. The origin of its axis, its orientation, and its scale are specified by a set of points in space. The purpose of it is to provide a standardized means of measuring and describing the coordinates of objects within that frame at any instant of time. An inertial reference frame, is a reference frame that is not undergoing any acceleration. You can tell if you are being accelerated, as you will feel a force, for example, in an accelerating car you will feel the chair being accelerated into you, with your body slightly lagging behind in the acceleration, and it is this basic principle for how accelerometers measure acceleration as shown in the diagram.

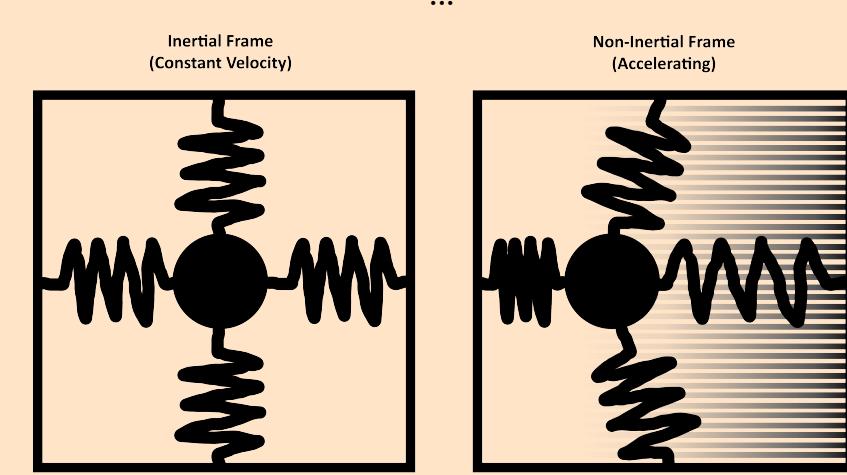


Figure 1.3: Diagram of a ball, attached to the walls of a box by springs, with the ball centered in the box in the inertial frame, i.e. with no acceleration (left), and in a non-inertial frame (right) where the box is now accelerating to the right, The ball lags behind as the box accelerates.

1.3 Classical Reference Frames

To see how to swap between reference frames in special relativity we will first have to introduce what the classical swapping between two reference frames looks like. To help us understand this, we will use a system of a moving platform, like we find on a treadmill. The two frames of reference will be that of the treadmill's platform and the room in which it is in.

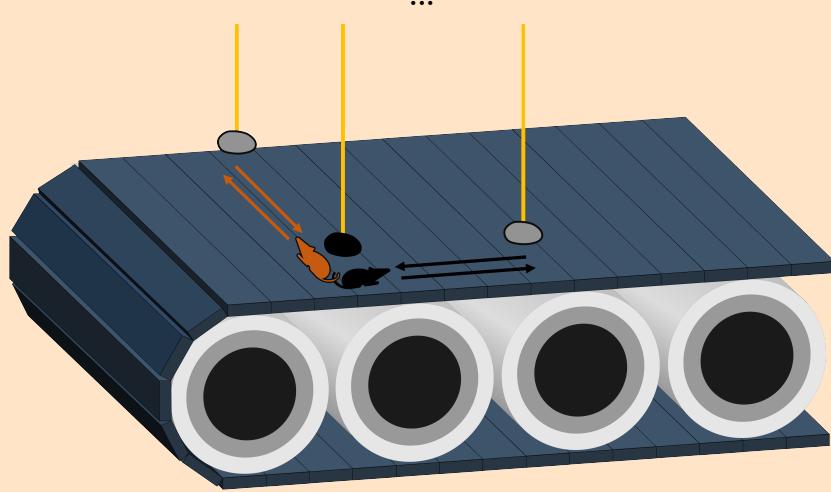
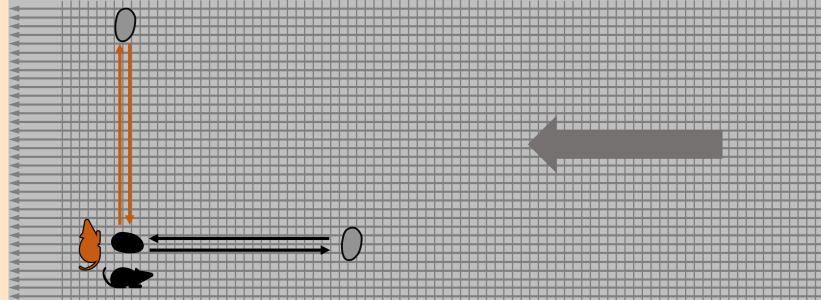


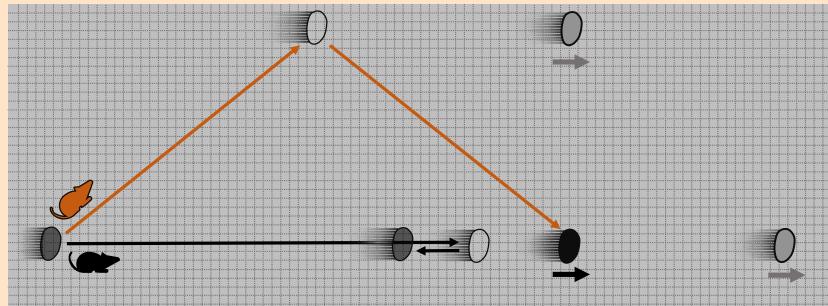
Figure 1.4: Diagram of 3d view of treadmill platform, with two rats and three hanging rocks.

To illustrate swapping between frames classically, We will look at a setup of two rats on a treadmill as shown in Figure 1.4, with three rocks hanging above the platform, at rest relative to the room. Both rats start under the same rock, one runs to a rock positioned in the forward direction of the treadmill, and the other runs to an equally distanced rock to the side, they then return to the starting rock. Here the platform can be seen as the medium in which the rats move. Both rats travel at the same constant speed relative to the platform, with the platform at a lower speed than the rats to allow them to get to the rocks.

If the platform is at rest, they will return to the starting rock at the same time. But if the treadmill is turned on and the platform is now moving, the rats will now have to also work against the movement of the platform to get to the rocks, this will lead to different distances the rats have to travel, and as a result, the rats get back to the starting rock at different times. The figures 1.5a and 1.5b show what the direction of movement of the rats and rocks look like in each reference frame.



(a) Room's reference frame.



(b) Platform's reference frame.

Figure 1.5: Rats running on a moving platform of a treadmill from two different reference frames

In the room's reference frame we have the platform moving backward, the rat moving straight ahead will have impedance to its movement but after it turns around it now has a boost from the platform to get back to the starting rock, whereas the sideways moving rat will have to go sideways and also balances out the backward pull of the platform to keep a perpendicular movement towards the hanging rocks. This gives different lengths of paths depending on the reference frame, as shown in figures 1.5a and 1.5b. It also gives different times for the rats to return to the starting rock. This is what we would expect and see with this classical example of frame swapping, but next, we will look at whether we can have the same sort of frame swapping when it comes to light traveling in two different frames of reference, and the notion of a universal rest frame referred to as the aether.

1.4 The Aether

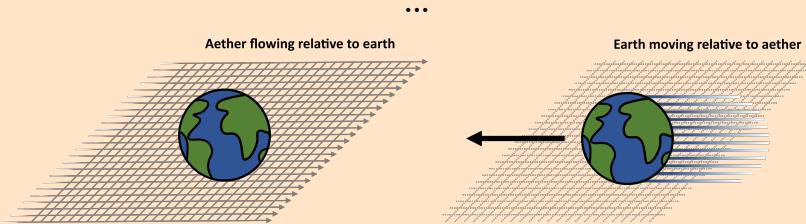


Figure 1.6: A diagram showing the aether's and earth's movement relative to each other.

In the 1800s, the theory was that light was a wave and therefore would need a medium for it to travel through that filled the vacuum of space, called "the luminiferous aether", and that light would travel at a constant speed relative to this aether, like how the rats in the previous section moved at a constant speed relative to the medium of the treadmill's platform.

An experimental setup by Michelson and Morley, shown in Fig. 1.7, was devised to measure earth's movement through the aether [1], by measuring how it affected the speed of light in different directions when observed in earth's reference frame.

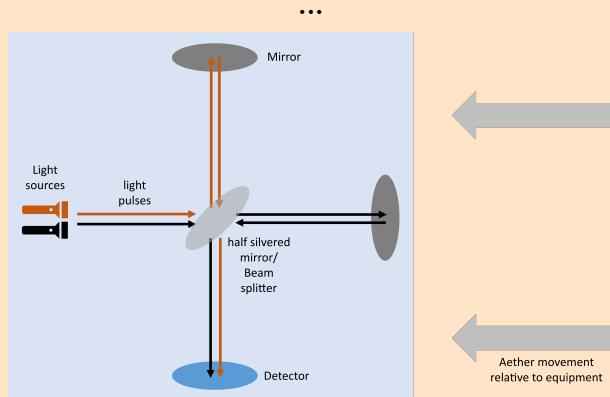


Figure 1.7: A diagram showing the Michelson-Morley experiment, (we can take the part of the paths between the the beam splitter and two mirrors to be analogous to the case of the paths in the previous rat and treadmill diagrams)

It did this by splitting a single light beam into two perpendicular paths, that are then reflected back to be recombined and sent towards a light detector. By rotating the whole interferometer setup, the two light paths could be aligned either parallel or perpendicular to the Earth's motion through the presumed aether. They reasoned that if the speed of light was constant with respect to the proposed aether, just like in the rat experiment from the previous section, the split light beams would recombine at different times. From the previous section's system of rats on a treadmill, we have the treadmill's platform as being analogous to the aether, and the rats as being analogous to the light, with both having the room that these experiments are done in as being the other inertial reference frames. However, when Michelson and Morley performed the experiment, they found no difference in travel time to the detector for both paths, indicating that there was no difference in the speed of light in any direction in earth's reference frame. hence no dependence of light's speed on the supposed aether. This null result seriously discredited the aether theories and ultimately led to the proposal by Einstein in 1905 that the speed of light (in a vacuum) is a universal constant and is independent of the motion of the observer or source. And to allow for us to have this universal speed of light (in a vacuum), it will require us to change our ideas of how time and positions are perceived by different observers.

1.5 Speed of Light

The experiments showed light does not have a medium that it travels at a constant velocity relative to, but instead travels at a constant velocity in a vacuum in all reference frames, independent of how fast the source of that light is moving in the frame, e.g. a moving truck's headlights. Light only moves slower in objects such as glass due to being impeded by the interactions between it and the material it is traveling through.

Light itself moves extremely fast compared to any other everyday speeds we are used to (roughly 300 000 000 m/s, it can travel the diameter of the world in the

blink of an eye). This is why we do not notice any delay in things we see around us in everyday life, though we should keep in mind that there is this delay, e.g. the light we see from the sun was emitted by it eight minutes ago for it to reach us now, and the further an object is located from us, the further back in time we are seeing it because of this delay.

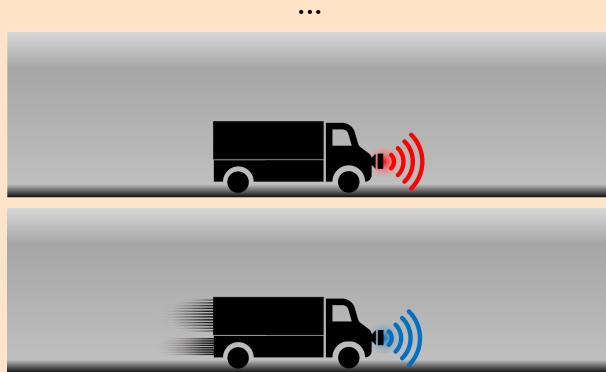


Figure 1.8: A diagram, showing light emitted from a truck in two different reference frames, with the emitted light having the same speed in each frame, though with different frequencies/energies due to the Doppler effect, which will be explained later.

When we look at the same truck setup as the previous sections, but in a vacuum this time, where the cannon firing a cannonball is swapped for the headlights emitting light, it gives the same speed of light when measured relative to the truck or the road, but how can this be true?

For this to be true, we need a new way of thinking about velocity addition, since the velocities of objects have to be added in a way that is consistent with the requirement that the speed of light is constant, but also gives the approximate classical addition at speeds of objects at much less than the speed of light, like how we observe with the cannon and truck. Since the speed of the light depends only on the units of time and positions, the only way to correct for this is to have the measured positions and times of objects transform differently when swapping between reference frames than the classical way. This is what we will talk about next. For the curious, it was the experiment by Ole Rømer that showed that light traveled at a finite speed rather than being instantaneously

emitted and received.

1.6 Positions and Time

1.6.1 Time Dilation

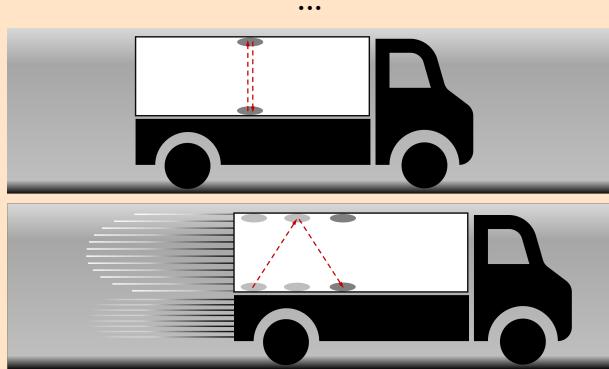


Figure 1.9: Diagram showing the extra distance light travels in the road's frame with the truck moving.

Let's imagine a simple clock, as shown in Figure 1.9, made of a light pulse moving back and forth between two mirrors on a moving truck, with a mirror on the truck's floor and on the roof directly above it, we will keep time by taking a tick of this clock to be when the light travels from one mirror to the other mirror and back again. For an observer in the truck, they will see the light go straight up and down between the mirrors, but an observer stationary relative to the road will see not just the light traveling up and down but also with the direction of the truck as well, hence it will travel a longer distance. But since the distance traveled by the light in the moving frame is longer and the speed of light is constant, we are only left with the possibility that the perceived travel time of the light has to be longer in this reference frame, i.e. the clock ticks slower to the observer on the road watching the moving truck. How much it is moving slower, can be solved using the ratio of the lengths of paths in each frame as they are shown in the diagram. This difference in the flow of time will later lead to more odd consequences. The slowing of the ticking is also the same

for any type of clock, and if you were to play a movie on the truck, it would take a longer amount of time for it to play through from the road's reference frame, it is the actual time itself being slowed down.

Also the closer the truck goes to the speed of light, the bigger the difference in how fast time flows relative to the truck's time (this can be seen as a longer horizontal stretch in the path taken for the light from the perspective of someone standing still on the road).

1.6.2 Simultaneity

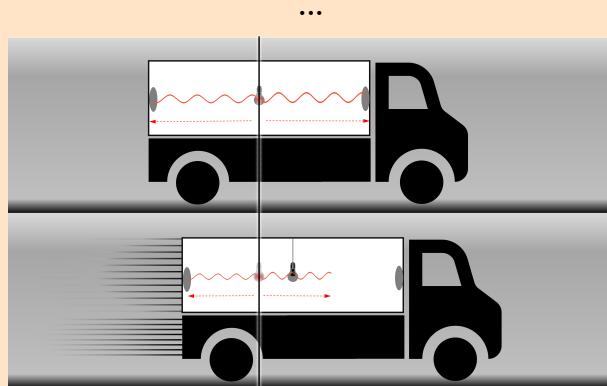


Figure 1.10: A diagram showing how the two events of light reaching the two walls of the truck are simultaneous in one frame but happen at different times in another, due to truck's movement in the second frame and the speed of light remaining the same.

Let's imagine a truck in its rest frame, with a light bulb in the middle and mirrors on the front and back walls, now if the light bulb gives off a pulse of light, the light will travel from the center of the truck to reach the mirrors simultaneously and back to the light bulb also simultaneously. But the observer on the road watching the truck drive past will see the light bulb simultaneously emit light in both directions and also simultaneously return to the bulb again. However, to do this the light reflects off each wall's mirrors at different times in this frame as shown in the diagram. This is because the speed of light is the same for both directions but the truck is moving, meaning the back of the truck

is moving towards where the bulb was when the pulse was emitted making the distances traveled shorter, and due to the front wall moving away from it, the light travels a longer distance to get to the front wall. Due to this and the speed of light being the same in both reference frames, the observer on the road will see the light hit the mirrors at different times. So, times of events (like when the light reaches either mirror) are different for observers in different reference frames, there is no one true order of events, e.g. an observer in a faster-moving truck moving in an opposite direction would see the light reach the front wall first. However for all observers the light will return to the central light bulb simultaneously. If two events happen in the same position at the same time (i.e. the light returning to the center of the light bulb simultaneously) then this happens simultaneously in all frames of reference.

I have left out any mention of length contraction, which will be introduced in the next section, and why it needs to be equal in both directions from the lightbulb, i.e. for the light to simultaneously return to the light bulb they need to maintain equal distances of the walls from the central light bulb.

1.6.3 Length Contraction

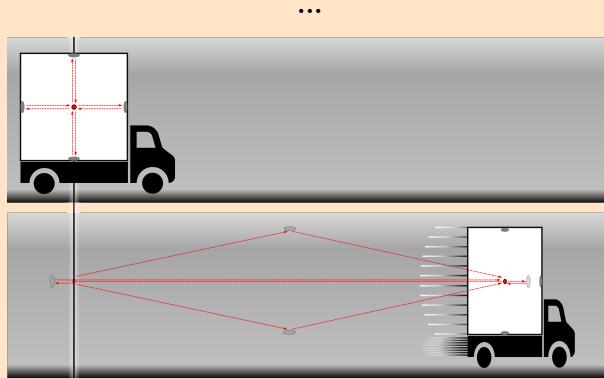


Figure 1.11: A diagram, showing a truck with a square container in its rest frame (top), emitting light from a central bulb in the four directions, with all light being reflected by the mirrored sides back to the center. In the second frame with the truck now moving (bottom) we see the truck's length contracted and what the light paths would be in this frame

If we have a truck in its rest frame with a square container with light emitted in all four directions from the center so that it will bounce off the mirrored sides of the truck and return to the center simultaneously. we require that in the moving frame they also all return to the center simultaneously, as multiple events that happen at a single point simultaneously in one frame, must happen simultaneously in all other frames. This time between the light being emitted and absorbed will be the dilated time, that was described in the time dilation section.

To achieve this simultaneity in the return of the light to the bulb, the length of the path of light in each direction has to be the same (as its speed is the same in all directions), we can work out the length of the upward path from the time dilation section, and this is the length the path needs to be in the horizontal directions as well, the paths can only have this length if the truck length is contracted when moving, the exact amount of contraction can be worked geometrically, and it turns out that the ratio of the increase in the amount of time that passes before the light is reabsorbed is inversely proportional to the amount the truck is contracted in the direction of its movement.

i.e. time passes by a factor slower in the moving frame, and the change in distances in the direction of movement are inversely proportional to the same factor.

The reason that we do not worry about contraction in the distance between the floor and ceiling mirrors, will be explained in a later section.

3 car system way of explaining length contraction (may delete)

Say we have 3 cars that are equally distanced and at rest on the road, the middle car sends out a pulse of light that reaches the front and back car at the same time to the road observer and the observers in the car when it reaches them both, all the cars accelerate for a predetermined fixed amount of time, so that after this acceleration period the cars are now moving but still equally spaced relative to the observers in the cars and the observer on the road. As for all

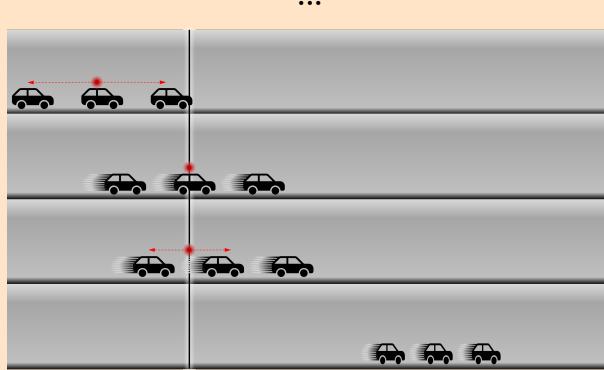


Figure 1.12: Diagram that illustrates an experiment with three cars initially at rest and equally spaced on a road. The middle car emits light that reaches the front and back cars simultaneously, triggering all cars to accelerate for a predetermined amount of time. Now another light is emitted from the middle car after the acceleration has finished, the road observer sees the light reach the back car first, as it is moving towards where the light had been emitted, this causes it to begin accelerating before the front car. This results in the cars being closer together after acceleration, demonstrating Lorentz length contraction. However, to the car observer, the distance between the cars remains unchanged from the initial distances.

observers the light reaches the front and back cars at the same time.

Now if we do this a second time, with the middle car releasing another pulse, we have in the frame of the observers in the cars: again the front and back drivers receiving the pulses at the same time meaning that they again will start accelerating simultaneously, and again after the acceleration there will be the same equal distance between the cars. But for the observer on the road, they see the back car receive the signal first as that car is moving towards the point where the light pulse was emitted and the front car is moving away from it. This would mean the back car would start accelerating first to get to the final constant velocity and get closer to the front car before the front car begins to accelerate, hence the cars end up closer together after all of this. The observers in the car and the road observer, do not agree with what the distances between the cars are. This contraction of length between cars, in the frame in which the cars are moving, is called Lorentz length contraction, which means that objects

that are moving faster become shorter, and the distances between the objects also become shorter.

1.6.4 No Width Contraction

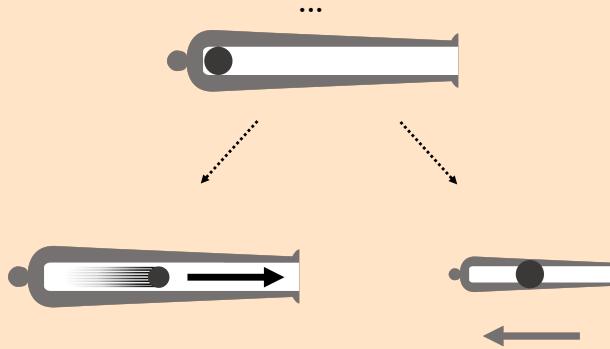


Figure 1.13: A diagram showing why there must not be a contraction in the perpendicular direction to the frame's relative motion. The top figure shows a ball and cannon at rest, the bottom figures show the cannon ball being fired in the frame of the cannon (left) where the moving ball has a contracted width, and the frame of the ball (right) where now the canon is now moving and has a contract width, with the ball at rest being the same size, both frames would contradict each other, if there was a change in width of moving objects, as the walls of the cannon and the surface of the ball would overlap in one frame and have a gap between them in the other. So we require that there is no change in size perpendicular to the movement of the object.

1.7 Doppler Effect

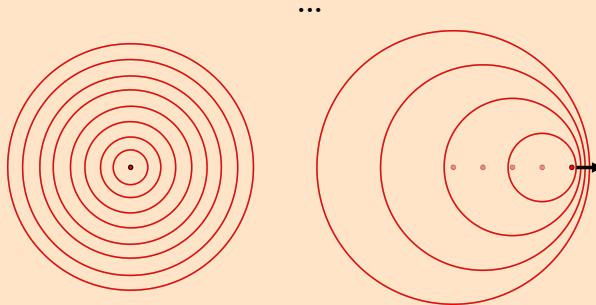


Figure 1.14: A diagram showing (left) a central source at rest emitting several circular pulses of light with equal time between each pulse, (right) the same source in a frame where it is now moving and emitting circular pulses of light, but each subsequent pulse is emitted from a different position as the source is moving, marked by a faded dot

If we have a source at rest, emitting circular pulses of light with equal times between each pulse, we will have concentric circular pulses in this frame, but if we move to a frame where the source is now moving to the right, each circular pulse is now being emitted from a different position as the source is moving. Due to the source moving, each pulse will be emitted closer to the right-hand side of the previous pulses, creating a bunching up of the pulses (increase in frequency) in the direction of movement and a spreading out (decrease in frequency) of the pulses in the opposite direction, that is what happens in the classical version of the Doppler effect, that you will notice that an ambulance or police car sounds different driving towards you and when driving away, this is due to the bunching up of the sounds waves in the direction of the moving vehicle. In special relativity, we also have to take the time dilation of the pulses into account, as there will be a longer time between each subsequent circular pulse, Due to perceived time moving more slowly for objects moving relative to the observer, this has a decreasing effect on the frequency in all directions, but directly in the direction of movement of the source this is outweighed by the frequency increase from the previous bunching up effect. Since the energy of the light is proportional to the frequency it is increased in the direction of motion

of the source due to the Doppler effects and decreased in the opposite direction. One thing not mentioned yet in this picture so far is how the light is also affected by what is called the aberration, which is the change in the direction of each part of the emitted light pulse, which leads to light not being circularly symmetric in each of the spherical pulses, which will be explained in the next section.

1.8 Aberration

The picture painted of the Doppler effect in the previous section has yet to show the effect on the angular distribution of the light in each of the spherical pulses. Here we will show that there is also a higher concentration of light in the direction of the source's movement, as shown in Figure 1.15.

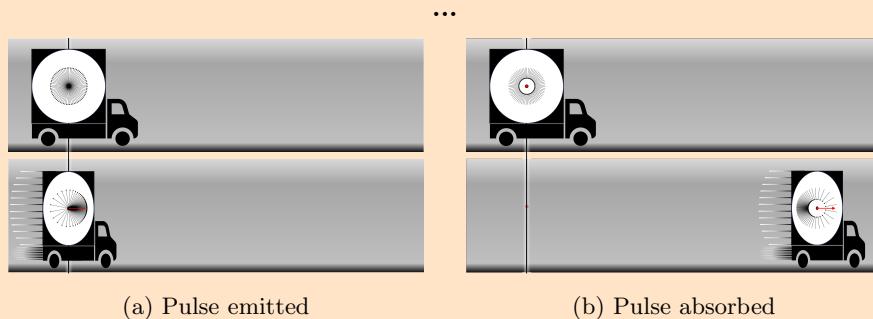


Figure 1.15: Two diagrams of a truck with a spherical container in two different frames, a light pulse is emitted from the center of the container (left) and reflected back by the mirrored wall to be absorbed in the center again (right), the truck in its rest frame (top of both figures) shows an evenly distributed light outward pulse, but in the frame with the moving truck (bottom of both figures) the emitted light is now more concentrated in the direction of the truck, which is required to allow the light to be reflected and returned to the center of the truck simultaneously.

With the help of the previous length contraction section, let's imagine a truck with a spherical container with mirrored walls, a central bulb emits light in all directions, in the rest frame all light reaches the spherical wall at the same time and returns to the center bulb simultaneously. Then in the moving frame, we have the spherical container length contracted and the light moving at the

same speed but reflecting off the walls at different times, returning to the center bulb simultaneously, for this to be true, the directions of the light have to be aberrated in the way shown in the diagram to allow for this simultaneous return to the bulb.

From the diagrams, you can see how light is aberrated when emitted from a moving source relative to its rest frame and also when being absorbed by a source, the faster the source is moving, the more the direction of each part of the light pulse is aberrated.

If the speed of the source was to approach the speed of light, all emitted light would approach the direction of movement of the source, if the source was theoretically able to reach the speed of light, all light would be emitted in the direction of the source but also move at the same speed of the source, i.e. the source would move with the emitted light and it would not leave the vicinity of the source. Though the rate at which it emits it would also tend to zero. If a photon theoretically had mass and its influence of the gravitational force moved at the speed of light, then we would not be able to feel any gravitational effects outside the vicinity of it, due to all of its gravitational field being propagated in the direction of its movement and at the same speed as the photon.

Also in astrophysics, it has to be taken into account that the earth's view of the universe is distorted by aberration, as shown in diagram ref... *** diagram of earth's view of the universe is distorted from the aberration (example on wiki)

1.9 Relativistic Beaming

If we now take both the Doppler picture and the aberration picture from the previous sections together, so that if we have a source that in its rest frame is emitting a spherical pulse of light with evenly distributed angles. Then in a frame with the source now moving, we will have this spherical pulse's wavelength bunched up in the direction of movement (giving a change in the color of the light), with the angular spread of the light at a higher concentration in the

direction of the movement of the source.

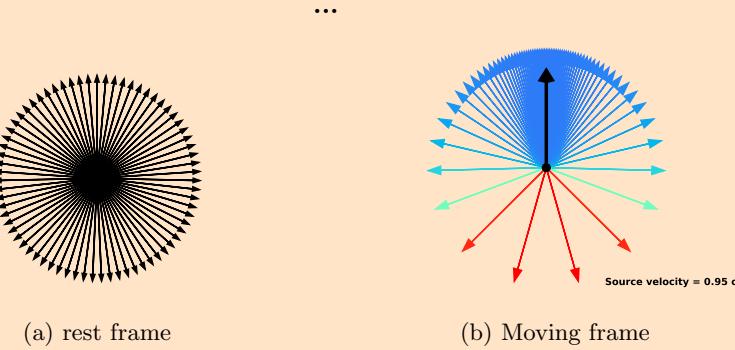


Figure 1.16: Outward pulse of light from source in a rest and moving frame

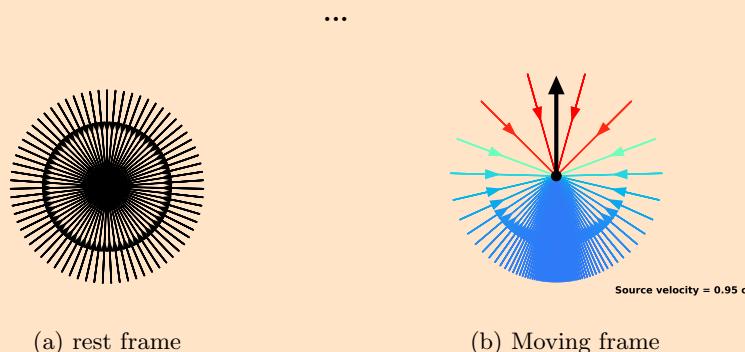


Figure 1.17: Inward pulse of light from source in a rest and moving frame

1.9.1 Retarded Field

If we were to have multiple pulses with a constant time between each pulse, and remembering that each pulses origin was from a point in the source's past position (the retarded position of the source), we can work out what the full view of all the pulses in the rest frame is in the moving frame, shown in Figure ???. From this, we can work out the distribution and hence concentration of the light for any moving source relative to that of it at rest.

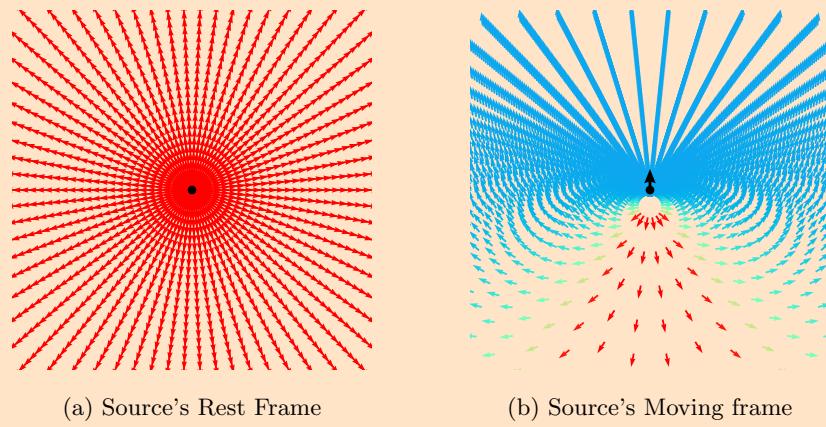


Figure 1.18: A 2D diagram showing multiple spherically symmetrical pulses of light from a source at rest (left), and in a corresponding primed inertial frame where the source is moving (right), the colors show the Doppler effect on lights light *** describe it as spherical waves with them being closer together in the direction of movement, could have separate diagram showing the concentric circles with different origins

1.9.2 Perceived vs Actual Speed

Imagine that a ball one light year (the distance it takes light to travel in a year) away, is fired directly towards you. It will take one year for the first light from the oncoming ball to reach you. If the ball moves at three-quarters of the speed of light, the ball will hit you in four-thirds of a year (a year and four months) after it is fired. The last light from the ball will reach you just as the ball hits you. As you see it, the time between the first and last light from the ball is four months. During those four months, you will see the ball start at its initial position and travel a distance of one light-year. So the ball appears to you to have been moving three times faster than light. This is just how it appears to the observer due to the delay in the light signal/the retarded view of the system, leading to latency in how the system is observed.

This shows how important it is to take the delay in the light from objects into account when observing relativistic systems, this view is called the delayed/re-tarded view, i.e what we see now are objects in their past positions, and the

further things are from you the further into the past we are currently seeing them.

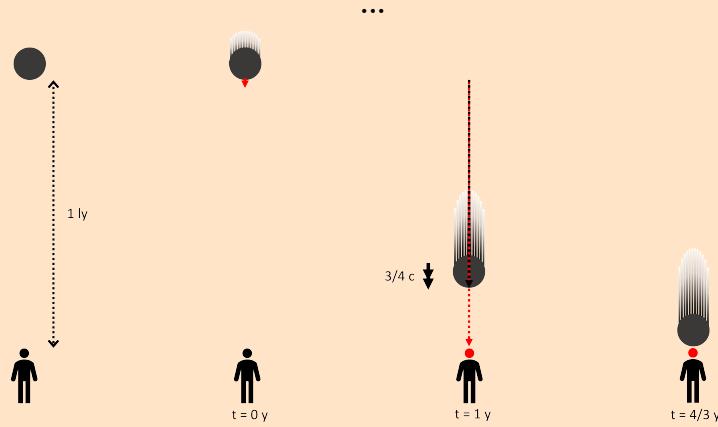


Figure 1.19: A diagram demonstrating the perceived speed of a ball vs its actual speed

1.10 Velocity Transformation

Previously, we explored how velocities at the speed of light change between different frames of reference, that is their directions are aberrated. However, for objects moving at speeds less than light, we need the transforming speeds to be different as well as aberrated due to the transform needing to tend towards the classical transform as the speeds tend towards zero. To understand these velocity transformations, we will need to first delve into the mathematical details which will be covered in later chapters.

1.11 The More Abstract Properties...

The transformation of more abstract quantities that depend on position and time, such as momentum, energy, and mass do not have a simple visual interpretation, and this is the point at which the mathematics is needed for a proper understanding of how they transform. With the maths, you will find

new conservation of momentum and energy laws.

1.12 Summary

A key early idea was that light traveled through a medium called the luminiferous aether. The Michelson-Morley experiment attempted to measure Earth's motion relative to this aether. Surprisingly, no difference was found in the speed of light regardless of direction. This conflicted with the intuitive addition of velocities, discredited the aether theory, and led Einstein to propose light's constant speed in all frames as a fundamental principle.

Special relativity emerged from the insight that light's speed in a vacuum is constant for all observers, regardless of the light source's speed relative to each observer. This required rethinking the concepts of time intervals and distances between points to accommodate light's fixed speed and led to the requirement that a clock moving relative to an observer ticks slower from the observer's perspective while also being contracted in the direction of its motion. This was shown in figure 1.9 and figure 1.11.

It also led to the simultaneity of events not being absolute, e.g. light emitted from the middle of a moving truck reaches the front and back walls at the same time as someone at rest in the truck. But someone standing still on the road sees the light hit both walls at different times, as shown in figure 1.10. There is no universal "now" at a distance - observers relate events differently.

Classical physics provides an excellent approximation of reality at low speeds. Relativistic effects only become readily noticeable at speeds approaching that of light. We don't normally experience them because motions in everyday life are very slow relative to that of light which is roughly three hundred million metres per second. The faster an object moves relative to an observer, the more its lengths contract and its time dilates, to the observer.

When it comes to velocities, they combine differently than in classical physics. However, we must ensure that the transforms reduce to the intuitive addition of velocities of classical physics at low speeds while accommodating light's constant

rate in all frames of reference. This leads to initially bizarre outcomes like **apparent** faster-than-light motion emerging from how relativistic optical effects play out across distances, due to the delay in light signals from objects at a distance. Nevertheless, special relativity gives a truer, more accurate picture of reality.

We also have seen that light's frequencies shift when emitted from a moving source due to space-time effects on wave propagation (the Doppler effect) and that light concentrates in the direction of the source's motion, the aberrational effect. Both together intensely focus light and radiation in the direction of rapidly moving objects, relativistic beaming.

These effects will all be explained more fully with math in the coming chapters, along with other effects that require the use of math to fully understand, as well as new conservation laws of energy, mass, and momentum. We will derive everything step by step starting with the transformation of positions and times between reference frames.

*** Explain what a photon sees, i.e. when it is emitted it sees a universe that is infinitely thin in direction of movement and the time between emission and absorption is then zero

1.13 Definitions

[Mathematical Framework]

Reference frame: An abstract coordinate system that is taken to be at rest consisting of three spatial and one time axis. The origin of its axis, its orientation, and its scale are specified by a set of points in space. The purpose of it is to provide a standardized means of defining the position and orientation of objects at any instant in time.

Inertial reference frame: A reference frame that is not undergoing any acceleration. An observer's reference frame is inertial if they do not feel a force acting on them.

Proper/Rest frame: The Proper/Rest frame of an object, is the reference frame of the object itself.

Primed frame: A reference frame that is moving relative to the current frame.

frame velocity: The velocity of a reference frame relative to the current reference frame.

Frame transform: The changing of coordinates and other quantities in a reference frame into their corresponding values in another reference frame.

Galilean/Classical frame transform: Transformation of coordinates and physical quantities according to Newtonian physics between two frames of reference.

Relativistic frame transform: The changing of the spatial coordinates, the time coordinate, and other physical quantities into their corresponding values in a different inertial reference frame using equations from special relativity.

Observer: An entity/thing that has an associated frame of reference/system of coordinates in which to measure the position, orientation, and other properties of objects with respect to.

Event: An occurrence that can be specified by a unique combination of the three spatial and one-time coordinates. These four coordinates, together, form a point in spacetime.

Simultaneity: Within a reference frame, two or more events are simultaneous if they occur at the same time within it. Classically if any two events are simultaneous in one reference frame, they are also simultaneous in any other reference frame. In special relativity, this is not true, and the order of events can depend on the frame of reference.

Retarded position: The previous position of a source of a field that an observer sees it at, due to the time delay between when a source emits a field and when that field is detected by the observer

[The Mechanics]

Aether: A proposed medium that filled the vacuum of space that light propagated through.

Vacuum: Empty space with no matter and negligible amount of effects from fields.

Lorentz length contraction: The shortening of an object along the direction of its motion relative to an observer (relative to the object at rest).

Time dilation: The slowing of the passing of time for objects in motion relative to an observer, compared to the passing of time for the observer.

Doppler effect: The change of a wave's frequency due to the relative motion of its source to an observer. This is due to the bunching up of the wave in the direction of movement of the source and the stretching of the wave in the opposite direction, as well as the dilation of time between two parts of the emitted wave.

Aberration of light: The changing of the direction of the propagation of light when moving into another reference frame.

Relativistic beaming: The accumulation of the Doppler effect and aberration of light, and how it affects the amount of light and its frequency emitted at different angles from a moving source. This effect can be extended to any wave or field being emitted at the speed of light from a moving source.

Retarded field: A moving source's field, where at each point has been propagated at the speed of light from the retarded position of the source...

Field:

[Defs for later]

Invariant: ...

Chapter 2

Positions, Time and Velocity

Now that we have learned about the basic concepts of special relativity, we can let the mathematics begin, starting with classical relativity as a starting point, it will give us an understanding of how to transform between reference frames, we can then compare special relativity to this to help with our intuition.

2.1 Classical Relativity

For classical (Galilean) relativity, if we have two observers moving relative to each other at constant velocity and want to find how each observer would describe the positions and motion of a particle relative to them, we need to know the positions and times in one observer's frame and how to transform them into the other observers frame. i.e. if we have how one observer would describe the coordinates of a particle over time relative to them and we want to now know the coordinates of the particle with respect to the other observer, we need to know how the transformation of coordinates works.

Before special relativity, it was assumed that distances between two positions

were the same for all observers, i.e. the first observer would see a stick and measure it to be the same length as another observer moving relative to them would measure it to be. long story short, lengths are the same for all observers in Galilean transform. Another assumption is that time is constant for everyone, i.e. clocks move at the same rate for each observer and if a particle was to emit light at a particular time. This light would be emitted at the same time for the other observers too.

So let's see how the mathematics works for the classical coordinate system transform between two observers. If we have the coordinate system in one inertial reference frame $\langle S \rangle$ to describe an event E , which is given by the set of spatial and time coordinates (x, y, z) and t , and we want to find the coordinates of the event in another "primed" inertial reference frame/coordinate system that is moving at a relative speed v to the first frame in the z -axis, that was initially overlapping with the first frame at time $t = 0$, with the events coordinates now described by x', y', z' at time t' . The direction of the second frame is in the upwards z -direction instead of being to the right in the case of the truck diagrams from the previous chapter, this is to help us when we use spherical polar coordinates later and to help see the symmetries.

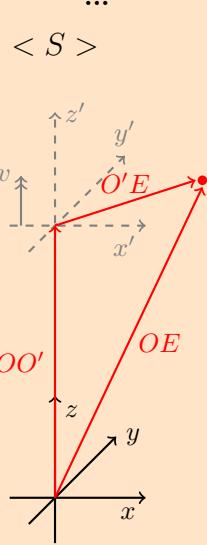


Figure 2.1: Diagram, of a reference frame $\langle S \rangle$ with its associated (x, y, z) coordinate system at rest in this frame, it also shows a primed coordinate axis (x', y', z') at time $t = t'$ from when the axis where overlapping, associated with another inertial reference frames moving at velocity v in the z -axis, it shows an event with position vector \overrightarrow{OE} in this frame and the vector $\overrightarrow{OO'}$ to this primed axis origin and the vector $\overrightarrow{O'E}$ which is the vector from this primed origin to the event in the shown frame.

For frame $\langle S \rangle$ we have its origin O and an event E located at

$$\overrightarrow{OE} = (x, y, z) \quad (2.1)$$

at time t from axis overlap, and a prime frame $\langle S' \rangle$'s origin O' moving at relative velocity v so that in the original frame we have after a time

$$\overrightarrow{OO'} = (0, 0, vt) \quad (2.2)$$

in $\langle S' \rangle$ we have the event's coordinates as

$$\overrightarrow{O'E} = (x', y', z') = \overrightarrow{OE} - \overrightarrow{OO'} = (x, y, z - vt) \quad (2.3)$$

for frame $\langle S \rangle$ we have

$$\overrightarrow{OE} = (x, y, z) = \overrightarrow{OO'} + \overrightarrow{O'E} = (x', y', z' + vt') \quad (2.4)$$

so we have $x = x'$, $y = y'$, $t = t'$ (time is assumed to be the same across reference frames in classical relativity), and for the z-components

$$\begin{aligned} z' &= z - vt \\ z &= z' + vt' \end{aligned} \quad (2.5)$$

This shows the symmetry between the frames, as if we were originally in the $\langle S' \rangle$ frame we would have the $\langle S \rangle$ frame moving at $-v$ relative to it, so we would expect that the inverse transformation would replace one frame's coordinates with the other and have the opposite frame velocity instead. So the inverse of going from primed frame to original is just a matter of swapping all primed coordinates with original coordinates and changing the sign on the frame velocity, this can be seen from starting at time $t = 0$ we have the frame's origins overlapping and that there is a symmetry from here, we have the freedom to choose either frame to be the original, as no frame is particularly special and depending which one you pick the relative velocity is either v or $-v$.

The full description of the position and time transform for an event is given as

$$\begin{aligned} x' &= x \\ y' &= y \\ z' &= z - vt \end{aligned} \quad (2.6)$$

at time

$$t' = t$$

2.2 Postulates of Special Relativity

1. The laws of physics are the same in all inertial frames of reference (i.e. if an observer carries out an experiment in an inertial reference frame, there would be no difference if the experiment was instead carried out within a different inertial reference frame)
2. The speed of light in a vacuum is the same in all inertial/non-accelerating frames of reference.

They are not the only assumptions, you may notice some small jumps in the logic when it comes to the derivation. But be aware small miss understandings have thrown people down the special relativity rabbit hole when trying to find inconsistencies or prove it wrong, throughout I will try to leave as few logical gaps as possible, and explain as simply and as in-depth as possible so that you do not do this yourself.

2.3 Derivations

For the derivation of the coordinate transformation between two inertial frames, we will follow a similar line of thought as we did in the first chapter on the conceptual introduction. For those who are curious, there is an alternative derivation using spherical light pulses given in Appendix ??.

2.3.1 Time Dilation

We can have a clock that keeps time using a light pulse being emitted from a bulb towards a mirror and back again, recording a tick on its clock each time it returns to where it was emitted, taking a time $\Delta\tau$ in the rest frame to return.

The distance the light travels to the mirror and back in this time is

$$2W = c\Delta\tau, \quad (2.7)$$

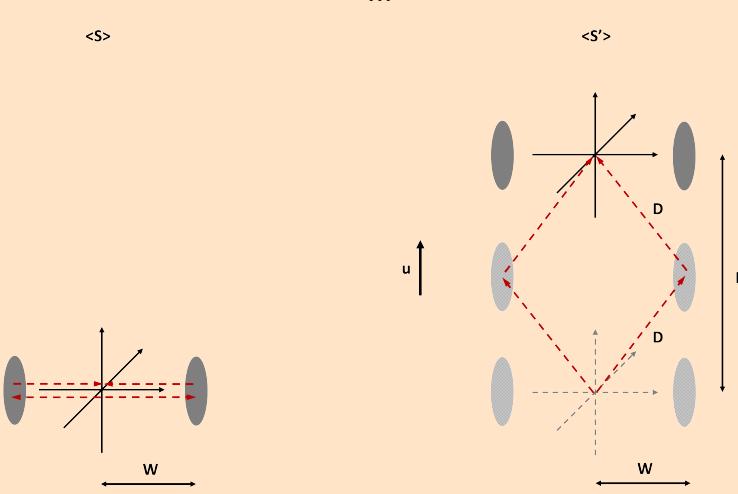


Figure 2.2: A diagram showing a clock keeping time using light, light is emitted from the origin and reflected off a mirror back to the origin, shown here in two different frames (describe what things are in the diagram)(need axis labels)

the distance the bulb travels during this in the moving frame is

$$L = u\Delta t' = -v\Delta t', \quad (2.8)$$

where $\Delta t'$ is the time it took light to return in moving frame and u is the speed of the bulb which is in the opposite direction of the primed frame relative to the proper frame, i.e. $u = -v$, the distance light travels in this time is

$$2D = c\Delta t', \quad (2.9)$$

this forms a triangle, so we have the square of the hypotenuse $2D$ equal to the square of the lengths W and L :

$$c^2\Delta t'^2 = v^2\Delta t'^2 + c^2\Delta\tau^2 \quad (2.10)$$

rearranging to find the time it took in the primed frame relative to the time in the proper frame, we find

$$\Delta t' = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \Delta \tau = \gamma \Delta \tau \quad (2.11)$$

with

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}. \quad (2.12)$$

which means

$$\Delta t' \geq \Delta \tau \quad (2.13)$$

So we take this outcome as the time taken between two events in the same position in the moving frame taking longer than in the rest frame of the clock, due to the longer distance it has to travel at the same speed.

2.3.2 Length Contraction TODO

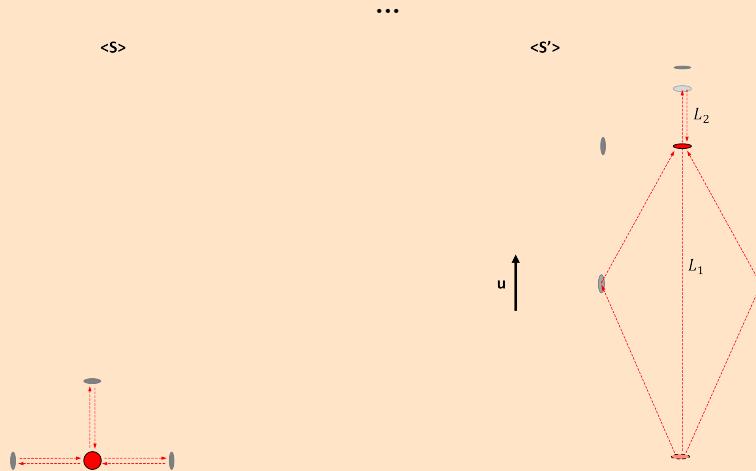


Figure 2.3: A diagram, showing a central bulb emitting light in the three directions, with all light being reflected by the mirrors back to the center, this is shown at rest on the left and moving in another inertial frame with the system now moving, with its length slightly contracted in the direction of movement.

If we have a system in its rest frame with light emitted in three directions from the center, one directly in front and the other two to the sides, with equally distanced mirrors in the three directions so that they will be reflected back and return to the center simultaneously. we require that in the moving frame they also all return to the center simultaneously, as multiple events that happen at a single point simultaneously in one frame, must happen simultaneously in all other frames. This time between the light being emitted and absorbed will be the dilated time from the previous section. To achieve this simultaneity in the return of the light to the bulb, the length of the full path of light from there center to the mirrors and back in each direction has to be the same (as its speed of light is the same in all directions and the return time has to be the same), we can work out the length of the sideways paths from the time dilation section, and this is the length the path needs to be in the vertical directions as well, the paths can only have this length if the system length is contracted when moving in the direction of movement. Note that there is no contraction in the perpendicular direction to the movement due to the paradox explained in chapter 1.6.4.

The pulse sent forward has a distance

$$L_1 = ct'_1 = D' + ut'_1 \quad (2.14)$$

to travel to get to the mirror that is moving away from the point it was emitted, where t'_1 is the time taken to get there and D' is the distance between the mirror in front and the bulb and u being the speed the bulb and mirrors are moving in the primed frame. After this, the light then has to travel back to the bulb a distance

$$L_2 = ct'_2 = D' - ut'_2 \quad (2.15)$$

where t'_2 is the time taken to return to the bulb from the mirror, both of these give the two times as

$$t'_1 = \frac{D'}{c-u} \quad (2.16)$$

and

$$t'_2 = \frac{D'}{c+u} \quad (2.17)$$

adding both times together gives the total primed time

$$t' = D' \left(\frac{1}{c+u} + \frac{1}{c-u} \right) = \frac{D'}{c} \frac{1}{1-u^2/c^2} = \frac{D'}{c} \frac{1}{1-v^2/c^2} = \gamma^2 \frac{D'}{c} \quad (2.18)$$

from the time dilation equation, we have

$$ct' = c\gamma t = \gamma D = \gamma^2 D' \quad (2.19)$$

and finally leading to the contracted distance to the front mirror in the primed frame being

$$D' = \frac{1}{\gamma} D \quad (2.20)$$

so we have

$$D' \leq D \quad (2.21)$$

we take from this that distances in the direction of movement of an object are contracted in its direction of movement compared to the rest frame of the object.

2.3.3 Position and Time Transformation

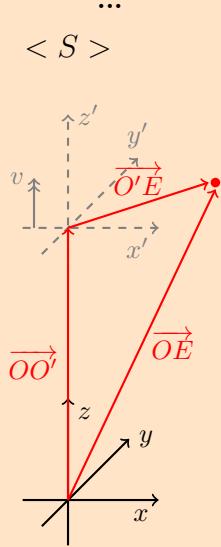


Figure 2.4: A diagram showing an event in a frame with another primed frame's axis shown moving in the z -direction.

If we are in a frame $\langle S \rangle$ with origin O at position $(0, 0, 0)$ and a primed frame $\langle S' \rangle$ moving in the z -direction at velocity v such that the position of the primed origin in this frame is

$$\overrightarrow{OO'} = (0, 0, vt) \quad (2.22)$$

at time t from when the origins overlapped, if there is an event (let's say light or another object reaching a mirror that is moving with the primed frames axis) described in this frames coordinates as

$$E = (x, y, z, t) \quad (2.23)$$

with the spatial coordinates shown in the diagram as

$$\overrightarrow{OE} = (x, y, z) \quad (2.24)$$

in the primed frame we have the event described as

$$E' = (x', y', z', t') \quad (2.25)$$

now in the current frame, we have the primed frame's axis moving with the mirror, from the previous chapter we should have this primed distance in the z-direction to the event contracted, as the primed z-coordinates are contracted in this frame, as seen from length contraction section; the events z-coordinate is $z = z'/\gamma$, so that distance between the primes origin and the event in this frame (not the primed frame) is

$$\overrightarrow{O'E} = (x, y, z'/\gamma) \quad (2.26)$$

noting that we have the x and y-components unaffected due to the cannon and ball thought experiment from the first chapter. From the diagram, we can see

$$\overrightarrow{OE} = \overrightarrow{OO'} + \overrightarrow{O'E} \quad (2.27)$$

subbing each component into this equation for the z-direction and rearranging, we have

$$z' = \gamma(z - vt) \quad (2.28)$$

If we were instead in frame $\langle S' \rangle$ we would have $\overrightarrow{OE} = (x, y, z/\gamma)$, $\overrightarrow{OO'} = (0, 0, vt')$, and $\overrightarrow{O'E} = (x, y, z')$, doing the same for these we have

$$z = \gamma(z' + vt') \quad (2.29)$$

This should be expected as this shows the expected symmetry between the frames because if we were originally in the $\langle S' \rangle$ frame we would have the $\langle S \rangle$ frame moving instead at $-v$ in the z-direction relative to the $\langle S' \rangle$ frame, so we would expect that the inverse transformation would replace one frame's coordinates with the other and have the negative of the velocity instead.

Now substituting the z transforms into each other we have

$$z = \gamma(\gamma(z - vt) + vt') \quad (2.30)$$

and rearranging gives

$$t' = \gamma \left(\left(\frac{1}{\gamma^2} - 1 \right) \frac{z}{v} + t \right) \quad (2.31)$$

with $\gamma^{-2} = 1 - v^2/c^2$ then giving

$$t' = \gamma \left(t - \frac{vz}{c^2} \right) \quad (2.32)$$

leaving full description of the location and time of an event as

$$\begin{aligned} x' &= x \\ y' &= y \\ z' &= \gamma(z - vt) \\ \text{at time} \\ t' &= \gamma \left(t - \frac{vz}{c^2} \right) \\ \text{with} \\ \gamma &= \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \end{aligned} \quad (2.33)$$

For the time transform the $\frac{vz}{c^2}$ term is the cause of simultaneous events in one inertial frame no longer being simultaneous in another inertial frame, and the gamma factor γ , is the term that causes the overall change in how fast time flows relative in each frame. In the spatial coordinate transforms the gamma factor is the cause of the change of lengths in the direction of the frame's movement. How the transformation works is shown in figure 2.5.

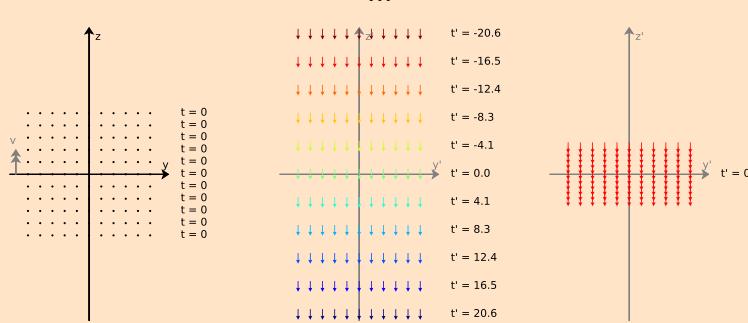


Figure 2.5: A diagram, showing a grid of points at time $t = 0$ that are at rest in their proper frame (left), giving the direct transformation into the moving primed frame (middle) where the points are now moving in the negative z -direction with the colors and labels showing the desynchronization of the primed times t' , the next figure (right) shows what these points would be when they are synchronized with the origin's time $t' = 0$. *** split plots into several svgs and for other set later too, also call first picture "initial" frame of reference and others primed

2.3.4 Infintesimal Changes of Coordinates

$$\begin{aligned}
 x'_2 - x'_1 &= x_2 - x_1 \\
 y'_2 - y'_1 &= y_2 - y_1 \\
 z'_2 - z'_1 &= \gamma(z_2 - vt_2) - \gamma(z_1 - vt_1) \\
 t'_2 - t'_1 &= \gamma\left(t_2 - \frac{v}{c^2}z_2\right) - \gamma\left(t_1 - \frac{v}{c^2}z_1\right)
 \end{aligned} \tag{2.34}$$

leads to

$$\begin{aligned}
 \Delta x' &= \Delta x \\
 \Delta y' &= \Delta y \\
 \Delta z' &= \gamma(\Delta z - v\Delta t) \\
 \Delta t' &= \gamma\left(\Delta t - \frac{v}{c^2}\Delta z\right)
 \end{aligned} \tag{2.35}$$

when taken the changes are infinitesimal this gives the differential quantities

$$\begin{aligned}
dx' &= dx \\
dy' &= dy \\
dz' &= \gamma(dz - vdt) \\
dt' &= \gamma \left(dt - \frac{v}{c^2} dz \right)
\end{aligned} \tag{2.36}$$

2.3.5 Transforms with Low Relative Frame Speed Change

For the coordinate transforms in the previous section we get back the classical coordinates transform for frame velocities much smaller than the speed of light, i.e. with $v \ll c$ we have $\gamma \approx 1$ we get

$$t' \approx t \tag{2.37}$$

$$z' \approx z - vt \tag{2.38}$$

which is the expected result for the classical coordinates transform, and what we are used to seeing things work in everyday life. The influence of the gamma factor is shown in figure 2.6.

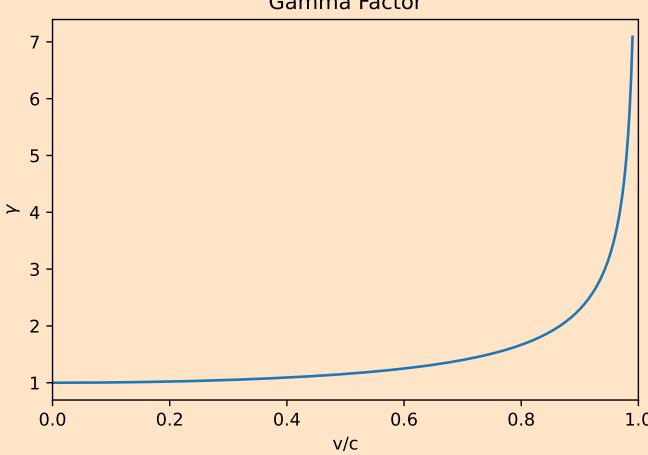


Figure 2.6: A diagram, showing the magnitude of the gamma factor with increasing ratio of frame speed to the speed of light, with γ tending to infinity as $v/c \rightarrow 1$.

2.4 Velocity

from taking the infinitesimals of equations from eq. (2.33), we have the velocities as

$$\begin{aligned}
 u'_x &= \frac{dx'}{dt'} = \frac{dx}{\gamma \left(dt - \frac{vdz}{c^2} \right)} \\
 u'_y &= \frac{dy'}{dt'} = \frac{dy}{\gamma \left(dt - \frac{vdz}{c^2} \right)} \\
 u'_z &= \frac{dz'}{dt'} = \frac{\gamma(dz - vdt)}{\gamma \left(dt - \frac{vdz}{c^2} \right)}
 \end{aligned} \tag{2.39}$$

now dividing the top and bottom of the fraction by dt we can give the equations in their vector form as

$$\mathbf{U}' = \frac{1}{\gamma \left(1 - \frac{v}{c^2} u_z\right)} \begin{pmatrix} u_x \\ u_y \\ \gamma(u_z - v) \end{pmatrix} \quad (2.40)$$

at position \mathbf{R}' and time t'

the generalised velocity transform derived in appendix (todo) is

$$\mathbf{U}' = \frac{1}{\gamma} \frac{\mathbf{u}_s + \left[\frac{\gamma-1}{\|\mathbf{V}\|^2} (\mathbf{u}_s \cdot \mathbf{V}) - \gamma \right] \mathbf{V}}{1 - \frac{\mathbf{u}_s \cdot \mathbf{V}}{c^2}} \quad (2.41)$$

with \mathbf{V} being the frame's velocity in any generalized direction, instead of just being in the z-direction

2.5 Acceleration

To find how the acceleration of a particle transforms for an observer in two different frames we take the differential of the velocity transform eq. (2.40) and use the differentiation rules for two generic functions f and g : $d(gf) = f dg + g df$, and $d[f(g(x))] = dg(x) * df(g(x))$

$$\begin{aligned} d\mathbf{U}' &= \frac{1}{\gamma \left(1 - \frac{v}{c^2} u_z\right)} \begin{pmatrix} du_x \\ du_y \\ \gamma du_z \end{pmatrix} + \frac{\gamma \frac{v}{c^2} du_z}{\gamma^2 \left(1 - \frac{v}{c^2} u_z\right)^2} \begin{pmatrix} u_x \\ u_y \\ \gamma(u_z - v) \end{pmatrix} \\ &= \dots \\ &= \frac{1}{\gamma \left(1 - \frac{v}{c^2} u_z\right)^2} \begin{pmatrix} du_x + \frac{v}{c^2} (u_x du_z - u_z du_x) \\ du_y + \frac{v}{c^2} (u_y du_z - u_z du_y) \\ \frac{1}{\gamma} du_z \end{pmatrix} \end{aligned} \quad (2.42)$$

Now, divide by the time differential from the previous chapter to get the acceleration transform

$$\mathbf{a}' = \frac{1}{\gamma^2 \left(1 - \frac{v}{c^2} u_z\right)^3} \begin{pmatrix} a_x + \frac{v}{c^2} (u_x a_z - u_z a_x) \\ a_y + \frac{v}{c^2} (u_y a_z - u_z a_y) \\ \frac{1}{\gamma} a_z \end{pmatrix} \quad (2.43)$$

at position \mathbf{R}' and time t'

2.5.1 Generalised Acceleration Transform

3 acceleration (check this)

$$\mathbf{a}' = \frac{1}{\gamma^3 \left(1 - \frac{(\mathbf{u} \cdot \mathbf{v})}{c^2}\right)^3} \left[\gamma \left(1 - \frac{(\mathbf{u} \cdot \mathbf{v})}{c^2}\right) \mathbf{a} + \frac{\gamma(\mathbf{a} \cdot \mathbf{v})}{c^2} \mathbf{u} + (1 - \gamma)(\mathbf{a} \cdot \hat{\mathbf{v}}) \cdot \hat{\mathbf{v}} \right] \quad (2.44)$$

at position \mathbf{R}' and time t'

E : an event given by the set of spatial and time coordinates.

$\langle S \rangle$: an inertial reference frame.

$\langle S' \rangle$: a second inertial reference frame moving relative to inertial frame $\langle S \rangle$.

v : is the velocity of second reference frame relative to the first reference frame

c : the speed that light travels in empty space when it is not impeded by matter and fields.

x, y, z, t : are the position and time coordinates of event in current inertial reference frame

x', y', z', t' : are the corresponding position and time coordinates the event in another inertial reference frame, referred to as the primed frame, and denoted with a' .

τ : the time of an event in the rest frame of an observer.

$\delta t, \delta t'$: ...

$\delta tau, \delta tau'$: ...

γ : a relativistic factor given by $1/\sqrt{1 - v^2/c^2}$ which determines how much of

an effect the transform of the coordinates deviates from the classical transform. It is only significant at relative frame speeds close to that of light.

Chapter 3

Observers Delayed World View

Light travels extremely fast (roughly 300,000,000 m/s) compared to other everyday speeds we are used to. In the blink of an eye, light could travel around the world twice. We don't need to normally worry about the delay in the light signal, as the time to propagate the distances in our everyday world is very short, so the signal can be approximated as being instantaneous, therefore ignoring the more complex modeling of the delayed signal. But when it comes to particle physics or astronomy, we might be dealing with speeds close to that of light or also huge distances, e.g. light has taken 152,000 years to travel from the Andromeda galaxy, so we are currently seeing this galaxy where it was and how it looked 152,000 years ago. These past positions that we currently see are called the retarded positions/delayed view positions.

3.1 Relativistic Observers Delayed view

When looking at moving objects we see them where they were in their past positions, due to the delay in the light signal, so let's talk about the transform

showing a source particle emitting a photon to an observer, showing the retarded position the observer sees.

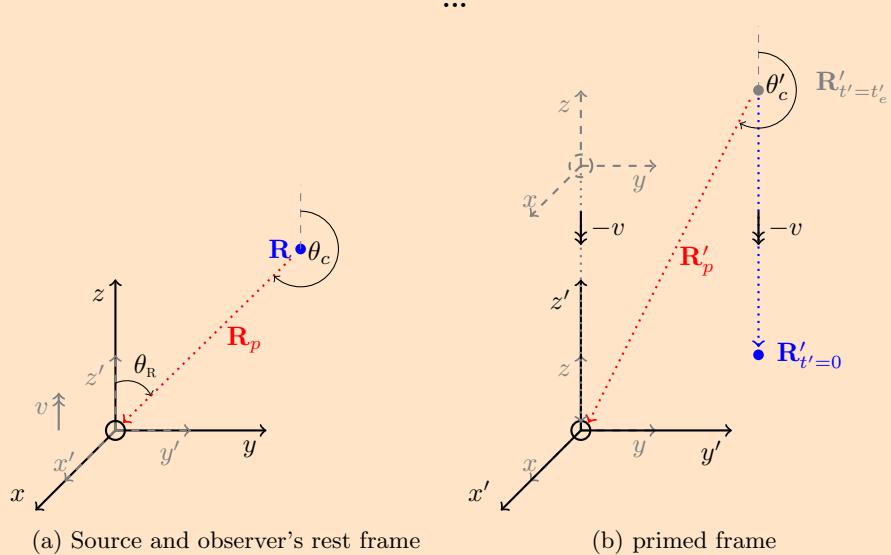


Figure 3.1: Diagram of a source particle (blue) emitting light (red), which is absorbed by an observer positioned at the origin at time $t = t' = 0$ with the observer and source at rest in the first inertial frame (left) and then shown moving in a primed frame (right).

system:

If we take the rest frame of an observer and a particle that emits a pulse of light from its position at \mathbf{R} so that it is received by the observer at the origin, at time $t = 0$.

The time the pulse of light was emitted is $T_e = -\frac{|\mathbf{R}|}{c}$ (where $\frac{|\mathbf{R}|}{c}$ is the time it takes light to propagate along \mathbf{R} to the origin), and the velocity of that light that is propagating towards the origin is

$$\mathbf{c} = c \begin{pmatrix} 0 \\ \sin \theta_c \\ \cos \theta_c \end{pmatrix} \quad (3.1)$$

giving the emitted position as

$$\mathbf{R} = cT_e \begin{pmatrix} 0 \\ \sin \theta_c \\ \cos \theta_c \end{pmatrix} \quad (3.2)$$

the position vector and light's velocity are in opposite directions which can be seen due to the emitted time being negative. If you wish to have this in terms of the position vector angle from the z-axis you can use the equation for the angles $\theta_R = \theta_c - \pi$. Now inputting the delayed positions and times (with origin time $t = 0$) in the first frame into the coordinate transform we get the delayed position in the primed frame as

$$\mathbf{R}' = cT_e \begin{pmatrix} 0 \\ \sin \theta_c \\ \gamma (\cos \theta_c - \frac{v}{c}) \end{pmatrix} = \begin{pmatrix} 0 \\ R_y \\ \gamma (R_z - vT_e) \end{pmatrix} \quad (3.3)$$

at time

$$T'_e = \gamma(1 - \frac{v}{c} \cos \theta_c)T_e$$

we can see that this does agree with the light's velocity transform as well, as it is in the direct opposite direction to \mathbf{R} as it should be. With the light being received at the origin of the primed frame axis at time $t' = 0$, this is the delayed position, if you wanted to find the position of where the particle actually is, we would need to propagate this to the current time of zero of the observer, given by

$$\mathbf{R}'_{t'=0} = \mathbf{R}' - \mathbf{v}T'_e = cT_e \begin{pmatrix} 0 \\ \sin \theta_c \\ \frac{1}{\gamma} \cos \theta_c \end{pmatrix} \quad (3.4)$$

3.2 Delayed View of Transformed Coordinates

How the delayed view of coordinates is taken into account after a transformation is shown in the figure below.

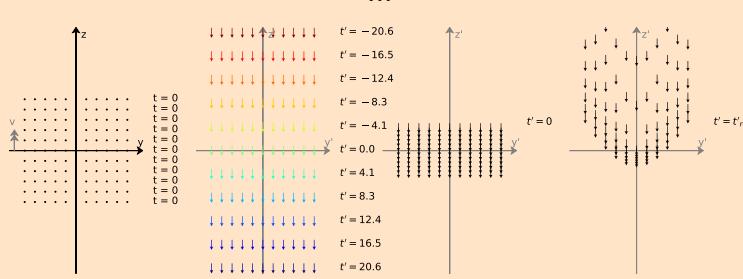


Figure 3.2: A diagram, showing an initial grid of points, representing position coordinates, at time $t = 0$ that are at rest in their proper frame $\langle S \rangle$, with an observer at the origin (left), giving the direct transformation into the moving primed frame $\langle S' \rangle$ (second) where the points are now moving at velocity $-v = -0.9$ which is in the negative z -direction with the colors and labels showing the desynchronization of the primed times $t' = -\gamma \frac{zv}{c^2}$, the next figure (third) shows what these points would be in the same primed frame when they are synchronized with the origin's time $t' = 0$, the last figure (fourth) shows the same primed frame $\langle S' \rangle$ but now with the particles positioned where they would be seen to be at by an observer at the origin with the observers time $t' = 0$, these positions are due to the delay in the light from these particles taking time to get to the observer, also known as the retarded positions of the particles.

*** todo: could include the plot of a larger grid of points to show the distribution of delayed view of particles better, or give the heat map of the density of delayed view coordinates

3.3 Steps for Changing Frames

This section outlines the steps that should be taken to change between two inertial frames:

Firstly let's talk about going from the retarded coordinates of objects in one frame to a second frame. So if in your initial frame, you have an observer at rest at the origin, and different objects are moving around in this frame, if we take the retarded positions \mathbf{R}_{ret} of the objects (i.e. as the observer sees them positioned due to the delay in the light from the object), and use the retarded times of these objects which is $t_{ret} = -\frac{|\mathbf{R}_{ret}|}{c}$ with the coordinate transform,

then these retarded positions in the initial reference frame will be transformed to their retarded positions in the second frame at their retarded times within it.

Another option is that you want to go from the object's positions that have times that are synchronized with the observer to positions that are also synced with the observer in the primed frame. Then we need to start with all objects where they currently are at $t = 0$ then the coordinate transformation transforms the positions to the primed positions in the second frame, but at desynchronized times, so that we will have to propagate these objects forward or back in time until they are all at the time of the observer at the origin which is $t' = 0$ if the velocities of the objects are constant then you will only need to find their primed velocities, but if the initial frames velocities are not constant this is then a bit trickier, as you will have to worry about the transformations of acceleration to find its position at $t' = 0$.

Chapter 4

Spherical light pulses

4.1 Spherical Polar Coordinates

This math section is just for reference in later sections,

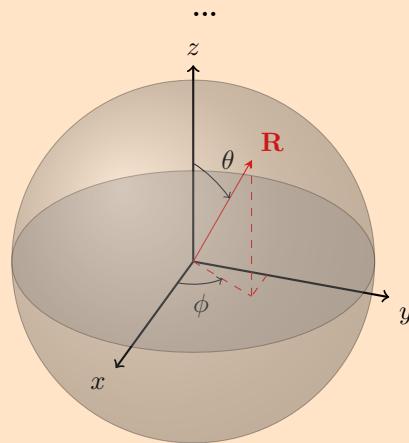


Figure 4.1: Diagram of an axis showing spherical polar coordinates

For a given coordinate \mathbf{R} we have the Cartesian coordinates (x, y, z) , but we can also write this in terms of spherical polar coordinates (r, θ, ϕ) where the angles from the Z-Axis and X-axis respectively are θ (between 0 and π) and ϕ (between 0 and 2π), and r is the distance to the point from the origin, i.e. the

magnitude $|\mathbf{R}|$. So that the coordinate can be written as

$$\boxed{\mathbf{R} = \begin{pmatrix} x \\ y \\ z \end{pmatrix} = r \begin{pmatrix} \cos \phi \sin \theta \\ \sin \phi \sin \theta \\ \cos \theta \end{pmatrix}} \quad (4.1)$$

with $r = \sqrt{x^2 + y^2 + z^2}$, $\cos \phi = \frac{x}{\sqrt{x^2+y^2}}$, and $\cos \theta = \frac{z}{r}$. Some useful formulae for later are the unit vectors

$$\hat{r} = \begin{pmatrix} \sin \theta \cos \phi \\ \sin \theta \sin \phi \\ \cos \theta \end{pmatrix}, \hat{\theta} = \begin{pmatrix} \cos \theta \cos \phi \\ \cos \theta \sin \phi \\ -\sin \theta \end{pmatrix}, \hat{\phi} = \begin{pmatrix} -\sin \phi \\ \cos \phi \\ 0 \end{pmatrix} \quad (4.2)$$

and the differentials with respect to each of the coordinates are

$$\begin{aligned} \frac{\partial \mathbf{R}}{\partial r} &= \hat{r} \\ \frac{\partial \mathbf{R}}{\partial \theta} &= \begin{pmatrix} r \cos \theta \cos \phi \\ r \cos \theta \sin \phi \\ -r \sin \theta \end{pmatrix} = r \hat{\theta} \\ \frac{\partial \mathbf{R}}{\partial \phi} &= \begin{pmatrix} -r \sin \theta \sin \phi \\ r \sin \theta \cos \phi \\ 0 \end{pmatrix} = r \sin \theta \hat{\phi} \end{aligned} \quad (4.3)$$

4.2 Light Speed Velocity Transform

The speed of light remains the same after transforming between frames, but the velocity is not as the direction changes. The velocity transformation in the case where the speed is that of light, $|u| = c$, with the direction in the initial frame at an angle α from the z-axis and φ from the x-axis, so that in spherical polar coordinates the velocity of light is

$$\mathbf{c} = c \begin{pmatrix} \cos \varphi \sin \alpha \\ \sin \varphi \sin \alpha \\ \cos \alpha \end{pmatrix}, \quad (4.4)$$

then from the velocity transform in eq. (2.40) we get

$$\mathbf{c}' = \frac{c}{\gamma(1 - \frac{v}{c} \cos \alpha)} \begin{pmatrix} \sin \alpha \cos \varphi \\ \sin \alpha \sin \varphi \\ \gamma(\cos \alpha - \frac{v}{c}) \end{pmatrix}$$

at position \mathbf{R}' and time t'

(4.5)

This still has the same speed of light which can be shown by taking the magnitude of the velocity vector (i.e. the square root of the squares of all of its elements). However, the angle of propagation has been changed.

4.3 Doppler Effect



Figure 4.2: Diagrams of Light being emitted by a source in two different frames.

In the rest frame of a source s , if it emits a wave at the speed of light in any direction, we have the length the wave has traveled from the source in the time $d\tau$ as being

$$d\lambda = cd\tau \quad (4.6)$$

If we change to a frame where the source is now moving at speed u'_s in the z -direction and we have the wave now emitted at an angle α' from the z -axis, traveling distance cdt' in a time dt' , with the particle moving a distance $u'_s \cos \alpha' dt'$

in the direction of this emitted part of the wave, this leads to a bunching up of the wave in the direction of the movement of the particle, as the same amount of the wave has been emitted in this primed time, but the distance it has traveled from source is

$$d\lambda' = cdt' - u'_s \cos \alpha' dt' = \left(1 - \frac{u'_s}{c} \cos \alpha'\right) cdt' \quad (4.7)$$

we have $dt' = \gamma d\tau$, so the ratio of space the wave takes up in that direction compared to in the rest frame is

$$\text{Ratio} = \frac{d\lambda'}{d\lambda} = \gamma \left(1 - \frac{u'_s}{c} \cos \alpha'\right) \quad (4.8)$$

since wavelength is inversely proportional to the frequency, and this bunching up at the front will proportionally decrease the wavelength

$$\Phi_r = \frac{1}{\gamma \left(1 - \frac{u'_s}{c} \cos \alpha'\right)} \quad (4.9)$$

4.4 Aberration in Special Relativity

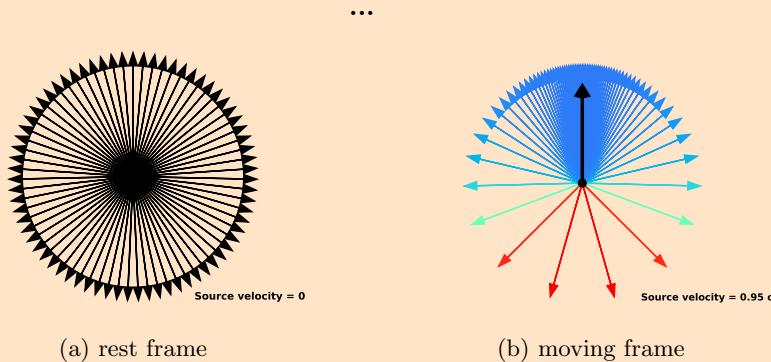


Figure 4.3: Diagram showing the directions of propagation of an evenly distributed light pulse in the rest frame of a source and the corresponding distribution in a frame in which the source is now moving, with the colors showing the Doppler effect on the light's frequency.

The propagation directions of the light as seen in figure 4.3, are transformed such

that the angles between them and the Z-Axis, have the aberrational relationship

$$\cos \alpha' = \frac{\mathbf{c}'_z}{\|\mathbf{c}'\|} = \frac{\cos \alpha + \frac{u'_s}{c}}{1 + \frac{u'_s}{c} \cos \alpha}. \quad (4.10)$$

This is the relativistic aberration formula [2], it shows how the field's propagation direction transforms. Figure 4.4 shows what would happen in the case of incoming light from all directions.

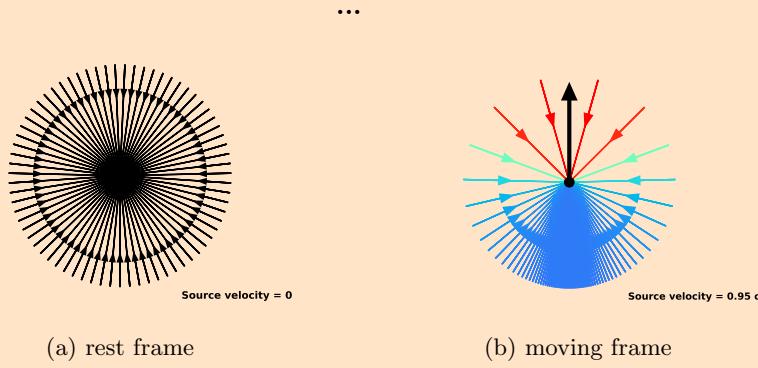


Figure 4.4: Diagram showing the directions of evenly distributed light being absorbed by a particle in its rest frame and the corresponding distribution in a frame in which the particle is now moving, with the colors showing the Doppler effect on the light's frequency.

It can be seen from eq. (4.10) that as a source's speed tends to the speed of light, $\cos \alpha'$ tends to 1 as the top and bottom of the fraction cancel each other out, hence all emitted or absorbed light in the source's rest frame tends to the direction of the source's movement. But we also have to remember that if there was any time between the start and finishing of the emitting the light, this time would tend to infinity as well, basically any emitted light would remain in the vicinity of the source.

4.5 Transform of Light Pulse from Origin

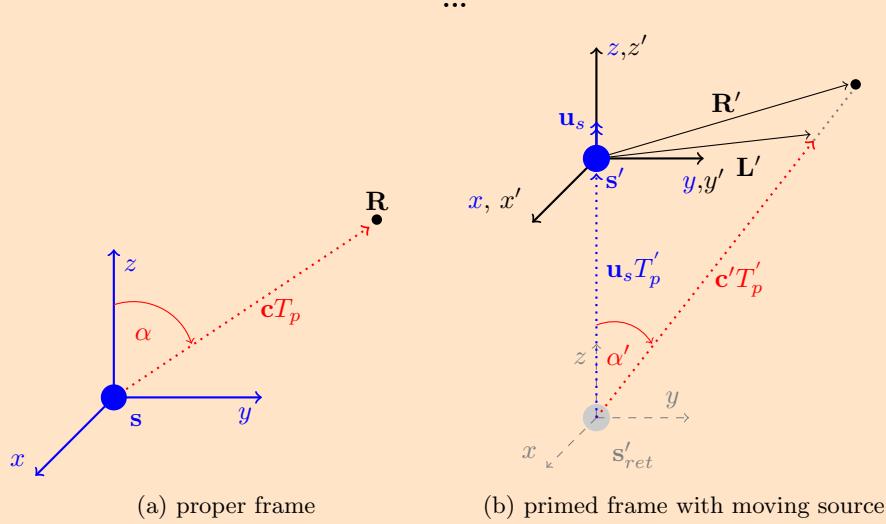


Figure 4.5: The diagram shows a source particle s in blue at rest in the proper frame, Fig (a), with a light pulse that has propagated along cT_p shown in red, at an angle α from the z -axis, to coordinate \mathbf{R} shown in black, with the current time for whole system being $t = 0$. Fig (b), shows the system transformed to a primed frame moving at velocity $\mathbf{v} = (0, 0, v)$ in the z -direction relative to the proper frame, such that the particle and its attached proper axis in this frame are moving at velocity $\mathbf{u}_s = -\mathbf{v}$ and is currently at the primed origin with time $t' = 0$. when the light was emitted from the particle, the particle and attached proper axis were at the retarded position \mathbf{s}'_{ret} shown in grey, the emitted light was then propagated at an angle α' from the z -axis, along $c'T'_p$ to reach \mathbf{L}' at $t' = 0$. it has not yet reached the corresponding Lorentz transformed coordinate \mathbf{R}'

*** in proper frame lights propagation is in same direction as the coordinate from the origin, but in i.e. vec c in same direction as vec r , but vec c' not in same direction as vec r'

As visualized in figure (4.5). If we have a light source particle s , in its rest frame that had emitted a pulse of light in one direction such that it propagates along cT_p at an angle α from the z -axis and is currently at position \mathbf{R} at time $t = 0$, where \mathbf{c} is the velocity of light and T_p is the time it took to propagate to \mathbf{R} . Using the coordinated transform from eq. (2.33) on \mathbf{R} with $t = 0$ we get

$$\mathbf{R}' = \begin{pmatrix} x \\ y \\ \gamma z \end{pmatrix}, \quad (4.11)$$

at time

$$t' = -\gamma \frac{vz}{c^2}$$

the light pulse's velocity using the velocity transform from eq. (4.5) in Cartesian coordinates and particles speed in the primed frame with $v = -u_s$ is

$$\mathbf{c}' = \frac{c}{\gamma \left(1 + \frac{z}{\|\mathbf{R}\|} \frac{u_s}{c} \right)} \begin{pmatrix} \frac{x}{\|\mathbf{R}\|} \\ \frac{y}{\|\mathbf{R}\|} \\ \gamma \left(\frac{z}{\|\mathbf{R}\|} + \frac{u_s}{c} \right) \end{pmatrix} \quad (4.12)$$

The time that the particle emits the light in the proper frame is $T_{ret} = -T_p$ which corresponds to the contracted time

$$T'_{ret} = \gamma T_{ret} = -\gamma T_p \quad (4.13)$$

So that in the primed frame, the retarded position of the source particle emmits the light pulse is

$$\mathbf{s}'_{ret} = -\gamma T_p \mathbf{u}_s \quad (4.14)$$

to find the position the light pulse is currently at, at time $t' = 0$, we have to rewind its time from the position \mathbf{R}' that it is at when using the Lorentz transform at time $t' = -\gamma \frac{vz}{c^2} = \gamma \frac{u_s z}{c^2}$ taking away the displacement it propagated in this time to get the current position of the light to be ($T'_p = -T'_{ret}$).

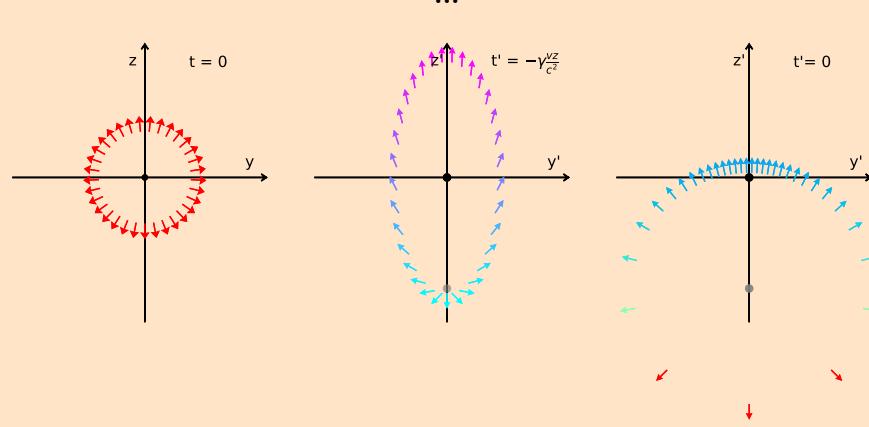
$$\begin{aligned}
\mathbf{L}' &= \mathbf{s}'_{ret} + \mathbf{c}' T'_p \\
&= (\mathbf{c}' - \mathbf{u}_s) T'_p \\
&= \frac{c T'_p}{\gamma \left(1 + \frac{z}{\|\mathbf{R}\|} \frac{u_s}{c}\right)} \begin{pmatrix} \frac{x}{\|\mathbf{R}\|} \\ \frac{y}{\|\mathbf{R}\|} \\ \gamma \left(\frac{z}{\|\mathbf{R}\|} + \frac{u_s}{c}\right) \end{pmatrix} - \begin{pmatrix} 0 \\ 0 \\ u_s T'_p \end{pmatrix} \\
&= \frac{c T'_p}{\gamma \left(1 + \frac{z}{\|\mathbf{R}\|} \frac{u_s}{c}\right)} \begin{pmatrix} \frac{x}{\|\mathbf{R}\|} \\ \frac{y}{\|\mathbf{R}\|} \\ \gamma \left(\frac{z}{\|\mathbf{R}\|} + \frac{u_s}{c}\right) - \gamma \left(1 + \frac{z}{\|\mathbf{R}\|} \frac{u_s}{c}\right) \frac{u_s}{c} \end{pmatrix}
\end{aligned} \tag{4.15}$$

which leads to

$$\mathbf{L}' = \frac{c T'_p}{\gamma \left(1 + \frac{z}{\|\mathbf{R}\|} \frac{u_s}{c}\right)} \begin{pmatrix} \frac{x}{\|\mathbf{R}\|} \\ \frac{y}{\|\mathbf{R}\|} \\ \frac{1}{\gamma} \frac{z}{\|\mathbf{R}\|} \end{pmatrix} \tag{4.16}$$

4.6 Relativistic Beaming

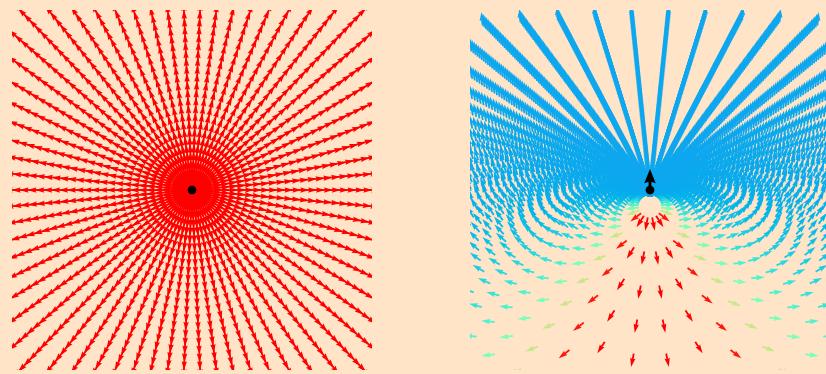
The previous section can be used to show us what happens to the angles of propagation in a spherical pulse of light with evenly distributed angles. Substituting $\frac{x}{\|\mathbf{R}\|} = \cos \varphi \sin \alpha$, $\frac{y}{\|\mathbf{R}\|} = \sin \varphi \sin \alpha$ and $\frac{z}{\|\mathbf{R}\|} = \cos \alpha$ for spherical polar coordinates, which leads to



(a) Source's Rest Frame (b) Source's Moving frame (c) Source's Moving frame

Figure 4.6: Diagram showing, the position of light pulse after a certain amount of time, first in the rest frame and next in the primed frame, the direct coordinate transform gives the ellipsoid with coordinates having different times, and the circular pulse in the primed frame is when you account for this and propagate each part of the pulse, to back or forward in time so that each part has the same time as the origin. *** could have the proper axis different color

If we then did the same for multiple pulses to see how they transform, noting that the time between each pulse is the dilated time in the primed frame.



(a) Source's Rest Frame

(b) Source's Moving frame

Figure 4.7: A 2D diagram showing multiple spherically symmetrical pulses of light from a source at rest (left), and in a corresponding primed inertial frame where the source is moving (right), the colors show the Doppler effect on lights light *** describe it as spherical waves with them being closer together in the direction of movement, could have separate diagram showing the concentric circles with different origins

4.7 Variables...

...

Chapter 5

The Field of a Moving Source

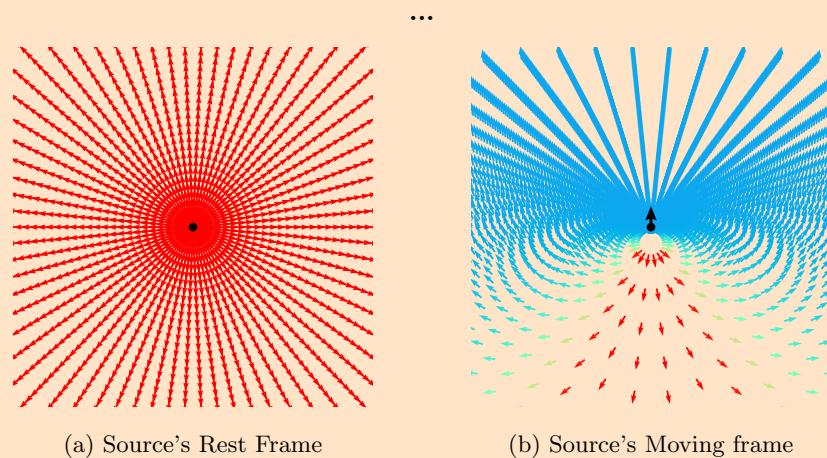


Figure 5.1: A 2D diagram showing multiple spherically symmetrical pulses of light from a source at rest (left), and in a corresponding primed inertial frame where the source is moving (right), the colors show the Doppler effect on lights light *** describe it as spherical waves with them being closer together in the direction of movement, could have separate diagram showing the concentric circles with different origins

*** What is the change in the rate of change of direction at a coordinate of the light $\frac{d\alpha(y',z',t')}{dt'}$ (maybe try $\frac{d\cos \alpha(y',z',t')}{dt'}$ and $\frac{d\sin \alpha(y',z',t')}{dt'}$ instead) and find at $t' = 0$ what is the change in the angle with change in position $\frac{d\alpha(y',z',t')}{dy'}$, $\frac{d\alpha(y',z',t')}{dz'}$ and $\frac{d^2\alpha(y',z',t')}{dy'dz'}$

*** What is the density of the field at a certain point, using $k \frac{1}{R^2}$ in a proper frame, and then taking into account the distance from coordinate to retarded position of the source and then the aberration and Doppler effect

5.1 Flux from Point Source

- flux describes the rate at which a quantity flows across a surface (flowrate per unit area)

5.1.1 Surface element

The surface element spanning from θ to $\theta + d\theta$ and ϕ to $\phi + d\phi$ on a spherical surface at (constant) radius r is then

$$dS = \left\| \frac{\partial \mathbf{R}}{\partial \theta} \times \frac{\partial \mathbf{R}}{\partial \phi} \right\| d\theta d\phi = \left| r\hat{\theta} \times r \sin \theta \hat{\phi} \right| = r^2 \sin \theta d\theta d\phi \quad (5.1)$$

The differential solid angle is

$$d\Omega = \frac{dS}{r^2} = \sin \theta d\theta d\phi \quad (5.2)$$

A solid angle Ω is a measure of the amount of the field of view from some particular point that a given object covers. The fraction of the field of view covered from a point is given by $\Omega/4\pi$

5.1.2 Field Flux

A proper frame of a particle's differential solid angle element (definition given in previous chapter (maybe move that subsection to here))

Differential Element in Spherical Polar Coordinates

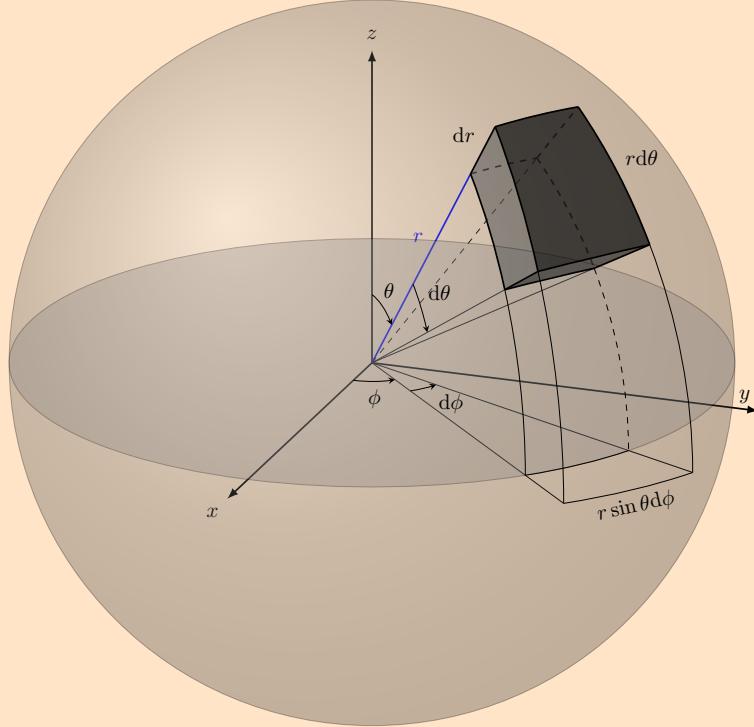


Figure 5.2: ...

$$d\Omega = \sin \theta d\theta d\phi, \quad (5.3)$$

encompasses a certain amount of the wavefront, this is the same amount of the wavefront that is encompassed by the coinciding aberrated differential solid angle

$$d\Omega' = \sin \theta' d\theta' d\phi'. \quad (5.4)$$

We can calculate this element by differentiating both sides of eq. (4.10) with respect to θ [3], which gives

$$\sin \theta' d\theta' = \frac{1 - \frac{u_s^2}{c^2}}{\left(1 + \frac{u_s}{c} \cos \theta\right)^2} \sin \theta d\theta = \frac{1}{\gamma^2 \left(1 + \frac{u_s}{c} \cos \theta\right)^2} \sin \theta d\theta \quad (5.5)$$

as $v = -u_s$ now using this and $d\phi' = d\phi$ (as the angle ϕ is always perpendicular to the motion of the particle and hence unaffected by transformation) we have the solid angle in the primed frame given as

$$d\Omega' = \frac{1}{\gamma^2 \left(1 + \frac{u_s}{c} \cos \theta\right)^2} \sin \theta d\theta d\phi. \quad (5.6)$$

The relative primed wavefront strength at a given angle is taken as being proportional to the amount of the wavefront per solid angle in the primed frame relative to that in the proper frame, referred to here as the aberrational wavefront strength weighting, given as

$$\Phi_\Omega = \frac{d\Omega}{d\Omega'} = \text{Å}^2. \quad (5.7)$$

(define what I mean by strength weighting)

*** This flux doesn't take into account that the source keeps changing position, so at a change ds from position field is coming from a different retarded position of the source, so there's a rotation of direction and a different flux, there's gradients and rotations

*** This flux transform is the change in flux in the emitted direction from the source only not what another particle would feel

Chapter 6

Invariant/Conserved Quantities

Quantities that are conserved (do not change) (frame-independent quantity) when transformed between inertial reference frames are called lorentz invariant. These conserved quantities give us equations to help us model complex systems and experiments

in this chapter, we will see the conservation principle related to space and time coordinates, velocities, and acceleration, but will also look at the more abstract conservation relation of momentum energy and mass. The reason we care about the concept of energy is because it has a related conserved quantity that can be used as a tool to help us make calculations and model more complex systems. We will find a quantity with the units of energy that differs from the previously known energy quantity and its conservation laws, but of course, will approximate to the previous laws at low velocities.

... the derivations to get the famous $E = mc^2$ use approximations, assumptions, and redefined definitions of energy and momentum, and the energy-momentum relation used quantum mechanics to derive, which would require another handbook to go into

... so instead we use the most important thing we need for energy to be and that is to have an associated conservation rule, which means finding a quantity that is unchanged between frames and has a quantity with the units of energy
... to do this we will start from the space-times' associated invariant quantity.

6.1 Space-Time Interval

We defined the space-time interval ΔS between two events in a generic frame $\langle S \rangle$ as:

$$\Delta S^2 = (c\Delta t)^2 - \Delta x^2 - \Delta y^2 - \Delta z^2 \quad (6.1)$$

All inertial frames of reference agree on the interval as it is conserved between frames $\Delta S = \Delta S'$ which can be shown when we sub in the primed version of this equation

$$\begin{aligned} \Delta S'^2 &= (c\Delta t')^2 - \Delta x'^2 - \Delta y'^2 - \Delta z'^2 \\ &= c^2 \gamma \left(\Delta t - \frac{v}{c^2} \Delta z \right)^2 - \Delta x^2 - \Delta y^2 - \Delta (\gamma(\Delta z - v\Delta t))^2 \quad (6.2) \\ &= (c\Delta t)^2 - \Delta x^2 - \Delta y^2 - \Delta z^2 \end{aligned}$$

if a particle is moving in frame $\langle S \rangle$ we can define a space-time interval in its positions in this frame as

$$\Delta S_p^2 = (c\Delta t_p)^2 - \Delta x_p^2 - \Delta y_p^2 - \Delta z_p^2 \quad (6.3)$$

and we have the proper time that passes for the particle as $\Delta\tau$, so taking the limit as the change in time goes to zero we can get the defined abstract space-

time velocity

$$\begin{aligned}
\left(\frac{dS_p}{d\tau}\right)^2 &= \left(c \frac{dt_p}{d\tau}\right)^2 - \left(c \frac{dx_p}{d\tau}\right)^2 - \left(c \frac{dy_p}{d\tau}\right)^2 - \left(c \frac{dz_p}{d\tau}\right)^2 \\
&= c^2 \gamma_{u_p}^2 \left(1 - \frac{u_x^2 + u_y^2 + u_z^2}{c^2}\right) \\
&= c^2
\end{aligned} \tag{6.4}$$

as can be seen, this is a conserved quantity as the speed of light is constant in all frames, taking the square root of both sides and multiplying by c and the particle's mass in its rest frame m_0 we have a conserved quantity with the units of energy, we will call this the rest energy of the particle

$$E_0 = m_0 c \frac{dS_p}{d\tau} = m_0 c^2 \tag{6.5}$$

$$m_0^2 c^2 = -\frac{E^2}{c^2} + (\gamma_u p_x)^2 + (\gamma_u p_y)^2 + (\gamma_u p_z)^2 = -\frac{E^2}{c^2} + (\gamma_u p)^2 \tag{6.6}$$

giving

$$E^2 = m_0^2 c^4 + \gamma^2 p^2 c^2 \tag{6.7}$$

for stationary particles $u = 0$

$$E = m_0 c^2 \tag{6.8}$$

this does not work for massless particles like photons, we have to use a different method of getting this, we know from before that $\mathbf{P}^2 = m_0^2 c^2 = 0$ leading to the components giving

$$\frac{E^2}{c^2} = |p_{light}|^2 \tag{6.9}$$

and finally

$$E = cp_{light} \quad (6.10)$$

.... not needed: definition of momentum for light borrowed from quantum mechanics

$$p = \frac{hf}{c} \quad (6.11)$$

where h is Planks constant and f is the frequency of the light so that the the reason we use this abstract 4-momentum is because it is conserved and the Newtonian momentum is not conserved when it comes to special relativity

The proper time between two events $\Delta\tau$ of an object is also agreed upon in all inertial frames and will also therefore agree on derivatives with respect to the proper time of the object τ (where it is the time between two events in the same location?), because they will be invariant we define the 4-position vector of an object in space from its position and time as, where we multiplied time by c

$$\mathbf{S} = \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} \quad (6.12)$$

taking a derivative with respect to the proper time of the object τ gives us the 4-velocity of an object with $\frac{d}{d\tau} = \frac{dt}{d\tau} \frac{d}{dt} = \gamma_u \frac{d}{dt}$

$$\mathbf{U} = \frac{d\mathbf{S}}{d\tau} = \gamma_u \frac{d}{dt} \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} = \gamma_u \begin{pmatrix} c \\ u_x \\ u_y \\ u_z \end{pmatrix} \quad (6.13)$$

know to show that the magnitude of this is invariant by showing all primed

4-velocities give

$$\|\mathbf{U}\| = c \quad (6.14)$$

... mention proper time and proper position

6.2 Energy-Momentum Relation

The 4 momentum is defined as

$$\mathbf{P} = m_0 \mathbf{U} = \begin{pmatrix} \gamma_u m_0 c \\ \gamma_u m_0 u_x \\ \gamma_u m_0 u_y \\ \gamma_u m_0 u_z \end{pmatrix} \quad (6.15)$$

we have the Taylor series expansion of γ_u factor is

$$\gamma_u = 1 + \frac{1}{2} \frac{u^2}{c^2} + \frac{3}{8} \frac{u^4}{c^4} + \dots \quad (6.16)$$

when subbed into 4-momentum, we have the time component and multiply by c , we have

$$c|\mathbf{P}_t| = m_0 c^2 \gamma_u = m_0 c^2 \left(1 + \frac{1}{2} \frac{u^2}{c^2} + \frac{3}{8} \frac{u^4}{c^4} + \dots\right) \quad (6.17)$$

$$c|\mathbf{P}_t| = m_0 c^2 + \frac{1}{2} m_0 u^2 + \frac{3}{8} \frac{u^4}{c^2} + \dots \quad (6.18)$$

since this

all terms here are in units of energy, the second term is the classical kinetic energy, and the higher-order terms may be seen as higher-order kinetic energy terms that are very small for $u \ll c$, the first term does not rely on the velocity of a mass, but only the mass itself, and this can be taken as a new type of energy all masses have even when at rest, i.e. the energy at rest for a mass is

$$E_{rest} = m_0 c^2 \quad (6.19)$$

we will write this new energy into the 4-momentum

$$\mathbf{P} = \begin{pmatrix} \frac{E}{c} \\ \gamma_u m_0 u_x \\ \gamma_u m_0 u_y \\ \gamma_u m_0 u_z \end{pmatrix} \quad (6.20)$$

the magnitude of the 4-momentum can easily be found to be

$$\|\mathbf{P}\| = m_0 c \quad (6.21)$$

another conserved quantity that does not depend on the frame and using this and the 4-momentum vector

*** this equation can be brought to earlier

*** it is the best equations to explain it all, its all a conserved quantity on both sides

$$m_0^2 c^2 = -\frac{E^2}{c^2} + (\gamma_u p_x)^2 + (\gamma_u p_y)^2 + (\gamma_u p_z)^2 = -\frac{E^2}{c^2} + (\gamma_u p)^2 \quad (6.22)$$

giving

$$E^2 = m_0^2 c^4 + \gamma^2 p^2 c^2 \quad (6.23)$$

for stationary particles $u = 0$

$$E = m_0 c^2 \quad (6.24)$$

this does not work for massless particles like photons, we have to use a different method of getting this, we know from before that $\mathbf{P}^2 = m_0^2 c^2 = 0$ leading to the

components giving

$$\frac{E^2}{c^2} = |p_{light}|^2 \quad (6.25)$$

and finally

$$E = cp_{light} \quad (6.26)$$

.... not needed: definition of momentum for light borrowed from quantum mechanics

$$p = \frac{hf}{c} \quad (6.27)$$

where h is Planks constant and f is the frequency of the light so that the
the reason we use this abstract 4-momentum is because it is conserved and the
Newtonian momentum is not conserved when it comes to special relativity

6.2.1 Energy-Mass Equivalence

*** E=mc² is derived using an approximation for vjjc

Imagine we have a particle at rest emitting light in all directions so that afterward it is still at rest, but in another frame we have this particle moving and the emitted light aberrated, but we still need the particle to remain at a constant velocity after emission. (maybe include a rotation in the impulse/effect of light in each direction on the particle of has)

- if we have a ball in its rest frame that emits a flash of light evenly in all directions, we will have the ball's energy decrease by the amount of energy of the emitted light, and it will remain at rest after due to the light being emitted evenly in all directions - its kinetic energy is zero in this frame

6.2.2 Energy-Momentum Equivalence

*** Dirac derived this and used quantum mechanics to do so, which is out of this handbook's domain, still consequences of this equation to be explored, such as the "zitterbewegung"

Momentum in special relativity is defined slightly differently, to ensure that the law of conservation of momentum holds true in all inertial frames, as required by the first postulate of relativity (it is still approximately unchanged at low velocities). momentum of a particle with rest mass m_0 is given the

$$\mathbf{p}' = m_0 \frac{dR'}{d\tau} = \gamma'_u m_0 \frac{dR'}{dt'} = \gamma m_0 u' \quad (6.28)$$

where u' is the velocity of the particle in this primed frame and $\gamma_{u'} = 1/\sqrt{1 - u'^2/c^2}$.

you will sometimes see the mass of an object referred to as $m = \gamma_{u'} m_0$ which is the apparent mass of a moving particle.

(why not use the 3d time differential equation here?)

6.3 Acceleration

...proper acceleration.

... this is more complex and one difference in this 4-vector here is that the spatial part of the 4-acceleration is not parallel to the 3-acceleration in a primed frame

$$\mathbf{A} = \gamma_u^4 \begin{pmatrix} -c \frac{\mathbf{a} \cdot \mathbf{u}}{c^2} \\ \frac{1}{\gamma_u^2} a_x - u_x \frac{\mathbf{a} \cdot \mathbf{u}}{c^2} \\ \frac{1}{\gamma_u^2} a_y - u_y \frac{\mathbf{a} \cdot \mathbf{u}}{c^2} \\ \frac{1}{\gamma_u^2} a_z - u_z \frac{\mathbf{a} \cdot \mathbf{u}}{c^2} \end{pmatrix} \quad (6.29)$$

this is not really the same thing as what we think of when we think of acceleration, as it does not give the rate of change of velocity in the normal sense in 3d space

$$A = \frac{dV}{d\tau} = \frac{dt}{d\tau} \cdot \frac{dV}{dt} \leftrightarrow \frac{dt}{d\tau} \cdot \frac{d}{dt} \begin{pmatrix} \gamma_u c \\ \gamma_u u \\ 0 \\ 0 \end{pmatrix} = \gamma_u \begin{pmatrix} \frac{u}{c} \gamma_u^3 a \\ \frac{u^2}{c^2} \gamma_u^3 a + \gamma_u a \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{u}{c} \gamma_u^4 a \\ \gamma_u^2 \left(\frac{u^2}{c^2} \gamma_u^2 + 1 \right) a \\ 0 \\ 0 \end{pmatrix} = \gamma_u^4 a \begin{pmatrix} \frac{u}{c} \\ 1 \\ 0 \\ 0 \end{pmatrix} \quad (6.30)$$

6.4 Force

$$\mathbf{F} = \frac{d\mathbf{p}}{d\tau} = m_0 \mathbf{A} \quad (6.31)$$

6.4.1 Force Transform

*** one of the most confusing topics in special relativity, from the way it is defined, how it is derived, and how it relates to other things

*** (a note that the 3-force in special relativity is specifically defined as the time derivative of momentum $\frac{d\mathbf{p}}{dt}$ and not mass times the acceleration $m\mathbf{a}$, this is due to newtons 2nd law holding only for the first definition in special relativity)

*** This means that acceleration turns out not to be in the same direction as the force in special relativity

force derivation is confusing as most sources give the time derivative of \mathbf{u} and \mathbf{v} as acceleration, (are they using that each force has opposite and equal force, but that would not take into account the retardedness)

$$\mathbf{a} = \frac{1}{m_0 \gamma(\mathbf{v})} \left(\mathbf{F} - \frac{(\mathbf{v} \cdot \mathbf{F})\mathbf{v}}{c^2} \right) \quad (6.32)$$

$$\mathbf{F} = \gamma^3 m_0 \mathbf{a}_{||} + \gamma m_0 \mathbf{a}_{\perp} \quad (6.33)$$

6.4.2 not sure what this is

The force parallel to the velocity of the primed frame is unchanged, $\mathbf{F}_{\parallel(\prime)} = \mathbf{F}_{\parallel}$, and the force perpendicular to this is $\mathbf{F}_{\perp(\prime)} = \gamma \mathbf{F}_{\perp}$, with total force, $\mathbf{F} = \mathbf{F}_{\parallel} + \mathbf{F}_{\perp}$ and

$$\mathbf{F}_{\parallel} = \frac{(\mathbf{F} \cdot \mathbf{v})}{|\mathbf{v}|^2} \mathbf{v} \quad (6.34)$$

and substituting this into the formula for total force:

$$\mathbf{F}_{\perp} = \mathbf{F} - \frac{(\mathbf{F} \cdot \mathbf{v})}{|\mathbf{v}|^2} \mathbf{v} \quad (6.35)$$

since $\mathbf{F}_{(\prime)} = \mathbf{F}_{\parallel(\prime)} + \mathbf{F}_{\perp(\prime)}$ we can use equations above to get:

$$\begin{aligned} \mathbf{F}_{(\prime)} &= \frac{(\mathbf{F} \cdot \mathbf{v})}{|\mathbf{v}|^2} \mathbf{v} + \gamma \left(\mathbf{F} - \frac{(\mathbf{F} \cdot \mathbf{v})}{|\mathbf{v}|^2} \mathbf{v} \right) \\ &= \gamma \mathbf{F} + (1 - \gamma) \frac{(\mathbf{F} \cdot \mathbf{v})}{|\mathbf{v}|^2} \mathbf{v} \end{aligned} \quad (6.36)$$

this is the same for the electric field as it only differs by some magnitude from the force.

this is similar to how we get the reverse transform: (is it the primed velocity used in this?)

$$\mathbf{F} = \frac{1}{\gamma} \mathbf{F}_{(\prime)} + (1 - \frac{1}{\gamma}) \frac{(\mathbf{F}_{(\prime)} \cdot \mathbf{v}_{(\prime)})}{|\mathbf{v}_{(\prime)}|^2} \mathbf{v}_{(\prime)} \quad (6.37)$$

6.5 Variables...

...

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